Nonlinear transient heat conduction in a stationary medium

Where k(T) is a function of temperature, In terms of Enthalpy, the equation can be modified to give

Where $H_2 - H_1 = \int_{T_1}^{T_2} \rho c dT$, the boundary conditions $T = T_b$ on Γ_b

$$k(T)\frac{\partial T}{\partial x}I_{x} + k(T)\frac{\partial T}{\partial y}I_{y} + k(T)\frac{\partial T}{\partial z}I_{z} + q + h(T - T_{\infty}) = 0 \quad \text{On } \Gamma_{q}$$

Where I_x , I_y and I_z are direction cosines of outward normal

h = Heat transfer coefficient, T_{∞} = ambient temperature, the initial condition for the problem is $T = T_0$ at t = 0

Application of Galerkin's method for nonlinear heat conduction problems

The solution domain is divided in to finite elements in space. The temperature is approximated within each element by

$$T(x, y, z, t) = \sum_{i=1}^{m} N_i(x, y, z) T(t)$$

Where N_i is the usual shape functions defined piecewise or element by element, T(t) is the nodal temperatures considered to be functions of time and m is the number of nodes in the element considered.

The Galerkin representation for the heat conduction problem is

$$\int N_{i} \left[\frac{\partial}{\partial x} \left(k_{x}(T) \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k_{y}(T) \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left(k_{z}(T) \frac{\partial T}{\partial z} \right) + Q - \rho c \frac{\partial T}{\partial t} \right] dx dy dz = 0$$
....(3)

Use integration by parts on the first three terms of equation (3) simplifies to

Inserting the temperature approximation equation (4) will simplify to

$$-\int \left[k_x(T) \frac{\partial N_j}{\partial x} \frac{\partial N_i}{\partial x} \{T\} + k_y(T) \frac{\partial N_j}{\partial y} \frac{\partial N_i}{\partial y} \{T\} + k_z(T) \frac{\partial N_j}{\partial z} \frac{\partial N_i}{\partial z} \{T\} + k_z(T) \frac{\partial N_j}{\partial z} \frac{\partial N_i}{\partial z} \{T\} \right] dx dy dz$$
$$-\int N_i N_j h \{T\} dT_h + \int N_i Q dx dy dz - \int N_i N_j \rho c \, dx dy dz \frac{\partial \{T\}}{\partial t}$$
$$-\int N_i q d\Gamma_q + \int N_i h T_\infty d\Gamma_h z = 0 \tag{5}$$

Equation (5) can be put in to more convenient forms as

$$\begin{split} M & \frac{dT}{dt} + kT = f \text{, where} \\ M_{ij} &= \int \rho c N_i N_j dx dy dz \\ K_{ij} &= \int \left[k_x(T) \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} + k_y(T) \frac{\partial N_i}{\partial y} \frac{\partial N_j}{\partial y} + k_z(T) \frac{\partial N_i}{\partial z} \frac{\partial N_j}{\partial z} + k_z(T) \frac{\partial N_i}{\partial z} \frac{\partial N_j}{\partial z} \right] d\Omega + \int N_i N_j h d\Gamma_h \\ f_i &= \int N_i Q d\Omega - \int N_i q d\Gamma_q + \int N_i h T_\infty d\Gamma_\infty \end{split}$$

This nonlinear equation set requires an iterative solution. Following the simplest form of iteration method we could start from some initial guess:

$$T = T^{0} = \left(T_{1}^{0}, T_{2}^{0}, T_{3}^{0}, \dots, T_{m}^{0}\right)$$

And obtain an improved solution T by solving the equation

$$M \frac{dT'}{dt} + k \left(T^0\right) T' = f^0$$

The general iteration scheme

$$M \frac{dT^n}{dt} + k \left(T^{n-1}\right) T^n = f^{n-1}$$

is then repeated until convergence. To within a suitable tolerance, is obtained.