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# NMR of Proteins and Small Biomolecules

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# NMR of Proteins and Small Biomolecules

Volume Editor: Guang Zhu

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## Aims and Scope

The series *Topics in Current Chemistry* presents critical reviews of the present and future trends in modern chemical research. The scope includes all areas of chemical science, including the interfaces with related disciplines such as biology, medicine, and materials science.

The objective of each thematic volume is to give the non-specialist reader, whether at the university or in industry, a comprehensive overview of an area where new insights of interest to a larger scientific audience are emerging.



Thus each review within the volume critically surveys one aspect of that topic and places it within the context of the volume as a whole. The most significant developments of the last 5–10 years are presented, using selected examples to illustrate the principles discussed. A description of the laboratory procedures involved is often useful to the reader. The coverage is not exhaustive in data, but rather conceptual, concentrating on the methodological thinking that will allow the non-specialist reader to understand the information presented.

Discussion of possible future research directions in the area is welcome.

Review articles for the individual volumes are invited by the volume editors.

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# Preface

Nuclear Magnetic Resonance (NMR) is one of only a few scientific techniques that have been widely applied in many different areas such as physics, chemistry, biology and medicine. Today, NMR plays a crucial role in structure-function studies of organic and inorganic compounds and of large biomolecules, particularly proteins and DNA. In medicine it is one of the most important imaging techniques available to physicians and it also has extensive applications in pharmaceutical research. In this special issue of NMR of Proteins and Small Biomolecules, review chapters cover a wide-range of topics on some of the latest developments in NMR techniques and their applications.

A review of recent developments in structure-based drug discovery begins this volume by showing how combining 1D and 2D NMR techniques with molecular docking can efficiently screen and identify novel “druggable” leads. The success of such an approach relies on the selection of worthwhile therapeutic targets; presumably proteins that are critical in particular diseases. In that context, studying protein-ligand interactions with NMR techniques that have unique powers for the study of weak protein-ligand interactions, as described in chapter 2, followed by the use of residual dipolar coupling approaches, reviewed in chapter 3, may be essential. Chapter 4 is dedicated to the study paramagnetic metalloproteins, giving an overview of paramagnetic NMR and  $^{13}\text{C}$  directly detected protonless NMR spectroscopy. Since NMR is especially powerful for the study of proteins dynamics, more recent developments in transverse relaxation dispersion experiments, which have extended the range of NMR relaxation studies to the milli-micro second timescale, are illustrated in chapter 5. Solid and liquid state NMR techniques for studying membrane proteins represent another “hot” area of contemporary research. Chapters 6 and 7 systematically describe these approaches. This volume concludes with chapter 8, a review of the considerable sensitivity enhancement attainable in magnetic resonance spectroscopic imaging (MRSI) through the use of dynamic nuclear polarization transfer. This non-invasive technique can be applied to the measurement of metabolites *in vivo* to allow early diagnosis and assessment of diseases in personalized medicine.

I hope that the reviews of these unique topics in NMR techniques and their applications as presented in this volume *NMR of Proteins and Small Biomolecules* are informative and fun to read; and I thank all the authors who made this project possible.

Hong Kong

Guang Zhu

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# Application of NMR and Molecular Docking in Structure-Based Drug Discovery

Jaime L. Stark and Robert Powers

**Abstract** Drug discovery is a complex and costly endeavor, where few drugs that reach the clinical testing phase make it to market. High-throughput screening (HTS) is the primary method used by the pharmaceutical industry to identify initial lead compounds. Unfortunately, HTS has a high failure rate and is not particularly efficient at identifying viable drug leads. These shortcomings have encouraged the development of alternative methods to drive the drug discovery process. Specifically, nuclear magnetic resonance (NMR) spectroscopy and molecular docking are routinely being employed as important components of drug discovery research. Molecular docking provides an extremely rapid way to evaluate likely binders from a large chemical library with minimal cost. NMR ligand-affinity screens can directly detect a protein-ligand interaction, can measure a corresponding dissociation constant, and can reliably identify the ligand binding site and generate a co-structure. Furthermore, NMR ligand affinity screens and molecular docking are perfectly complementary techniques, where the combination of the two has the potential to improve the efficiency and success rate of drug discovery. This review will highlight the use of NMR ligand affinity screens and molecular docking in drug discovery and describe recent examples where the two techniques were combined to identify new and effective therapeutic drugs.

**Keywords** Drug discovery, FAST-NMR, *In silico* screening, Ligand affinity screens, Molecular docking, Nuclear magnetic resonance, Virtual screening

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## 1 Introduction

The completion of the human genome project [1] coupled with an increase in R&D investments was widely anticipated to be the cornerstone of personalized medicine with a corresponding explosion in new pharmaceutical drugs targeting a range of diseases. Nearly a decade later, the rate at which new drugs enter clinical development and reach the market has declined dramatically despite the influx of novel therapeutic targets and R&D investments. In the past 5 years the number of new molecular entities (NMEs) receiving FDA approval has decreased by 50% from the previous 5 years [2]. There are several reasons for this decline, but most stem from the fact that drug discovery is a complex and costly endeavor. Approximately 80–90% of drugs that reach the clinical testing phase fail to make it to market [3, 4]. Efforts to reduce costs often lead pharmaceutical companies to invest their time and money in proven therapies, like “best-in-class” drugs, instead of “first-in-class” drugs that target new mechanisms of action or diseases. As a result, many diseases are “orphaned” and lack any therapeutic compounds in the discovery pipeline. Addressing these issues will require fundamental changes to create a more efficient drug discovery process.

The enormous costs and high failure rates inherent to the pharmaceutical industry are clearly contributing factors to the declining number and diversity of new therapeutics. Efforts that minimize costs without restricting research endeavors will evidently benefit the development of drugs for various human diseases. The availability of hundreds of whole-genome sequences for numerous organisms provides an invaluable data set for drug research [1, 5, 6]. Identifying a novel “druggable” protein target is a critical first step for a successful and efficient drug discovery effort. Unfortunately, bioinformatics analysis alone does not generally provide enough information to justify embarking upon an expensive drug discovery program [7, 8]. Instead, knowing the three dimensional structure of a protein greatly

enhances the value of the bioinformatics analysis. Protein structures often provide insights into the molecular basis of the protein's biological function and its relationship to a particular disease. A protein structure also provides detailed information on the sequence and structural characteristics that govern ligand binding interactions. Building a drug discovery effort based on structural information promises to help in the identification of novel therapeutic targets, in the discovery of new lead compounds, and in the optimization of drug-like properties to improve efficacy and safety. Currently, the drug discovery process within the pharmaceutical industry employs high-throughput screening (HTS) as the primary method for identifying lead compounds. However, the high false positive rate [9–12] combined with a significant cost in time and money has encouraged the development of alternative methods to drive the drug discovery process [13, 14].

Nuclear magnetic resonance (NMR) spectroscopy is uniquely qualified to assist in making the drug discovery process more efficient [15, 16]. NMR is useful for several reasons: (1) it directly detects the interaction between the ligand and protein using a variety of techniques, (2) samples are typically analyzed under native conditions, (3) hundreds of samples can be analyzed per day, and (4) information on the binding site and binding affinity can be readily obtained. These features allow NMR to be an effective tool at multiple steps in the drug discovery pathway, which includes verifying HTS and virtual screening hits [15, 17–19], screening fragment-based libraries [15, 20–22], optimizing lead compounds [15, 17, 23, 24], evaluating ADME-toxicology [25–27], and identifying and validating therapeutic targets [28, 29]. Nevertheless, there are still intrinsic costs to maintaining an NMR instrument, screening a compound library, and producing significant quantities of a protein. One way to significantly reduce experimental costs is to utilize *in silico* methodologies to supplement the lead identification and optimization steps of the drug discovery process [30].

Molecular docking is a computational tool that predicts the binding site location and conformation of a compound when bound to a protein [30–32]. This approach has been found to be fairly successful in redocking compounds into previously solved protein–ligand co-structures [33], where more than 70% of the redocked ligands reside within 2 Å root mean squared deviation (RMSD) of the actual ligand pose. During the prediction of protein–ligand co-structures, molecular docking programs calculate a binding score that allows for the selection of the best ligand pose. The binding score is typically based on a combination of geometric and energetic functions (bond lengths, dihedral angles, van der Waals forces, Lennard-Jones and electrostatic interactions, etc.) in conjunction with empirical functions unique to each specific docking program [34–39]. A large variety of docking programs are available that include AutoDock [40], DOCK [41], FlexX [42], Glide [43], HADDOCK [44], and LUDI [45, 46].

Binding energies are also routinely used to rank different ligands from a compound library after being docked to a protein target. The virtual or *in silico* screening of a library composed of thousands of theoretical compounds can be accomplished in a day with minimal cost [47–49]. Thus, a virtual screen can significantly accelerate the hit identification and optimization process while



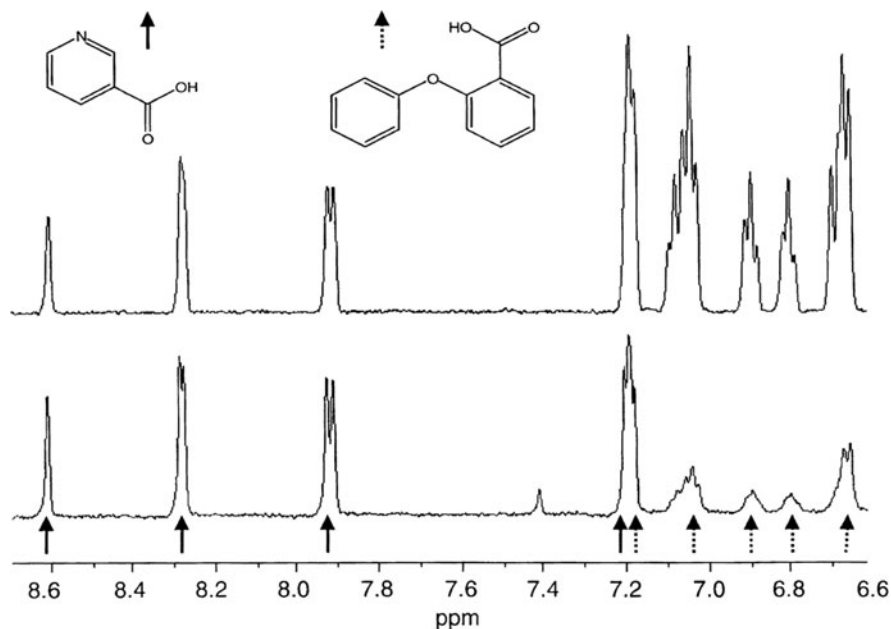
reducing the amount of experimental effort. However, a virtual screen does have significant limitations that prevent it from completely replacing traditional HTS [50–52]. These limitations include inaccurate scoring functions, use of rigid proteins, and simplified solvation models. In essence, a virtual screen only increases the likelihood that a predicted ligand actually binds the protein target, experimental verification is essential. Despite the individual drawbacks, NMR ligand affinity screens and molecular docking are complementary techniques. This review will highlight the use of NMR ligand affinity screens and molecular docking in drug discovery and describe recent examples where the combination of the two techniques provides a powerful approach to identify new and effective therapeutic drugs.

## 2 NMR Ligand Affinity Screens

NMR ligand affinity screening is a versatile technique that is useful for multiple stages of the drug discovery process [15, 17, 22, 53]. This versatility arises from the ability of NMR to directly detect protein–ligand binding based on changes in several NMR parameters. A binding event is detected by the relative differences between the protein or ligand NMR spectrum in the bound and unbound states. However, the specific type of information obtained about the binding process depends on whether a ligand-based or target-based NMR experiment is used.

### 2.1 *Ligand-Based NMR Screens*

Ligand-based NMR screens typically monitor the NMR spectrum of a ligand under free and bound conditions. Distinguishing between a free ligand and a protein–ligand complex is generally based on the large molecular weight difference that affects several NMR parameters. Small molecular weight molecules have slow relaxation rates ( $R_2$ ), negative NOE cross-peaks, and large translational diffusion coefficients ( $D_t$ ). If a protein–ligand binding event occurs, the ligand adopts the properties of the larger molecular-weight protein, increasing  $R_2$ , producing positive NOE cross-peaks, and decreasing  $D_t$ , all of which can be observed by NMR [54]. Most ligand-based NMR screens use one-dimensional (1D)  $^1\text{H}$ -NMR experiments to monitor these changes, which provide significant benefits for a high-throughput screen. 1D NMR experiments are typically fast (2–5 min) and routinely use mixtures without the need to deconvolute [55]. The deconvolution of mixtures is avoided by ensuring that NMR ligand peaks do not overlap in the NMR spectrum (Fig. 1). The application of mixtures allows for hundreds to thousands of compounds to be screened in a single day. Another advantage of ligand-based NMR methods is the minimal amount of protein required ( $<10\ \mu\text{M}$ ) for each experiment. Additionally, isotopically labeled proteins are not needed for the



**Fig. 1** An example of the use of a ligand-detect NMR experiment to observe the line broadening (increase  $R_2$ ) that occurs when one compound, in a mixture of two compounds, binds a protein target. The <sup>1</sup>H-NOESY spectra of nicotinic acid (*left* structure) and 2-phenoxybenzoic acid (*right* structure) in a mixture without protein (*top* spectrum) and with the protein, p38 MAP kinase, added (*bottom* spectrum). The *solid* and *dashed* arrows represent the resonances of nicotinic acid and 2-phenoxybenzoic acid, respectively. In this case, the resonances corresponding to 2-phenoxybenzoic acid are broadened, indicating binding of this compound to the protein. (Reprinted with permission from [178], copyright 2001 by Academic Press)

NMR ligand affinity screen and protein molecular weight is not a limiting factor [21]. In fact, higher molecular-weight proteins enhance the observation of a binding event in a ligand-based NMR screen. All of these characteristics make ligand-based NMR screens a routinely used drug discovery technique.

There are several screening techniques created from ligand-based NMR experiments: line broadening [56], STD NMR [57], WaterLOGSY [58], SLAPSTIC [59], TINS [60], transferred NOEs [61], FAXS [62, 63], FABS [64, 65], and diffusion measurements [66, 67]. Each of these methods utilizes a specific NMR parameter that indicates ligand-binding, such as a change in ligand NMR peak width or diffusion, a saturation transfer from the protein or solvent to the ligand, an NOE transfer between the free and bound ligand, a spin-label induced paramagnetic relaxation, or fluorine chemical shift anisotropy. The choice of which method to use typically depends upon the protein target and the compound library being screened. In addition, line broadening and STD, among other techniques, can be used to measure dissociation constants ( $K_D$ ) [68, 69]. Conversely, ligand-based NMR screens don't provide any structural information about the protein-ligand complex.

## 2.2 Target-Based NMR Screens

A target based screen focuses on changes in the protein (or other target) NMR spectrum to identify a binding event. Typically, chemical shift perturbations (CSPs) occur in the protein NMR spectrum upon ligand binding. The complexity and severe peak overlap in a protein 1D  $^1\text{H}$  NMR spectrum makes it impractical to observe subtle CSPs for weak binding ligands. Instead, two-dimensional (2D) heteronuclear NMR [70–72] experiments are typically used for target-based NMR ligand affinity screens [73]. 2D  $^1\text{H}$ - $^{13}\text{C}/^{15}\text{N}$  HSQC/TROSY NMR experiments require a significant increase in experiment time ( $>10$  min) due to the additional dimension and the need to collect a reference spectrum for the ligand-free protein. Also, the protein needs to be  $^{15}\text{N}$  and/or  $^{13}\text{C}$  isotopically labeled. Importantly, 2D  $^1\text{H}$ - $^{13}\text{C}/^{15}\text{N}$  HSQC/TROSY NMR experiments provide additional information about the ligand binding site.

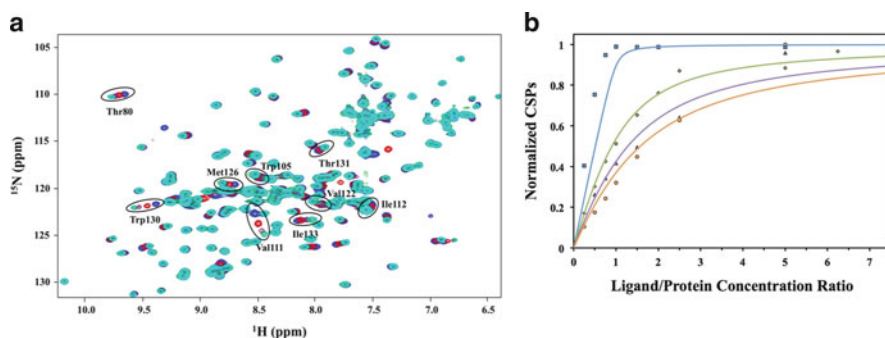
A binding ligand often results in the observation of CSPs of the resonances in a 2D  $^1\text{H}$ - $^{15}\text{N}$ - or  $^1\text{H}$ - $^{13}\text{C}$ -HSQC spectrum (Fig. 2a). These CSPs are usually caused by a change in the chemical environment for residues proximal to the bound ligand or residues undergoing ligand-induced conformational changes. The availability of the protein structure and the NMR sequence assignments (correlation of an NMR resonance with a specific amino acid residue) allows for the CSPs to be mapped onto a three-dimensional (3D) representation of the protein’s surface. A cluster of residues on the protein surface with observed CSPs often identifies the ligand-binding site.

The ligand binding affinity or  $K_D$  is also routinely determined from CSPs measured from a series of 2D  $^1\text{H}$ - $^{13}\text{C}/^{15}\text{N}$  HSQC/TROSY NMR experiments. The magnitude of the CSPs at varying ligand concentrations is correlated to the  $K_D$  for the protein–ligand complex using the following equation [74, 75]:

$$\text{CSP}_{\text{obs}} = \text{CSP}_{\text{max}} \frac{(K_D + [L] + [P]) - \sqrt{(K_D + [L] + [P])^2 - 4[L][P]}}{2[P]}, \quad (1)$$

where  $[P]$  is the protein concentration,  $[L]$  is the ligand concentration,  $\text{CSP}_{\text{max}}$  is the maximum CSP observed for a fully bound protein, and  $\text{CSP}_{\text{obs}}$  is the observed CSP at a particular ligand concentration. A least squares fit of (1) to the experimental CSP data is used to calculate a  $K_D$  (Fig. 2b).

As previously mentioned, since target-based screens require the use of multidimensional NMR experiments, data collection is significantly longer relative to ligand-based NMR screens. Also target-based screens require higher protein concentrations ( $>50$   $\mu\text{M}$  compared to  $<10$   $\mu\text{M}$ ). This severely limits the utility of target-based NMR screens for the high-throughput analysis of large compound libraries. Instead, the approach is typically used to validate hits from a high-throughput screen or the analysis of relatively small fragment-based libraries [76–78]. A fragment-based library consists of low molecular-weight compounds ( $<250$ – $350$  Da) that are fragments of known drugs or have drug-like properties



**Fig. 2** (a) An overlay of the 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC spectra for the protein YndB titrated with increasing amounts of chalcone. The perturbed residues can be used to identify a consensus binding site. (b) NMR titration data for YndB bound to chalcone (blue), flavanone (green), flavone (purple), and flavanol (orange). The magnitude of the chemical shift perturbation can be used to calculate the dissociation constants for each compound. (Reprinted with permission from [112], copyright 2010 by John Wiley and Sons)

[79]. Recent advances like the SOFAST-HMQC experiment [80, 81] and the Fast-HSQC experiment [82] have decreased the time and amount of protein necessary for a target-based screen. Nevertheless, NMR ligand affinity screens are still very resource intensive, requiring a significant amount of time and material. Also, since any high-throughput screen produces a significant amount of negative data (most ligands don't bind or inhibit a protein), a more efficient approach is to screen a library of compounds with a higher probability of binding the protein target. In effect, a virtual or *in silico* screen can be used to enrich a library with likely binders.

### 3 Molecular Docking

An accurate prediction of the interactions between two molecules requires an in-depth understanding of the energetics that led to a stable biomolecular complex. Unfortunately, a model that correctly accounts for all the factors involved in a productive protein–ligand interaction is currently unknown. Further, the problem is exponentially more complex than just modeling the specifics of a protein–ligand interaction. A protein contains thousands of atoms that have specific interactions with each other, with the solvent, and with other ions; in addition to the bound ligand. Because of this complexity, computational efforts that attempt to model protein–ligand interactions require significant amounts of processing power and time. Many efforts that utilize molecular dynamics and distributed computing [83, 84] are generally limited to a detailed analysis of a *single* system. These methods are generally not practical for the majority of researchers interested in conducting a virtual screen of a library containing upwards of millions of compounds. To make molecular docking computationally feasible and easily accessible, many simplifications and trade-offs in the process are necessary.

Many computer programs are available to perform or assist with molecular docking. The vast number of docking programs makes it impractical to describe them all in detail within a single review (for other reviews please see [85–89]). Each docking program does have some unique features that make them particularly useful for a given situation or problem. However, nearly all the docking programs consist of two primary components: docking (or searching) and scoring [30, 31]. Docking refers to the sampling of the ligand’s conformation space and its orientation relative to a receptor. Scoring is used to evaluate and rank the current pose of the ligand.

### 3.1 Docking

The docking process requires, at a minimum, two inputs: the three-dimensional structures of the receptor (protein) and the ligand. The most common simplification to the docking process is to keep the structure of the receptor rigid and stationary. Only the ligand is typically allowed to be flexible as it is docked to the protein. Keeping the protein rigid significantly minimizes the complexity of the calculation. Sampling the conformations and orientations of the ligand is done using systematic or stochastic methods [30, 31].

Systematic search methods attempt to sample all of the possible conformations of a ligand by incrementing the torsional angles of each rotatable bond. Unfortunately, this technique is computationally expensive due to the exponential increase in the number of possible conformations ( $N_{\text{conf}}$ ) as the number of rotatable bonds increases:

$$N_{\text{conf}} = \prod_{i=1}^N \prod_{j=1}^{n_{\text{inc}}} \frac{360}{\theta_{i,j}}, \quad (2)$$

where  $N$  represents the number of rotatable bonds,  $n_{\text{inc}}$  is the number of incremental rotations for each rotatable bond, and  $\theta_{i,j}$  is the size of the incremental rotation for each rotatable bond. As a result, purely brute force systematic approaches are generally not used. Instead, most systematic searches require the use of efficient shortcuts. As an illustration, MOLSDOCK [90] uses mutually orthogonal Latin squares (MOLS) to identify optimal ligand conformations. Latin squares are an  $N \times N$  matrix, where each parameter (torsion angle value) occurs only once in each row and column. Orthogonal Latin squares are two or more superimposed  $N \times N$  matrices, where each parameter still only occurs once in each row and column. MOLS are used to identify the  $N^2$  subset of ligand conformations used to calculate binding energies. Simply, only a small subset of the possible ligand conformations is sampled to construct the potential surface and identify the minima.

Perhaps the most commonly utilized systematic search method is incremental construction, which is used by DOCK [41], FlexX [42], E-Novo [91], LUDI [45, 46], ADAM [92], and TriX [93]. In this particular method, the ligand is

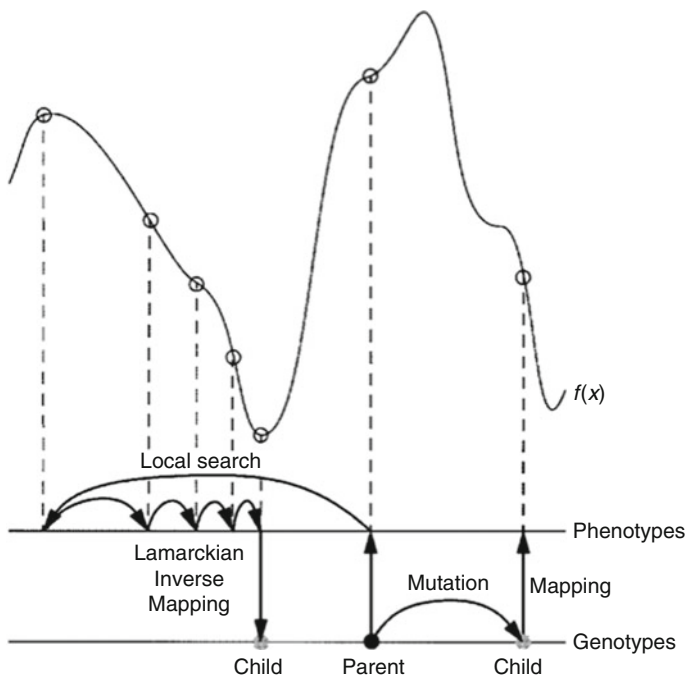
split into fragments. The most rigid fragments are often used as the core or anchor and are docked first into the receptor binding pocket. The remaining fragments are incrementally added back onto the core fragment, where each addition is systematically rotated to evaluate the most optimal conformation. Thus, incremental construction drastically reduces the number of possible conformations that need to be searched in order to identify the optimal pose.

Another systematic approach uses rigid docking in combination with a pre-defined library of ligand conformations, which is implemented in OMEGA [94], FLOG [95], Glide [43], and the TrixX Conformer Generator [96]. This technique generates several low energy conformers for a ligand that are clustered by RMSD. A representative conformer from each cluster is then docked into the receptor. The approach is very fast because the docking process keeps the ligand rigid, eliminating the need to spend computation time on searching torsional space. A tradeoff for this increase in speed is a potential loss in accuracy, since the binding potential for all possible conformers may not be explored. Conversely, a major benefit of the technique is the fact that the library of structural conformers only needs to be generated once. This is a significant savings in time for the pharmaceutical industry, where screening libraries may consist of millions of compounds.

Unlike systematic approaches that attempt to sample all possible ligand conformations, stochastic searches explore conformational space by making random torsional changes to a single ligand or a population of ligands. The structural changes are then evaluated using a probability function. There are three types of stochastic searches: Monte Carlo algorithms [97], genetic algorithms [98], and tabu search algorithms [99]. The most basic stochastic method is the Monte Carlo algorithm, which utilizes a Boltzmann probability function to determine whether to accept a particular ligand pose:

$$P \sim \exp \left[ \frac{-(E_1 - E_0)}{K_B T} \right], \quad (3)$$

where  $P$  is the probability the conformation is accepted,  $E_0$  and  $E_1$  are the ligand's energy before and after the conformational change,  $K_B$  is the Boltzmann constant, and  $T$  is the temperature. The simple scoring function used by the Monte Carlo algorithms is more effective than molecular dynamics in avoiding local minima and finding the global minimum. Alternatively, genetic algorithms utilize the theory of evolution and natural selection to search ligand conformation space. In this case, the conformations, orientations, and coordinates of a ligand are encoded into variables representing a "genetic code." A population of ligands with random genetic codes is allowed to evolve using mutations, crossovers, and migrations. The new population is evaluated using a fitness function that eliminates unfavorable ligand poses. Eventually, a final population converges to ligands with the most favorable "genes" or conformations (Fig. 3). Tabu searches, like other stochastic methods, randomly modify the conformation and coordinates of a ligand, score the conformer, and then repeat the process for a new conformation. Tabu searches



**Fig. 3** An illustration of the genetic algorithm approach, where the states of the ligand (translation, orientation, and conformation relative to the protein) are interpreted as the ligand genotype and the atomic coordinates represent the phenotype. A plot of the change in the fitness function ( $f(x)$ ) as the ligand population is allowed to mutate, crossover, and migrate. The genetic evolution of the ligand effectively samples conformational space where the best conformer is identified by a minimum in the fitness function (Reprinted with permission from [179], copyright 1998 by John Wiley and Sons)

utilize a tabu list to remember previous ligand states. A pose is immediately rejected if it is close to a prior conformation. The tabu list encourages the search to progress to unexplored regions of conformational space.

### 3.2 Scoring

While docking algorithms are generally efficient at generating the correct ligand pose, it is important for the docking program to actually select the correct ligand conformation from an ensemble of similar conformers. In essence, the scoring function should be able to distinguish between the true or optimal binding conformation and all the other poses. The scoring function is also used to rank the relative binding affinities for each compound in the library. Ideally, the scoring function should be able to calculate the free energy ( $\Delta G_{\text{binding}}$ ) of the protein–ligand binding interaction, which is directly related to the  $K_D$ :

$$\Delta G_{\text{binding}} = -RT \ln \frac{1}{K_D}. \quad (4)$$

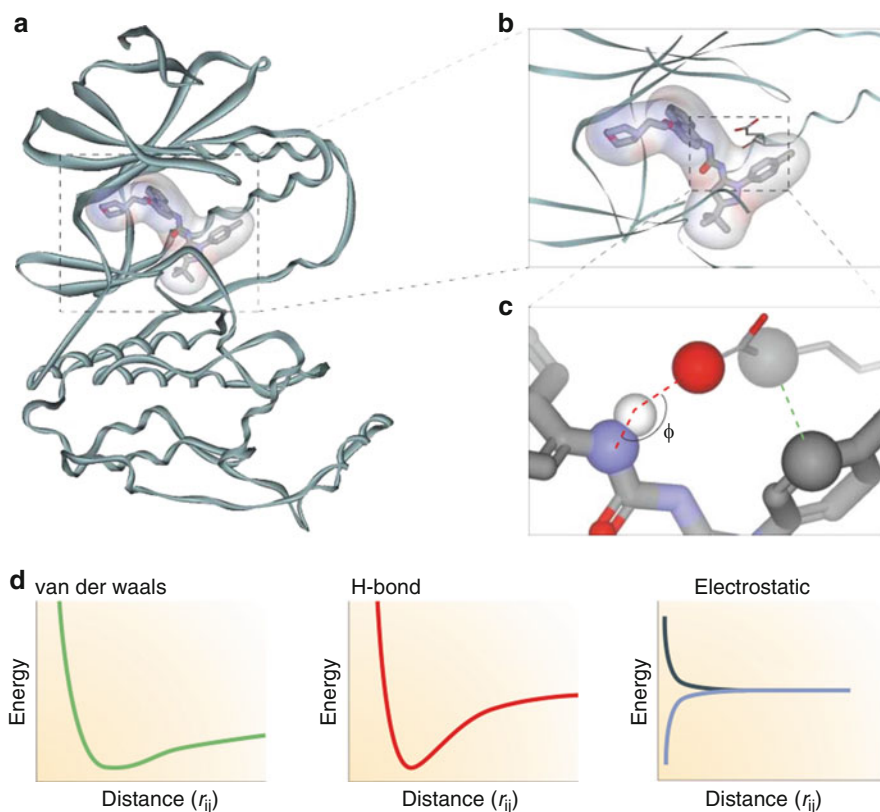
Unfortunately, accurately calculating the binding free energy is very challenging due to the many forces that influence binding. In molecular docking, there are five primary types of scoring functions: force field-based, empirical, knowledge-based, shape-based, and consensus [100–102].

Force field-based scoring functions [30, 31] are used to calculate the free energy of binding by combining the receptor–ligand interaction energy and the change in internal energies of the ligand based on its bound conformation (Fig. 4). The internal energy of the receptor is usually ignored since the receptor is kept rigid in most docking programs. The protein–ligand binding energies are typically defined by van der Waal forces, hydrogen bonding energies, and electrostatic energy terms. The van der Waals and hydrogen bonding terms often utilize a Lennard-Jones potential function, while the electrostatic terms are described by a coulombic function. Unfortunately, these interaction energies were originally derived from measuring enthalpic interactions in the gas phase. Of course, receptor–ligand binding interactions actually occur in an aqueous solution, which introduces additional interactions between the solvent molecules, the receptor, and the ligand. Protein–ligand binding energies are also dependent on the entropic changes that occur upon binding, which include torsional, vibrational, rotational, and translational entropies. Most entropy and solvation-based energy terms can't be calculated using force field-based scoring functions. As a result, force field-based scoring functions are incomplete and inaccurate.

Empirical scoring functions [103–106] are similar to force field-based scoring functions since they use a summation of individual energy terms. But empirical scoring functions also attempt to include solvation and entropic terms. This is typically achieved by using experimentally determined binding energies of known ligand–receptor interactions to train the scoring system using regression analysis. Empirical scoring functions are fast, but the accuracy is completely dependent upon the experimental data set used to train the scoring function. In general, empirical scoring functions are reliable for ligand–receptor complexes that are similar to the training set.

Knowledge-based scoring functions [107–109] are fundamentally different from force field-based and empirical scoring functions. Knowledge-based scoring functions don't attempt to calculate the free energy of binding. Instead, these scoring functions utilize a sum of protein–ligand atom pair interaction potentials to calculate a binding affinity. The atom pair interaction potentials are generated based upon a probability distribution of interatomic distances found in known protein–ligand structures. The probability distributions are then converted into distance-dependent interaction energies. In this manner, knowledge-based scoring functions allow for the modeling of binding interactions that are not well understood. The approach is also very simple, which is useful for screening large compound libraries. Unfortunately, knowledge-based scoring functions are designed





**Fig. 4** (a) A representation of p38 mitogen-activated protein kinase structure bound to BIRB796 and (b) an expanded view of the binding site. (c) A representation of the hydrogen-bonding (*red*) and electrostatic interactions (*green*) between the atoms of the protein and the atoms of the ligand. (d) A representation of three force-field energy terms (van der Waals, hydrogen-bonding, and electrostatic) as distance between the interacting atom pairs change. (Reprinted with permission from [30], copyright 2004 by the Nature Publishing Group)

to reproduce known experimental structures, and the binding score generated has little relevance to an actual binding affinity. This is an issue similar to empirical scoring functions; the accuracy of the scoring function is strongly dependent on the similarity of the protein–ligand complex to the training data set.

As implied, shape-based scoring functions are based on a shape match between the ligand and the ligand binding site [110]. These scoring functions are typically used as prefilters to eliminate compounds that are unable to fit into the ligand binding site [111, 112]. Shape-based scoring functions are very fast, but are limited relative to more accurate scoring functions that calculate binding affinities. Shape-based scoring functions typically generate smooth energy surfaces using Gaussian functions [111], which are more tolerant to atomic variations and make protein

clash interactions “softer.” This essentially helps minimize the effect of small structural variations that may occur during ligand binding.

While the above scoring methods are generally useful in describing protein–ligand interactions, the simplifications used in each approach limits the overall accuracy in predicting the correct docked ligand pose [113, 114]. The major weakness of most docking programs has been shown to be the scoring function. One approach to compensate for this deficiency is to use a consensus score from a combination of scoring functions to rescore a docked pose. Consensus scoring [31, 115] has been shown in several examples to improve docking results compared to a single scoring function. However, like individual scoring functions, the improvement is not consistent and the proper choice of scoring functions to calculate a consensus score is typically based on trial and error.

### 3.3 *Protein Flexibility*

Proteins are inherently flexible and undergo a range of motions over different time scales, and thus the use of rigid protein structures by molecular docking is problematic [116, 117]. This is especially troublesome for therapeutic targets where only an apo-structure is available. Conformational changes upon ligand-binding may range from small perturbations in side chain conformation at the site of ligand binding to large rearrangements of the entire protein structure. Not accounting for such structural changes during ligand docking can drastically alter the ability to identify reliable protein–ligand models correctly [118–122]. Conversely, attempting to dock a large library of flexible ligands to a completely flexible protein structure using molecular dynamics is too computationally expensive to be practical.

Several approaches to “solve” the protein flexibility problem have been explored. The first generally applicable approach utilized soft docking in the scoring function, which reduces the van der Waals repulsion terms in the empirical scoring function [123, 124]. This allows for some overlap between ligand and protein atoms. While this approach is simple and fast, it can only accommodate very small changes in side chain conformations. Other approaches attempt to implement protein structural changes into the docking process. For example, a library of side chain rotamers for residues only in the ligand binding site is routinely used [40, 125]. This dramatically reduces the number of active rotatable bonds during the docking process and has a lower computational cost compared to molecular dynamics. However, the inclusion of a library of rotamers in the docking protocol is significantly slower than rigid protein docking. Furthermore, the approach is limited to local side chain conformational changes.

The most common docking technique that attempts to account for protein flexibility uses multiple protein structures. The ensemble of structures is expected to represent the range of conformations sampled by the protein and has the benefit of being able to evaluate both small and large conformational changes. The molecular docking is repeated for each individual protein conformation, which

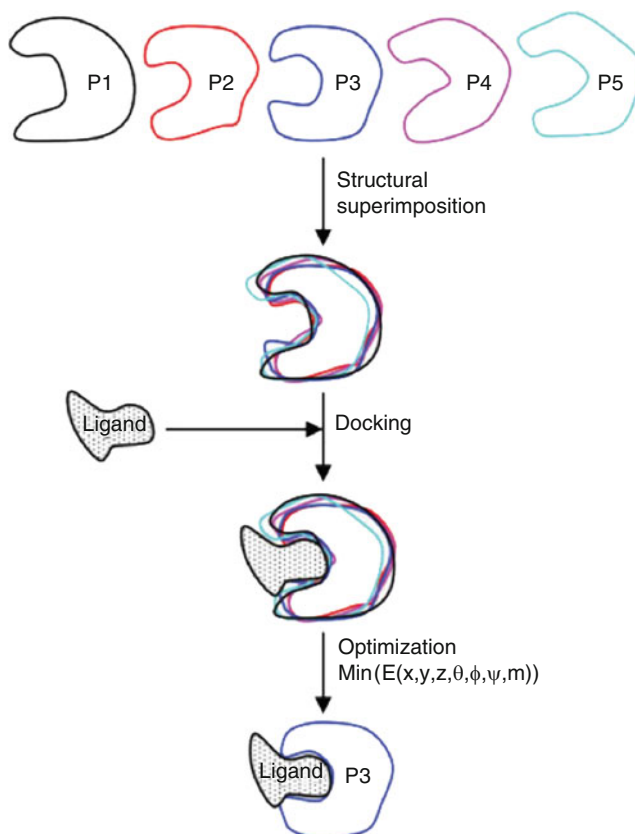
results in a proportional increase in computational time. Also, the results may be ambiguous, since there may be several equally valid ligand poses for each different protein conformation. This is especially apparent in virtual screening approaches where enrichment factors suffer when docking to multiple structures (please see Sect. 3.4). This is likely due to an increase in the number of false positives among the top hits [126]. Ensemble docking is an alternative to docking multiple structures that removes the ambiguity [118]. All the protein structures from the ensemble are superimposed in order to generate an average structure or an average receptor grid. The docking is then performed against the average structure or average receptor grid (Fig. 5). The ensemble docking approach allows for a single docking at a significantly lower computational cost; however, it may suffer from accuracy problems if the ensemble is biased towards the unbound form of the protein. Effectively, a biased ensemble may negate the goal of incorporating protein flexibility if it represents a single conformation.

### 3.4 *Virtual Screening and Assessment*

Using molecular docking to identify lead candidates is an attractive approach for the pharmaceutical industry; it allows for the rapid evaluation of millions of chemical compounds while using minimal resources compared to traditional HTS. The process by which molecular docking is used to rank compounds within a library based on a predicted binding affinity is known as virtual screening [127, 128]. The potential benefit to drug discovery has inspired the development and evaluation of numerous virtual screening methodologies. A virtual screen requires a balance between optimizing speed and maximizing accuracy. Specifically, the goal of a drug discovery virtual screen is the rapid and efficient separation of a small subset of active compounds from a relatively large random library of inactive compounds. Unfortunately, determining the effectiveness of a specific virtual screening process is challenging, where independent evaluators routinely generate inconsistent results [87, 129–131].

The ambiguous nature of the results from a virtual screen requires additional methods to evaluate its success. Typically, a virtual screening process is evaluated against a protein target with a set of known binders. Assessing the performance of a virtual screen is primarily based on the accuracy of the predicted ligand pose and binding affinity. The correct binding pose is often evaluated by calculating the RMSD between the docked and experimental ligand structures. The evaluation of binding affinity is typically based on the accurate ranking of known binders instead of the absolute scores because of the known limitations with calculating a binding energy. Other modes of performance assessment involve evaluating enrichment and generating diverse hit lists.

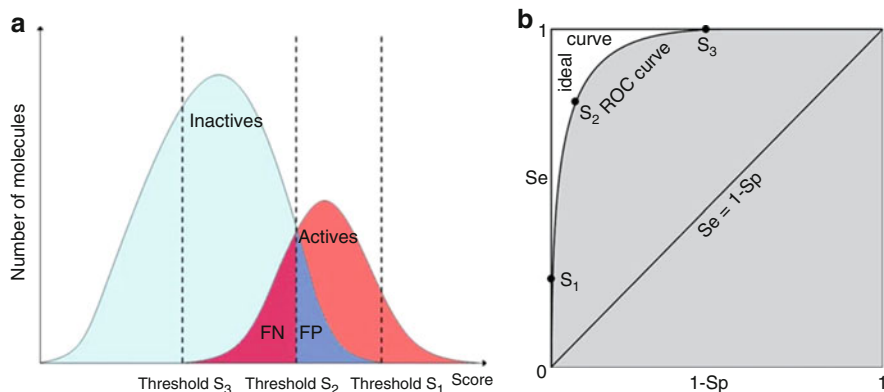
In a virtual screening protocol, every compound in a library ( $N_{\text{tot}}$ ) is docked to the protein and a corresponding binding score is calculated. The binding score for the ligand's best docked pose is used to rank the ligand relative to the entire library.



**Fig. 5** A cartoon illustration of ensemble docking, where five individual protein structures are superimposed to create a single scoring parameter for the docked ligand. Ensemble docking minimizes the computational effort since a single docking occurs to select the best conformer instead of five separate molecular docking simulations. (Reprinted with permission from [118], copyright 2007 by John Wiley and Sons)

A virtual screen never results in all the truly active compounds being top ranked. Instead, most virtual screening protocols set a binding score or ranking threshold to identify the predicted active compounds or “hits.” In general, top ranked compounds are expected to be enriched with active compounds compared to a random selection (Fig. 6a). A high enrichment factor ( $EF > 10$ ) is considered the benchmark of success for a virtual screening [132]. Enrichment is dependent on sensitivity ( $Se$ ) and specificity ( $Sp$ ). Sensitivity represents the true positive rate, which is the ratio of true positives ( $TP$ ) found by the virtual screening vs the total number of actives ( $A$ ) in the library. The number of actives corresponds to both true positive ( $TP$ ) and false negative ( $FN$ ):

$$Se = \frac{TP}{TP + FN}. \quad (5)$$



**Fig. 6** (a) A theoretical distribution of compounds in a virtual screen based upon the docking score. The overlap between active and inactive compounds indicates that the scoring threshold used to identify a hit by virtual screening is critical. (b) A ROC curve is used to evaluate the enrichment of a virtual screen and select a scoring threshold. A ROC curve that approaches  $Se = 1$  and  $1-Sp = 0$  represents perfect enrichment. The area under the ROC curve (AUC) represents the probability that a true active is identified. (Reprinted with permission from [131], copyright 2008 by Springer)

Specificity is the measure of the true negative rate, which represents the ratio of true negatives ( $TN$ ) to the total number of inactive compounds. The number of inactive compounds corresponds to both true negatives ( $TN$ ) and false positives ( $FP$ ):

$$Sp = \frac{TN}{TN + FP}. \quad (6)$$

The enrichment factor is a common method for evaluating the enrichment capabilities of a virtual screen:

$$EF = \frac{\left(\frac{TP}{TP+FP}\right)}{\left(\frac{TP+FN}{N_{tot}}\right)}. \quad (7)$$

The enrichment factor is dependent upon the ratio of active compounds to the total number of compounds in the library. As a result, enrichment scores are difficult to compare between virtual screens with different libraries. Also, the enrichment factor does not distinguish between high and low ranking compounds.

Perhaps the more popular approach for evaluating enrichment is to generate a receiver operating characteristic (ROC) curve [133]. The ROC curve is a plot of the true positive rate ( $Se$ ) against the false positive rate ( $1-Sp$ ) at varying thresholds for determining a hit. A ROC curve allows for the evaluation of a virtual screening method without using an arbitrary scoring threshold. Enrichment occurs when the resulting data point at a particular threshold resides above the diagonal ( $Se = 1-Sp$ ), which corresponds to a random selection of compounds. In a perfect virtual

screen where every active compound is identified as a hit and every inactive compound falls below the threshold, the ROC curve approaches the top left corner ( $Se = 1$  and  $1 - Sp = 0$ ) (Fig. 6b).

Hit list diversity is also an important consideration for the success of a virtual screen since there is more value in identifying a few unique compounds instead of many compounds all based on the same chemical scaffold. One way that diversity can be determined is by comparing the structural similarities of hits from a virtual screen by using the Tanimoto index [134] and then clustering the results. Basically, a Tanimoto index is calculated based on the fraction of similar chemical substructures present in two structures. Generally, 1,365 chemical substructures are used to describe a structure. The substructures include individual elements, two-atom substructures, single rings, condensed rings, aromatic rings, other rings, chains, branches, and functional groups:

$$TI = \frac{C}{A + B + C}, \quad (8)$$

where  $A$  represents the substructural features present in the first structure,  $B$  represents the substructural features present in the second structure, and  $C$  represents the substructural features common to both structures. Identical structures have a  $TI$  score of 1, where completely dissimilar structures have a  $TI$  value of 0.

## 4 Combining Molecular Docking with NMR Ligand Affinity Screens

The vast majority of initial leads in drug discovery are identified from HTS [13, 135, 136]. Pharmaceutical companies have invested heavily in developing and maintaining large chemical libraries (>1,000,000 compounds), which are screened using automated, biological assays intended to monitor a specific response or biological effect [136]. Unfortunately, HTS is extremely inefficient due to the high cost of developing, maintaining, and screening such large libraries of compounds. Furthermore, the random search for an effective drug in the vastness of chemical space ( $\sim 10^{60}$  compounds) [137] is almost guaranteed to fail. Thus, HTS hit rates are typically very low, where <0.5% of compounds exhibit any inhibitor activity in an assay [138]. Correspondingly, HTS assays are highly inefficient since most of the screening effort is spent on the analysis of negative data. Additionally, HTS assays, by nature, are mechanistic “black boxes,” and a response does not provide any information on the mechanism of inhibition. This often leads to numerous false positives from undesirable interactions [11, 12, 139] that may lead the drug discovery project astray. Improving the efficiency of drug discovery requires the implementation of advanced techniques that better guide the selection of lead candidates without sacrificing speed.

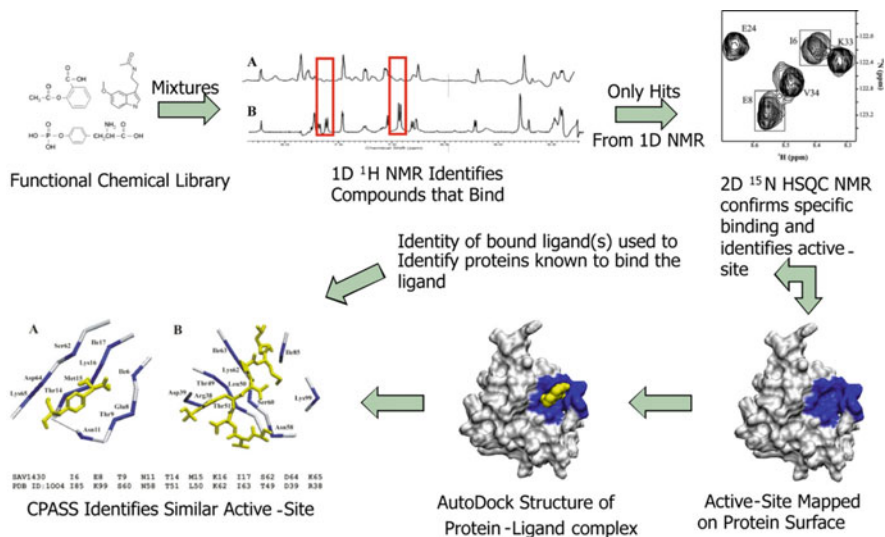
Ideally, an entirely *in silico* approach to screening a large compound library would significantly improve efficiency and reduce costs [140, 141]. However, several assessments of virtual screens have concluded that, without prior in-depth analysis of the protein's ligand binding site, only a marginal improvement in finding successful leads is observed relative to standard HTS [32]. NMR can complement a virtual screen by providing rapid experimental validation of lead compounds. NMR allows for a ligand-binding event to be directly observed instead of relying on false-positive prone activity assays. Also, NMR provides detailed structural information about the ligand binding site and the orientation of the bound ligand. An NMR ligand affinity screen can be used to validate upwards of thousands of predicted hits from a virtual screen [142]. Thus, combining NMR with virtual screens may provide a more efficient approach to lead identification and drug discovery.

#### 4.1 Identification of New Therapeutic Targets

The functional assignment of unannotated proteins is essential to the drug discovery process. Greater than 40% of protein sequences encoded in eukaryotic genomes consist of proteins of unknown function and represent an important opportunity to identify new therapeutic targets [143]. Assigning a function to an uncharacterized protein is an arduous and time-consuming task. The process often requires detailed biochemical studies that may include analyzing cell phenotypes through knockout libraries, monitoring of gene expression levels, or utilizing pull-down assays [144–147].

Since the interactions of proteins with other biomolecules or small molecules is the basis of a functional definition or classification, identifying the functional ligand, the functional epitope or ligand binding site, and the 3D structure of the protein–ligand complex are invaluable for a functional annotation. A functional epitope or ligand binding site is evolutionarily conserved relative to the rest of the protein structure in order for the protein to maintain its biological function. Therefore, proteins that share similar binding site structures are expected to be functional homologs and bind a similar set of ligands [28, 29]. Correspondingly, numerous *in silico* approaches attempt to infer a function for an uncharacterized protein by predicting ligand binding sites using geometry-based, information-based, and energy-based algorithms [148–150]. Unfortunately, unambiguously identifying the ligand binding site on a protein can be challenging without experimental evidence, especially for proteins with no known function.

Functional Annotation Screening Technology using NMR (FAST-NMR) [28, 29] is one approach that combines HTS by NMR with molecular docking and bioinformatics analysis in order to assign a function to a protein (Fig. 7). In this process, a compound library that contains approximately 430 biologically relevant compounds [151] is screened by NMR using a multistep approach [152]. First, a ligand-based screen using 1D NMR<sup>1</sup>H line-broadening experiments identifies



**Fig. 7** A flow diagram of the FAST-NMR process. Mixtures of biologically active compounds are first assayed in a ligand-based 1D line broadening screen against the protein of interest. Compounds that are identified as hits are then verified using CSPs from a 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC experiment that define a binding site on the protein surface. The CSPs are used to guide and filter an AutoDock molecular docking calculation to generate a protein–ligand co-structure. The ligand binding site defined by the co-structure is then compared to other experimental binding sites in the PDB using CPASS. (Reprinted with permission from [28], copyright 2008 by Elsevier)

potential binders. These hits are then verified in a target-based screen using a 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC experiment, where the occurrence of CSPs allows for the identification of the ligand binding site. Molecular docking is used to generate a rapid protein–ligand co-structure [121] that serves as input for the Comparison of Protein Active-Site Structures (CPASS) program [153]. CPASS compares the sequence and structure of this NMR modeled ligand binding site to ~36,000 unique experimental ligand binding sites from the RCSB Protein Databank [143]. Thus, a protein of unknown function can be annotated from a protein with a known function that shares a similar ligand binding site [154]. The FAST-NMR and CPASS approach has been used for the successful annotation of two hypothetical proteins, SAV1430 from *S. aureus* [29] and PA1324 from *P. aeruginosa* [155]. It has also been used to identify a structural and functional similarity between the bacterial type III secretion system and eukaryotic apoptosis [156].

The FAST-NMR approach was recently applied to protein YndB from *Bacillus subtilis* to generate a functional annotation [112]. FAST-NMR was augmented by the inclusion of a virtual screen using the Nature Lipidomics Gateway library that contains ~22,000 lipids. Eight major categories of lipids are represented in the library (fatty acyls, glycerolipids, glycerophospholipids, sphingolipids, sterol lipids, prenol lipids, saccharolipids, and polyketides), which are further divided into a total of 538 distinct subclasses. The initial goal was to identify lipid scaffolds that



preferentially bound YndB to infer the natural ligand. OMEGA [94] was used to generate a database of ~10,000,000 conformers from the lipid library. The program FRED was then used to dock the lipid conformer library to YndB. FRED [111] used rigid docking based on shape complementarity and a consensus scoring system to rank the ligands. The relative enrichment for each lipid class was calculated at different thresholds. Only one lipid category, the polyketides, had a positive relative enrichment, where all of the polyketides identified belonged to the flavonoid class of lipids. Within the flavonoids, three subclasses emerged as favorable hits from the virtual screen, where chalcones/hydroxychalcones, flavanones, and flavones/flavonols accounted for 44.9%, 28.6%, and 14.3% of the top 50 hits, respectively. *trans*-Chalcone, flavanone, flavone, and flavonol were selected to represent each class. The compounds were titrated into YndB to confirm binding and to measure  $K_D$ . The titrations were followed using a series of 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC NMR experiments, where CSPs were measured to calculate  $K_D$ s (Fig. 2). *trans*-Chalcone ( $K_D < 1 \mu\text{M}$ ), flavanone ( $K_D 32 \pm 3 \mu\text{M}$ ), flavone ( $K_D 62 \pm 9 \mu\text{M}$ ), and flavonol ( $K_D 86 \pm 16 \mu\text{M}$ ) were all shown to bind YndB in the same ligand binding site with  $K_D$ s that mimicked the virtual screen ranking. Chalcones and flavonoids have not been identified among the natural products of *Bacillus* organisms, but are important precursors to plant antibiotics. The screening results are consistent with the symbiotic relationship between *B. subtilis* and plants. *B. subtilis* YndB is proposed to be part of a stress-response network that senses chalcone-like molecules during a plant's response to a pathogen infection. The stress-response may induce *B. subtilis* sporulation or the production of antibiotics to assist in combating the plant pathogens.

## 4.2 Rapid Protein–Ligand Structure Determination

A protein–ligand complex is instrumental to a structure-based approach to drug discovery. A new protein–ligand structure is required for each iteration of the lead modification process, until the compound has been evolved into a drug candidate. As a result, rapid protein–ligand structure determination benefits the drug discovery process. There are several methods that utilize NMR CSPs from a protein–ligand binding interaction with molecular docking to generate a corresponding co-structure. Some recent techniques include the McCoy and Wyss method [157], LIGDOCK [158], NMRScore [159], AutoDockFilter [121], QCSP-Steered Docking [160], and HADDOCK [44]. Basically, the CSPs are used to guide the docking process qualitatively and then to steer or filter the docking quantitatively. The docked model is validated by an agreement with the experimental CSPs.

AutoDockFilter (ADF) utilizes a post-filtering approach for rapidly (~35–45 min) generating a co-structure. First, CSPs from the 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC spectrum are mapped onto the protein surface to define the AutoDock 4.0 3D search grid. A 100 docked ligand poses are generated within the CSP defined search grid. Second, the CSPs are used to filter the ligand conformers and select the best pose

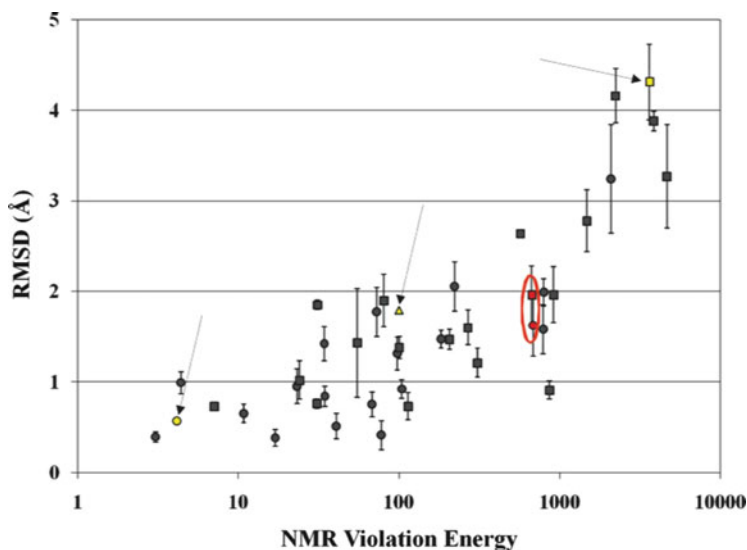
with the AutoDockFilter (ADF) program. ADF calculates a pseudodistance ( $d_{\text{CSP}}$ ) based on the magnitude of the CSPs and compares it to the shortest distance ( $d_s$ ) between any atom in the residue that incurred the CSP with any atom in the docked ligand pose. A violation energy is attributed to each protein residue that is further from the docked ligand pose than predicted by the CSP pseudodistance. The sum of these violation energies generates an overall NMR energy ( $E_{\text{NMR}}$ ) for the docked ligand conformer:

$$E_{\text{NMR}} = k \sum_{i=1}^n (\Delta_{\text{Dist}})^2 \Delta_{\text{Dist}} = \begin{cases} d_{\text{CSP}} - d_s & d_{\text{CSP}} < d_s \\ 0 & d_s \leq d_{\text{CSP}} \end{cases}. \quad (9)$$

The conformer with the lowest NMR energy corresponds to the best protein-ligand co-structure based on a consistency with the experimental CSPs. The NMR energy also provides a qualitative way to evaluate the reliability of the co-structure, with high NMR energies correlating to unreliable co-structures (Fig. 8).

NMRScore [159] is very similar to ADF. NMRScore uses poses generated by AutoDock and seven other docking programs. CSPs are calculated for each pose using DivCon, where a CSP RMSD is determined between the calculated and experimental CSPs. The best pose corresponds to the conformer with the lowest CSP RMSD. The McCoy and Wyss method [157] also uses simulated chemical shift changes. But, unlike the NMRScore approach, the docked ligand is replaced by a number of randomly placed amino-acid probes within the ligand binding site. Proton chemical shifts, primarily from ring-current effects, are calculated for the protein with and without the docked amino-acid probes. The proton chemical shifts are calculated using the SHIFTS program [161], where CSPs are determined based on the difference between the two sets of calculated proton chemical shifts. The best pose for the amino-acid probe is chosen based on a minimal difference between the experimental and calculated proton CSPs. The ligand is then docked to the protein by aligning the ligand with the amino-acid probes.

Instead of simulated chemical shifts, the HADDOCK [44] and LIGDOCK [158] programs use CSPs to define ambiguous interaction restraints (AIRs) [162]. AIRs are an intermolecular distance restraint between all atoms of the residue with the CSP and all atoms of the ligand. Importantly, other experimental information (STDs, mutational data, etc.) can also be used to define AIRs. HADDOCK and LIGDOCK employ a three-tiered approach to refining the protein-ligand complex. First, the ligand is docked to a rigid protein structure. Next, the protein-ligand structure is refined with simulated annealing in torsional space [163]. Finally, the structure is optimized with explicit solvent to remove any remaining structural problems. HADDOCK and LIGDOCK are particularly beneficial since the protein-ligand co-structure is directly refined against the experimental CSPs. The methods do suffer from long computation times and potential difficulties with proper parameterization of the ligand. HADDOCK was initially developed to dock protein-protein interactions and was later modified to accommodate ligands, whereas LIGDOCK was specifically designed to generate protein-ligand co-structures.



**Fig. 8** A comparison of the NMR docking energy from AutoDockFilter to the rmsds between the best docked ligand conformers and the experimental protein–ligand co-structure. An improved correlation is observed for the docking of ligands to the bound form of the protein (*circles*) compared to the apo-protein structure (*squares*). The *red* data points correspond to AutoDockFilter docking results using experimental CSPs for staphylococcal nuclease (PDB-ID: 1EY0, 1SNC) [180–182]. The *yellow* data points correspond to a docking to the apo-structure of acetylcholinesterase (PDB-ID:1ACJ, 1QIF) that resulted in a high rmsd. However, the inclusion of side chain flexibility for residues in the ligand binding site resulted in an improved docking and lower rmsd. (Reprinted with permission from [121], copyright 2008 by the American Chemical Society)

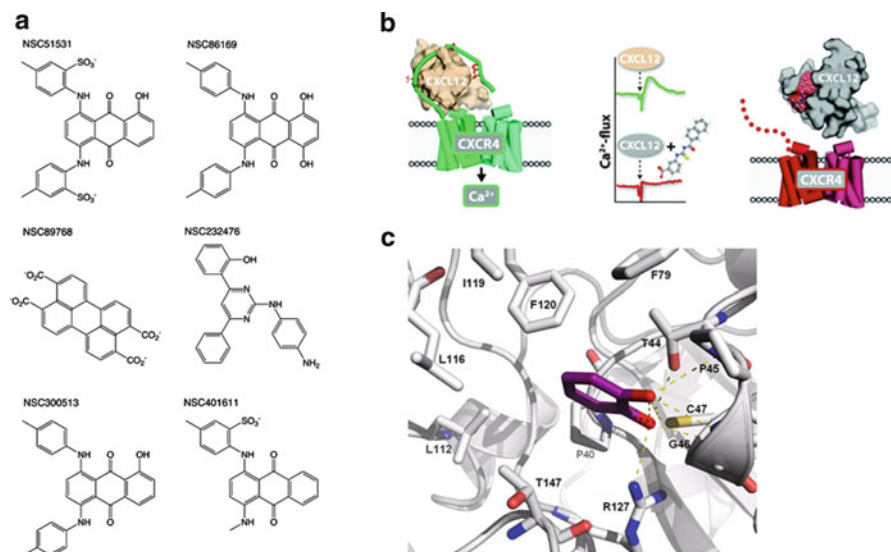
Gonzalez-Ruiz and Gohlke describe a conceptual hybrid (QCSP-Steered Docking) of the AutoDockFilter and the HADDOCK/LIGDOCK procedures, effectively combining the best features of both methods [160]. AutoDock 3.0.5 was modified to incorporate a new hybrid scoring scheme utilizing the DrugScore target function [164] with an amended CSP energy function. Basically, AutoDock is used to generate poses similar to AutoDockFilter, but when an energetically acceptable pose is obtained, CSPs are calculated for the pose. The calculated CSPs are based only on ring current effects [165] from aromatic rings in the ligand. A comparison between the calculated and experimental CSPs is used to calculate an energy violation. Instead of an absolute difference, a Kendall’s rank correlation coefficient is used to account for magnitude differences between the experimental and calculated CSP values. The pose with the lowest DrugScore and CSP energy is chosen. Thus, QCSP-Steered Docking is as fast as AutoDockFilter, but allows for direct refinement against the experimental CSPs like HADDOCK/LIGDOCK.

### 4.3 Lead Identification

Several recent approaches have investigated the combination of NMR and molecular docking for identifying inhibitors for specific proteins. Typically, these approaches apply one of two methodologies: (1) a virtual screen of a large compound library followed by validation of potential binders by NMR or (2) a fragment-based screen using NMR followed by the use of molecular docking to generate a protein–ligand co-structure for optimization.

Virtual screening followed by NMR validation is perhaps the most commonly used combination of these two techniques. Several recent studies have highlighted the use of this approach [166–169]. Branson et al. [166] used a virtual screen with NMR to identify inhibitors of lupindiadenosine 5',5'''-P<sup>1</sup>,P<sup>4</sup>-tetrphosphate (Ap<sub>4</sub>A) hydrolase. These proteins are found in eukaryotes, prokaryotes, and archaea and have been proposed to be involved in several biological functions, ranging from apoptosis, DNA repair, to gene expression. In bacteria, it has also been shown to be involved in pathogenesis, which makes this a potential target for developing antimicrobial agents. There is also a significant difference in sequence between the bacterial and animal forms of the protein, which makes this even more attractive as a drug target. In this study, a virtual screen using DOCK 4 [41] was performed on Ap<sub>4</sub>A hydrolase from *Lupinusangustifolius* with a database of ~120,000 compounds. The docked poses from DOCK were reranked according to consensus scoring using six different scoring functions, where the top 100 ranked ligands were selected and then filtered again to remove all compounds with a logP of 3 or greater in order to select for compounds likely to be water soluble. The result was seven compounds, of which six were commercially available. These six compounds were then subjected to isothermal titration calorimetry to identify any inhibition of hydrolase activity. From that analysis, one compound (NSC51531), which contains a 1,4-diaminoanthracene-9,10-dione core, showed significant binding affinity (~1 μM  $K_D$ ) and was chosen for analysis by 2D <sup>1</sup>H-<sup>15</sup>N HSQC. The NMR analysis showed CSPs consistent with the ATP binding site of the protein. In addition, introducing NSC51531 to the human Ap<sub>4</sub>A hydrolase showed non-specific binding and had no apparent toxic effects against human fibroblasts. This is likely due to structural differences between the binding sites of the lupin and human forms of Ap<sub>4</sub>A hydrolase. Potentially, a scaffold based upon NSC51531 could result in an inhibitor with specificity towards the bacterial form of the protein leading to an effective microbial agent (Fig. 9a).

Veldcamp and coworkers [169] utilized a similar method that targeted the chemokine CXCL12, which activates the CXCR4 receptor shown to be involved with cancer progression. In this approach, nearly 1.5 million compounds from the ZINC database [170] were screened using DOCK 3.5 [171] against the region of CXCL12 that interacts with CXCR4. Specifically, a sulfotyrosine (sY21) was targeted since it was anticipated to be an important residue for the CXCL12-CXCR4 interaction. The top 1,000 hits were manually inspected to identify five compounds with a favorable interaction with sY21. These five compounds were



**Fig. 9** (a) The inhibitors to lupin Ap<sub>4</sub>A hydrolase, where NSC51531, NSC232476, and NSC89768 were identified by the virtual screen and NSC86169, NSC300513, and NSC401611 were structural analogs of NSC51531. (Reprinted with permission from [166], copyright 2009 by the American Chemical Society). (b) A representation of the interaction between the three sulfotyrosine groups of chemokine CXCL12 and the N-terminal region of the G-protein-coupled receptor CXCR4. Virtual screening and NMR identified 3-(naphthalene-2-carbonylthiocarbonylamino)benzoic acid (ZINC 310454) as a possible inhibitor of the binding between CXCL12 and CXCR4, which was verified with a calcium flux assay. (Reprinted with permission from [169], copyright 2010 by the American Chemical Society). (c) The docked pose of fragment F152 (*magenta*) in the active site of human peroxiredoxin 5 with the hydroxyl groups oriented towards catalytic cysteine (C47). (Reprinted with permission from [174], copyright 2010 by PLoS)

then screened using 2D <sup>1</sup>H-<sup>15</sup>N HSQCs, which showed that four of the compounds bound weakly, but specifically, to CXCL12 in the region of interest. The strongest binder, ZINC 310454, had a *K<sub>D</sub>* of ~64 μM. Additional NMR screens with analogs to ZINC 310454 showed the importance of the carboxylic acid and naphthyl group, since analogs lacking these features showed no binding in the 2D <sup>1</sup>H-<sup>15</sup>N HSQC experiments. Furthermore, a calcium flux assay demonstrated that 100 μM ZINC 310454 inhibited CXCL12-mediated signaling. Correspondingly, ZINC 310454 may be a useful scaffold for drug development (Fig. 9b). The results also reinforced the validity of chemokines as a target for drug discovery.

Using molecular docking to screen a large compound library does reduce the time and resources relative to an HTS assay, but it still suffers from an unfocused approach. In general, virtual screens or HTS assays don't efficiently sample chemical space or improve the diversity of hits. Molecular modeling also requires a priori knowledge of the binding site to guide the virtual screen, which may be difficult when dealing with new potential therapeutic targets. One approach to these problems may be to utilize NMR as the primary screening tool and molecular

docking to generate protein–ligand co-structures. Since it is not practical to use NMR to screen the large library of compounds typically utilized by HTS or virtual screening, a more focused approach with a smaller compound library is employed.

Fragment-based screening utilizes a significantly smaller library consisting of simple, low molecular-weight (<250–350 Da) molecules [15, 20–22]. These fragment-like molecules typically have weaker binding affinities (millimolar range) compared to hits found in high-throughput screens (micromolar range), but NMR is sensitive enough to detect these weak protein–ligand interactions. Importantly, fragment-based libraries are more efficient in covering chemical space. Simply, the number of possible compounds decreases drastically as the number of atoms is reduced. Thus, a smaller chemical library actually covers a larger percentage of chemical space. An even greater structural diversity can be achieved by chemically linking multiple fragments. This also results in an additive improvement in binding affinity. Evolving a drug from smaller fragments in this manner has the added benefit of improving ligand efficiency, which typically results in a more bioavailable compound that minimizes non-specific and unfavorable interactions [172, 173].

A recent study [174] by Barelier and colleagues utilized fragment-based screening by NMR and molecular docking in the investigation of the human peroxiredoxin 5 (PRDX5) ligands. Peroxiredoxins are important enzymes that catalyze the reduction of hydroperoxides through a conserved cysteine. However, very few ligands have been identified that bind these proteins despite the availability of crystal structures for PRDX5 bound with benzoate (PDB ID: 1HD2, 1H40) [175]. A compound library of 200 fragment compounds was screened by NMR using STD and WaterLOGSY experiments, where six fragments were identified as binders. STD experiments were also used to calculate the binding affinities for the six fragment molecules, which were in the 1–5 mM range. Since the 1D experiments did not provide information about the location of the binding site, AutoDock 4 [40] was used to dock the fragments to the PRDX5 protein structure. The docking was done against the entire protein structure; a grid search focusing on the benzoate ligand binding site was not used. Not surprisingly, ambiguous results were obtained. The molecular fragments bound to several locations on the PRDX5 structure that were indistinguishable based on binding energies.

Of necessity, the NMR backbone assignments for PRDX5 were obtained to enable the identification of the ligand binding site by monitoring CSPs in 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC experiments. All the fragments were shown to generate a similar set of CSPs consistent with a binding site that included the proposed catalytically active cysteine. The docked binding conformation was also further confirmed from CSPs for derivatives of these fragments. Analysis of the PRDX5 structure with the docked fragments identified the presence of a potentially important hydroxyl functional group that was pointed towards the catalytic cysteine (Fig. 9c). Interestingly, the benzoate compound found in the PRDX5 crystal structure did not show binding by NMR. But, derivatives of benzoate that included a hydroxyl functional group showed improved affinity, further indicating the importance of this hydroxyl group in ligand binding to PRDX5. These results provide further validation of the

value of combining fragment-based NMR screens with molecular docking to generate chemical leads.

While fragment-based screens have been shown to be an effective approach to drug discovery, NMR ligand affinity screens require more time and material than a virtual screen. However, fragment-based screens are extremely helpful for new therapeutic targets with unknown binding sites. Also, the approach has the added benefit of providing information about the druggability of the protein target. There is a correlation between the hit rate of a fragment-based NMR screen and the ability of the protein target to bind drug-like compounds with high affinity [176, 177].

## 5 Concluding Remarks

Significant advances continue to be made in the fields of molecular docking and NMR ligand affinity screens that are benefiting drug discovery. Molecular docking provides an extremely rapid way to evaluate likely binders from a large chemical library with minimal cost. Unfortunately, limitations in the accurate ranking of true binders by molecular docking programs require further experimental validation. Conversely, NMR ligand-affinity screens can directly detect a protein–ligand interaction, can measure a corresponding  $K_D$ , and can reliably identify the ligand binding site. However, NMR-ligand affinity screens are resource intensive and are generally limited to relatively small chemical libraries. Thus, the strengths and weakness of virtual screens and NMR ligand affinity screens are perfectly complementary. Combining the two screening techniques has the potential of significantly improving the efficiency of drug discovery. The combination of NMR and molecular modeling techniques has been shown to enable the rapid determination of reliable protein–ligand co-structures, the identification of new therapeutic targets, and the successful discovery of new drug leads.

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# NMR as a Unique Tool in Assessment and Complex Determination of Weak Protein–Protein Interactions

Olga Vinogradova and Jun Qin

**Abstract** Protein–protein interactions are crucial for a wide variety of biological processes. These interactions range from high affinity ( $K_d < \text{nM}$ ) to very low affinity ( $K_d > \text{mM}$ ). While much is known about the nature of high affinity protein complexes, our knowledge about structural characteristics of weak protein–protein interactions (wPPIs) remains limited: in addition to the technical difficulties associated with their investigation, historically wPPIs used to be considered physiologically irrelevant. However, emerging evidence suggests that wPPIs, either in the form of intact protein complexes or as part of large molecular machineries, are fundamentally important for promoting rapid on/off switches of signal transduction, reversible cell–cell contacts, transient assembly/disassembly of signaling complexes, and enzyme–substrate recognition. Therefore an atomic-level elucidation of wPPIs is vital to understanding a cornucopia of diverse cellular events. Nuclear magnetic resonance (NMR) is famous for its unique abilities to study wPPIs and, by utilization of the new technical developments combined with sparse data based computational analysis, it now allows rapid identification and structural characterization of wPPIs. Here we present our perspective on the NMR methods employed.

**Keywords** Chemical shifts · NMR · NOE · PCS · PRE · RDC

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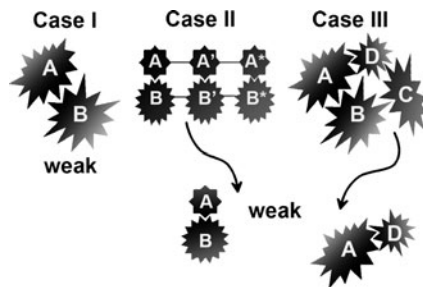


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## 1 Introduction

Living organisms are very complex, highly structured, tightly regulated systems with precisely orchestrated communications at every level of organizational hierarchy. These communications are largely mediated by protein–protein interactions (PPIs). The complete genome sequencing now reveals that there are thousands of potential PPIs that may function as building blocks for these communication networks [1,2]. PPIs can be classified based upon the strength of interaction, which is often rendered by the equilibrium dissociation constant ( $K_d$ ) equal to  $k_{\text{off}}/k_{\text{on}}$ , where  $k_{\text{off}}$  is the rate constant of the complex dissociation reaction and  $k_{\text{on}}$  is the rate constant of the association reaction. The window of biologically relevant  $K_d$  values is extremely wide and can cover 12 orders of magnitude [3]. PPIs can be very loosely divided into three major subclasses [4]: (1) strong, with  $K_d < 10^{-9}$  M, and permanent association, (2) strong and transient, where the change in the quaternary state can be triggered, for example, by ligand binding, and (3) weak, with  $K_d > 10^{-4}$  M, and transient association with  $k_{\text{off}}$  rates of up to  $10^4 \text{ s}^{-1}$ , which results in lifetimes as short as 100  $\mu\text{s}$  [5]. Decades of extensive studies have illuminated structural and functional features for the PPIs from the first two subclasses characterized by strong binding with  $K_d < 10^{-6}$  M, which are summarized in numerous reviews [6–8]. The weak PPIs and their physiological importance (wPPIs, with  $K_d > 10^{-4}$  M), on the other hand, are less well understood. This could be attributed in part to the technical difficulties encountered during attempts to characterize them directly *in vitro* or *in vivo*. The other reason relates to a common prejudice that wPPIs might not be found in living cells, especially considering low ( $< 10^{-7}$  M) protein concentrations estimated by the whole cell volumes. However, it is now being increasingly appreciated that wPPIs are crucial for promoting diverse biologically important processes such as reversible cell–cell contacts, transient assembly and/or disassembly of large signaling complexes, and dynamic regulation of enzymes [9]. Figure 1 provides three possible scenarios of wPPIs: (1) wPPI between two intact proteins, (2) wPPI as part of multi-domain interactions between two intact proteins, and (3) wPPI as part of a multi-protein complex. Conventional methods such as X-ray crystallography, surface Plasmon resonance (SPR), and isothermal titration calorimetry (ITC) often fail to study these wPPIs accurately. In contrast, nuclear magnetic resonance (NMR) has been proven as a particularly powerful tool to examine them at atomic level resolution



**Fig. 1** Three representative cases of wPPIs. Case I: a weak protein–protein interaction found in a locally highly crowded manner. Case II: a weak domain–domain interaction, exemplified by A–B pair, as part of a tight multi-domain complex. Such weak binary domain–domain interaction may be undetectable by many conventional methods including deletion mapping, yeast-hybrid approach, immunoprecipitation, etc., but become apparent when the tertiary structure of the tight complex is challengingly determined. However, NMR may be able to pick this interaction at early stage of the characterization. Case III: a weak protein–protein interaction as a part of multi-protein complex. Similar to (II), a weak A–D pair may not be detectable in isolated manner by any conventional techniques except NMR

and at near physiological conditions [7,8,10,11]. In this chapter we present the various NMR approaches to assess and characterize these three types of wPPIs structurally.

### 1.1 Chemical Shift Perturbation Mapping

The resonance frequencies, also known as chemical shifts, of individual atoms in a particular protein strongly depend on the local environment and, because of that, are often considered as fingerprints of its structure. The chemical shift patterns of  $^{15}\text{N}$  and directly attached amide  $^1\text{H}$  are especially sensitive in this respect. Thus their perturbation, as the result of complex formation, provides a highly sensitive tool for mapping the binding interface. Binding surface on the target protein is identified by titrating the unlabeled target into the solution of the  $^{15}\text{N}$ -labeled protein and monitoring the resultant spectral changes in its  $^1\text{H}$ - $^{15}\text{N}$  HSQC (heteronuclear single quantum correlation) spectrum or, for the larger proteins, its TROSY-based (transverse relaxation-optimized spectroscopy) version [12]. The utility and popularity of this experiment are based upon its straightforward nature and high sensitivity – the spectrum can be recorded in 5–40 min on a typical protein sample (~0.1–1 mM). The spectral changes, also denoted as chemical shift perturbations (CSP), are usually associated with a particular set of amino acids that either (1) directly participate in interaction or are situated very closely to the binding site or (2) reflect binding-induced conformational rearrangements (e.g., a disorder–order transition). The former happens more frequently for wPPIs, which have characteristic small CSP with little or no conformational change, at least within the backbone of well folded domains. Hence the binding interface can

be qualitatively deduced from the spectrally perturbed residues. In addition, for wPPIs with fast exchange, characterized by high  $k_{\text{off}}$  rates,  $K_d$  can be estimated from concentration-dependent titrations [13]. However, the potential problem associated with this analysis relates to the low affinity of the complex: an interfacial residue might not necessarily undergo a significant CSP upon binding, meaning that interfaces derived from CSP data alone are not always complete [6]. One way to avoid this caveat is to increase the ratio of the unlabeled titrant, which can drive the equilibrium towards the bound form with bigger CSP. However, the ratio cannot be too high since it may cause some non-specific effects. The weakest PPI analyzed by CSP demonstrated a  $K_d$  of  $10^{-2}$  M for the flavodoxin/flavodoxin reductase complex [14]. Both (1) and (2) might occur for strong PPIs, which often undergo significant conformational rearrangements upon binding, especially if a disorder–order transition occurs. In such cases, analysis of binding interfaces is more challenging and less straightforward, and one has to rely on additional techniques, such as incorporation of a cross-saturation [15] or inter-molecular Nuclear Overhauser Effects/Enhancements (NOEs) (described in detail in the next section). To conclude, although for wPPIs the CSP method provides fast and robust assessment of the residues forming an intermolecular interface, the mutual orientation of the two partners remains elusive. Thus, if the goal is to generate the structural model of the complex, additional experiments have to be performed and/or novel computational approaches have to be employed.

## 1.2 Nuclear Overhauser Effect

NOE, a relaxation mechanism based upon magnetic dipole–dipole interactions of the nuclei, allows measurement of interproton distances with the basic  $r^{-6}$  distance proportionality. This provides major distance restraints for structural calculations. Supplemented with additional data, such as original dihedral angle restraints obtained from J-couplings or more recent information about the orientation of the bond vectors connecting magnetically active nuclei with respect to the external magnetic field, this approach has been the foundation for NMR-based protein structure determination since its dawn in 1984 [11].

### 1.2.1 NOE in the Determination of the Structure of wPPIs

From the perspective of wPPIs structural characterization, two particular applications of NOE are proven most fruitful.

#### Transferred NOE Experiments

Transferred NOE Experiments (trNOESY) is a quick two-dimensional  $^1\text{H}$  NMR experiment that allows the observation of intramolecular proton contacts ( $<5$  Å)

for the small peptide bound to its target protein [16]. In a nutshell, if the affinity of the interaction in question is low due to the fast dissociation rate, cross relaxation between protons of the peptide in the bound state, which is governed by the large correlation time of the complex, is transferred to the free state through chemical exchange. This phenomenon is manifested by the appearance of additional peaks at the original (corresponding to free state) frequencies in the NOESY spectrum of the peptide when it is mixed with a small molar portion of the large, typically over 50 kDa in molecular weight, target protein. Protein-peptide ratios for trNOESY may vary from 1:10 to 1:200 with mixing times ranging from 50 to 500 ms. Both parameters need to be optimized for each particular case. Substantially increased number of NOEs should be observed for the peptide in the presence of its target protein comparatively to the peptide free form. The method requires no isotope-labeling and is suitable for examination of protein-ligand interactions over a wide range of  $K_d$ s (micromolar-millimolar) [7]. The method can be applied to study initial lead compounds weakly bound to the target protein, which allows the structure determination of the bound compounds for further optimization leading to high affinity binding.

#### Half-Filtered NOESY (Intermolecular NOEs)

Half-filtered NOESY approach was developed to detect inter-molecular contacts in the form of NOEs only between protons pairs in which one of the protons is attached to  $^{15}\text{N}$  or  $^{13}\text{C}$  nuclei while the other is attached to magnetically inactive  $^{14}\text{N}$  or  $^{12}\text{C}$  nuclei. Thus it requires the preparation of  $^{15}\text{N}$ - $^{13}\text{C}$ -labeled protein mixed with its unlabeled binding partner, and/or vice versa. Two types of half-filtered NOESY experiment are commonly used [17]: (1) three dimensional  $^{13}\text{C}$ -separated- $^{15}\text{N}$ ,  $^{13}\text{C}$ -filtered NOESY, which detects NOEs between protons attached to  $^{13}\text{C}$  atoms of the doubly labeled protein and those attached to  $^{12}\text{C}$  and  $^{14}\text{N}$  on the unlabeled protein, and (2) three-dimensional  $^{15}\text{N}$ -separated- $^{15}\text{N}$ ,  $^{13}\text{C}$ -filtered NOESY, which detects NOEs between  $^{15}\text{N}$ -attached and  $^{12}\text{C}$ -,  $^{14}\text{N}$ -attached protons. However, there are pulse sequences available with smart combinations of both, when separation in  $^{13}\text{C}$  and  $^{15}\text{N}$  dimensions can be achieved simultaneously, significantly reducing the experimental time. The sensitivity of this experiment crucially depends on the lifetime of a protein complex. For weak PPIs characterized by high  $k_{\text{off}}$  the actual portion of the complex within the sample might not be high enough for observing the intermolecular NOEs. However, in favorable cases, when the concentrations of both binding partners are high [18], the complexed state could be detectable with the help of high-sensitivity NMR instruments, such as those equipped with cryo-probes. Another relatively more sensitive experiment is the  $^{15}\text{N}$ -edited NOESY on a  $^{15}\text{N}$ /100% deuterated protein bound to the target, which will detect the NOEs between the amide proton of the  $^{15}\text{N}$ -labeled protein and any nearby protons of the target [19]. This experiment can be complementary to those in (1) and (2) but provides a very unambiguous assignment of the amide protons, which are usually well resolved in the HSQC spectrum, to the protons of the

unlabeled target. This experiment can detect very weak NOEs, possibly up to 7 Å in distance, due to the deuteration effect.

### 1.3 Residual Dipolar Couplings

Although NMR-based applications on dipolar coupling have been mainly associated with solid-state NMR or NMR of oriented samples, they have been recently applied to solution NMR for studying macromolecular structure and function in an aqueous solution [20,21]. The dipolar coupling describes a through-space interaction that arises between any two magnetically active nuclei. It depends upon the distance between two atoms, which is constant for the nuclei connected by covalent bonds such as  $^1\text{H}$ - $^{15}\text{N}$  or  $^1\text{H}$ - $^{13}\text{C}$ , and the orientation of the connecting vector with respect to the external magnetic field. In solution, dipolar interactions are averaged because of fast isotropic tumbling. However, if the macromolecules experience obstacles in some directions due to partial alignment with orienting media, for example, bacteriophage, bicelles or polyacrylamide gels, the dipolar couplings are not completely canceled out and whatever is left is designated as residual dipolar couplings (RDC). By measuring RDC the orientation of the molecular alignment tensor could be defined, providing the information about mutual orientation between the domains within single macromolecule or between binary units of the complex. Thus the quest began to find robust ways to orient the media weakly without significant increase in viscosity of the system or generation nucleation points for aggregation. The general idea is based upon a fact that certain media, possessing sufficiently large magnetic susceptibility anisotropy, can be aligned spontaneously by high magnetic field. In earlier 1990s, bicelles, disk-shaped particles made from the lipid/detergent mixtures, were introduced [22] for this purpose, followed by rod-shaped viruses [23], mechanically orientable systems [24,25], and G-tetrad DNA [26]. As compared to the more conventional NOEs approach, RDC carry complementary information: while NOEs provide only local distance restraints, RDC contain long range orientational information (e.g., see review by [27]), thus delivering powerful long-range geometric constraints for proper subunits orientation during the structure determination of the complex. In the case of wPPIs that may undergo fast exchange between the bound and free forms of the binding partners, measured RDC will represent the population weighted average values of those in bound and free form. Theoretically, knowing the molar ratio of bound and free forms (from  $K_d$  and molar concentrations), and after measuring RDC in the free form, one can calculate RDC in bound form [27]. From the RDC of weakly bound subunits, their alignment tensors can be calculated and matched for defining the structure of a weak complex. One example using this strategy to determine the structure of weak complex is  $\alpha$ -methyl mannose bound to mannose-binding protein with a  $K_d \sim 1$  mM [28]. In practice, however, it is not always straightforward for small ligands to determine the accurate fraction of bound form and, thus, full saturation could be beneficial to utilize the RDCs of the fully bound form.

## 1.4 Paramagnetic NMR

As NMR spectroscopists are constantly on the quest to improve the line-shape and to reduce the width of the peaks in their spectra, elimination of the paramagnetic species, often causing significant line-broadening effects, has been considered as paramount in sample preparations. However, the usually undesirable line-broadening effect can provide unique structural long range information [29] when the effect is specific and paramagnetic center is localized to a particular site of the macromolecules. Historically, this understanding has been mainly applied to proteins containing metal-binding sites (reviewed in [30]). The idea to utilize surface exposed cysteines for introduction of paramagnetic tags came later with a cornucopia of chemical agents and procedures developed for reliable conjugation at specific sites (reviewed in [31]). This approach not only provides information about intramolecular distances but can also help in defining alignment of binding subunits within complexes characterized by wPPIs [32], although actual quantification of the distance dependence in such systems is not always straightforward as we discuss below. Two distinguished NMR phenomena, based upon the specific nature of the magnetic moment of an attached paramagnetic center, present the basis for structural investigation. These are paramagnetic relaxation enhancement (PRE) characteristic for all paramagnetic moieties and the pseudocontact shifts (PCS) effect specific for the subclass of paramagnetic ions with an anisotropic electron  $g$ -tensor.

### 1.4.1 The PRE Effect

The large magnetic dipolar interaction of the unpaired electron from a paramagnetic atom with the neighboring NMR-active nucleus results in an increase of the relaxation rate of the above nucleus [33]. Similar to NOE, this effect has basic  $r^{-6}$  distance proportionality, but, because of the larger magnetic moment of the electron, PRE effect is observable at longer, up to 25–35 Å, distances depending upon the nature of the particular paramagnetic group [34,35]. Thus PRE measurements can provide much longer range distance restraints for structural calculations in comparison to the classical NOE approach. The caveat of PRE application for short distances determination comes from the same original source and relates to the fact that nuclei situated very close to a paramagnetic center are often broadened beyond detection. However, the data sets acquired by NOE and PRE approaches are at least complementary. The PRE measurements are based on the correlation of the increased transverse relaxation rate with the distance between the introduced paramagnetic moiety and the affected nucleus [36]. Simply speaking, we are measuring the distance-dependent reduction in peak intensities in a  $^{15}\text{N}$ -HSQC spectrum of the target protein when a single paramagnetic tag is attached to it at the specific site, usually through a thioether bond formed with the side chain of the cysteine residue. Nitroxide stable radicals or metal chelators, such as EDTA- $\text{Mn}^{2+}$ , which are characterized by an unpaired electron with an isotropic  $g$ -tensor, are especially useful since they do not give rise

to pseudo-contact shifts, and Curie-spin relaxation is negligible [37]. The advantage of their employment for pure PRE measurements is the fast and straightforward nature of the method: the resonance assignments in the correlation spectra, known from through-bond scalar triple resonance experiments, are not perturbed by paramagnetic modification and the high sensitivity results in an experimental time of 5–40 min on a typical protein sample (the same advantages we discussed for CSP method, as both are based on the same types of data acquisition – HSQC experiments). The potential problems are associated with intrinsic flexibility/rotation of either the paramagnetic tag itself or its attachment to the protein, resulting in the time average distances sampled over all possible conformations. Thus certain caution is required for incorporation of the derived distances as the restraints for structure calculations, where a paramagnetic center, for example, can be treated as an ensemble average rather than a fixed point [38]. The other possibility for highly flexible systems is to use PRE data loosely, as a guide, rather than major geometric restraints, for example, in structure determination of the complex when the orientation of the peptide, which could be labeled by a paramagnetic tag, in a particular binding site needs to be addressed [39,40]. The potential ability to study transient low population intermediates in macromolecular interactions is conceivably one of the most exciting PRE implementations in structural biology. These illusive species are rarely accessible by other than NMR biophysical techniques. In an exchanging system the observed PRE measured on the resonance of the major species can be modulated by the minor species to the extent depending upon the rate of exchange [41], with the fast exchange allowing one to characterize structurally populations comprising as low as 1%. The example illustrating PRE potential to demonstrate the existence and visualize the distribution of an ensemble of transient non-specific intermediates in addition to specific complex formation has been presented by Clore and colleagues for a bacterial phosphotransferase system [42]. Thus, it has been proven that PRE data is highly sensitive asserting weak interactions characterized by large  $k_{\text{off}}$  rate and is salutary for structural analysis of weak PPIs.

### 1.4.2 The PCS Effect

PCS is a phenomenon that is only observable for paramagnetic systems with anisotropic unpaired electrons such as those found in  $\text{Dy}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Fe}^{3+}$  lanthanide ions, characterized by an anisotropic electron  $g$ -tensor. In general, if the  $g$ -tensor is anisotropic, then the magnetic susceptibility tensor (usually referred to as the  $\chi$ -tensor) is anisotropic as well. The magnitude of the PCS depends on the orientations of the vectors connecting the lanthanide ion and affected nuclei with respect to direction of the external magnetic field. These orientations are not averaged because tumbling in aqueous solution appears to be non-isotropic due to the effect of large  $\chi$ -tensor intrinsic for these paramagnetic species. This large magnetic susceptibility tensor provides enough energy to overcome random Brownian motion and to generate

preferable orientation or alignment of the macromolecule containing the paramagnetic tag. This phenomenon causes changes in chemical shift of the affected nuclei which are sufficiently close to the paramagnetic center (reviewed in [43]). Importantly, the PCS displays basic  $r^{-3}$  distance proportionality in contrast to the  $r^{-6}$  dependence for the PRE or NOE. In theory, this will define a relatively longer experimentally attainable distance range, extending it up to, for example,  $\sim 40$  Å for  $\text{Dy}^{3+}$  in metalloproteins. In practice, the principal axis of the  $\chi$ -tensor is not rigidly fixed within the molecular frame when an extrinsic metal ion is attached to a macromolecule using a chelator with a flexible linker, causing significant reduction in the magnitude of the PCS because of  $\chi$ -tensor principal axes fluctuations within the frame of the macromolecule. From the perspective of studying wPPIs, PCS restraints can be generated using a  $^{15}\text{N}$ -labeled and/or  $^{13}\text{C}$ -labeled protein bound to an unlabeled but paramagnetically tagged partner.  $^{15}\text{N}$  and/or  $^{13}\text{C}$ -HSQC experiments then need to be recorded for both the paramagnetic and diamagnetic states of a sample and chemical shift changes should be extracted from the spectra [31]. However, to use PCS data, one first has to define the tensor describing the anisotropic magnetic moment of the paramagnetic center [44]. When the structures of the individual proteins are known, PCS data can be combined with rigid-body docking to produce a model for a protein complex. This approach has been proven successful in determination of a 30-kDa complex between the  $\theta$  and  $\varepsilon$  subunits of *Escherichia coli* polymerase III [45], where the active-site bound  $\text{Mg}^{2+}/\text{Mn}^{2+}$  pairs were exchanged with paramagnetic  $\text{Dy}^{3+}$  or  $\text{Er}^{3+}$  and corresponding  $^{15}\text{N}$ -HSQC spectra of the diamagnetic apo-complex and paramagnetic-ion-bound complexes were compared. An analogous approach taking advantage of the intrinsic iron-binding capability of cytochrome f has been used earlier to define the structure of its transient complex with plastocyanin: conveniently, iron in its oxidized  $\text{Fe}^{3+}$  form is paramagnetic while in the  $\text{Fe}^{2+}$  form it displays a diamagnetic nature [46].

## 2 Conclusions

While tight protein interactions can be addressed experimentally by many techniques, including X-ray crystallography, the vast majority of these approaches fail or become unreliable when the interactions are weak. Solution NMR spectroscopy is unique among the structural techniques, permitting the characterizing of weak interactions and providing structures of weak protein-target complexes. If such interactions involve small molecules, NMR can be employed for optimization and development of drug-leads. In the current post-genomic era, the NMR methods we have highlighted in combination with functional and computational approaches hold significant promise for characterizing the plethora of weak protein complexes that regulate cellular events, thereby providing an unbiased and comprehensive view of how proteins function in living cells.



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# The Use of Residual Dipolar Coupling in Studying Proteins by NMR

Kang Chen and Nico Tjandra

**Abstract** The development of residual dipolar coupling (RDC) in protein NMR spectroscopy, over a decade ago, has become a useful and almost routine tool for accurate protein solution structure determination. RDCs provide orientation information of magnetic dipole–dipole interaction vectors within a common reference frame. Its measurement requires a nonisotropic orientation, through a direct or indirect magnetic field alignment, of the protein in solution. There has been recent progress in the developments of alignment methods to allow the measurement of RDC and of methods to analyze the resulting data. In this chapter we briefly go through the mathematical expressions for the RDC and common descriptions of the alignment tensor, which may be represented using either Saupe order or the principal order matrix. Then we review the latest developments in alignment media. In particular we looked at the lipid-compatible media that allow the measurement of RDCs for membrane proteins. Other methods including conservative surface residue mutation have been invented to obtain up to five orthogonal alignment tensors that provide a potential for de novo structure and dynamics study using RDCs exclusively. We then discuss approximations assumed in RDC interpretations and different views on dynamics uncovered from the RDC method. In addition to routine usage of RDCs in refining a single structure, novel applications such as ensemble refinement against RDCs have been implemented to represent protein structure and dynamics in solution. The RDC application also extends to the study of protein–substrate interaction as well as to solving quaternary structure of oligomer in equilibrium with a monomer, opening an avenue for RDCs in high-order protein structure determination.

**Keywords** RDC · Alignment medium · Ensemble · Dynamics · Oligomer

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## 1 Introduction

Solution NMR spectroscopy is a powerful technique to study protein structure and dynamics at the atomic level. The method relies on a variety of experimental restraints to determine protein structure. These include the nuclear Overhauser effects (NOEs) that provide interproton distances, the  $J$ -coupling constants that depend on dihedral angles [1–3], the paramagnetic relaxation enhancement (PRE) that is distance dependent with respect to the paramagnetic center [4], and the residual dipolar couplings (RDC) that report on internuclei vector orientations. NOE typically measures interproton distances of less than 5 Å and  $J$ -coupling probes spin nuclear interactions within a few bonds away and they are therefore local in nature. In contrast, PRE can measure distances up to 20–30 Å from the paramagnetic center. In this respect RDC is unique. It can provide relative orientations among internuclei vectors irrespective of their distance separations. This unique property of RDC has opened up new possibilities in using NMR to study phenomena that were previously unattainable.

In the presence of a magnetic field, RDCs arise when the proteins in solution weakly align relative to the field, thus creating an anisotropic condition. The direction of the alignment of the protein molecules in the magnetic field is commonly referred to as the alignment tensor frame. Under such anisotropic condition, with the presence of an external field, a magnetic dipole–dipole interaction does not average to zero and yields a measurable dipolar coupling. The magnitude of the dipolar coupling depends on the angle between the internuclei vector and the external magnetic field as well as the internuclei distance. If the dipolar interaction is between two covalently bonded nuclei, then the internuclear distance is fixed and only the orientation dependence remains. A typical measurement may report hundreds of

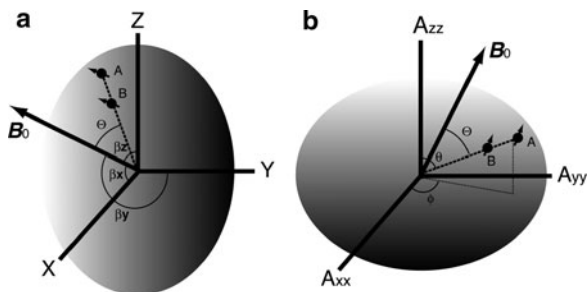
RDCs within a protein that correspond to bond directions within the alignment tensor frame, providing orientation restraints for protein structure determination.

There are at least two approaches to create weak alignment conditions for measuring RDCs. One is to take advantage of the large magnetic susceptibility of a protein where its interaction with the magnetic field could produce a weak alignment [5, 6]. The other is to mix the protein sample with a medium that can be mechanically manipulated to create an anisotropic matrix or one containing supramolecules with substantially large susceptibility anisotropy that can be aligned under an external magnetic field. The interaction between the proteins and the media in turn could induce weak alignment of the proteins [7]. The latter approach creates a degree of alignment that is roughly one order of magnitude stronger than the former one, significantly larger than the experimental error, and thus is more practical for general applications.

A common usage of RDCs is to include them with other NMR restraints in refining a protein structure. For studies of large or membrane associated proteins, where high level of deuteration is required to achieve narrower line-widths, the number of observed NOEs will be reduced greatly. Therefore RDC restraints are necessary [8]. In this chapter we will focus on a short description on how RDC was developed, its practical mathematical expressions, and novel methods used in creating different alignment conditions that would allow more proteins to be studied using RDCs. We will describe RDC data interpretations and some common approximations. Finally we will discuss the most recent RDC applications in ensemble structure refinement, protein–ligand, and protein high-order structure determinations.

## 2 Theoretical Expressions

Dipolar coupling measures the interaction between two magnetic nuclei in an external field  $\mathbf{B}_0$ . If the vector connecting nuclei  $A$  and  $B$  is parallel to the field  $\mathbf{B}_0$ , the coupling is at its strongest with a magnitude  $D_{\max}$ , which is given by Eq. (1), where  $\mu_0$  is the vacuum permeability,  $h$  is Planck's constant,  $\gamma_A$  and  $\gamma_B$  are the gyromagnetic ratios of nuclei  $A$  and  $B$ , respectively, and  $r_{AB}$  is the distance between nuclei  $A$  and  $B$ . In some expressions vacuum permeability  $\mu_0$  was assumed and therefore omitted in the  $D_{\max}$  expression, and there will be a factor of  $4\pi$  difference in the denominator of Eq. (1).  $D_{\max}$  is bond type dependent and usually on the order of  $10^3$  Hz, e.g.,  $D_{\max}$  for protein backbone amide  $^1\text{H}$ – $^{15}\text{N}$  spin nuclei pair is 21.66 kHz with an assumed bond length of 1.04 Å. Because of diffusive motion the direction of the internuclei vector fluctuates relative to the  $\mathbf{B}_0$  direction and therefore the dipolar coupling has to be evaluated with respect to every possible orientation. This orientation dependency follows the second order Legendre polynomials. Shown in Eq. (2) is the expression for the dipolar coupling  $D$  of an internuclei vector with a fixed distance (which is the case for bonded nuclei), where  $\Theta$  is the instantaneous angle between the dipole–dipole or bond vector and  $\mathbf{B}_0$  (Fig. 1), and the angular bracket indicates time or population average. The average



**Fig. 1** Illustrations of relationships between RDC internuclei vector  $AB$  and an arbitrary molecular frame (a) and the alignment tensor frame (b). A protein molecule, carrying spin nuclei  $A$  and  $B$ , is represented using an ellipsoid.  $B_0$  is the external field.  $\Theta$  is the instantaneous angle between the internuclei vector  $AB$  and  $B_0$ .  $\beta_{x,y,z}$  specify the projection angles of  $B_0$  onto each axis of a molecular frame. Polar angle  $\theta$  and azimuth angle  $\phi$  are spherical coordinates of the vector  $AB$  in the alignment tensor frame

for  $D$  will be zero if Eq. (2) is integrated over the variable  $\Theta$  that covers the surface of a sphere, corresponding to all possible orientation for an isotropic diffusing molecule. However, due to a weak alignment causing a nonisotropic sampling of orientations, the average of  $D$  is not zero. In practice, adjusting the concentration of the alignment media can allow between an equivalent of 0.1–1% fractions of protein molecules being aligned and it results in a value for  $D$  to be within  $\pm 20$  Hz for  $^1\text{H}$ – $^{15}\text{N}$  vectors, referred to as RDC.

$$D_{\max} = \frac{-\mu_0 h \gamma_A \gamma_B}{8\pi^3 r_{AB}^3}, \quad (1)$$

$$D = D_{\max} \left\langle \frac{3\cos^2\Theta - 1}{2} \right\rangle, \quad (2)$$

$$D = D_{\max} \begin{pmatrix} x & y & z \end{pmatrix} \begin{pmatrix} S_{xx} & S_{xy} & S_{xz} \\ S_{yx} & S_{yy} & S_{yz} \\ S_{zx} & S_{zy} & S_{zz} \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}, \quad (3)$$

$$S_{ij} = \left\langle \frac{3\cos\beta_i \cos\beta_j - \delta_{ij}}{2} \right\rangle, \quad ij = \{x, y, z\}. \quad (4)$$

The equation for RDC, Eq. (2), could be rewritten in the form of the expectation value for a vector in a Saupe order matrix, i.e., Eqs. (3, 4) [9]. The  $x, y, z$  in Eq. (3) are directional cosines of a bond vector in an arbitrary molecular frame (Fig. 1a), the most convenient would be an existing PDB coordinate frame; Saupe element  $S_{ij}$  is the averaged projection of axes of the molecular frame onto the direction of  $B_0$  with  $\beta_{x,y,z}$  (Fig. 1a) specifying the projection angle for each axis;  $\delta_{ij}$  is Kronecker delta. The Saupe matrix is symmetric and traceless and contains five unknown variables,

i.e.,  $S_{xx}$ ,  $S_{yy}$ ,  $S_{xy}$ ,  $S_{xz}$ , and  $S_{yz}$ . The product on the right side of Eq. (3) is a scalar, the value of which equals the average of the second order Legendre polynomials in Eq. (2). The aforementioned  $10^{-3}$  to  $10^{-4}$  (0.1–1%) scaling in the dipolar coupling is contained within all Saupe order elements. Since the five unknowns in the Saupe matrix are common to all bond vectors or RDC measurements in one aligned protein sample, in theory with more than five RDC bond vectors pointing in different directions one could solve the Saupe matrix. Prestegard and coworkers have written a protocol for solving the five Saupe unknowns using the singular-value-decomposition (SVD) method to obtain alignment parameters [10].

$$D = \frac{3}{2}D_{\max}(x \ y \ z)T^*(\alpha, \beta, \gamma) \begin{pmatrix} A_{xx} & 0 & 0 \\ 0 & A_{yy} & 0 \\ 0 & 0 & A_{zz} \end{pmatrix} T(\alpha, \beta, \gamma) \begin{pmatrix} x \\ y \\ z \end{pmatrix}, \quad (5)$$

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = T(\alpha, \beta, \gamma) \begin{pmatrix} x \\ y \\ z \end{pmatrix}, \quad (6)$$

$$D = D_a \left( 3\cos^2\theta - 1 + \frac{3}{2}R\sin^2\theta \cos 2\phi \right). \quad (7)$$

One can diagonalize the Saupe matrix to obtain the alignment parameters. Diagonalization of the Saupe matrix results in the principal order matrix and Euler rotation matrices  $T(\alpha, \beta, \gamma)$  and  $T^*(\alpha, \beta, \gamma)$ , where  $\alpha$ ,  $\beta$ , and  $\gamma$  are Euler angles and \* indicates conjugate transpose (Eq. (5)). The Euler rotation of Cartesian coordinates in the molecular frame  $xyz$  of Eq. (3) results in a set of new coordinates  $x'y'z'$  Eq. (6) for the bond vector within the principle order frame (Fig. 1b). The principal order matrix equals the alignment tensor matrix with its Eigenvalues,  $A_{xx}$ ,  $A_{yy}$ ,  $A_{zz}$  in Eq. (5), representing the alignment order in the corresponding tensor direction. There are different conventions in describing the alignment order parameters. For instance, one can keep the  $S$  representation for Eigenvalues in the alignment tensor matrix, i.e.,  $S_{zz}$ , etc. Alternatively, the  $A_{zz}$ , notation which is equivalent to  $2S_{zz}/3$  in the diagonalized Saupe matrix can be used [11].

Following the convention in Eq. (5), and in analogy to  $S_{xx}$  and  $S_{yy}$  of the diagonalized Saupe matrix, there are the following relationships:  $A_{xx} = A_{zz}(-1/2 + 3R/4)$  and  $A_{yy} = A_{zz}(-1/2 - 3R/4)$ , where  $R$  is the rhombicity which can be in the range of 0–2/3, with the convention that  $|A_{zz}| \geq |A_{yy}| \geq |A_{xx}|$ . Because  $A_{zz}$ , on the order of  $10^{-3}$  to  $10^{-4}$ , is not a convenient number,  $A_{zz}$  can be replaced with a more convenient  $D_a$  ( $= 3D_{\max}A_{zz}/4$ ) representation, which allows for an easier comparison among different alignment conditions. Here is an alternative to using the five Saupe unknowns. We use  $D_a$  to specify the alignment order of a sample, rhombicity  $R$  to describe asymmetry of the alignment tensor, and three Euler angles to define the tensor directions instead. With Euler angles one could conveniently visualize the tensor within a molecular frame. Further simplifications can be made by using spherical coordinates, i.e., polar angle  $\theta$  and azimuth angle  $\phi$ , to replace

$x'y'z'$  (Fig. 1b). After substituting  $D_a$ , rhombicity  $R$ , and the spherical coordinates in the alignment frame into Eqs. (5, 6), one will end up with a more familiar expression for RDC Eq. (7). It is essential to keep in mind that Eq. (2) and Eq. (7) are the similar formula, and the only difference is that the latter carries no ensemble averages.

### 3 Alignment Methods

Dipolar coupling manifests itself as an additional coupling to the scalar or  $J$ -coupling. RDC measurements normally require two NMR samples prepared in parallel, one with and the other without the presence of an aligning medium, corresponding to an isotropic and an anisotropic samples, respectively. An identical NMR pulse sequence is applied to both samples to measure the  $J$  and  $J + D$  values. The difference yields the dipolar coupling  $D$ . Therefore any pulse sequences developed to measure  $J$ -coupling are applicable for measuring dipolar coupling. There are two general methods to determine  $J + D$  value under anisotropic conditions, measuring direct splitting on NMR spectra and fitting a  $J$ -modulated intensity [12]. Methods for measuring  $D_{H-N}$ ,  $D_{H\alpha-C\alpha}$ ,  $D_{C\alpha-C'}$ ,  $D_{N-C'}$ ,  $D_{H-C'}$ , and  $D_{H-H}$  were extensively reviewed [11, 13, 14]. Here we focus on ways of adjusting alignment conditions and improvement in measurement accuracy.

A number of liquid-crystalline media have been employed to generate weak alignment for solution protein samples. The utilization and mechanisms of the commonly used alignment media have been reviewed previously [11, 15] and they include bicelles made of dimyristoylphosphatidylcholine (DMPC) and dihexanoylphosphatidylcholine (DHPC) [16, 17], filamentous phages *Pf1* [18] or *fd* [19], stretched [20, 21] or compressed [22] polyacrylamide gel, and poly(ethylene glycol)/hexanol mixture [23]. In addition, one can also add charged molecules into the alignment media to generate different alignment tensors for protein molecules [24, 25]. Though the pool of alignment media seems large, the demand for new media still exists for several reasons. One reason is that when there are more choices of alignment media the chance of finding a compatible one for challenging proteins or membrane proteins increases. The second reason is that any additional orthogonal alignment tensors from new alignment media add new information and can potentially better define bond vector orientation and dynamics. This has a promising potential for de novo protein structure determination. Recent additions in alignment media include charged gel [26], novel DNA based media [27, 28], and collagen gel [29]. The physical interactions between proteins and the medium are almost exclusively to be either steric and/or electrostatic, which could potentially limit the ability to obtain a complete set of five orthogonal alignment tensors. However, recent developments such as conservative mutation on protein surface [30] and use of composite media [31] may overcome these limitations. To improve measurement accuracy by providing more consistent reference (isotropic) data, approaches such as extracting RDCs from two samples at different



concentrations of alignment media [32] and application of a new two-stage NMR tube [33] have been proposed.

### 3.1 *Charged Polyacrylamide Gel*

Mechanically stretched [21] or compressed [22] polyacrylamide gel medium is a promising medium for measuring RDCs on membrane proteins that are reconstituted in micelles or bicelles. This is due to its inertness and inability to react with detergents. However, to establish weak alignment that is practical, it is necessary to polymerize at least 7% (w/v) acrylamide in a sample [34]. At such concentration the narrow pore size of the gel matrix limits protein diffusion, resulting in peak line-width broadening. Meier et al. [35] initially showed a sufficient alignment order was achieved by copolymerizing only 2% (w/v) acrylamide and acrylic acid, leading to an anionic polymer. Cierpicki and Bushweller [26] used acrylamide (<5% w/v) with different charged polymer units to generate alignment order. In addition to anions, positive charges were introduced by addition of (3-acrylamidopropyl)-trimethylammonium chloride or *N*-(2-acryloamidoethyl)-triethylammonium iodide. With such a charged gel, satisfactory sample stability and NMR spectra quality were obtained using integral membrane protein OmpA dissolved in dodecylphosphocholine (DPC) micelles. The RDCs obtained from the charged gel were directly used for membrane protein structure determination.

### 3.2 *DNA Based Media*

Nucleic acids carry a relatively stronger magnetic susceptibility than proteins and seldom react with detergent. In the continuous effort to develop detergent compatible alignment media, Douglas et al. [27] initially exploited DNA nanotubes as alignment media. Two bundles of six DNA strands in 7,000-base length were linked through base-pairing sticky ends to form a micron-long DNA filament. The trans-membrane helices of a T-cell receptor reconstituted in DPC/SDS bicelles were aligned in a cosolvent of DNA nanotubes at a concentration of 28 mg/mL. The measured  $D_{H-N}$  and  $D_{H\alpha-C\alpha}$  were shown to be consistent with the protein structure determined without RDC.

To ease the DNA nanotube medium preparation and reduce the cost, Lorieau et al. [28] used potassium salt of dinucleotide (d(GpG)) that would tetramerize through guanosine hydrogen bonds at a concentration of 10 mg/mL. The G-tetrad DNA, similar to bacteriophage *Pf1*, is strongly negatively charged. However, different from phage *Pf1*, its liquid crystalline phase is chiral nematic and the director can run perpendicular to the external field. Analysis of RDCs collected on a mutant of protein GB3 aligned in G-tetrad DNA showed the obtained alignment tensor carried the same directions as those in phage *Pf1*, but the sign of  $D_a$

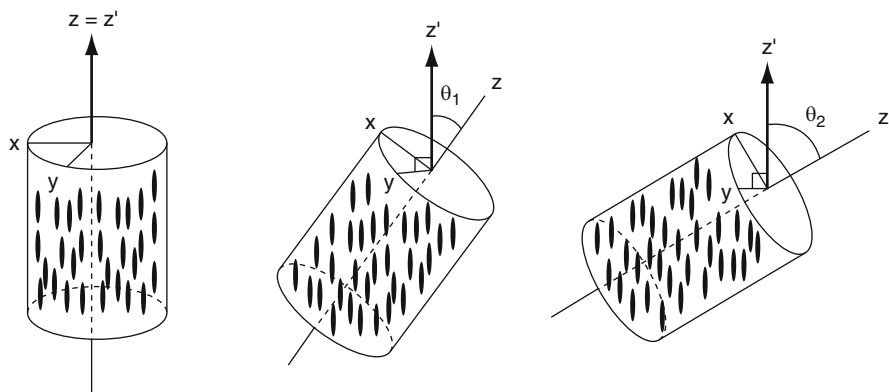
switched [28]. In this case, it is the electrostatic interaction between alignment medium and protein that determine the uniqueness of the alignment tensor. Nevertheless, G-tetrad DNA allowed RDC measurements on the fusion peptide of the influenza viral hemagglutinin solubilized in DPC micelles [36].

### 3.3 Collagen Gel

Collagen proteins, abundant in mammals, are made of trimer of polypeptide chains. Each chain is rich in proline or hydroxyl-proline, glycine, and others, which together form an extended left-handed polyproline-II structure [37]. Three copies of such chains form a right-handed helix, carrying a weak magnetic moment. For in vitro use the rat tail tendon type I collagen monomers were prepared [38] and stored in acidic buffer at 4 °C. At pH 6–8 and temperatures over 30 °C the collagen helices will polymerize and cross-link through amino acid side chains, e.g., lysine. When such reactions occur in the presence of magnetic field, weak alignment can be achieved [29]. At a concentration of 13 mg/mL, a sample with collagens polymerized in the presence of magnetic field yielded deuteron splitting of 20 Hz, usually sufficient for dipolar coupling [29].

### 3.4 Composite Media

As mentioned before, the nature of alignment forces, either steric or electrostatic, could limit the number of orthogonal alignment tensors we can observe experimentally. However, Ruan and Tolman [31] showed the interference of the two alignment forces could produce one additional orthogonal tensor. They polymerized 5% (w/v) polyacrylamide gel together with 3–4 mg/mL of bacteriophage *Pf1* in the presence of the magnetic field. A special tube with a cross section size of 7 × 5 mm was filled with both media and positioned at a maximal angle of 55° to field  $B_0$  (Fig. 2). The gel was dried, later soaked with protein sample, and stretched into a 5-mm NMR tube. Phages were assumed to be field aligned and trapped at the direction because their motions were inhibited by polymerized acrylamide. Essentially, the direction of phages formed a tilting angle with the direction of gel stretch. Proteins were subject to interfered alignment forces, resulting in both strong steric and electrostatic interactions. Surprisingly, the measured alignment tensors at different tilting angles are not a simple linear combination of those obtained from individual alignment using phage or acrylamide gel alone. As reported [31], at least three orthogonal tensors were solved for a ubiquitin sample subjected to composite alignment media. The same group also reported that the unambiguous bond direction might be determined with only three orthogonal alignment tensors [39].



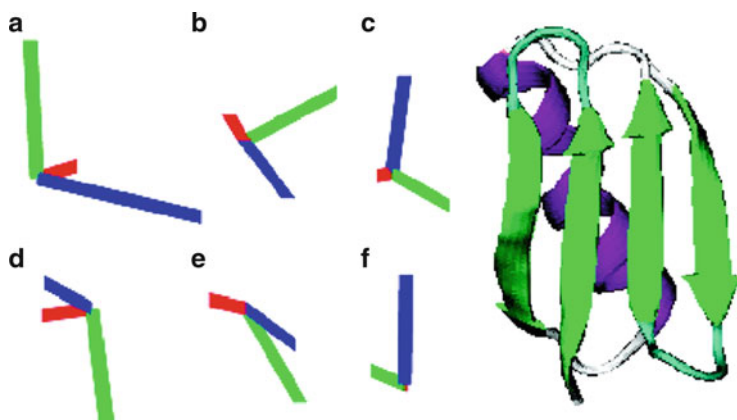
**Fig. 2** Variation of alignment using stretched polyacrylamide gels (SAG) and bacteriophage *Pf1*, which has been embedded and aligned (along  $z'$ ) at different angles relative to the long axis ( $z$ ) of the sample. The gels were cast in an approximately ellipsoidal (squashed cylinder) geometry, with dimensions of 5, 7, and 10 mm along the  $x$ ,  $y$ , and  $z$  axes, respectively. The gels were then dried, rehydrated, and stretched to fit within a 4.2 mm i.d. NMR tube. (Reprinted with permission from [31])

### 3.5 Conservative Mutation

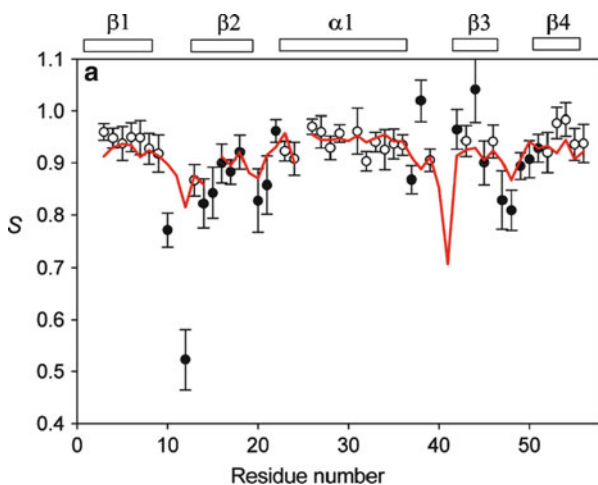
The surface charge of *Pf1* phage is heavily negative and proteins will be aligned according to its surface electrostatic potentials [40]. Yao and Bax [30] carefully modified the surface charge distribution of the 6-kDa protein GB3 by either conservatively mutating one or two residues at each time, e.g., K to E, or keeping the histidine-tag at terminus of the native GB3. The backbone structure was later found to be unperturbed in those modified proteins. A total of six mutant proteins were found to align quite differently with respect to the  $B_0$  field and five singular values were obtained (Fig. 3). With this well-defined system, the amplitude and direction of RDC bond dynamics up to the millisecond (ms) were unambiguously obtained and compared to nanosecond to picosecond (ns–ps) dynamics from spin relaxation measurements (Fig. 4) [41]. Interestingly, Yao et al. found both RDC and relaxation measurements showed the same amount of flexibility for residues in regular secondary structures. Bond vectors within loop regions, however, were shown by RDC to have larger amplitude of motions compared to what were suggested from the relaxation data [41].

### 3.6 RDC/RCSA Accuracy Improvement

In addition to RDC, protein alignment creates slight chemical shift changes for the aligned sample relative to the isotropic sample. The chemical shift difference, named residual chemical shift anisotropy (RCSA), comes from the projection of the chemical shift tensor, which is not averaged to zero, onto the alignment tensor. RCSA is also a long range structural restraint, providing orientation dependences



**Fig. 3** Alignment tensor orientations relative to the ribbon backbone structure of GB3 for six mutants, all in liquid-crystalline *Pfl* medium. The six tensors are for A – K19AD47K; B – K19ED40N; C – K19EK4A-C-His-tag; D – K19EK4A-N-His-tag; E – K19AT11K; F – K19EK4A. Diagonalized tensor elements, Dxx (red), Dyy (green), and Dzz (blue) have magnitudes proportional to the length of the corresponding lines. (Reprinted with permission from [30])



**Fig. 4** Experimental order parameters,  $S$ , of NH bond vectors in GB3 derived from iterative Direct Interpretation of Dipolar Couplings (DIDC) using all six sets of RDCs. The red line marks the order parameters derived from  $^{15}\text{N}$  relaxation. Filled symbols represent residues for which the fully anisotropic model was required to get a satisfactory fit to the data, while, for open symbols, the isotropic internal motion model was able to fit the RDCs to within the experimental noise. (Reprinted with permission from [41])

generally complementary to RDCs [42–44]. The solvent used in the isotropic protein sample usually differs from that in the anisotropic sample. This can introduce errors in measuring both dipolar couplings and chemical shift differences. Such error may be below the measurement accuracy for RDCs, on the order of 0.1–1 Hz. However,

RCSA, which is on the order of parts per billion (ppb), may be greatly affected. The solvent effect can be removed by measuring several aligned samples at increasing alignment orders and extrapolating the RCSA values to zero concentration [32]. The methods were applied to backbone  $^{15}\text{N}$  [32] and  $^{13}\text{C}'$  [45] RCSA measurements.

Another way to overcome this is to measure isotropic sample under the same unaligned medium. Liu and Prestegard [33] developed a two-stage NMR tube with two different internal diameters (I.D.) at upper and lower parts of an open-ended NMR tube. For isotropic condition the protein sample was soaked into polyacrylamide gel in 5 mm I.D. part, then the same piece of gel was pulled and stretched into the 3 mm I.D. part by vacuum created using a syringe at the other end of the two-stage tube, resulting anisotropic condition. The method allows for higher accuracy measurements for both RDC and RCSA.

Aside from using both aligned and unaligned samples for RDC and RCSA, one may also keep a single sample that contains alignment media for both isotropic and anisotropic measurements. Upon the application of magic angle spinning, similar to what is used in solid state NMR, the alignment effect for bicelles [46] or phage [47] was removed. In this way the measurement will not contain any solvent effects.

## 4 Interpretation and Implementation

Measured RDC values are representative averages of the whole ensemble of dipolar interactions within protein molecules in solution. Such an ensemble should include all protein conformers interconverting at time scales faster than the inverse of RDC values ( $1/D$ ). For instance, the observed dipolar coupling is affected by the internuclei or bond vectors that stretch and vibrate on a femtosecond to nanosecond (fs–ns) time scale, the protein domain reorientation on a nanosecond to microsecond (ns– $\mu\text{s}$ ) time scale, and conformational change that ranges from nanoseconds, e.g., unstructured terminus, to milliseconds. It is nearly impossible to describe protein structure and dynamics using RDC values without any assumptions. Some approximations have to be made in interpreting RDC measurements.

### 4.1 Approximations

Most RDCs except for  $D_{\text{H-H}}$  [48] are measured on fixed internuclei distances such as bond vectors or geometries within the peptide plane so that in Eqs (1) and (2) a constant internuclei distance  $r_{AB}$  in  $D_{\text{max}}$  is assumed. This in itself is an approximation because the effective bond lengths vary due to dynamic processes. For instance, Yao et al. [49] determined the average protein backbone H–N bond length to be 1.01–1.02 Å. However, for deriving the true alignment order an effective value of 1.04 Å was proven proper and used extensively to account for H–N bond libration dynamics [50]. A slight increase in effective bond length for RDC analysis also applies

to  $H\alpha-C\alpha$  bonds. On the other hand, bond lengths for heavy nuclei, e.g.,  $C\alpha-C'$ , are consistent with values reported from crystallography studies [50].

By expressing RDC using Eq. (3), where the Saupe matrix is multiplied by bond directional cosines within a local molecular frame, we assume protein conformational change does not affect the alignment tensor. Thus Eq. (3) is an approximation to Eq. (2), in which the average bracket over the second order Legendre polynomial is applied. This approximation is valid when RDCs from the structured part of the proteins are analyzed. For small proteins GB3 [41, 51] and ubiquitin [52] the model free (MF) order parameter from  $^{15}\text{N}$  spin relaxation analysis and the order parameter derived from RDC analysis were strongly correlated. These indicate that the effect of slowmotion dynamics on microsecond to millisecond time scales on RDC may be negligible and a single structure representation is sufficient for the structured part of proteins within the current experimental uncertainty [53].

When the dynamics involve large scale amplitudes of motion, e.g., the MF order parameter  $S^2 < 0.6$ , multiple structures may exist and each one of them is subjected to its own alignment tensor that may vary significantly due to different steric, and electrostatic interactions of each conformer to the alignment media. Specifically for unfolded proteins it becomes challenging to separate intrinsic bond dynamics from Saupe order parameters in the laboratory frame because local motion and the overall perturbed diffusive motions are coupled [54]. Monte-Carlo simulation of an ensemble of conformers and the following comparison to measurements remain the only option to interpret RDCs in a flexible system [55, 56]. The application of RDCs to study protein in an unfolded state is an active research field [57–59]. For instance, efforts have been made to represent the conformational space of urea denatured ubiquitin with as few as 200 conformers, which seem to reproduce measured RDCs [59]. However, additional specific bond type RDC scaling factors had to be applied for H–N,  $H\alpha-C\alpha$ , etc., which indicate different amounts of motion along bond vectors within a peptide plane [57, 59].

## 4.2 Common Applications of RDC

The most common use of RDCs is for structure validation and refinement. Given structure coordinates determined either from X-ray crystallography or solution NMR, one can readily fit RDC data to the corresponding bond vector directions within the molecular frame. Any programs that utilize Eqs (3) or (5) in a chi-square minimization routine or formal software packages such as PALES [60, 61] and REDCAT [62] which implement the SVD algorithm can be employed to carry out the numerical fitting. The fitting will result in the optimized alignment tensor and values that best matched the measured RDCs. To quantify the agreement between a structure and measured dipolar couplings, Cornilescu et al. [63] proposed quality factor  $Q$ . This factor can better determine the quality of the fit than Pearson's correlation  $R$ . Shown in (8) is the expression for the  $Q$  factor where rms refers to root mean square. It provides an estimate of average disagreement in percentage

between measured and calculated dipolar couplings. A  $Q$  factor of 40%, roughly corresponding to Pearson's  $R$  of 0.9, are commonly found on structures with 2–3 Å resolution; and a  $Q$  value of 20% or less indicates the structure is at high resolution (1.5 Å) and accuracy. Clore and Garrett [64] suggested an alternative form  $[2D_a^2(4 + 3R^2)/5]^{1/2}$  to replace the denominator of (8) in the case of a limited RDC sampling over all of the possible orientations. A thorough discussion on  $Q$  factor can be found elsewhere [65].

$$Q = \frac{\text{rms}(D_{\text{meas}} - D_{\text{calc}})}{\text{rms}(D_{\text{meas}})}, \quad (8)$$

$$E = k(D_{\text{meas}} - D_{\text{calc}})^2. \quad (9)$$

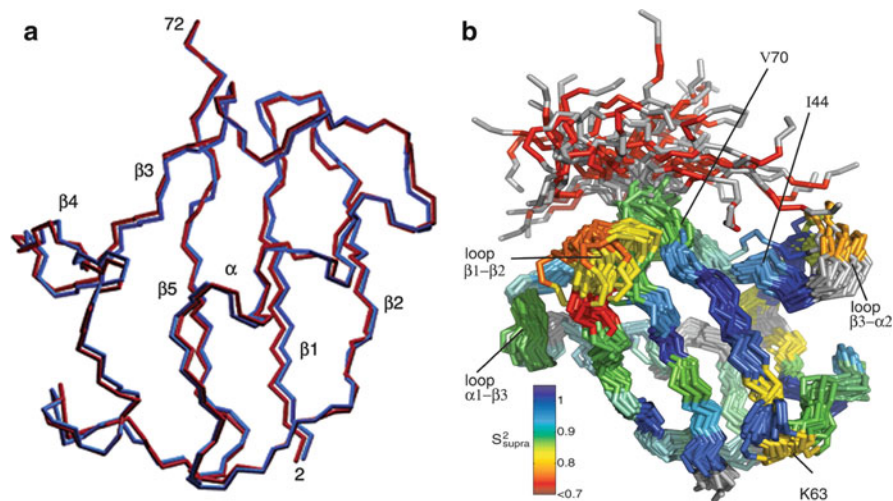
For solution structure determination, RDC is normally not included as a potential term (9) during the initial structure calculation in a simulated annealing protocol in a program such as Xplor-NIH [66, 67]. Due to directional degeneracy associated with RDC restraints, as described in the theoretical expression section, multiple local potential minima will hinder the successful search for the right conformation. Typically a rough tertiary fold of protein is obtained first with the use of NOE restraints and such a fold can be used as a starting structure for the next simulated annealing procedure with the combined NOE and RDC restraints. The procedure to include RDC restraints in a program such as Xplor-NIH uses four pseudo-atoms (OXYZ) [68] to represent the alignment tensor directions [69]. The RDC force constant,  $k$  in (9), can be increased gradually as the temperature is being lowered in the simulated annealing. In the end it is ideal to adjust the force constant  $k$  so that the deviation from the measured RDCs matches the experimental error.

### 4.3 Ensemble Minimization

As discussed before, RDC reflects the ensemble averaged dipolar coupling and incorporates a wide range of time scales up to milliseconds. During RDC restrained structure calculations we made approximations by assuming a single conformer. The question remains whether it is reasonable to keep this assumption in all cases. Clore and Schwieters [52] initially employed a two-member ensemble minimization to test whether a better agreement between measured and calculated RDC data could be achieved, and whether it was statistically significant. The ensemble algorithm proposed by Clore and Schwieters [52] kept a user-defined number of conformers during the course of a simulated annealing, and any evaluated physical quantity, e.g.,  $D_{\text{calc}}$ , is linearly averaged among individual conformers. After refining the ubiquitin structure with  $D_{\text{H-N}}$  data sets collected in 11 different alignment media, and a couple of other heteronuclear RDCs, they found a single conformer in most cases is sufficient to yield a good quality factor  $Q$  less than 20%. The structure was validated against the  $D_{\text{H}\alpha\text{-C}\alpha}$

data set that was not being applied as a restraint. A 15% reduction in  $Q$  value was obtained with an ensemble size of two, which was statistically significant. In addition, several residues were found to be undergoing anisotropic motion and can be better represented by the two-member ensemble (Fig. 5a) [52]. Later a simulated annealing with an ensemble size of four to eight on protein GB3 had been found to be optimal [71]. The consensus is that other than residues having large amplitude or anisotropic motions, most structured residues can be represented with a single structure for monomeric small proteins because the accuracy of NMR structural coordinates is well within the measurement uncertainty of the RDCs.

Other studies on ensemble minimization reveal a relatively larger conformational sampling for ubiquitin [72]. In another study, an ensemble of around 100 ubiquitin structures was generated using NMR restrained simulated annealing and molecular dynamics. Interestingly, such an ensemble covers structural heterogeneity observed in 40 or so ubiquitin–ligand complex crystal structures (Fig. 5b) [70]. RDC restraints used in such ensemble minimizations were measured in over 30 different alignment media [73]. The biological implication is that ligand induced ubiquitin conformers preexist in its structural ensemble. Further, the ubiquitin N–H order parameters derived from RDC analysis [73–76] were overall lower by 0.1–0.2 than that of MF order parameter derived from relaxation data [77], indicating an appreciable amount of microsecond to millisecond motion that was not observed in spin relaxation measurement. It is nevertheless an interesting and different conclusion from other findings on ubiquitin [52] and GB3 [41].



**Fig. 5** Structure ensemble of Ubiquitin. (a) Two members (shown in *red* and *blue*) of a typical ensemble from the two-member size calculation. (Reprinted with permission from [52].) (b) Backbone trace of 40 randomly chosen structures from the ensemble. Residues are colored by the amount of additional (slower than- $\tau_c$ ) mobility as compared with the Lipari-Szabo order parameters. (Reprinted with permission from [70])

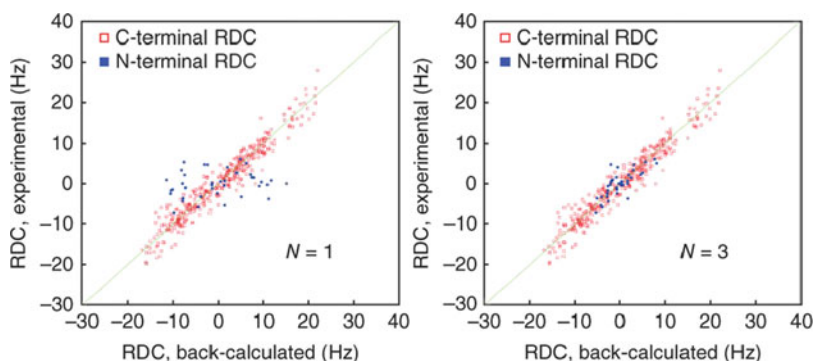


Besides backbone dynamics for monomeric proteins, ensemble minimization may be better suited for studying relatively slower and larger scale protein domain dynamics and structure. Liu et al. [78] implemented ensemble minimization to characterize the dynamic structure of a membrane-anchored ADP ribosylation factor (Arf). The yeast Arf is composed of flexibly joined C-terminal domain and the short N-terminal helix that interacts with the membrane. The Arf protein anchored on DMPC/DHPC bicelles were aligned in a negatively charged polyacrylamide gel, and the resulting alignment order parameters  $D_a$  and  $R$  for C- and N-terminal domains are 8.6 Hz and 0.29, and 4.7 Hz and 0.61, respectively. The differences in alignment order indicate significant interdomain motions. A three-member ensemble minimization significantly increased the linear correlation coefficient between measured and back-calculated RDC (Fig. 6) and PRE data on the N-terminal helix. Such dynamics could be crucial for Arf to function as a regulator of effectors GTP/GDP.

#### 4.4 Structure of a Ligand in a Bound State

The conformational ensemble described previously represents a continuous distribution of conformers and no significant protein fold or high-order structural change within the ensemble members was taken into account. Other systems may also be composed of discrete conformers in equilibrium; for instance, large structural differences could exist between free and bound states of ligand molecules in the presence of their receptors. Generally the NMR signals from the ligands in this type of equilibrium are dominated by the ligands in the free states. RDCs can play a role in solving ligand structure in the bound form not generally accessible by other techniques.

Photo-activated retina rhodopsin, a G-protein coupled receptor (GPCR), is a major component in purple membranes (PM). Due to the high structural content of helices in GPCR, GPCR rich PM disks align in the presence of a magnetic field



**Fig. 6** Ensemble structural fitting to RDCs. Agreement between experimental and back-calculated RDCs for a one-state ensemble (*left*) and a three-state ensemble (*right*). Data include NH,  $\text{NC}'$ , and phenyl CH (the latter two are normalized to NH). RDCs are collected in positive and negative gels. (Reprinted with permission from [78])

with the membrane normal being parallel to the field [79, 80]. The binding of the GDP-bound form of the heterotrimeric G protein transducin (Gt) to the light activated MII intermediate of the rhodopsin activates the visual signal transduction pathway. A peptide consisting of the C terminal ten residues of the  $\alpha$ -subunit of transducin (Gt $\alpha$ ) is capable of competing with Gt for binding and was studied using NMR spectroscopy [81]. In the presence of light, Gt $\alpha$  transiently binds activated GPCR that is embedded in membrane disks of bovine retina. Since these disks were aligned in the magnetic field, RDCs for Gt $\alpha$  could be observed. The measured RDCs followed a time-decay course after light activation, while GPCR returned to its inactive state. The decay time constant is on the order of 1 h that allowed rapid acquisition of 2D  $^1\text{H}$ - $^{15}\text{N}$  and  $^1\text{H}$ - $^{13}\text{C}$  spectra on the Gt $\alpha$  peptide sample. The RDC values at zero time were obtained from extrapolation of multiple RDCs collected along a time decay curve. The structure and orientation of the peptide in the bound state to GPCR were derived from RDCs and transferred NOEs [81]. The validity of such use of RDCs relies on two conditions. The first condition is the fast exchange between the free and bound states of Gt $\alpha$  that allows the observation of only one set of resonances of the free state. The other condition is that the free Gt $\alpha$  peptides are not aligned in GPCR enriched PM media so that alignment contribution all comes from the bound state. Indeed the alignment order of the system was weak with a  $D_a$  value of only 1.6 Hz.

In the above case, the ligand receptor GPCR is naturally buried in bicelles, which by itself is the alignment medium. This ensures strong alignment order for bound ligand molecules. However, for a general study protein receptors are not necessarily buried in any disks that could be aligned; thus the alignment order for ligand molecules in equilibrium with its free form would be too weak to be observed. Seidel et al. [82] demonstrated a method of anchoring polyhistidine-tagged protein receptors onto the bicelles that were doped with histidine-tag binding lipid molecules. This increased the alignment order of the protein–ligand complex and prevented measurements of the protein RDCs. The RDCs of the bound ligand in fast exchange with free ligand molecules, however, were observed and they could be used to map the ligand configuration on the protein receptors.

#### 4.5 Structure of Oligomeric State

RDCs are very useful in improving the structure determination of a monomeric protein. In addition, they can be utilized to establish quaternary structure of symmetric oligomers. Prestegard and coworkers have developed methods to determine dimer structure using RDC data [83, 84]. The underlying principle is that the rotational symmetric  $C_2$  axis of any protein dimer should be parallel to one axis of the alignment tensor of the dimer. By solving alignment tensor axes from different alignment media, one could identify the common tensor axis that should be parallel to the symmetric axis [85, 86]. *Bacillus subtilis* proteins YkuJ dimerize with strong affinity and its  $K_d$  is on the order of  $10^{-9}$  M [83]. When working at the typical NMR

protein concentration ( $\sim 0.1$ – $1.0$  mM) the measured RDC would primarily reflect the dimer structure. The crystal structure of YkuJ is a tetramer and there are two possible configurations for a dimer structure. Wang et al. [83] aligned the protein using two alignment media, 5% poly(ethylene glycol)/hexanol and 10 mg/mL phage *Pf1*. Both yielded different alignment tensor directions. However, the  $x$  axes from the two tensors were close to each other and were believed to be the symmetric axis of the dimer. The further rotation and docking procedure with energy minimization yielded dimer structure very close to one possible dimer structure observed in the crystal structure of the tetramer. Prestegard and colleagues further extended this approach to study structures of weakly associated dimer. The *Staphylococcus epidermidis* proteins SeR13 weakly dimerize at a  $K_d$  on the order of  $10^{-3}$  M [84]. The measured RDCs were the weighted average of monomer and dimer states and the resulting alignment tensor from direct fitting would be the average of both monomer and dimer alignment tensors. Lee et al. [84] circumvented this by deriving the exact  $K_d$  value from concentration dependent chemical shift values, then extrapolating the RDCs as a function of protein concentration together with the  $K_d$  to yield pure RDC data set for the dimer. The rest of the work to determine the dimer structure was similar to the previous approach [83], but with the additional restraints from chemical shift perturbation and paramagnetic surface perturbation data that helped identifying newly buried surface residues upon dimerization.

## 5 Conclusion

Some additional reviews [11, 65, 87] can be found on RDC topics with different emphasis such as methods [69, 88, 89], theories [13, 14, 90], and dynamics [15]. In this chapter we have briefly summarized the theoretical expressions for the widely used RDCs and common descriptions for magnetic alignment under either Saupe order matrix or alignment tensor representations. RDCs have become notably more relevant in solution structure and dynamics studies for larger system due to the fact that the number of NOEs diminishes. In some cases RDCs may be the sole major experimental NMR restraints for determining domain positions in large systems of over 100 kDa [8]. With its essentiality in mind, more alignment media are being explored. The more choices available, the higher the chance biologically important proteins and membrane proteins reconstituted in micelles or bicelles can be aligned without any interference with the media. More RDC measurements on large and multidomain proteins can reveal their quaternary structure and dynamics manifested through differences in alignment order. However, caution should be exercised in RDC data interpretations because a single RDC data set may not be able to separate differences in domain positions and dynamics simultaneously. Cross-validating the results with different RDCs under another orthogonal alignment tensor, rotational diffusion tensors from spin relaxation, and small angle X-ray scattering that are sensitive to the protein overall shape should help overcome this problem.

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# NMR Studies of Metalloproteins

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**Abstract** Metalloproteins represent a large share of the proteomes, with the intrinsic metal ions providing catalytic, regulatory, and structural roles critical to protein functions. Structural characterization of metalloproteins and identification of metal coordination features including numbers and types of ligands and metal-ligand geometry, and mapping the structural and dynamic changes upon metal binding are significant for understanding biological functions of metalloproteins. NMR spectroscopy has long been used as an invaluable tool for structure and dynamic studies of macromolecules. Here we focus on the application of NMR spectroscopy in characterization of metalloproteins, including structural studies and identification of metal coordination spheres by hetero-/homo-nuclear metal NMR spectroscopy. Paramagnetic NMR as well as  $^{13}\text{C}$  directly detected protonless NMR spectroscopy will also be addressed for application to paramagnetic metalloproteins. Moreover, these techniques offer great potential for studies of other non-metal binding macromolecules.

**Keywords** C direct detection · Metal coordination · Metalloprotein · NMR spectroscopy · Paramagnetic metalloprotein

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## 1 Introduction

Metal ions play important roles in life science and the molecular mechanism of metal-dependent life processes and all matters regarding metal within a cell or tissue/organ are topics of the emerging fields of metallomics and metalloproteomics [1–3]. Metal ions must usually associate with proteins (and other biomolecules), i.e., metalloproteins, to prevent the toxic effects of metal excess. Metalloproteins are one of the most diverse classes of proteins with the intrinsic metal ions providing catalytic, regulatory, and structural roles critical to protein function, and are found in plants, animals, and many microorganisms. It has been estimated that metalloproteins account for approximately one-quarter to one-third of all the proteins in the human body [4]. A recent study revealed that many of metalloproteomes still remain uncharacterized [5]. A systematic bioinformatics survey of 1,371 metalloenzymes with known structures showed that about 40% of enzyme-catalyzed reactions involve metal ions, e.g., magnesium, zinc, and iron [6]. Enormous effort has been devoted toward understanding the structure and function of metalloproteins and such knowledge has been used to design a new functional metalloprotein [7] and to rationalize and to search for new metalloproteins by a bioinformatics approach [6, 8]. However, it is impossible currently from gene sequences to predict the numbers and types of metals an organism assimilates from its environment or uses in its metalloproteins because the geometry and composition of metal binding site are diverse and poorly recognized [9, 10]. Therefore, understanding of the function of metalloproteins comes from individual characterization of the structures of the proteins and chemical states of the metal centers by various spectroscopic techniques including NMR spectroscopy, circular dichroism (CD), electronic absorption spectroscopy (UV), small angle X-ray absorption, as well as extended X-ray absorption fine structure (EXAFS).

Over the last three decades, NMR spectroscopy has been developed into a very important and versatile analytical technique both in the chemical and biological sciences. It has been used within the framework of Structural Genomic (SG) projects worldwide for determination of structures of proteins at the atomic level under physiologically relevant conditions [11–13]. Moreover, NMR spectroscopy is applicable to study of the interactions of proteins with other molecules including proteins, nucleic acids, and even small molecules which are mainly based on the sensitivity of the chemical shifts towards changes in chemical environments [14, 15]. Application of this technique to structural studies is limited to small proteins (30–35 kDa) even with the aid of isotopic labeling (<sup>13</sup>C, <sup>15</sup>N and <sup>2</sup>H) although backbone assignments of a 723-residue enzyme with a molecular weight

of 81.4 kDa were achieved [16]. Recently, a few breakthroughs have been made in this field. This includes the first NMR structure of a seven-helix transmembrane protein determined in membrane-mimetic environments [17] and the first structure determined in living cells by in-cell NMR [18]. Apart from its application in structural studies, NMR spectroscopy is also able to monitor the internal motion of biomolecules ranging over times from subnanoseconds to beyond seconds. Characterization of dynamics of biomolecules, such as folding transition, will be a great help for our understanding of the biological function of biomolecules.

Application of NMR spectroscopy in studies of metalloproteins is in principle the same as other proteins if the proteins contain diamagnetic metal ions. In the case of paramagnetic metalloproteins, things become more challenging since paramagnetic metals affect longitudinal and transverse relaxation rates of observed nuclei. However, recent advances in the hardware and methodology have enable structures of such metalloproteins to be determined, these advances including  $^{13}\text{C}$ -detected experiments, solid-state NMR, and the discovery of paramagnetic relaxation enhancement (PRE) [19–21]. In this review, we will address mainly the application of NMR spectroscopy in studies of metalloproteins including the contribution of NMR to structural characterization of metalloproteins, and special attention will be devoted to the utilization of NMR in characterization/identification of the metal binding sites and their coordination environments as well as to probing conformational changes of metalloproteins upon metal binding and release. The techniques used for paramagnetic metalloproteins will also be included briefly since a number of reviews in this field can be found [19–22]. A systematic review of the application of metal NMR spectroscopy will not be made and interested readers are directed elsewhere [23].

## 2 The Contribution of NMR to Structural Metalloproteins

### 2.1 *Conventional Method for Structure Determination of Metalloproteins*

NMR spectroscopy can be applied to characterize structurally diamagnetic metal containing metalloproteins similar to other proteins. Amongst the structures deposited in the Brookhaven Protein Data Bank (PDB), ca. 15% of the structures were resolved by NMR. Here, we will give a brief summary since the detailed methodology can be found in numerous references [11, 24–26].

In structural studies of proteins/metalloproteins, concentrations of about 1 mM are typically required and proteins must be soluble and stable over a period of time (weeks). For small proteins with several tens of amino acids, e.g. metallothionine [27, 28], it is sufficient to use  $^{15}\text{N}$ -labeled samples to determine structures of the proteins. However, if proteins can be overexpressed in a bacterial system (e.g., *Escherichia coli*), it is desirable to overexpress the protein with uniform enrichment

of  $^{15}\text{N}$  and  $^{13}\text{C}$  even for small proteins. This will make full use of multidimensional heteronuclear NMR experiments to increase the spectral resolution. Unlike other proteins, expression of metalloproteins in a bacterial system usually requires specific metal ions to be supplemented in the medium to induce overexpression of the targeted proteins [14, 29]. Alternatively, metal ions have to be incorporated into the proteins after purification, particularly if the metal ions play structural roles, as otherwise the proteins may not be stable for structural characterization. For example, HypA from *Helicobacter pylori* precipitates easily in the absence of zinc which serves a structural role [15]. However, caution has to be taken during metal incorporation since excess metal ions may also cause protein aggregation.

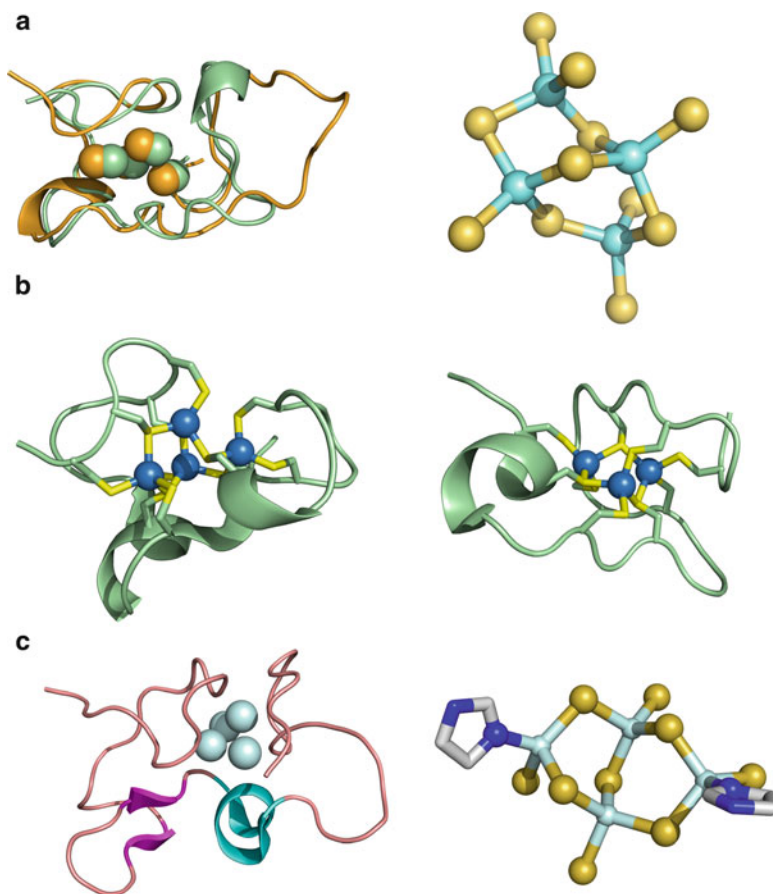
A major bottleneck in solving protein structures by NMR is the highly peak-picking and assignment of chemical shifts and NOEs. The strategy of the assignment process and structure calculation can be found in an excellent review [30]. In general, for a  $^{15}\text{N}/^{13}\text{C}$ -labeled protein, a series of double/triple resonance experiments are recorded for resonance assignments. Backbone assignments are derived from HNCA, HNCOC, HNCACB, CBCACONH, HNCO, and HNCACO, whereas side-chain protons and carbon atoms are assigned from HCCH-TOCSY, HCCH-COSY, HBHACONH, C(CO)NH, and H(CCO)NH [24, 25]. The chemical shifts of backbone and side-chain are then used to assign NOEs ( $^{15}\text{N}$ -/ $^{13}\text{C}$ -HSQC-NOESY) to derive inter-protein distance restraints. Usually the structure determination process goes through several iterations of compiling a NOESY peak list, assignment of NOE cross-peaks to sequence-specific interactions, structure generations and assessment, refinement of NOESY peak lists, and reassignment of the cross-peaks, which can be carried out automatically [31, 32]. In addition to distance restraints, dihedral angle restraints are usually obtained from several experiments, e.g., HNHA [33] or HNHB [34], or predicated from TALOS, a program that empirically predicts backbone angles ( $\phi, \psi$ ) based on the chemical shifts of  $\text{H}^\alpha$ ,  $\text{C}^\alpha$ ,  $\text{C}^\beta$ ,  $\text{C}'$ , and N [35], as well as the H-bond restraints derived from H-D exchange experiments. For elongated macromolecules, residual dipolar couplings (RDC) as additional restraints are necessary for structure determination. Information about RDC can be found in this book. In order to get relatively good quality for the structures, the numbers of NMR restraints used for structural determination are usually of the order of 10–20 independent interatomic distances per amino acid plus some dihedral restraints and as well as atom–atom vector directions. The quality of calculated structures has to be evaluated using programs PROCHECK, WHATIF, etc., and detailed description can be found in a recent review [13].

In addition to the general strategies described above, metal-based NMR parameters are also of great help in the evaluation of structures of metalloproteins, especially for those metalloproteins whose folding is highly metal-dependent. To incorporate metal cluster constraints into structural calculation, residues that coordinate to metal ions (e.g.,  $\text{Zn}^{2+}$ ) must be identified first either by mutagenesis studies or by physical characterization such as UV absorption spectroscopy, EXAFS, and NMR spectroscopy [36–38]. Providing that metal coordination residues and geometries are unveiled, metal cluster restraints can be obtained based on relevant crystal structures of either macromolecules or small molecules. Usually, metal cluster restraints are not

used in initial structure generation to avoid bias, especially during the automated assignment procedure in CYANA [39]; instead, metal ligand restraints are incorporated in the last step of structural calculation as well as in the refinement stage [40, 41].

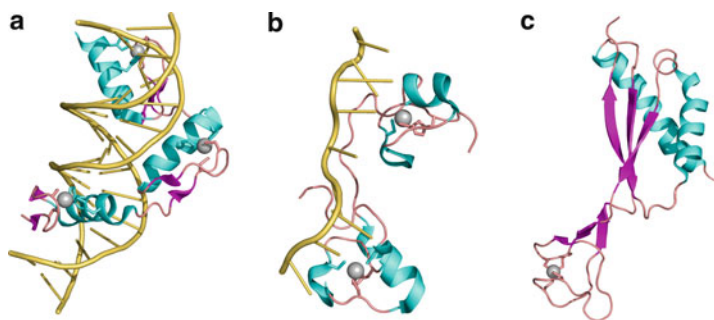
The application of NMR spectroscopy in protein structure determination actually started with a small metalloprotein, metallothionine (MT) [42]. Metallothionines are a class of low molecular weight (typically 6–7 kDa) cysteine-rich proteins. The proteins lack a well-defined secondary structure and their folding is dictated mostly by a clustered network of cysteine residues and metal ions usually represented by  $\text{Zn}^{2+}$ ,  $\text{Cu}^+$ , and  $\text{Cd}^{2+}$  [43, 44]. Since the first solution of the structure of rabbit liver  $\text{Cd}_7\text{MT}_2$  [42], numerous three dimensional structures of metallothionines from different isoforms (MT1/MT2/MT3) or different species such as blue crab and mammalian (rabbit, rat, and human) have been resolved by NMR spectroscopy [45] with only one structure (rat liver  $\text{Cd}_5\text{Zn}_2\text{MT}_2$ ) determined by X-ray crystallography [46]. The protein consists of two dynamic metal-thiolate clusters folded into two domains ( $\alpha$ ,  $\beta$ ) and the structural mobility of the protein makes it difficult to be crystallized. The metal cluster restraints, e.g., Cd–S bond lengths, as well as Cd–S–Cd, S–Cd–S, and  $\text{Cys}^{\beta}\text{–S–Cd}$  bond angles from the X-ray crystal structures of model cadmium complexes and rat liver  $\text{Cd}_5\text{Zn}_2\text{MT}_2$  were often incorporated with other distance and angle restraints in structure calculation. Recently, a new member of metallothionine MT3 with the conserved CPCP motif in the N-termini has been involved in the growth inhibitory activity and is down-regulated in the brain of Alzheimer's patients [47]. The solution structures of both human [27] and mouse MT3 [48] resolved by NMR spectroscopy for the C-terminal  $\alpha$ -domain, Fig. 1a, revealed a similar  $\text{Cd}_4\text{Cys}_{11}$  cluster as well as very similar tertiary folds to MT1/2. However, a loop in the acidic hexapeptide insertion is found and is slightly longer in human MT3 than in mouse MT3. The first solution structure of  $\text{Cd}_7\text{MT-nc}$  of the Antarctic fish *Notothenia coriiceps* was also determined [28]. The position of the ninth cysteine of  $\text{Cd}_7\text{MT-nc}$  is different from mammalian MT which results in a structural change of the domain, in particular in the orientation of the loop (Lys50–Thr53), Fig. 1b, and in turn a different charge distribution with respect to mammalian MT [28]. Interestingly, an intriguing class of histidine-containing metallothionines has also been identified in fungi and bacterial [49]. The histidine residue has been thought to be able to modulate zinc affinity and reactivity. Solution structure of one of this class of MTs,  $\text{Zn}_4\text{SmtA}$  from cyanobacterium *Synechococcus* PCC 7942 was determined [41], Fig. 1c, revealing a  $\text{Zn}_4\text{Cys}_9\text{His}_2$  cluster with a topology similar to that of the  $\text{Zn}_4\text{Cys}_{11}$  cluster of the  $\alpha$ -domain of mammalian MT. However, the two  $\text{ZnCysHis}$  sites and one  $\text{ZnCys}_4$  site readily exchange  $\text{Zn}^{2+}$  for exogenous  $\text{Cd}^{2+}$ . Moreover, SmtA contains a short  $\alpha$ -helix and two small antiparallel  $\beta$ -sheets surround the inert zinc site, which resemble zinc finger portions of GATA and LIM proteins. Such a structure of SmtA probably produces its function of specific protein and/or DNA recognition [41].

NMR spectroscopy makes an enormous contribution to structural biology of metalloproteins, particularly in zinc-binding proteins. Zinc, the second most abundant metal found in eukaryotic organisms, plays important catalytic and structural roles in a variety of biological processes. Binding of zinc is able to stabilize the folded



**Fig. 1** (a) Overlaid view of structures of human MT3 (PDB: 2f5h) in *orange* and mouse MT3 (PDB: 1JI9) in *pale green* with the cadmium ions shown as spheres (*left*), and the four metal-thiolate cluster of human MT3 in the  $\alpha$ -domain with the metal ions and sulfur atoms from cysteines shown in *cyan* and *yellow* respectively (*right*). (b) NMR structures of the *Notothenia coriiceps* Cd<sub>7</sub>MT-nc with the  $\alpha$ -domain (PDB: 1m0g) (*left*) and the  $\beta$ -domain (PDB: 1m0j) (*right*). The metal-thiolate clusters are also shown with the cadmium ions shown as spheres in *sky blue* and sulfurs as sticks in *yellow*. (c) Solution structure of bacterial SmtA (PDB: 1jjd) with the zinc ions shown in *pale cyan* as spheres (*left*) and the metal cluster Zn<sub>4</sub>Cys<sub>9</sub>His<sub>2</sub> (*right*) with the zinc ions shown in *pale cyan*, sulfurs in *yellow*, and the nitrogens of histidines in *blue*

conformations of domains, which renders a proper function of the protein [50]. Zinc finger proteins are the most abundant class of zinc-binding proteins that contain conserved cysteines and histidines coordinated to zinc. Intensive structural and functional studies have established the invariance of the  $\beta\beta\alpha$  framework of the Cys<sub>2</sub>His<sub>2</sub> zinc finger module and provided a sound basis for understanding the nature of DNA recognition [40, 51–55]. Diverse structures of zinc fingers also account for their diverse functions such as DNA recognition, RNA packaging,



**Fig. 2** (a) Ribbon diagram of the Gfzf35-DNA complex (PDB: 2kmk) with the zinc ions shown as *gray spheres* and the side-chains of two coordinated histidines and two cysteines shown in sticks and DNA shown in *yellow sticks*. (b) Solution structure of the RNA complex of TIS11d (PDB: 1rgo) with the zinc ions in *gray spheres* coordinated to three cysteines and one histidine and RNA shown as *yellow sticks*. (c) Solution structure of HypA from *Helicobacter pylori* (PDB: 2kdx) with the zinc ions in *gray sphere* coordinated to four cysteine sulfurs

transcriptional activation, regulation of apoptosis, protein folding and assembly, as well as lipid binding [52, 56, 57]. Comprehensive reviews in this area can be found [56–58]. Here, we will only highlight some of the recent studies. Solution structure of Gfi-1 zinc finger 3–5 complex with a 16-mer consensus DNA (Fig. 2a) demonstrated that zinc fingers 3–5 bind into the major groove of the target DNA, reminiscent of canonical Cys<sub>2</sub>His<sub>2</sub> zinc-finger domains, which provide valuable insight into the structure determinants for DNA binding specificity as well as molecular rationales for a naturally occurring mutation that causes acute myeloid leukemia [59]. Poly(ADP-ribose) (PAR) is an important post-translational modification in higher eukaryotes. Solution structures of two PBZ modules (PAR-binding zinc finger) of PNK-like factor (APLF) and the PDB domain of *Drosophila melanogaster* CG1218-PA reveal a novel type of Cys<sub>2</sub>His<sub>2</sub> zinc finger and provide a structural basis for PBZ-PAR recognition. Intriguingly, Cys<sub>2</sub>His<sub>2</sub> zinc coordination of the PBZ modules is structurally and functionally dissimilar from canonical double stranded DNA-binding TFIIIA-type zinc fingers; rather they resemble single-stranded RNA-binding Cys<sub>3</sub>His<sub>1</sub> tandem zinc fingers (TZFs). Both of them lack secondary structures but have rigid backbone conformations as a result of zinc binding [60, 61]. Zinc finger proteins are also able to bind to RNA. The NMR structure of tandem zinc finger (TZF) domain of the protein TIS11d bound to the RNA sequence 5'-UUAUUUAUU-3' (Fig. 2b) reveals a pair of novel Cys<sub>3</sub>His<sub>1</sub> fingers which independently recognizes the four nucleotide sequence UAUU and the sequence specificity in RNA recognition is achieved by a network of intermolecular hydrogen bonds [62]. This structure provides insights into RNA-binding function of this family of Cys<sub>3</sub>His<sub>1</sub> zinc finger proteins [62]. The Cys<sub>3</sub>His<sub>1</sub> zinc finger motif is also found in the structure of SAP30 polypeptide of the Sin3 corepressor complex which adopts a novel fold comprising two  $\beta$ -strands and two  $\alpha$ -helices with the zinc organized center. Such a structure may also function as a double-stranded DNA-binding motif [63]. The zinc finger CW (zf-CW) domain

with a motif of about 60 residues is frequently found in proteins involved in epigenetic regulation. Interestingly, NMR structure of human zf-CW domain and PWWP domain containing proteins reveal a new fold in which a zinc is coordinated tetrahedrally by four conserved Cys residues [64]. Such a structure partially resembles the plant homeo domain (PHD) finger bound to the histone tail, implicating a similar function of zf-CW domain [64]. This kind of Cys<sub>4</sub> motif is widely found in other metalloproteins such as [NiFe] hydrogenases accessory protein HypA. Solution structure of HypA from *H. pylori* (Fig. 2c) showed that zinc coordinated to four cysteines donated from loops and no apparent secondary structure found in the zinc-domain [15]. The X-ray structure of HypA from *Thermococcus kodakaraensis* KOD1 further confirmed such a zinc coordination sphere [65].

## 2.2 Utilization of Chemical Shifts to Generate Structures

Protein NMR chemical shifts are highly sensitive to local structure and reflect a wide array of structure factors including backbone and side-chain conformation, secondary structure, hydrogen bonds, and the orientation/position of aromatic rings. Chemical shift data can be used in conjunction with protein sequence information and reasonable force field to generate 3D structure models using the method of CHEMSHIRE or CS-ROSETTA [66–68]. The Chemical-Shift-ROSETTA (CS-ROSETTA) is a robust protocol available for de novo protein structure generation. The method uses experimental chemical shifts of <sup>13</sup>C<sup>α</sup>, <sup>13</sup>C<sup>β</sup>, <sup>13</sup>C<sup>γ</sup>, <sup>15</sup>N, <sup>1</sup>H<sup>α</sup>, and <sup>1</sup>H<sup>N</sup> as an input to select polypeptide fragments in existing protein structures (e.g., PDB data bank) in conjunction with the standard ROSETTA Monte Carlo fragment assembly and energy minimization protocol [67, 68]. The CS-ROSETTA has been further combined with CYANA using unassigned NOESY data to direct Rosetta trajectories toward the native structure and produces a more accurate models than CS-ROSEAAR alone [69]. Moreover, chemical shifts have been further extended in determination of protein–protein complex structures via the CamDock method [70]. The method that utilizes chemical shifts to generate structures may potentially provide a new direction for high-throughput NMR structure determination of proteins including metalloproteins, although such a method has not yet been applied in metalloproteins so far.

## 3 Identification of Metal Coordination

### 3.1 Homonuclear and Heteronuclear Metal NMR Spectroscopy

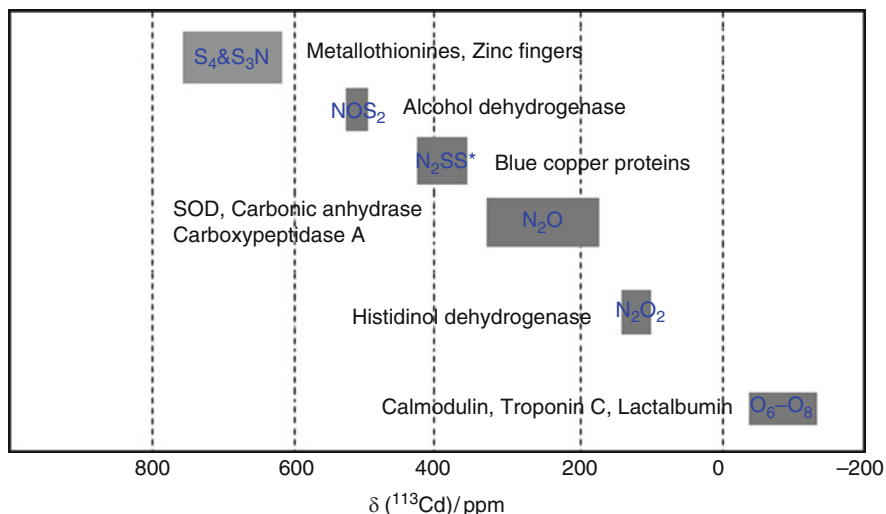
Metalloprotein functionality depends on subtle interaction between properties of the metal ion, dictated by its coordination chemistry. Our present knowledge in terms of structure-function of metalloproteins, in particular the role of metal ions

involved, varies considerably from protein to protein. Therefore, identification of metal coordination parameters including numbers and types of ligands and metal-ligand geometry and mapping the structural and dynamic changes upon metal binding are of significance towards understanding biological functions of metalloproteins. Today, NMR spectroscopy is one of the leading techniques for this purpose. Applicability of homonuclear metal NMR and heteronuclear  $^1\text{H}$ -metal HMQC to monitor directly protein-metal interactions rely greatly on the properties of the nuclei. Some of the metal (e.g.,  $^{113}\text{Cd}$ ) NMR has been used extensively to identify coordination spheres of the metal ions. Moreover, the coupling constants between the NMR active metals and nuclei of the protein provide insight into the identity and geometry of the metal ligands [71]. Many metal (e.g.,  $^{43}\text{Ca}$  and  $^{67}\text{Zn}$ ) NMRs are less powerful and hardly used owing to the fact that these nuclei have the spin quantum number  $I$  greater than  $1/2$  which leads to lower sensitivity, poor resolution, and broadening due to large quadrupolar moments although ultrahigh fields improve it. Several reviews have systematically summarized the application of heteronuclear NMR spectroscopy in biological and medicinal chemistry as well as in the study of metalloproteins [23, 71–73]. Here, we will highlight some of the recent progresses as a snapshot of using metal NMR to identify metal coordination.

Cadmium is one of the most widely used metal nuclei for probing metal-protein interactions, despite its toxic properties. It has two NMR active nuclei,  $^{113}\text{Cd}$  and  $^{111}\text{Cd}$  (spins of  $1/2$ ), with the former being slightly more sensitive and therefore usually used as a preferred nucleus. At natural abundance, the sensitivity of  $^{113}\text{Cd}$  is very low (ca. 7.6-fold of  $^{13}\text{C}$ ), and therefore isotopic enrichment (ca. 96%) of  $^{113}\text{Cd}$  is usually needed to ensure reasonable quality of spectra to be acquired in a relatively short period of time (a few hours for ca. 0.5 mM samples).  $^{113}\text{Cd}$  or  $^1\text{H}$ - $^{113}\text{Cd}$  NMR spectroscopy has been utilized in the study of a variety of metalloproteins where the native  $\text{Zn}^{2+}$ ,  $\text{Ca}^{2+}$ ,  $\text{Mg}^{2+}$ ,  $\text{Mn}^{2+}$ ,  $\text{Fe}^{2+}$ , and  $\text{Cu}^{2+}$  can be substituted by  $^{113}\text{Cd}$ , given that the adaptable ligand coordination number and geometry of  $\text{Cd}^{2+}$  is similar to  $\text{Zn}^{2+}$  and the ionic radius of  $\text{Cd}^{2+}$  (0.97 Å) is similar to that of  $\text{Ca}^{2+}$  (0.9 Å) [27, 74–77]. Moreover, the substitution of the native zinc from metalloenzymes and DNA-binding proteins by cadmium caused almost no changes in their structures and functions [78, 79].

$^{113}\text{Cd}$  chemical shifts are very sensitive to the nature, number, and geometric arrangement of the coordinated ligands [71], as shown in Fig. 3. Such wide chemical shift dispersion not only provides information about the types and numbers of ligand at a particular metal site, but also discriminates multiple metal sites with identical ligand coordination environments.  $^{113}\text{Cd}$  NMR and  $^1\text{H}$ - $^{113}\text{Cd}$  HMQC have been employed exclusively in identification of metal-thiolate clusters in a family of small proteins, e.g., metallothionines [27, 48, 71, 80]. Both homonuclear 1D  $^{113}\text{Cd}$  decoupling studies (Fig. 4a) and 2D  $^{113}\text{Cd}$ - $^{113}\text{Cd}$  COSY (Fig. 4b) of  $^{113}\text{Cd}_7$ -MTs established the existence of two metal-thiolate clusters in this protein, while  $^1\text{H}$ - $^{113}\text{Cd}$  HMQC (Fig. 4c) was used to identify sequence-specific cysteine-cadmium coordination bonds. The chemical shift patterns for the two clusters  $\text{Cd}_3\text{Cys}_9$  and  $\text{Cd}_4\text{Cys}_{12}$  of human MT3 as shown in Fig. 4a showed seven resonances at analogous positions compared with MT1/2 with chemical shift ranging from 600 to 690 ppm [27].

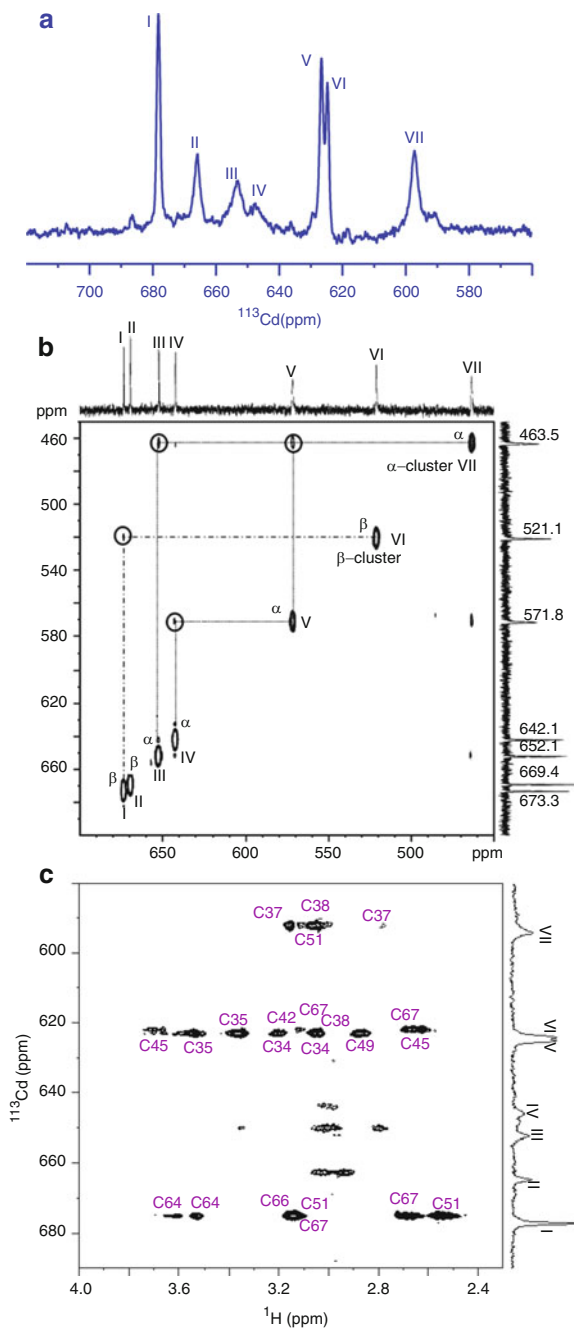




**Fig. 3** Chemical shifts of  $^{113}\text{Cd}$  for structurally characterized  $^{113}\text{Cd}$ -substituted metalloproteins relative to external 0.1 M  $[\text{Cd}(\text{ClO}_4)]$ . The chemical shift positions are represented by gray bars and coordinating atoms are highlighted in blue with typical proteins listed besides. Here S represents sulfur from cysteine, S\* represents sulfur from methionine, O represents oxygen from carboxylate or water, and N represents nitrogen from histidine [71]

However, the resonances from the  $\alpha$ -domain (I, V, VI, and VII) are slightly shifted, particularly for resonance VII, probably attributable to the hexa-peptide insertion in this domain. The homonuclear 2D  $^{113}\text{Cd}$ - $^{113}\text{Cd}$ , Fig. 4b, clearly shows the correlation of cadmium signals which confirms the existence of two Cd-thiolate clusters [81]. The 2D  $^1\text{H}$ - $^{113}\text{Cd}$  HMQC as shown in Fig. 4c is normally used to obtain detailed metal-thiolate connectivity within each of these clusters [27]. Recently, metallothionines from different species, such as sea mussel *Mytilus galloprovincialis* ( $\text{Cd}_7\text{MT10}$ ) [81] and blue crab *Callinectes sapidus* (MTC) [82], have been studied by combined use of  $^{113}\text{Cd}$  NMR,  $^{113}\text{Cd}$ - $^{113}\text{Cd}$  COSY, and  $^1\text{H}$ - $^{113}\text{Cd}$  HMQC. Both proteins have 21 cysteine residues with the position of cysteines distinct from mammalian MTs. The unique structure and dynamic features of the metal-thiolate cluster in these proteins are obviously seen from their distinct NMR parameters of  $\text{Cd}_7\text{MTC}$  [81, 82].  $^{111}\text{Cd}$  NMR has also been applied in studies of His-containing metallothionine, e.g., SmtA. The metal cluster of  $\text{CdS}_4$  and  $\text{CdN}_2\text{S}_2$  were identified [83], and  $^1\text{H}$ - $^{111}\text{Cd}$  HMQC of  $\text{Cd}_7\text{SmtA}$  unequivocally demonstrated couplings of two  $\text{Cd}^{2+}$  to both  $\text{H}^{\epsilon 1}$  or  $\text{H}^{\delta 2}$  protons of two histidine residues [41]. Apart from metallothionines, cadmium NMR has also been employed to identify metal coordination environments in various metalloproteins, such as zinc finger [40] and  $[\text{NiFe}]$  hydrogenase accessory protein HypA [15], where, in the latter case, zinc ions (substituted by  $^{113}\text{Cd}$ ) are coordinated to four cysteine side-chains tetrahedrally. Moreover, it was also used to investigate major zinc binding sites on human albumin [74, 84]. The chemical shifts of  $^{111}\text{Cd}$  of human albumin

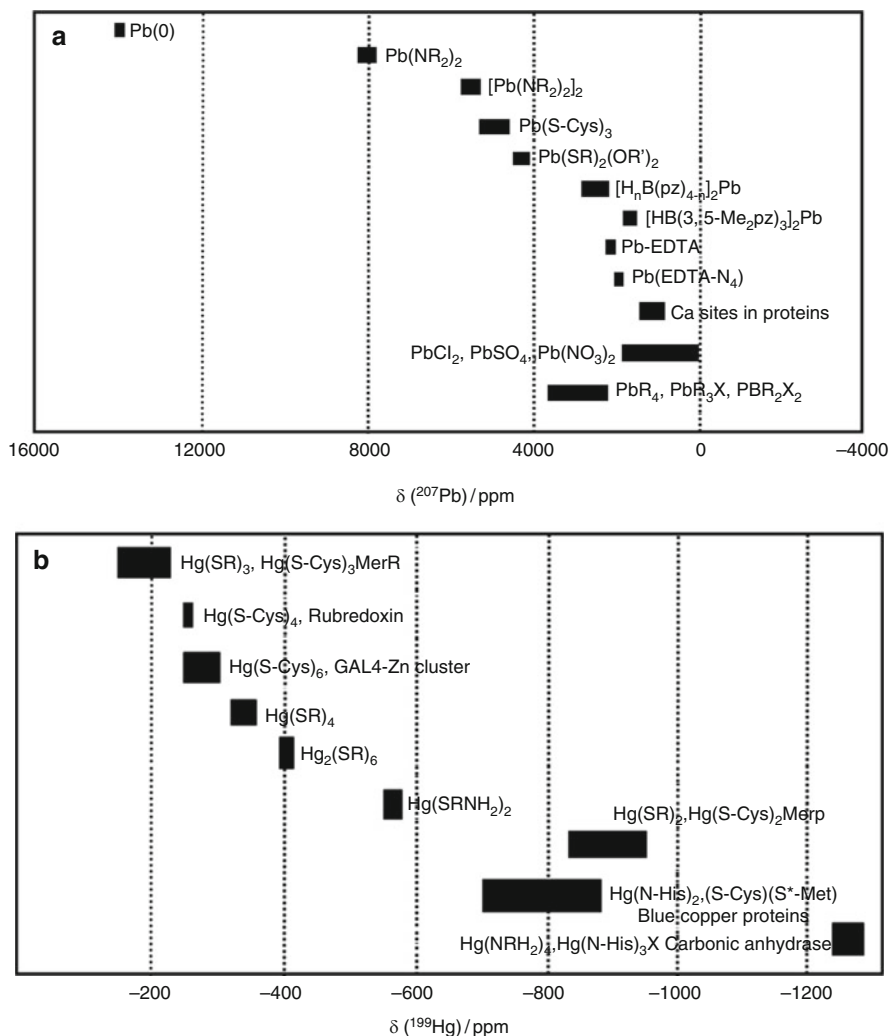
**Fig. 4** (a) The 500-MHz NMR spectrum of human  $^{113}\text{Cd}_7\text{-MT3}$  in 15 mM phosphate buffer, pH 7.3 and 310 K. Resonances I, V, VI, and VII are assigned to the  $\alpha$ -domain and II, III, and IV to the  $\beta$ -domain [27]. (b) 2D  $^{113}\text{Cd}$ - $^{113}\text{Cd}$  COSY NMR of  $\text{Cd}_7\text{-MT10}$  in 17 mM Tris- $\text{d}_{11}$  buffer, pH 7.0, 298 K with the one-dimensional proton-decoupled  $^{113}\text{Cd}$  NMR spectrum at the top. The cadmium connectivity in the  $\alpha$ - and  $\beta$ -domains is also shown [81]. (c) The 500 MHz two-dimensional  $^1\text{H}$ - $^{113}\text{Cd}$  HMQC spectrum of human  $^{113}\text{Cd}_7\text{-MT3}$  in 15 mM phosphate buffer with  $^3J$  ( $^1\text{H}$ ,  $^{113}\text{Cd}$ ) of 30 Hz. The assignments were also labeled with one-letter amino acid code [27]



(ca. 130 and 30 ppm) in combination with other techniques (EXAFS and mutagenesis studies) clearly demonstrated that the major zinc is a five-coordinate site with residues of histidine and aspartate.

Similarly, heteronuclear magnetic resonance spectroscopy with nuclei of spins 1/2 (e.g.,  $^{109}\text{Ag}/^{107}\text{Ag}$ ,  $^{207}\text{Pb}$ , and  $^{199}\text{Hg}$ ) has also been extensively employed in studies of metal coordination environments or the active site structures of metalloproteins.  $^{109}\text{Ag}/^{107}\text{Ag}$  NMR has so far found limited application in biological systems due to low sensitivity and extremely long spin-lattice relaxation times. The only known metalloprotein studied by 2D  $^1\text{H}$ - $^{109}\text{Ag}$  HMQC spectroscopy to date is the silver-substituted yeast metallothioneine [85]. The toxic lead ( $^{207}\text{Pb}$ ), although not directly biological relevant, is an excellent substitute for  $\text{Ca}^{2+}$  in calcium-binding, allowing them to retain similar structures and function.  $^{207}\text{Pb}$  has a moderate resonance frequency, a vast chemical shift range (over thousands), and potentially large spin-spin coupling to neighboring nuclei which limits its application in studies of metalloproteins. Although  $^{207}\text{Pb}$  has been used to study active sites of model compounds or peptides of metalloenzymes [86, 87], there appears to be only one application using  $^{207}\text{Pb}$  as a probe to study  $\text{Pb}^{2+}$  binding to the  $\text{Ca}^{2+}$  site of calcium-binding proteins including calmodulin (CaM) [88]. Binding of  $^{207}\text{Pb}$  to both carp and pike parvalbumins gave rise to two  $^{207}\text{Pb}$  signals from 750 to 1,260 ppm downfield relative to aqueous  $[\text{Pb}(\text{NO}_3)_2]$ . Similarly, four  $^{207}\text{Pb}$  signals, which fall in the same chemical shift window, could be observed for CaM. Both  $^{207}\text{Pb}$  and  $^1\text{H}$  have demonstrated that  $\text{Pb}^{2+}$  binds to all four sites simultaneously, in contrast to the behavior of the protein in the presence of  $\text{Ca}^{2+}$  [88]. The large chemical shift dispersion and remarkable sensitivity to the chemical environment of  $^{207}\text{Pb}$  signals promote extensive studies using model complexes or peptides on mining the relationship of chemical shifts of  $^{207}\text{Pb}$  with the numbers and types of coordinating ligands [86, 87, 89]. A recent study again demonstrated that chemical shifts of  $^{207}\text{Pb}$  can be used to discriminate between  $\text{PbS}_3$  (from Cys of thiolate-rich peptides) with other ligands such as  $\text{PbS}_3\text{O}$  and other O, S, and N donor ligands [89]. Therefore, based on  $^{207}\text{Pb}$  chemical shift maps, Fig. 5a, both homonuclear and heteronuclear  $^{207}\text{Pb}$  spectroscopy should provide a useful tool for investigation of  $\text{Pb}^{2+}$  coordination in more complex biological systems.

In spite of being widely known as a protein-modifying agent, the toxic metal  $\text{Hg}^{2+}$  ion can be used as a probe to substitute native metal ions for several metalloenzymes in a manner that preserves catalytic activity. Mercury has two NMR-active isotopes,  $^{199}\text{Hg}$  and  $^{201}\text{Hg}$ , with the latter having the spin of 1/2 and a natural abundance of 16.84%.  $^{199}\text{Hg}$  exhibits several favorable NMR properties for structural and functional study, such as large coupling constants, a wide range of chemical shift dispersions (ca. 5,000 ppm), and a relative sensitivity of 5.4 times that of  $^{13}\text{C}$  and 8 times that of  $^{113}\text{Cd}$  for an equal number of nuclei. Therefore,  $^{199}\text{Hg}$  NMR techniques ( $^{199}\text{Hg}$  and  $^1\text{H}$ - $^{199}\text{Hg}$  HMQC) have been used as useful tools to probe the metal coordination environment in biological relevant complexes [91, 92] as well as regulatory proteins, copper enzymes, and zinc transcription factor complexes as large as 50 kDa [90, 93–96]. The chemical shift of  $^{199}\text{Hg}$  is very sensitive to the primary coordination spheres including numbers, types of ligands, as well as



**Fig. 5** (a)  $^{207}\text{Pb}$  chemical shifts of various lead model complexes and  $^{207}\text{Pb}$ -substituted proteins relative to external 1 M  $[\text{Pb}(\text{NO}_3)_2]$  in 99.9%  $\text{D}_2\text{O}$   $\text{pH}^* 3.3$ . The pz represents the pyrazolyl ring and S-Cys represents a thiolate from cysteine (adapted from [87, 89]). (b)  $^{199}\text{Hg}$  chemical shifts of aliphatic amine/thiol model complexes and  $^{199}\text{Hg}$ -substituted proteins relative to  $\text{Hg}(\text{CH}_3)_2$  at 298 K. The  $\text{NRH}_2$  represents a primary amine and SR represents a thiolate, and the coordination environments of the Hg proteins include His imidazole nitrogen, Cys thiolate S-Cys, and Met thioether S\*-Met [90]

coordination geometry. For examples,  $^{199}\text{Hg}$  bound blue copper proteins gave rise to signals at ca.  $-880$  ppm for azurin,  $-749$  ppm for plastocyanin, and  $-706$  ppm for rusticyanin [93, 95]. These blue proteins have very similar coordination environments, e.g., two histidines and one cysteine and the fourth weakly

associated ligand (Met or others) which is slightly different for these proteins. Nevertheless, such a subtle difference can be faithfully reflected by the chemical shifts of  $^{199}\text{Hg}$ . The chemical shift map of  $^{199}\text{Hg}$  can be derived based on various model complexes and proteins; Fig. 5b. The large chemical shift dispersion for  $^{199}\text{Hg}$  allows clear differentiation between a variety of  $M(\text{SR})_n$  environments. Given that  $\text{Hg}^{2+}$  is readily exchanged for the native metal ion in many copper, zinc, and iron metalloproteins [96],  $^{199}\text{Hg}$  NMR methods can play an important role in structural, spectroscopic, and chemical studies of metalloproteins and metal binding domains where the tertiary structure of the folded proteins dictates the geometry of the metal ion.

### 3.2 Chemical Shift Perturbation

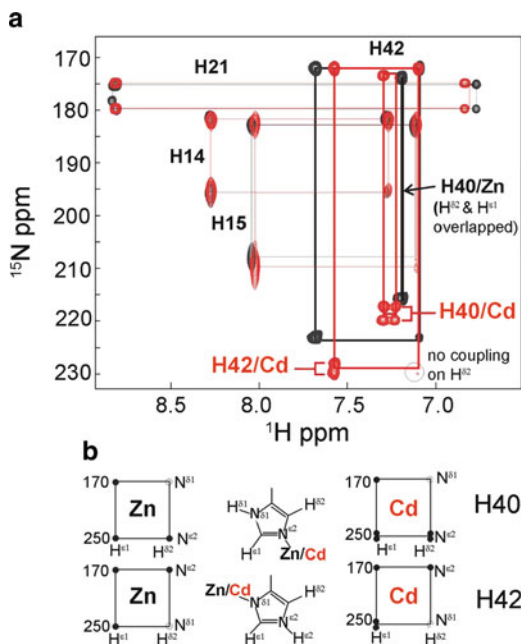
Protein NMR chemical shift is highly sensitive to the exact environment of the atom and can provide valuable insights into structural features including metal ligation. For examples, chemical shifts of  $^{13}\text{C}^\beta$  for the zinc bound cysteines (ca. 34 ppm) are significantly downfield shifted relative to those of non-metal bound cysteines (ca. 27 ppm) [97]. Such an index has often been used to discriminate zinc bound cysteine residues [15, 98]. The chemical shift perturbations upon metal ion binding can be used to estimate the affinity, stoichiometry, and kinetics of metal binding, and, moreover, they can be used to identify metal coordination environments. This approach is usually denoted as chemical shift mapping, which has been widely used to study protein–protein interaction [99, 100].  $^1\text{H}$  NMR and  $^1\text{H}$ - $^1\text{H}$  TOCSY have been used to identify the types of residues binding to metal ions for unlabeled small metalloproteins [101–103]. However, two dimensional HSQC, especially  $^1\text{H}$ - $^{15}\text{N}$  HSQC, is often employed in studies of metal coordination environments in metalloproteins due to the fact that it is well resolved in comparison with 2D  $^1\text{H}$ - $^{13}\text{C}$  HSQC. The identity of each cross-peak in the 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC spectra is assigned based on a series of triple resonance experiments. Binding of metal ions would lead to the appearance of new peaks or disappearance of original peaks depending on the exchange rates of the apo- and metal-bound forms on the NMR time scales [99, 104]. The chemical shifts perturbation (CSP) can be followed in titration experiments, where the concentration of diamagnetic metal ion is increased gradually. When heteronuclear data are available, the binding site is usually predicated by the combined chemical shift perturbation  $\Delta\delta_{\text{comb}}$ , which has been shown to be a more reliable approach to evaluate titration data quantitatively [105]. Although several approaches are available to obtain this value as summarized previously [105], in practice, the weighting for  $^1\text{H}$  and  $^{15}\text{N}$  is considered to be the same, and the  $\Delta\delta_{\text{comb}}$  is usually quantified by the following equation:

$$\Delta\delta_{\text{comb}} = \sqrt{\frac{\Delta\delta_{\text{HN}}^2 + \frac{1}{25}\Delta\delta_{\text{N}}^2}{2}}$$

Chemical shift mapping provides information about the location of the metal binding sites and has been used to identify the metal coordination for numerous metalloproteins [15, 106–108]. For example, addition of  $\text{Cu}^+$  to human Cox17 induced significant chemical shift variations over residues Lys20 and Ala24, and the appearance of the NH signals of Cys22–Cys23, which is thought to serve as the  $\text{Cu}^+$  binding motif [107]. However, this approach usually has to be used in combination with other physical techniques or biological approaches, e.g., mutagenesis to specify the metal binding sites, since the chemical shift perturbations mainly stem from either direct binding or conformational changes caused by the metal ions. Binding of  $\text{Ni}^{2+}$  to *H. pylori* HypA led to disappearance of signals of Glu3 and Asp40 in the 2D  $^1\text{H}$ - $^{15}\text{N}$  HSQC spectrum. When combining with mutagenesis study, side-chain 2D  $^1\text{H}$ - $^{15}\text{N}$  HMQC, UV absorption, as well as CD, it was proposed that  $\text{Ni}^{2+}$  coordinates with His2 (side-chain  $\text{N}^\delta$ ), His2 (backbone N), and the backbone nitrogens of Glu3, and Asp40 with a square-planar geometry [15]. Such a binding also induced structural changes which were thought to be important for its downstream receptor's recognition [15].

Histidine often serves as a metal binding ligand in metalloproteins and can provide both backbone and side nitrogens to coordinate with metal ions such as  $\text{Zn}^{2+}$  and  $\text{Ni}^{2+}$ . It has been shown that different tautomeric forms of histidine imadazole rings have different, distinguishable signal patterns in a long-range 2D  $^1\text{H}$ - $^{15}\text{N}$  HMQC spectrum [109], metal coordination often causes recognizable changes in the NMR spectrum of histidine side-chains, and imadazole nitrogen atoms involved in direct metal coordination have specific chemical shift [15, 40, 98, 110]. This technique has been extensively used to identify  $\text{Zn}^{2+}$  binding. The chemical shifts observed for the unprotonated imadazole nitrogen atoms of a zinc finger domain Hdm2(429–491) appeared at ca. 215 ppm. The cross-peak pattern in  $^1\text{H}$ - $^{15}\text{N}$  HMQC spectrum (with  $^2\text{J}_{\text{HN}}$ ) of zinc bound Hdm2 (429–491) showed that His452 and His457 assumed different tautomeric forms with the former being  $\text{N}^{\epsilon 2}$ -protonated and the latter  $\text{N}^{\delta 1}$ -protonated with chemical shifts around 170 ppm, which demonstrated that  $\text{Zn}^{2+}$  coordinated to both His452 and His457 via the  $\text{N}^{\epsilon 2}$  and  $\text{N}^{\delta 1}$  respectively [98]. More interestingly, a comparison of the 2D  $^1\text{H}$ - $^{15}\text{N}$  HMQC spectra of  $\text{Zn}^{2+}$ -bound proteins with  $^{113}\text{Cd}^{2+}$ -bound proteins, as shown in Fig. 6, allows one to observe the  $^{15}\text{N}$ - $^{113}\text{Cd}$  coupling, which assists identification of overlapping of histidine side-chains that bind to metal ions [40]. The HMQC spectrum of the zinc-bound domain, shown as black in Fig. 6a, clearly shows two of the  $^{15}\text{N}$  resonances His42 and His40 shifted downfield as a result of zinc coordination. Based on the pattern of the cross-peaks, His42 is in the  $\epsilon$  tautomeric state indicative of zinc coordination to the  $\text{N}^{\delta 1}$  of His42; Fig. 6b. Such a method cannot be used to assign the nitrogen atom of His40 due to overlapping of the  $\text{H}^{\delta 2}$  and  $\text{H}^{\epsilon 1}$  resonances of His40. The HMQC spectrum of  $^{113}\text{Cd}$  bound protein, shown as red in Fig. 6a, clearly shows coupling between  $^{113}\text{Cd}$  and the  $\text{N}^{\delta 1}$  of His42 observed on  $\text{H}^{\epsilon 1}$ , suggesting a covalent bond between them. Importantly, the pattern of the connectivities and the coupling observed to both  $\text{H}^{\delta 2}$  and  $\text{H}^{\epsilon 1}$  from the  $^{113}\text{Cd}$  allows the unambiguous assignment of the  $\delta$  tautomer for His40, with the metal coordination via the  $\text{N}^{\epsilon 2}$  of the side-chain; Fig. 6b [40]. The side-chain 2D  $^1\text{H}$ - $^{15}\text{N}$  HMQC spectrum of histidine has also been used to identify  $\text{Ni}^{2+}$  binding

**Fig. 6** (a) Comparison of the 500-MHz  $^1\text{H}$ - $^{15}\text{N}$  HMQC spectrum of Zn-bound ZZ domain in *black* with the  $^{113}\text{Cd}$ -bound form in *red*. The coupling of  $^{113}\text{Cd}$  to  $\text{N}^{\text{e}2}$  is observed on  $\text{H}^{\text{e}2}$  and  $\text{H}^{\text{e}2}$  for His40; whereas  $^{113}\text{Cd}$  coupling to  $\text{N}^{\delta 1}$  is observed on  $\text{H}^{\delta 1}$  for His 42. (b) Diagrams showing the connectivities of two zinc-ligand histidine residues H40 and H42 based on the above HMQC spectrum. The coupling of  $^{113}\text{Cd}$  to the attached  $^{15}\text{N}$  is useful for identification of metal bound histidine residues and the location of the metal coordination on the histidine rings (adapted from [40])



atoms in *H. pylori* HypA and significant downfield shift of His2 upon  $\text{Ni}^{2+}$  binding indicates that  $\text{Ni}^{2+}$  binds to His2 through the  $\text{N}^{\delta 1}$  atom [15].

## 4 NMR in Studies of Paramagnetic Metalloproteins

Metalloproteins represent a large share of a proteome. A large number of them contain paramagnetic metal ions, which possess unpaired electrons. The presence of a paramagnetic center causes pronounced effects in NMR spectra and reduces dramatically the intensity of NOEs and the efficiency in the transfer of scalar couplings both in homonuclear and heteronuclear experiments, which hampers spectrum assignment and structural determination through standard approaches. However, with the advances in novel experiment design and development of software protocol in recent years, the presence of paramagnetic centers has been used as a precious source of structural information [19, 20]. Extensive reviews regarding this topic can be referred to [19–22, 111–113]. Here we will give a very brief description of paramagnetism-based restraints as well as application of  $^{13}\text{C}$ -detected experiments in the structural examination of paramagnetic metalloproteins.

### 4.1 Paramagnetism-Based Structural Restraints

In paramagnetic systems, where paramagnetic metal ions are either intrinsic or extrinsic, there are three NMR experimental observables that yield long-range

structural information, i.e., paramagnetic resonance enhancement (PRE), pseudocontact shifts (PCS), and RDC induced by anisotropic paramagnetic centers. In addition, cross-correlated relaxation (CCR) effects between anisotropic paramagnetic centers and anisotropic parameters of the nuclear spins can also be exploited to generate long-range restraints [19, 111, 112]. Paramagnetic centers with isotropic electron spin distribution ( $\text{Mn}^{2+}$  and  $\text{Gd}^{3+}$ ) produce large PREs due to slow electron relaxation. In contrast, paramagnetic centers with anisotropic electron spin distribution for most paramagnetic metal ions, including most of the lanthanides, create all four long-range paramagnetic effects, which contain rich structural information [114]. Here, we will focus on PRE and PCS and their applications. Information about RDC can be found in this book.

The PRE arises from magnetic dipolar interactions between a nucleus and the unpaired electrons of the paramagnetic center, resulting in an increase in nuclear relaxation rates. In contrast to NOE, where the effects are limited to short range interaction ( $<6 \text{ \AA}$ ), the PRE effects are relatively large and can be detected up to  $35 \text{ \AA}$  owing to the large magnetic moment of an unpaired electron. There are two mechanisms, i.e., the Solomon mechanism and the Curie spin mechanism, that give the PREs, with the former being predominant for slowly tumbling molecules with long lifetimes of electronic spin state (such as  $\text{Mn}^{2+}$  and  $\text{Gd}^{3+}$ ). However, the Curie relaxation becomes important when the electronic relaxation is much faster than the rotational tumbling of the molecules, which is the case for the majority of paramagnetic metal ions. Theoretical and experimental aspect of PRE as well as its application in studies of structures of proteins and protein–protein complexes can be found in recent reviews [114–116].

At high magnetic fields (over 500 MHz for  $^1\text{H}$  frequency), the PRE rate,  $\Gamma_2$ , which arises from the dipole–dipole interaction between a nucleus and unpaired electrons with an isotopic  $g$ -tensor, is conventionally calculated by the Solomon–Bloembergen (SB) equation:

$$\Gamma_2 = \frac{1}{15} \left( \frac{\mu_0}{4\pi} \right)^2 \gamma_I^2 g^2 \mu_B^2 S(S+1) r^{-6} \left\{ 4\tau_c + \frac{3\tau_c}{1 + (\omega_H \tau_c)^2} \right\}$$

where  $r$  is the distance between the paramagnetic center and the observed nucleus,  $\mu_0$  the permeability of vacuum,  $\gamma_I$  the nuclear gyromagnetic ratio,  $g$  the electron  $g$ -factor,  $\mu_B$  the electron Bohr magneton,  $S$  the electron spin quantum number, and  $\tau_c$  the PRE total correlation time. In practice,  $\Gamma_2$  is measured as a difference in transverse relaxation rates between the paramagnetic ( $R_{2,\text{para}}$ ) and diamagnetic ( $R_{2,\text{dia}}$ ) states. A two-time point measurement is recommended as a simple approach for obtaining  $\Gamma_2$  rates and their corresponding errors without fitting procedures. In this approach, the  $\Gamma_2$  rates are determined from two time points ( $T = 0$  and  $\Delta T$ ) for transverse relaxation as shown by the following equation [116]:

$$\Gamma_2 = R_{2,\text{para}} - R_{2,\text{dia}} = \frac{1}{T_b - T_a} \ln \frac{I_{\text{dia}}(T_b) I_{\text{para}}(T_a)}{I_{\text{dia}}(T_a) I_{\text{para}}(T_b)}$$



where  $I_{\text{dia}}$  and  $I_{\text{para}}$  are the peak intensities for the diamagnetic and paramagnetic states, respectively. The choice of time points is important to minimize the error. For example, if the range of  $\Gamma_2$  rates is  $0\text{--}75\text{ s}^{-1}$ , a second time point  $T_b$  should be at ca.  $1.15/(R_{2,\text{dia}} + 50)$  s, representing a reasonable choice [115, 116].

For isotropic metal ions such as  $\text{Mn}^{2+}$  and  $\text{Gd}^{3+}$ , the Curie-spin relaxation that could potentially exhibit significant cross-correlation with other relaxation mechanisms is negligible for medium-size macromolecules [116], and the Solomon relaxation is predominant. PRE analysis in such a system is thus simple. The PRE has been used extensively in metalloproteins that possess a rigid intrinsic paramagnetic center [117–119]. Such a strategy has also been extended not only in the NMR structure determination of non metalloproteins [120–122], in which paramagnetic metal ions ( $\text{Mn}^{2+}$  or  $\text{Gd}^{3+}$ ) or nitroxide radicals were conjugated through appropriate chemical modification [123], but also in the characterization of protein–protein/nucleic acid complexes [124–126] and membrane-proteins [127], in particular in transient macromolecular interactions [115, 128–131].

PCSs are precious sources of structure information and are observed only in paramagnetic systems with anisotropic unpaired electrons, e.g.,  $\text{Dy}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Fe}^{3+}$ . The magnitude of the PCS,  $\delta^{\text{PCS}}$ , is calculated using the following equation [132]:

$$\delta^{\text{PCS}} = \frac{1}{12\pi} r^{-3} \left\{ \Delta\chi_{\text{ax}} (3\cos^2\theta - 1) + \frac{3}{2} \Delta\chi_{\text{rh}} \sin^2\theta \cos 2\varphi \right\}$$

$$\Delta\chi_{\text{ax}} = \Delta\chi_{\text{zz}} - \frac{1}{2}(\chi_{\text{xx}} + \chi_{\text{yy}}) \quad \text{and} \quad \Delta\chi_{\text{rh}} = \chi_{\text{xx}} - \chi_{\text{yy}}$$

where  $r$  is the distance between the metal ion and the nuclear spin,  $\theta$  and  $\varphi$  are the angles describing the position of the nuclear spin with respect to the principle axes of the magnetic susceptibility tensor  $\chi$ , and  $\Delta\chi_{\text{ax}}$  and  $\Delta\chi_{\text{rh}}$  are the axial and rhombic components, respectively, of the magnetic susceptibility tensor.

The PCS are manifested by large changes in chemical shifts of the nuclear spins that are exposed to the paramagnetic metal ions and arise from through-space dipolar interactions with rapidly relaxing unpaired electrons. The PCS displays an  $r^{-3}$  distance dependence, in contrast to the  $r^{-6}$  dependence for the PRE, which results in a relatively long distance range for the PCS to be detected (ca.  $40\text{ \AA}$  for  $\text{Dy}^{3+}$ ) [133]. In general, the  $\delta^{\text{PCS}}$  values can be measured after the complete assignment is obtained for the  $^1\text{H}\text{--}^{15}\text{N}$  HSQC spectra of the both the diamagnetic and the paramagnetic samples, and are calculated as the difference between the chemical-shift values observed for the nuclei in a paramagnetic system and in a diamagnetic analog. The  $\delta^{\text{PCS}}$ -derived restraints alone cannot be used to solve the structures. Instead, the PCSs have to be incorporated with NOEs and dihedral-angle restraints to determine structures of proteins or to refine protein structures. The first example of using PCS in a structure refinement was reported on a low-spin  $\text{Fe}^{3+}$  heme protein [134]. Such a strategy has been extended not only in studies of paramagnetic proteins [135, 136], but also in non-metal binding proteins [137, 138], which were labeled by paramagnetic metal ions such as lanthanides [114, 139] or

genetically encoded  $\text{Co}^{2+}$ -binding amino acid [140]. Moreover, PCSs can also be used as restraints in molecular dynamics [141, 142].

## 4.2 $^{13}\text{C}$ -Detected Protonless NMR

Direct detection of heteronuclei, in particular  $^{13}\text{C}$ , offers a valuable alternative to  $^1\text{H}$  detection in the study of biological macromolecules [143–145] as well as paramagnetic proteins [136, 145–147]. The recent development of high magnetic fields as well as the availability of cryogenically cooled probe heads has improved  $^{13}\text{C}$  sensitivity significantly, which has greatly stimulated research using  $^{13}\text{C}$ -detected experiments on enriched samples to study biological macromolecules [148]. The  $^{13}\text{C}$  detection takes advantage of the slower relaxing  $^{13}\text{C}$  spins and overcomes the drawbacks produced by the fast  $^1\text{H}$  transverse relaxation, which leads  $^1\text{H}$  signals to broaden beyond detection limits for large proteins. Such an approach is particularly useful in paramagnetic systems since  $^{13}\text{C}$  direct detection is less affected by the paramagnetic center than  $^1\text{H}$  owing to the lower  $^{13}\text{C}$  gyromagnetic ratio, which decreases the paramagnetic dipolar contributions to its relaxation by a factor of around 16  $(\gamma_{\text{C}}/\gamma_{\text{H}})^2$  [136, 146, 149].

In  $^{13}\text{C}$  direct detection experiments, several approaches were used to achieve “virtual” decoupling (to remove homonuclear one-bond carbon–carbon couplings) such as IPAP schemes (in-phase anti-phase) [150–152], in which two FIDs for each increment are recorded and stored separately, one for in-phase and another for anti-phase, the two components being combined to remove the splitting. An alternative is  $\text{S}^3\text{E}$  schemes [152, 153] (spin-state selective excitation), in which two different experiments are performed with one being absorptive and another dispersive. One or more of these building blocks (IPAP and  $\text{S}^3\text{E}$ ) can be implemented in any experiments based on  $^{13}\text{C}$  direct detection. A set of  $^{13}\text{C}$  based experiments, which can be used for the assignment of backbone and side-chains of  $^{13}\text{C}/^{15}\text{N}$  labeled proteins, is now available and summarized in a recent review [144]. The sequence-specific assignment was achieved by CACO and CANCO, which provide the correlation of each CO to the two neighboring  $\text{C}^\alpha$  nuclei; CACO, CBCACO, and  $^{13}\text{C}$ - $^{13}\text{C}$  TOCSY can provide spin-system assignment [143, 154, 155]. The CON-IPAP experiment is used to correlate backbone nitrogen with CO through the one-bond  $\text{C}'\text{-N}$  coupling [149]. The  $^{13}\text{C}$ - $^{13}\text{C}$  NOESY experiments based on dipole–dipole interaction with longitudinal magnetization transfer represent a valuable alternative to COSY experiments based on scalar couplings to detect C–C one-bond correlation for large macromolecules and paramagnetic metalloproteins [149, 156, 157]. The  $^{13}\text{C}$ - $^{13}\text{C}$  NOESY with direct  $^{13}\text{C}$  detection is exploited as a valuable tool to extend the assignment to side-chains in large molecules such as  $\text{C}'\text{-C}^\beta$  if mixing times is long enough (e.g., 800 ms) [156]. The correlations between nuclei not directly bound and not mediated by spin diffusion in  $^{13}\text{C}$ - $^{13}\text{C}$  NOESY spectra are identified, which would represent a breakthrough of structure determination of large macromolecules in solution by providing distance constraints.

However, the sensitivity for  $^{13}\text{C}$  direct detection is required to be improved before long-range correlations to obtain  $^{13}\text{C}$ – $^{13}\text{C}$  distance constraints can be used in structural characterization of large macromolecules.

$^{13}\text{C}$  direct detection has been successfully applied to paramagnetic proteins, where the contribution to line broadening coming from the paramagnetic center is so large that  $^1\text{H}$  signals around the metal ion are beyond detectable limits [136, 144, 147, 149, 158]. Such a technique can also be used in generation of paramagnetism-based restraints including PCS, PRE, and RDC [148, 159–161]. It has been demonstrated that  $^{13}\text{C}$  directly detected spectra provide an alternative method for the measurement of RDC with precision as good as that from  $^1\text{H}$  detection, but with additional advantage for measuring those broad resonances in  $^1\text{H}$  detection [148]. Direct detection of  $^{13}\text{C}$  intrinsically offers a way to detect resonances close to the metal ion where  $^1\text{H}$  resonances are too broad to be detected. Indeed, with the aid of a  $^{13}\text{C}$  direct detection approach,  $^{13}\text{C}$  resonances as close as 6 Å from the metal ion are detected for CopC, a  $\text{Cu}^{2+}$  binding protein involved in copper homeostasis, whereas no  $^1\text{H}$  resonance can be detected within a sphere of 11 Å from the metal due to fast relaxation caused by paramagnetic  $\text{Cu}^{2+}$  [136]. Incorporation of heteronuclear paramagnetism-based restraints, e.g., PCSs and longitudinal relaxation rate enhancement, allows CopC structures to be resolved with the RMSD of  $\text{Cu}^{2+}$  determined only by the paramagnetism-based constraints of 1.1 Å [136]. The  $^{13}\text{C}$  direct detection technique has also been used for residue-specific assignments of resonances, in particular those near paramagnetic centers (e.g.,  $\text{Ni}^{2+}$  and  $\text{Fe}^{3+}$ ) such as in a 20-kDa Ni-containing enzyme, acireductone dioxygenase (ARD) [162], and oxidized human [2Fe-2S] ferredoxin [146], as well as a 19-kDa  $\text{Fe}^{3+}$  hemophore HasA [147]. In many paramagnetic systems, the longitudinal relaxation rates are influenced to a smaller extent than the transverse relaxation rates. The  $^{13}\text{C}$ – $^{13}\text{C}$  NOESY experiments are therefore a useful approach to overcome the quench of scalar coupling based transfer, in particular for large macromolecules. The use of  $^{13}\text{C}$  direct detected experiments, e.g.,  $^{13}\text{C}$ – $^{13}\text{C}$  COSY,  $^{13}\text{C}$ – $^{13}\text{C}$  NOESY, and  $^{13}\text{C}$ – $^{13}\text{C}$  COCAMQ, allows  $^{13}\text{C}$  signals as close as 4 Å to  $\text{Cu}^{2+}$  to be detected in oxidized monomeric copper/zinc superoxide dismutase (SOD) [149]. The advantage of  $^{13}\text{C}$ – $^{13}\text{C}$  NOESY experiments for higher molecular weights was seen by comparison of the protein SOD [149, 156]. All of the expected  $\text{C}\alpha$ –CO connectivities were detected with higher intensity in the dimeric protein than in the monomeric state. In addition, most of the two bond CO– $\text{C}\beta$  cross-peaks were observable for the dimeric SOD when the long mixing times were used [156]. Interestingly, the intrinsic asymmetry of a  $^{13}\text{C}$ – $^{13}\text{C}$  COSY experiment allows the coordinating residues of paramagnetic metal ions to be identified easily, providing a unique method to distinguish between monodentate and bidentate coordinating side-chain carbonyls [163]. Significantly, the  $^{13}\text{C}$ -based strategy in combination with solid-state NMR led to partial sequence-specific (35%) and side-chain assignments for the iron storage protein, ferritin, a very large protein with a molecular mass of 480 kDa and 24 subunits [164, 165]. The solution  $^{13}\text{C}$ – $^{13}\text{C}$  NOESY spectra for side-chain observation has provided the identification of an iron channel that guides the direction transport of the multimeric

Fe<sup>3+</sup> products from the active site toward the nanocage. The interior of the four-helix bundle is identified as the functional channel based on the observed paramagnetic effects on residues lining the internal face of the four-helix bundle. The NMR data provide a basis for the pathway of iron from the ferrous/dioxygen oxidoreductase site to the central cavity of ferritin [164]. Such studies open new avenues for the application of <sup>13</sup>C direct detection experiments to systems with molecular assemblies larger than 100 kDa.

## 5 Perspectives

Since the first protein solution structure was determined by high resolution NMR spectroscopy about 25 years ago [166], NMR has been established as the only experimental method that provides both structural and dynamical information at atomic resolution close to physiologically relevant conditions. Protein structure determination in living cells has also been achieved recently by in-cell NMR [18]. However, the limitation of this technique in structural studies lies in low sensitivities and poor resolution when the size of macromolecules increases. Some proteins, in particular metalloproteins, might not be stable for a period of time (days or weeks) or have limited solubility. Moreover, new challenges in life science have also promoted development of new NMR methods which will improve sensitivities and reduce acquisition times to fulfil the requirement of characterization of these proteins and their complexes.

Enormous effort has been made to improve NMR instrumentation in terms of experimental sensitivity, which results in availability of high-field magnets, cryogenically cooled probes. In the meantime, tremendous advances in methodology have contributed to an increased interest in the study of molecular systems of increasing size and complexity. The introduction of a nonlinear sampling scheme (instead of conventional uniform sampling) allows a very fast acquisition of multi-dimensional NMR [167–169]. Many schemes have been developed to reduce the spectral dimensionality and thus to speed up the experiments, which enables quick assignment of large proteins [170, 171]. The examples include the G-matrix Fourier transform (“GFT”) NMR approach where sub-spectra from joint sampling of indirect dimensions are linearly recombined and analyzed [172]. In the projection reconstruction (“PR”) method, the corresponding full-dimensional spectrum is reconstructed [171, 173, 174]. Moreover, various ultrafast NMR techniques including SOFAST/BEST NMR [175, 176] and Hadmard NMR [177] are also available for studying biomolecules even in real time. All these new schemes deliver appreciable improvement in the speed of data acquisition and show promise for speeding up multidimensional NMR of normal size proteins [170, 178, 179] and very large proteins [180, 181] as well as sequence assignment for intrinsically unstructured proteins [178].

<sup>13</sup>C NMR spectroscopy is emerging as a powerful tool to complement <sup>1</sup>H NMR spectroscopy in the investigation of biomolecules, in particular for large molecules and paramagnetic metalloproteins and also for the study of short-lived molecules

[182]. However, data acquisition time is rather long even for samples with high concentrations (ca. mM). Implementation of fast NMR methods such as non-uniform sampling in the indirect dimension significantly reduced experimental times [183]. Such a strategy will open new avenues to applications of  $^{13}\text{C}$  NMR to take advantage of the favorable heteronuclear chemical shift dispersion in biological systems, especially for systems of increasing size.

Chemical shifts of selective nuclei ( $^{113/111}\text{Cd}$ ,  $^{209}\text{Pb}$ , and  $^{195}\text{Pt}$ ) are sensitive towards types of ligands (N, O, and S), numbers, and geometries, and will continue to play a role in characterization of metal–protein local coordination. The ultra-field NMR facilities direct observation of biologically important metal ions with half integer, quadrupolar nuclei (e.g.,  $^{67}\text{Zn}$  and  $^{25}\text{Mg}$ ) [184, 185].

In spite of the availability of all these new techniques, their applications to metalloproteins are currently sparse. There is an urgent need to promote these advanced techniques in the scientific community through introducing integrated software packages for experimental set-up, data processing, and analysis. This will enable protein chemists and bioinorganic chemists, who are not NMR experts, to employ the new techniques in their research. The combination of fast NMR techniques,  $^{13}\text{C}$  directly detected NMR with paramagnetic NMR, will offer great possibilities in tackling new challenges in life science and will open new avenues for NMR spectroscopy to be utilized not only in the characterization of single biomolecules, e.g., structural and dynamical studies of proteins/metalloproteins and paramagnetic proteins, short-lived macromolecules, and intrinsically unstructured proteins, but also in the investigation of more complex systems to give an integrated view of interacting molecular networks. In particular, ultra-fast NMR opens up new perspectives for NMR structural investigations of unstable protein/metalloprotein samples and real-time site-resolved studies of protein kinetics or monitoring folding/unfolding processes of proteins/metalloproteins caused by ligand or metal binding/release.

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# Recent Developments in $^{15}\text{N}$ NMR Relaxation Studies that Probe Protein Backbone Dynamics

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**Abstract** Nuclear Magnetic Resonance (NMR) relaxation is a powerful technique that provides information about internal dynamics associated with configurational energetics in proteins, as well as site-specific information involved in conformational equilibria. In particular,  $^{15}\text{N}$  relaxation is a useful probe to characterize overall and internal backbone dynamics of proteins because the relaxation mainly reflects reorientational motion of the N–H bond vector. Over the past 20 years, experiments and protocols for analysis of  $^{15}\text{N}$   $R_1$ ,  $R_2$ , and the heteronuclear  $^{15}\text{N}$ – $\{^1\text{H}\}$  NOE data have been well established. The development of these methods has kept pace with the increase in the available static-magnetic field strength, providing dynamic parameters optimized from data fitting at multiple field strengths. Using these methodological advances, correlation times for global tumbling and order parameters and correlation times for internal motions of many proteins have been determined. More recently, transverse relaxation dispersion experiments have extended the range of NMR relaxation studies to the milli- to microsecond time scale, and have provided quantitative information about functional conformational exchange in proteins. Here, we present an overview of recent advances in  $^{15}\text{N}$  relaxation experiments to characterize protein backbone dynamics.

**Keywords** Dispersion · Dynamics · NMR · Protein · Relaxation

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## 1 NMR Relaxation to Detect Protein Dynamics

In this chapter, we describe mainly  $^{15}\text{N}$  relaxation experiments to characterize protein backbone dynamics in solution. NMR spin-relaxation is a phenomenon in which perturbed magnetization is restored to statistical equilibrium by random fluctuations of local magnetic fields. The major local magnetic fields in diamagnetic proteins are generated by the amide  $^1\text{H}\text{-}^{15}\text{N}$  dipolar interaction and  $^{15}\text{N}$  chemical shift anisotropy (CSA). Interchange among different chemical shift environments by chemical exchange or conformational exchange also contributes to the spin relaxation. In this section, we review the types of  $^{15}\text{N}$  relaxation experiments that are used to characterize protein backbone dynamics.

The model-free approach is the most frequently applied protocol to extract information about overall and fast (faster than overall) dynamics in proteins. In this approach the spectral density function characterizing the randomly fluctuating local fields is written in terms of correlation times for overall and internal motion and a generalized order parameter (the model free parameters). These parameters are obtained from the model-free analysis using  $^{15}\text{N}$  longitudinal relaxation rate ( $R_1$ ), transverse relaxation rate ( $R_2$ ), and  $^{15}\text{N}\text{-}\{^1\text{H}\}$  nuclear Overhauser effect (NOE) measurements. Measuring at least four relaxation rates, at two or more static magnetic field strengths, improves determination of the model-free parameters. Alternatively,  $^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE alone may be used to evaluate the high-frequency spectral density function,  $J(\omega_{\text{H}} \pm \omega_{\text{N}})$ , instead of derivation of model-free parameters. However this approach is not straightforward at high magnetic field strength for proteins with significant mobility as described in Sect. 2.3.

Although relaxation of nuclei other than  $^{15}\text{N}$  can be used to characterize backbone dynamics,  $^{15}\text{N}$  relaxation experiments have been the most widely applied, and for this reason are the focuses of this chapter.

## 2 Recent Improvements in the $^{15}\text{N}$ $R_1$ , $R_2$ , and $^{15}\text{N}\text{-}\{^1\text{H}\}$ NOE Experiments

A set of  $^{15}\text{N}$   $R_1$ ,  $R_2$ , and  $^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE observations is typically analyzed using model-free analysis to obtain the generalized-order parameter,  $S^2$ , that characterizes degree of internal motion [1, 2]. Since a set of  $^{15}\text{N}$   $R_1$ ,  $R_2$ , and  $\{^1\text{H}\}\text{-}^{15}\text{N}$  NOE

observables mainly reflects values of the spectral density functions at zero,  $^{15}\text{N}$ , and  $^1\text{H}$  frequencies, (i.e.,  $J(0)$ ,  $J(\omega_{\text{N}})$ , and  $J(-\omega_{\text{H}})$ , respectively), the model-free analysis using the three experimental data sets is suitable to characterize fast internal motion in proteins [3–5]. In the model-free analysis, a correlation time for internal motion,  $\tau_i$ , is determined for each amide site in addition to  $S^2$ . Moreover, when the simple model-free spectral density function is unable to fit the data, an extended model that contains an order parameter for faster internal motion,  $S_f^2$ , an order parameter for slower internal motion,  $S_s^2$ , and a correlation time,  $\tau_s$ , for the slower time scale motion is tested. A chemical exchange term,  $R_{\text{ex}}$ , is added to test for the presence of slow (milli- to microsecond) motions as well. These parameter optimizations are conducted by initially assuming a spherical rigid body rotation of the molecule, i.e., assuming a single rotational correlation time,  $\tau_R$ . Subsequently, either an axially symmetric or fully asymmetric model of the molecular rotational diffusion may be tested. The principles, protocols, and verification of the parameterization derived by the model-free analysis have extensively been studied and described [6–34]. In this section we focus on the recent developments in applying the model-free approach.

## 2.1 Practical Aspects in $^{15}\text{N}$ $R_2$ Experiment

Since  $R_2$  is the only observable that provides information about the  $J(0)$  spectral density contribution, accurate measurement of this observable is of particular importance. Transverse relaxation rates are typically measured by either a spin-lock ( $R_{1\rho}$ ) or a Carr–Purcell–Meiboom–Gill (CPMG) experiment. In the following, advantages and disadvantages of the two experiments are described with particular consideration of (1) limitations on the applied  $B_1$  field strength in which  $\omega_1 = \gamma_{\text{N}}B_1$  ( $\gamma_{\text{N}}$  is gyromagnetic ratio of  $^{15}\text{N}$ ), (2) off-resonance error, and (3) suppression of cross correlation by  $^1\text{H}$ – $^{15}\text{N}$  dipolar interaction (DD) and  $^{15}\text{N}$  CSA.

### 2.1.1 Spin-Lock $R_{1\rho}$ Experiment

In the  $R_{1\rho}$  experiment, in which relaxation is measured in a rotating reference frame, an rf field,  $B_1$ , is applied during the relaxation period, during which time the magnetization is “locked” almost parallel to  $B_1$ .  $R_{1\rho}$  is a function of  $R_1$  and  $R_2$ , given by

$$R_{1\rho} = R_2 \sin^2 \theta + R_1 \cos^2 \theta. \quad (1)$$

Here,  $\theta$  is given by  $\tan(\omega_1/(\omega_0 - \omega_{\text{rot}}))$ , and  $\omega_0$  and  $\omega_{\text{rot}}$  are the Larmor frequency of the signal and the angular frequency of the rotation frame, respectively. To obtain  $R_2$  from  $R_{1\rho}$  most accurately, it is advantageous to increase the  $R_2$

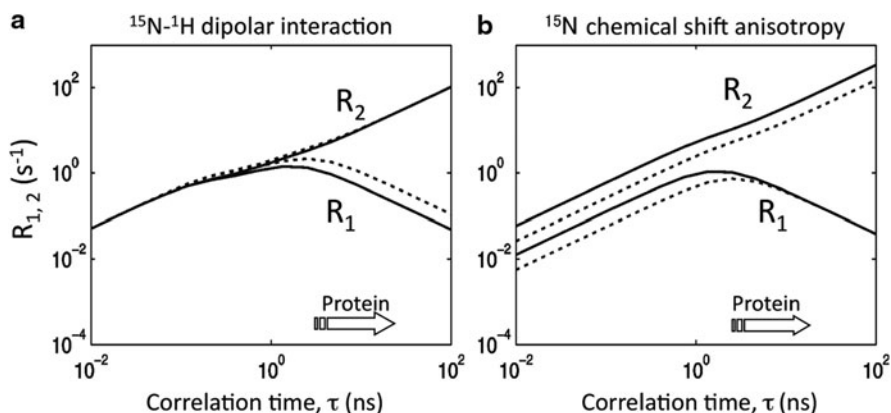


contribution to  $R_{1\rho}$ . For this purpose, a strong  $B_1$  field strength, which makes  $\theta$  close to  $90^\circ$ , is employed.

*Limitations of applicable  $B_1$  field strength* and the rf duty cycle depend on individual probes. In general, at  $^{15}\text{N}$  resonance frequency at 61 MHz, we recommend that a spin lock field  $\gamma_{\text{N}}B_1/2\pi > 2$  kHz be applied for 60–80 ms to determine  $^{15}\text{N}$   $R_{1\rho}$ . The  $\gamma_{\text{N}}B_1/2\pi$  was calculated assuming that the entire chemical shift range for amide backbone  $^{15}\text{N}$  signals in diamagnetic proteins is  $\pm 900$  Hz ( $\pm 15$  ppm) at 61 MHz, which corresponds to  $\sin \theta > 0.9$ . The spin lock duration was estimated based on approximate  $^{15}\text{N}$   $R_2$  of a folded 10–20 kDa protein at room temperature. Although the highest measurement accuracy is obtained when data is recorded until the magnetization decays sufficiently (typically for a time  $\geq 1/R_{1\rho}$ ), it may not be possible to satisfy this condition for  $^{15}\text{N}$  sites which relax slowly (as in unfolded proteins or small peptides) without reducing  $B_1$  which leads to a reduction in  $\sin \theta$ .

*Errors in  $R_{1\rho}$*  resulting from off-resonance effects may be significant but can be corrected. It is an advantage of the  $R_{1\rho}$  experiment that, even when a signal is off-resonance from the rf carrier frequency and for which  $\sin \theta$  is small, an accurate  $R_{1\rho}$  value can be obtained using equation (1). Although the correction requires an  $R_1$  value, this will be available when  $R_1$  data is recorded to characterize fast backbone dynamics using model-free analysis. As described above, the accuracy of  $R_{1\rho}$  measurements decreases for signals located far off-resonance. Compared to CPMG  $R_2$  that is described in Sect. 2.1.2,  $R_{1\rho}$  values do not need to be recorded at different carrier-frequencies without discarding any data.

*Cross-correlation* interference by  $^1\text{H}$ – $^{15}\text{N}$  dipolar interaction (DD) and  $^{15}\text{N}$  CSA has to be suppressed to detect accurate  $^{15}\text{N}$  transverse relaxation rates (Fig. 1). The cross term is suppressed by flipping the sign of the DD term by applying  $^1\text{H}$   $180^\circ$  pulses at a rate greater than the decay rate of the two  $^{15}\text{N}$ – $^1\text{H}$  J-coupled components [35–37]. However, in a weak  $B_1$  field, the two J-coupled components undergo



**Fig. 1** (a)  $^{15}\text{N}$ – $^1\text{H}$  dipolar and (b)  $^{15}\text{N}$  CSA contribution to longitudinal and transverse relaxation rates ( $R_1$ , and  $R_2$ ) as a function of a correlation time. The rates were calculated assuming a simple Lorentzian spectral density function,  $J(\omega) = \tau/(1 + \omega^2\tau^2)$ . Solid and dotted lines indicate rates calculated assuming at 900 MHz and 600 MHz instruments, respectively

precessions at frequencies  $\pm J/2$  relative to the rotating frame. Because this frame is significantly tilted from the rotating frame, application of frequent  $^1\text{H}$  pulses results in incomplete cancelation of the cross term. Therefore, infrequent  $^1\text{H}$  pulses are applied to suppress DD/CSA correlation when the off-resonance effect is significant [38, 39].

Other parameters that have to be specifically considered in the  $R_{1\rho}$  experiment are the spatial homogeneity of the  $B_1$  field strength and linearity of the power amplifier. Although  $B_1$  homogeneity has been improved in recent NMR probes, estimation of the inhomogeneity remains important to confirm the accuracy of the obtained  $R_{1\rho}$ . For this purpose, measurement of the  $B_1$  inhomogeneity using inverse detection is useful [40]. Amplifier linearity has also improved. However, since pulse power is switched for spin-lock, it is important to check that a phase shift accompanies a power change, and make a correction should one be needed.

### 2.1.2 Carr–Purcell–Meiboom–Gill $R_2$ Experiment

In CPMG  $R_2$ , *limitation of applicable  $B_1$  field strength* is in general smaller than that of the  $R_{1\rho}$  experiment. Since CPMG  $180^\circ$  pulses are applied with interpulse delays ( $2\tau_{\text{CP}}$ ), CPMG pulses with much stronger  $\gamma_{\text{N}}B_1/2\pi$  than that of spin-lock can be applied. However, in a protein, the RF field strength used in the  $^{15}\text{N}$  CPMG  $R_2$  experiment must be carefully considered because the interpulse delay is set short (typically,  $2\tau_{\text{CP}} \sim 1$  ms) to suppress generation of antiphase terms,  $\text{N}_{\text{X,Y}}\text{H}_{\text{Z}}$ , caused by  $^1\text{H}$ – $^{15}\text{N}$  J coupling. For example, if 6 kHz CPMG pulses are applied with  $2\tau_{\text{CP}} = 1$  ms, the RF power delivered to the probe is nine times stronger than that of a 2 kHz spin lock. Given that the duty cycle is ca. 10% in the CPMG experiment, about the same amount of energy is deposited in the probe in both CPMG and  $R_{1\rho}$  experiments.

In CPMG  $R_2$ , *off-resonance error* is negligible at low magnetic field strength but significant at high magnetic field strength [41, 42]. Although the stronger  $\gamma_{\text{N}}B_1/2\pi$  for each CPMG pulse inverts magnetization more uniformly than the spin-lock, CPMG pulse train accumulates error caused by a combination of pulse imperfections and off-resonance effects. The CPMG error of a signal located at off-resonance frequency  $f_{\text{off}}$  ( $f_{\text{off}}$  is the difference between the signal and carrier frequencies) is estimated with a function of  $2\tau_{\text{CP}}$  and  $B_1$ . Importantly, the off-resonance error is maximized at  $2\tau_{\text{CP}}f_{\text{off}} = n$  ( $n$  is an integer): when  $2\tau_{\text{CP}} = 1$  ms, the off-resonance error is significant at  $f_{\text{off}} = 1, 2,$  and  $3$  kHz [41]. Magnitude of the error depends on the  $B_1$  field strength the  $180^\circ$  pulses. This relationship indicates that  $R_2$  can be recorded without significant off-resonance error at 61 MHz for signals because the entire chemical shift range for amide backbone  $^{15}\text{N}$  signals spans approximately less than  $\pm 15$  ppm. However, if the same  $2\tau_{\text{CP}}$  is used at  $^{15}\text{N}$  91 MHz resonance frequency, signals located at  $\pm 11$  ppm off-resonance suffer from significant errors in measured  $R_2$  values.

It is a disadvantage of CPMG  $R_2$  experiment that there is no simple equation to correct for CPMG  $R_2$  off-resonance effects. In practice, it is recommended to discard  $R_2$  data obtained at  $f_{\text{off}} = n/2\tau_{\text{CP}}$  frequency, and record the data at two different carrier frequencies. As an alternative, a phase cycle to average out the off-resonance effect in CPMG  $R_2$  may be used [43, 44]. Using this method,  $R_2$  is

determined from the observed relaxation rate using known  $R_1$ ,  $\tau_{CP}$ , and the CPMG pulse width. Application of the sequence to relaxation dispersion has also been proposed [45]. Dissection of the relaxation contribution in the  $R_2$  dispersion determined using the sequence has been described [46].

*Cross-correlation* of  $^1\text{H}$ – $^{15}\text{N}$  dipolar interaction (DD) and  $^{15}\text{N}$  CSA in CPMG  $R_2$  experiments has been suppressed by applying  $^1\text{H}$   $180^\circ$  pulses [35–37]. Since CSA relaxation increases as a function of the magnetic field strength (Fig. 1), relative contribution of DD/CSA cross-correlation in the  $^{15}\text{N}$  transverse relaxation increases up to an external field strength.

### 2.1.3 Practical Relaxation Delay at High Magnetic Field Strength

Propagation of experimental noise,  $\Delta I$ , to the uncertainty,  $\Delta R$ , of a relaxation rate,  $R$ , in an experiment where signal is measured at only two relaxation time points is calculated using (2). Here, a two-point single exponential decay function,  $I(t) = I^0 \exp(-RT)$ , is assumed [46, 47]:

$$\Delta R/R = (\Delta I/I^0)[1 + \exp(2RT)]^{1/2}/(RT). \quad (2)$$

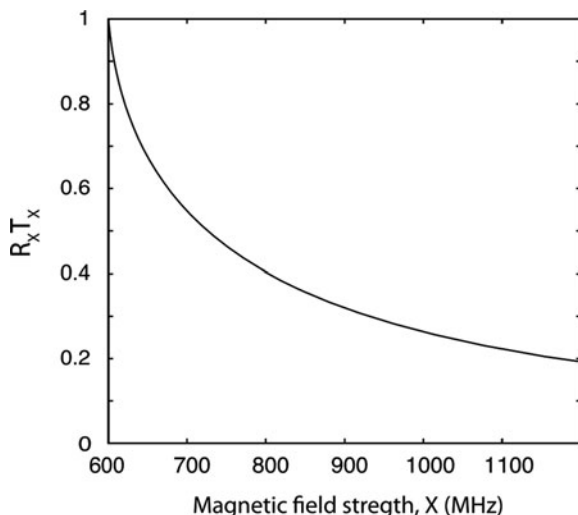
Signal-to-noise ratio,  $I^0/\Delta I$ , is proportional to  $(B_0)^{3/2}$  (the magnetic field strength is  $B_0$ ) [48]. By taking the advantage of the gain of the signal-to-noise ratio, the time  $T$  can be reduced in the experiments performed at a higher magnetic field strength when similar fractional error,  $\Delta R/R$ , to that obtained at a lower magnetic field strength is desired. Assuming that  $R$  at 600 MHz is determined by a two-point exponential fitting at  $T = 0$  and the optimal  $T (=1/R)$ , the  $R_X T_X$  at  $X$  MHz is determined by the following function:

$$(B_0^X/B_0^{600})^{3/2} = 2.86 R_X T_X / [1 + \exp(2R_X T_X)]^{1/2}. \quad (3)$$

This equation predicts how much  $T$  can be reduced to have  $R_X$  value with a similar uncertainty to that at 600 MHz (Fig. 2). For example, when  $R$  at 600 MHz is determined by a two-point exponential fitting at  $T = 0$  and  $T = 1/R$ , the  $R_X T_X$  at  $X = 900$  MHz can be reduced by a factor of 0.32 times the optimal value,  $T_X = 1/R_X$ . Since there are other factors that contribute to actual signal-to-noise ratio of relaxing magnetization and that determine sensitivity of two different magnetic fields, this calculation is a rough estimate.

## 2.2 Practical Aspects in $^{15}\text{N}$ $R_1$ Experiment

$^{15}\text{N}$   $R_1$  depends upon  $J(\omega_N)$  and  $J(\omega_H)$ , but not on either  $J(0)$  or chemical exchange. The pulse sequence used to measure  $R_1$  in proteins usually incorporates



**Fig. 2** Estimate of  $R_X T_X$  value at the magnetic field strength at  $X$  MHz to provide the uncertainty of the relaxation rate ( $R_X$ ) equivalent to that obtained at 600 MHz. The graph was calculated using (3), in which  $(I^0/\Delta I)$  is proportional to  $(B_0)^{3/2}$  and  $R$  at 600 MHz is determined by a two-point exponential fitting at  $T = 0$  and  $T = 1/R$ . For example, at  $X = 900$  MHz (*horizontal scale*),  $R_X T_X$  can be reduced to 0.32 (*vertical scale*) than the optimal  $T_X = 1/R_X$  to obtain equivalent  $R_X$  uncertainty at 600 MHz

a Freeman–Hill phase cycle in which measured magnetization decays from an initial value of  $I(0)$  to zero of magnetization starting from  $I(0)$  to  $z$  [49]. Since this phase cycle avoids recovery from  $-I(0)$  to  $I(0)$ , there is no need to record the magnetization recovery till equilibrium is attained. Otherwise, the total experimental time is enormously long [50]. In addition, because it is known that the  $I(t)$  approaches zero at infinite time  $t$ , offset is not required as an unknown parameter in the exponential fitting. Typically, only two unknown parameters ( $R_1$  and  $I^0$ ) are optimized in the exponential fitting.

The pulse scheme to suppress DD/CSA cross-correlation interference is similar to that used to measure  $R_2$  in which  $^1\text{H}$   $180^\circ$  pulses are applied during every 5–10 ms (at a rate greater than the decay rate of the faster-relaxing components of the  $^{15}\text{N}$ – $^1\text{H}$  J-coupled two components) [37]. As the magnetic field strength increases, the magnitude of the  $\tau_R(1 + \tau_R^2\omega^2)$  term decreases (Fig. 1). However, in the slow molecular tumbling limit ( $\tau_R\omega_N \gg 1$ ), the  $R_1$  contribution by the CSA term is field independent because  $\sigma_{\text{CSA}}B_0$  increases cancels by the reduction of the  $\tau_R(1 + \tau_R^2\omega^2)$  term. Thus, although the relative contribution of DD/CSA increases, the absolute cross-correlation effect on  $R_1$  is not necessarily increased with increase in the magnetic field strength. The application rate of the  $^1\text{H}$   $180^\circ$  pulses may depend more on the apparent relaxation rates of the two components, i.e., the  $^1\text{H}$  spin-flip rate, but not necessarily on the magnetic field strength.

To suppress DD/CSA cross-correlation in  $^{15}\text{N}$   $R_1$  experiment, application of accurate  $^1\text{H}$   $180^\circ$  inversion pulses is important. In principle, it is best to invert only

amide protons and not perturb water protons [51]. In this way, saturation of water proton magnetization, which reduces amide proton intensity by water-amide exchange at high-pH, is avoided. However, in practice this can be achieved only at high fields in which there is a sufficient chemical shift separation between amide and water proton chemical shifts. Otherwise, complete selective inversion is achieved at the cost of putting amide proton magnetization in the transverse plane for significant periods of time.

### 2.3 Practical Aspects in $^{15}\text{N}\text{-}\{^1\text{H}\}$ NOE Experiment

$^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE equals the ratio of steady state  $^{15}\text{N}$  signal intensities recorded with/without  $^1\text{H}$  saturation. Since the  $^{15}\text{N}$  signal intensities have to be accurately encoded in the  $t_1$  dimension, INEPT transfer from  $^1\text{H}$  to  $^{15}\text{N}$  is not used prior to the  $t_1$  evolution. As a result, the sensitivity of the  $^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE experiment is ca. ten times lower than that of  $^{15}\text{N}$   $R_1$  and  $R_2$ . To compensate for the low signal-to-noise ratio, experiment recording times are larger in the NOE experiment than in  $R_1$  and  $R_2$ . In addition, a long magnetization recovery time ( $>3\text{ s}^{-1}$ ) between each scan adds to the total time required to obtain data with adequate sensitivity [6, 52–54].

Sufficient magnetization recovery of  $^{15}\text{N}$  (in the experiment with  $^1\text{H}$  saturation) or both  $^1\text{H}$  and  $^{15}\text{N}$  (in the experiment without  $^1\text{H}$  saturation) is crucial to determine NOE values accurately. In proteins where  $\tau_R\omega \gg 1$ , dipolar longitudinal relaxation rates decrease as the resonance frequency,  $\omega$ , increases. Therefore, a longer recovery time will be required for the  $^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE experiment at higher magnetic field strength. In a rigid protein with a rotational correlation time,  $\tau_R$ , of 10 ns,  $^{15}\text{N}$   $R_1$  is  $1.1\text{--}1.4\text{ s}^{-1}$  at 61 MHz (in a 600 MHz NMR instrument) whereas  $^{15}\text{N}$   $R_1$  is ca.  $0.69\text{--}0.85\text{ s}^{-1}$  at 91 MHz (in a 900 MHz instrument). Thus, recovery times more than 3 and 5 s are required at 61 and 91 MHz, respectively. However, these are the recovery times estimated from  $^{15}\text{N}$   $R_1$ .  $^1\text{H}$  recovery times often become longer when there is not much surrounding  $^1\text{H}$  nuclei or less  $^1\text{H}$  spin-flip (such as deuterated proteins, unfolded proteins, or in a loop region of a folded protein). In this case, insufficient  $^1\text{H}$  Z-magnetization recovery is corrected using the equation derived by Bax and Grzesiek [52]. When the  $^{15}\text{N}$   $R_1$  recovery is not sufficient, another correction equation that counts both  $^1\text{H}$  and  $^{15}\text{N}$   $R_1$  recovery is used [55]. These corrections work reasonably once accurate  $^1\text{H}$  and  $^{15}\text{N}$   $R_1$  values are obtained.

When the amide proton magnetization does not recover in a single exponential manner, the correction equations do not give accurate results. In particular, when there is severe DD/CSA cross correlation in a deuterated protein in which the proton spin-flip rate is small at high-magnetic field strength, the decay of  $^1\text{H}$  magnetization of one of the two  $^{15}\text{N}$  coupled components becomes slow and nonexponential [56]. A simple solution will be to apply a sufficiently long recovery time. An alternative solution will be a pulse sequence that has recently introduced by Ferrage and coworkers [57, 58].

$^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE is thought to decrease monotonically as the rate of fast internal motion increases. However, the actual dependence of the NOE on correlation time is more complex. For example, when the model-free approach is used to express the spectral density using a correlation time for internal motion ( $\tau_i$ ) as well as an overall correlation time ( $\tau_R$ ), the NOE is a two-valued function of  $\tau_R$ , attaining a maximum value at one value of  $\tau_R$  (for example, see Fig. 19.10 in [59]).

### 3 Extending Relaxation Measurements Beyond $R_1$ , $R_2$ , and $^{15}\text{N}\text{-}\{^1\text{H}\}$ NOE

Although a set of  $^{15}\text{N}$   $R_1$ ,  $R_2$ , and  $^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE is commonly used to characterize backbone protein dynamics, other relaxation experiments are also useful to characterize protein dynamics. Use of other than three experimental data allows an application of a more detail dynamics model than the conventional model-free model. For example, cross-correlated longitudinal ( $\eta_Z$ ) and transverse ( $\eta_{XY}$ ) rates between  $^1\text{H}\text{-}^{15}\text{N}$  DD and  $^{15}\text{N}$  CSA have provided useful information about protein backbone dynamics [14, 60, 61]. Since cross-correlated relaxation occurs together with auto-relaxation, the rate is obtained by multiple exponential fitting [36, 62–65]. As an alternative approach, the  $\eta_{XY}$  rate has been more accurately determined by taking intensity ratios of the inphase and antiphase magnetization [66].

In the analysis of cross-correlated relaxation rates to detect protein dynamics, the relative orientation between the  $^{15}\text{N}\text{-}^1\text{H}$  dipole and the  $^{15}\text{N}$  tensor ( $\sigma_{\parallel}$ ) is required as an additional parameter for fitting the data [67, 68]. Numerous measurements of cross correlated relaxation have been used to estimate the  $^{15}\text{N}$  CSA in protein backbone in solution [14, 15, 66, 67, 69–73]. According to these results,  $^{15}\text{N}$  CSA is an axially symmetric CSA tensor,  $169 \pm 5$  ppm, with a relative orientation about  $21.4^\circ \pm 2.3^\circ$  tilted against the N–H dipolar tensor [73]. Thus, a disadvantage of the use of  $\eta_{XY}$  is that this tilted angle from the N–H vector has to be included as a fixed parameter in the model-free analysis. Since the relative contribution of  $\eta_{XY}$  in  $S^2$  depends on the degree of internal motion in each residue, use of  $\eta_{XY}$  may introduce an additional uncertainty in the model-free analysis. Therefore, it is important to clarify how much  $^{15}\text{N}$  CSA varies site-specifically.

Recently, relaxation rate products ( $2N_X H_X$ ,  $2N_Z H_X$ ,  $2N_X H_Z$ , and  $2N_Z H_Z$ ) have been measured to extract information about internal motion of proteins [74]. By addition and subtraction of these four terms there ideally remains only the relaxation rate, that contains the  $J(0)$  term, obtained by dipolar coupling. Using this relaxation rate as well as  $R_1$  and  $^{15}\text{N}\text{-}\{^1\text{H}\}$  NOE values,  $S^2$  values independent from the chemical exchange contribution were determined. The same set of data has also been applied to extract chemical exchange contribution [75].

Dynamics on a time scale much slower (i.e.,  $\sim 10$  ms) than can be measured by  $R_2$  relaxation dispersion is often characterized by measuring the exchange of the longitudinal  $^1\text{H}$  magnetizations among species undergoing chemical exchange

[76]. Detection of exchange crosspeaks in  $^1\text{H}$  spectra is often difficult in large proteins because of interference from the many  $^1\text{H}$  NOE cross peaks. In contrast, observation of such exchange through  $^{15}\text{N}$  spectroscopy has the advantage that only exchange peaks are observed. The  $^{15}\text{N}$  Z-exchange two-dimensional spectra are acquired using a pulse sequence similar to  $^{15}\text{N}$   $R_1$  experiment but with  $t_1$  chemical shift evolution period prior to the Z-mixing (relaxation) period [4, 77–79]. The measured rates of exchange provide information useful to characterize ligand–protein and protein–protein interactions.

## 4 Relaxation Dispersion Experiments

In contrast to  $^{15}\text{N}$   $R_1$ ,  $R_2$ , and  $^{15}\text{N}\{-^1\text{H}\}$  NOE experiments that characterize subnanosecond motions, CPMG and spin-lock relaxation experiments provide quantitative information about milli- to microsecond time scale motions. In this section, the relaxation dispersion is first defined, and, subsequently, CPMG and spin-lock  $R_2$  dispersion experiments that have recently been developed for applications to proteins are reviewed.

### 4.1 Relaxation Dispersion in General

Although currently the term “relaxation dispersion” or “ $R_2$  dispersion” often refers to CPMG or spin-lock (off-resonance or resonance  $R_{1\rho}$ ) measurements, more generally the term refers to the relaxation rates measured as a function of magnetic field strength. Typically either the static field,  $B_0$ , provided by the spectrometer magnetic or the radio-frequency (RF) field,  $B_1$ , generated by the probe transmitter coil is varied over a wide range.  $B_0$ -dependent dispersion studies are also known as “NMR relaxometry,” “field cycling,” or “nuclear magnetic relaxation dispersion (NMRD)” in the literature. In these experiments  $R_1$  of particular nucleus is measured as a function of  $B_0$  [80–90]. In these relaxation dispersion applications, the spectral density function  $J(\omega)$  is determined at numerous values of  $\omega$ , allowing various dynamical features of macromolecules, such as paramagnetic interaction with proteins and residence times of water molecules in proteins to be obtained [84, 87]. The greatest advantage of the  $R_1$  dispersion experiments is that the effective field strength varied is very wide. However, there are significant technical challenges to varying rapidly the static field strength of the samples and to increasing sensitivity [91–93]. In contrast, the dependence on  $R_2$  on  $B_1$  is readily measured, and has been used to study chemical exchange for a long time [94–101]. The  $R_2$  dispersion experiment can detect such low field effects of the exchange in chemical shifts whereas the range of variable effective field strength in the  $R_2$  dispersion is relatively small compared to that of the  $R_1$  dispersion.  $R_2$  dispersion experiments that are recently applied for biological systems will be described below.

## 4.2 $^{15}\text{N}$ CT-CPMG Relaxation Dispersion Experiment

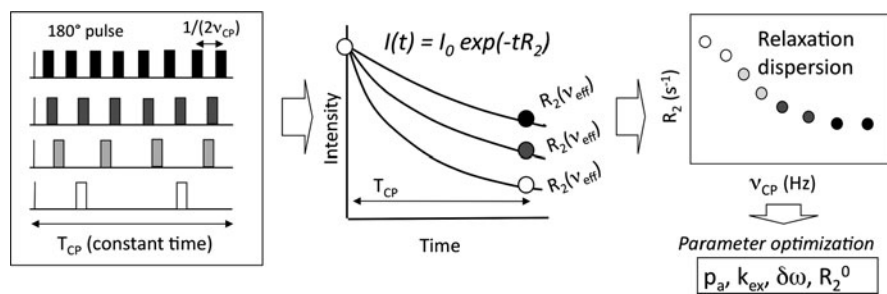
The CPMG  $R_2$  experiment has long been applied to detect chemical exchange phenomena [95, 96, 102]. A simple case is when the exchange rate is larger than the intrinsic relaxation rate,  $R_2^0$  (determined by dipolar and CSA relaxation, details described below), and the observed  $R_2$  is the sum of  $R_2^0$  and the chemical exchange contribution,  $R_{\text{ex}}$ , [98, 103]:

$$R_2\tau_{\text{CP}} = R_2^0 + R_{\text{ex}}(\tau_{\text{CP}}). \quad (4)$$

Here,  $\tau_{\text{CP}}$  is a half duration of CPMG interpulse delay, and the effective field strength is defined by  $\nu_{\text{CP}} = 1/(4\tau_{\text{CP}})$ . CPMG experiments have also been applied to characterize protein conformational equilibria and kinetics on the milli- to microsecond time scale [103–108]. The current form of the constant-time (CT) version of the experiment (Fig. 3) consists of the following two distinct steps.

One step averages the contributions to  $R_2^0$  from the inphase and the antiphase ( $N_{XY}$  and  $N_{XY}H_Z$ , respectively) components by having tandem CPMG periods linked by an rc-INEPT [109]. Averaging is achieved because one CPMG period starts from the inphase coherence and the other from the antiphase coherence for all values of  $\tau_{\text{CP}}$  [109]. This insures that  $R_2^0$  remains the same at all values of  $\tau_{\text{CP}}$ . Recently, use of a relatively strong  $^1\text{H}$  CW field ( $B_1 > 15$  kHz) that decouples the inphase and the antiphase terms has been introduced [110]. The approach is advantageous because the maximum value of  $\tau_{\text{CP}}$  that can be achieved increases twofold, using the same transverse relaxation duration. However, application of a strong  $^1\text{H}$  CW field may cause heating by dielectric or inductive losses for samples containing high salt at high magnet field strength [111–114].

The other step in the CPMG relaxation dispersion experiment is determination of  $R_2$  values by two-point intensity measurement: one is the magnetization at the initial time and the other is the magnetization at time  $T_{\text{CP}}$  [115, 116]. Although



**Fig. 3** Schematic flow chart of the experiment and analysis of constant-time CPMG  $R_2$  dispersion



such a two-point fitting has been studied [117, 118], it was not often applied for the conventional relaxation experiments for proteins, which record several data points to optimize the relaxation rate. For example, in the  $^{15}\text{N}$   $R_1$  and  $R_2$  experiments for the model-free analysis, it must be important to ensure that the signal decay is expressed as a single-exponential function so that the theoretical equations are applied for the analysis. This is not the case in the analysis of  $R_2$  dispersion because  $R_2^0$  is not dissected to extract parameters for internal motion. It will also be noteworthy that recent developments of the commercial NMR instrument have significantly increased signal-to-noise ratio, which enables two-point intensity measurements of  $^{15}\text{N}$   $R_2$  of protein samples in more practical.

By using the rc-INEPT and by applying the two-point exponential fitting, the CPMG period is held constant in the CT-CPMG relaxation dispersion experiment. The template of this experiment was initially applied to probe side chain dynamics of Asn and Gln  $\text{NH}_2$ -sites, and then applied to detect backbone amide  $^{15}\text{N}$  sites [115, 116]. There are three critical experimental parameters for the CT-CPMG experiment:  $180^\circ$  CPMG pulse width,  $p_{90}$ , the half duration between the CPMG pulses,  $\tau_{\text{CP}}$ , and the total CPMG relaxation delay,  $T_{\text{CP}}$ . To obtain the relaxation dispersion profile, one reference spectrum without a CPMG period and a series of spectra with a fixed CPMG period ( $T_{\text{CP}}$ ) but variable  $\tau_{\text{CP}}$  (a half duration between  $180^\circ$  CPMG pulses) are recorded. Two-point exponential fitting to determine  $R_2$  values is done for the entire set of CPMG spectra.

$R_2$  values are independent of  $\nu_{\text{CP}}$  ( $=1/(4\tau_{\text{CP}})$ ) when there is no chemical exchange ( $R_{\text{ex}} = 0$ ) on the time scale similar to that of  $\tau_{\text{CP}}$ . In contrast,  $R_2$  typically decreases as  $\nu_{\text{CP}}$  increases when there is chemical exchange. However,  $R_2$  may not be independent of  $\nu_{\text{CP}}$  even when there is no chemical exchange, if artifacts are introduced by off-resonance effects or CPMG pulse imperfections. To reduce these systematic errors, it is important to apply the strongest (shortest)  $180^\circ$  CPMG pulses possible within probe limits. Artifacts are maximized at  $2\tau_{\text{CP}}f_{\text{off}} = n$  (here, the  $n$  is integer) as discussed in Sect. 2.1 [41, 119]. Typically, the author's group employs  $90 \mu\text{s}$  or a shorter  $180^\circ$ -pulse at two different carrier frequencies, and records data up to  $1 \text{ kHz}$   $\nu_{\text{CP}}$ .

Sample heating is a more critical issue than regular CPMG  $R_2$  experiment because  $\tau_{\text{CP}}$  is shortened to achieve high  $\nu_{\text{CP}}$ . In principle, heating of the samples by CPMG pulses should be avoided because heating is not uniform at varying  $\tau_{\text{CP}}$  in the CT-CPMG  $R_2$  dispersion experiment. However, one may insert a  $^{15}\text{N}$  pulse scheme to compensate heating to perform experiments at uniform temperature [120, 121]. Such a compensation scheme for  $^1\text{H}$  pulses has also been implemented in the CT-CPMG  $R_2$  dispersion experiment with strong  $^1\text{H}$  CW irradiation [110]. Once the heating is so severe that the compensation sequence is required, the actual temperatures during the experiments have to be recorded, particularly when the  $R_2$  dispersion data are recorded at two or more static magnetic field strengths. From this aspect, it must be better to avoid heating as much as possible.

Another critical issue in  $^{15}\text{N}$  CPMG  $R_2$  dispersion is the magnitude of  $R_2^0$  in large molecules. Since  $R_{\text{ex}}$  is extracted from the measured  $R_2$ , the accuracy of  $R_{\text{ex}}$  decreases when  $R_2^0$  is large. In a pulse sequence that averages the inphase and antiphase components,  $^1\text{H}$   $R_1$  contributes to  $R_2^0$  as does  $^{15}\text{N}$   $R_2$  [109, 116]. Thus, as the molecular

tumbling becomes slower, both  $^{15}\text{N}$   $R_2$  and  $^1\text{H}$   $R_1$  increase (the later by proton spin-flip), resulting in increase in  $R_2^0$ . To overcome this problem, use of  $^1\text{H}$  CW decoupling scheme decrease eliminates the  $^1\text{H}$   $R_1$  term, resulting in  $R_2^0$  determined only by  $^{15}\text{N}$   $R_2$  [110]. Further reduction of  $R_2^0$  is achieved by a pulse sequence in which there is an additional  $\text{H}_2\text{N}_Z$ -relaxation period to set the  $^1\text{H}$   $R_1$  term to be time-independent, called the “constant relaxation time” scheme [106, 122] with combination of TROSY [123, 124]. Alternatively, use of a deuterated  $^{15}\text{N}$  protein is a simple way to reduce the  $^1\text{H}$   $R_1$  contribution in the  $R_2^0$  in  $^{15}\text{N}$  CPMG  $R_2$  dispersion.

### 4.3 $^{15}\text{N}$ Off-Resonance $R_{1\rho}$ Experiment

Characterization of milli- to microsecond motions in proteins based on chemical exchange can also be performed by the dispersion version of the  $^{15}\text{N}$  off-resonance  $R_{1\rho}$  experiment [39, 106, 125–128]. As described in the section on  $^{15}\text{N}$   $R_2$  (Sect. 2.1), a simple equation, (1), is applied to determine  $R_{1\rho}$  from  $R^{\text{obs}}$  and  $R_1$ . Although a strong  $B_1$  field strength that satisfies  $\omega_1 \gg (\omega_0 - \omega_{\text{rot}})$  is applied to minimize the  $R_1 \cos^2 \theta$  term in the standard (on-resonance)  $R_{1\rho}$  experiment, the  $(\omega_0 - \omega_{\text{rot}})$  term is significant and is varied in the off-resonance  $R_{1\rho}$  experiment [40, 106].  $R_{1\rho}$  is plotted as a function of the off-resonance field strength,  $\omega_e/2\pi$ , given by  $\omega_e = (\omega_0^2 + \omega_{\text{rot}}^2)^{1/2}$ .

In theory,  $R_{1\rho}$  is well-suited to record spectra at high  $\omega_e/2\pi$  values because high  $\omega_e/2\pi$  is achieved by increasing off-resonance field strength without increasing  $B_1$  field strength. However, since  $R^{\text{obs}}$  decreases as  $\omega_e/2\pi$  increases, a longer spin-lock period is required to record the reduced  $R^{\text{obs}}$  and this determines the limits of applicable  $\omega_e/2\pi$ . Although the sensitivity largely depend on sample concentration and intrinsic  $R_2^0$  at each site, the data in the literatures cited above have been mostly acquired for 140–200 ms with  $\omega_e/2\pi$  up to 2–4 kHz.

Most of the experimental parameters are the same as those of the on-resonance  $R_{1\rho}$  experiment except for a couple of points. (1) Constant relaxation-time scheme is useful to subtract part of the  $R_1$  component and to simplify the  $R^{\text{obs}}$  equation [106], and is also known as an “ $R_{1\rho} - R_1$ ” sequence [126]. (2) A scheme to spin-lock most of the magnetization uniformly is needed. For this, a scheme with periods of evolution due to chemical shift offset [38, 39, 106] or adiabatic rotation of magnetization by amplitude and phase modulated pulses [125, 127] is applied. When there is a large chemical shift dispersion, such as a high magnetic field strength, the latter has been recommended [125]. (3) Suppression of DD/CSA cross correlation and the artifact caused by the antiphase component is required. As described in the section on  $^{15}\text{N}$   $R_2$  (Sect. 2.1), sequences to take care of these effects have been used, in particular to record the data at weak off-resonance field strength [38, 39]. For application to large proteins, a TROSY-selected version has been proposed [128].

#### 4.4 Bloch–McConnell Equation and Related Equations

Parameter optimization of CPMG  $R_2$  dispersion data is carried out by minimizing chi squared as given by

$$\chi^2 = \sum_m \sum_i \left( \frac{R_{2a}^{i,exp} - R_{2a}^{i,cal}}{\sigma_{i,err}} \right)^2 \quad (5)$$

Here,  $R_{2a}^{i,exp}$  and  $R_{2a}^{i,cal}$  are experimental and calculated  $R_2$  values of  $i$ th  $\sigma$ CP value, respectively, and  $\sigma_{i,err}$  is experimental uncertainty of the  $i$ th  $R_2$  value.  $R_{2a}^{i,exp}$  is calculated as described below. The number,  $m$ , indicates the number of residues to be analyzed. When fitting each residue,  $m = 1$ .  $R_{2a}^{i,cal}$  is calculated by solving the Bloch-McConnell equation including the effects of  $180^\circ$  pulses iteratively or by using its analytical solutions [94, 95, 98, 103]. Typical analytical equations applied for CPMG  $R_2$  dispersion are (1) Luz-Meiboom equation that is suitable to analyze fast exchange and easy to incorporate in optimization programs because  $R_{ex}(\tau_{CP})$  is expressed by a single equation [95] or (2) Carver-Richards equation that is also suitable to analyze intermediate and fast exchange and when there are differences in  $R_2^0$  in two sites [98, 103]. Violation of these equations in the slow limit has been well described in the literature [116]. In contrast to these analytical solutions, the Bloch-McConnell equation is applicable to any exchange regime and any relaxation rates [94]. However, since intensity is calculated step by step for each  $\tau_{CP}$ , a relatively longer computation time is required.

When optimizing the parameters for on/off-resonance  $R_{1\rho}$  experiments, the same principles apply in the minimization of  $\chi_2$  in (5). Since  $R_{1\rho}$  is applied to investigate dynamics on the time scale faster than that of CPMG  $R_2$ , the fast exchange equation is often applied [39, 106]. Recently, the equation that is applicable for the slow time scale has become available, and has been applied to proteins undergoing slow conformational exchange [129–131].

#### 4.5 Practical Aspects Parameter Optimization

Two steps are involved in optimizing the fitted exchange parameters: first, parameters are optimized for each residue ( $m = 1$  in (5)) and, second, parameters are optimized for a group of residues ( $m > 1$  in (5)). Overall flow is described as follows. First, the data of individual residues are fit to get verification that the  $R_2$  dispersion data is in reasonable agreement with theory and to estimate the time scale of motions. This step may also select only the sites that exhibit significant dispersion profiles. Next, once this is done, one fits the data of a group of residues to determine global exchange parameters. When only one  $R_2$  dispersion is fit, the maximum number of unknown parameters in the two-site exchange model is four:

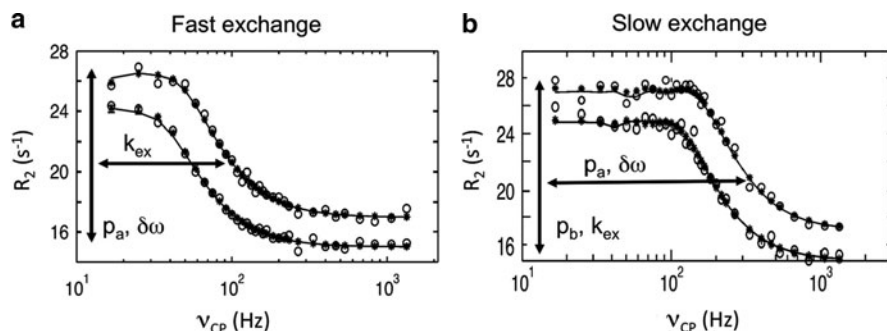
$p_a$ ,  $k_{\text{ex}}$ ,  $\delta\omega/2\pi$ , and  $R_2^0$ . In the group fits, the  $p_a$  and  $k_{\text{ex}}$  are assumed to be uniform in the group, while  $\delta\omega/2\pi$  and  $R_2^0$  are assumed to be residue specific. The total number of parameters in a group fit is therefore equal to  $2 + 2m$  ( $m$  is the number of residues, as described in (5)). As described below, there are variations in the number of parameters depending on the models and the kinds of experimental data applied for the analysis.

In the individual fit, there are a couple of practical points for better optimization. To explain it, the following simplified fast-exchange (6) and slow-exchange (7) [97, 98, 116] are useful (Fig. 4):

$$R_2 = R_2^0 + p_a p_b (\delta\omega)^2 k_{\text{ex}} / (k_{\text{ex}}^2 + (2\pi\nu_{\text{CP}})^2) \quad (6)$$

$$R_2^a = R_2^{a0} + p_b k_{\text{ex}} - p_a k_{\text{ex}} \{ \sin(\delta\omega/4\nu_{\text{CP}}) / (\delta\omega/4\nu_{\text{CP}}) \}. \quad (7)$$

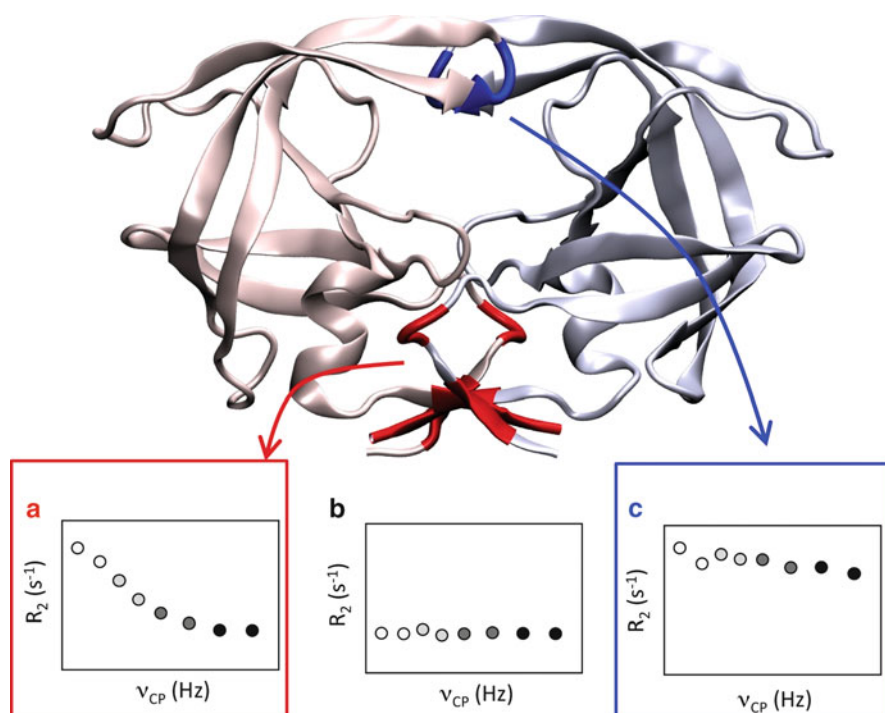
First, as seen in (6) and (7), the parameters are not independent of each other. For example, in the fast-exchange regime (6),  $p_a$  and  $\delta\omega$  are not independently determined from one relaxation dispersion profile (Fig. 4a). In such a fit for each  $R_2$  dispersion profile, the  $p_a p_b (\delta\omega)^2$  term is given as a single term,  $\Phi_{\text{ex}}$ , and the  $k_{\text{ex}}$  is determined [132, 133]. In particular, these parameters are extracted from an analysis of  $R_{1\rho}$ , in which a strong  $B_1$  field strength is applied and exchange is therefore assumed to be in the fast limit. In the slow exchange,  $(\delta\omega/4\nu_{\text{CP}})$  term is not separated (Fig. 4b). Moreover, since the dispersion profiles in the fast and slow exchanges are similar [134], the simplified equations for each time scale may be used when the time scale of exchange has been estimated by other experiments. Second, since the parameters are not such



**Fig. 4** Generated (a) fast-exchange and (b) slow-exchange  $R_2$  dispersion data (open circles) with 1% intensity noise, and their data points obtained by fit (asterisks), and the fit curve (solid lines). In (a), the noisy data points were generated assuming a two-site exchange with the following parameters:  $p_a = 0.9$ ,  $k_{\text{ex}} = 100 \text{ s}^{-1}$ ,  $\delta\omega/2\pi = 60 \text{ Hz}$ , and  $R_2^0 = 15 \text{ s}^{-1}$  at 61 MHz. In (b), the noisy data points were generated assuming a two-site exchange with the following parameters:  $p_a = 0.9$ ,  $k_{\text{ex}} = 100 \text{ s}^{-1}$ ,  $\delta\omega/2\pi = 200 \text{ Hz}$ , and  $R_2^0 = 15 \text{ s}^{-1}$  at 61 MHz. In both cases, the two sets of data with the noise were generated assuming 61 and 81 MHz. The fit data points were slightly different from the fit curves because of the noise. The vertical and horizontal bars and parameters besides indicate the parameters that determine magnitudes of  $R_{\text{ex}}$  and the effective field strength,  $\nu_{\text{CP}}$ , respectively

independent, it is safe to acquire multiple ( $>2$ )  $R_2$  dispersion data sets at different  $B_0$  field strength. When two sets of  $R_2$  dispersion data recorded at two magnetic field strengths are analyzed for each residue, the unknown parameters for optimization are  $p_a$ ,  $k_{ex}$ ,  $\delta\omega/2\pi$ , and  $R_2^0$  at one  $B_0$  field, and  $R_2^0$  at another  $B_0$  field strength. Third, even when  $R_2$  data recorded at two magnetic field strengths are available, the parameters may not be well optimized because the experimentally varied  $\nu_{CP}$  range may not be sufficiently wide. Thus, it is beneficial if  $R_2^0$  value(s) are determined by independent experiments so that the number of unknown parameters is decreased. Several reports of  $R_2^0$  determinations have been made [14, 74, 122, 135]. However, such determination of  $R_2^0$  by other methods may not be necessarily performed because of the added measurement time and when the group fit will be performed subsequently.

The group fit is useful to extract  $p_a$  and  $k_{ex}$  of the group of residues [116, 136–138]. As described above,  $p_a$  and  $k_{ex}$  may not be accurately determined by the individual fit.



**Fig. 5** Ribbon presentation of HIV-1 protease and the overview of the regions CPMG  $R_2$  dispersion profiles for (a, red) the terminal  $\beta$ -sheet region, (b, gray) the core of the protein, and (c, blue) the flap region. The terminal  $\beta$ -sheet residues exhibited significant chemical exchange ( $p_a \sim 0.94$  and  $k_{ex} \sim 650$  s<sup>-1</sup>) by the CPMG  $R_2$  dispersion experiments [137]. In contrast, the flap region exhibited too large  $R_2$  values in the CPMG  $R_2$  dispersion experiments to be analyzed [137]. However, the flap region had been found to undergo conformational exchange by the model-free analysis previously [141]. To prevent misinterpretation of data, since CPMG  $R_2$  dispersion experiments can detect exchange in a limited time scale, inspection of  $R_2$  is important as well as evaluation of the optimized parameters

The group fit may be performed for (1) a group of residues and/or (2) for a set of  $R_2$  dispersion data recorded by observing different nuclei [47, 139, 140]. The fit using a group of residues will be advantageous to identify regions that undergo conformational equilibrium in the same time scale and population, presumably indicating cooperative dynamics. Selection of the group of residues may not be straightforward because uncertainty of the optimized parameters at each residue is not necessarily Gaussian distributed [47]. In the fits of data of different types of nuclei, there may be systematic errors in individual experiments so that contour maps of the  $\chi_2$  values (5) consistent with data acquired for all types of nuclei may not be obtained [139].

Finally, a practical aspect of evaluation of regions that undergo conformational exchange is described. General technical limitation of spectroscopies, such as NMR relaxation, is that information of only a limited frequency range is obtained.  $R_2$  dispersion experiment detects chemical (conformational) exchange on or around the  $\nu_{\text{CP}}$  range. In the CT-CPMG  $R_2$  dispersion studies of Human Immunodeficiency Virus-1 (HIV-1) protease (as depicted in Fig. 5), significant  $R_2$  dispersion profiles were detected to optimize the exchange parameters only in the terminal  $\beta$ -sheet region [137]. Although the flap region exhibited very high  $R_2$  values, the data did not have the sensitivity to be fit, presumably because the time scale of motion is much faster than the studied  $\nu_{\text{CP}}$  range. Such fast dynamics (comparing to the  $\nu_{\text{CP}}$ ) of the HIV-1 protease has previously been characterized by the model-free analysis and amide  $^1\text{H}$   $R_{1\rho}$  experiments [141, 142]. Thus, application of the  $R_2$  dispersion with other experiments will be important to avoid misinterpretation of data.

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# Contemporary Methods in Structure Determination of Membrane Proteins by Solution NMR

Tabussom Qureshi and Natalie K. Goto

**Abstract** Integral membrane proteins are vital to life, being responsible for information and material exchange between a cell and its environment. Although high-resolution structural information is needed to understand how these functions are achieved, membrane proteins remain an under-represented subset of the protein structure databank. Solution NMR is increasingly demonstrating its ability to help address this knowledge shortfall, with the development of a diverse array of techniques to counter the challenges presented by membrane proteins. Here we document the advances that are helping to define solution NMR as an effective tool for membrane protein structure determination. Developments introduced over the last decade in the production of isotope-labeled samples, reconstitution of these samples into the growing selection of NMR-compatible membrane-mimetic systems, and the approaches used for the acquisition and application of structural restraints from these complexes are reviewed.

**Keywords** Detergents · Global folds · Isotope labeling · NMR structure restraints · Protein expression

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## Abbreviations

$\beta$ -OG	$\beta$ -Octyl glucoside
14- <i>O</i> -PC	1,2-Di- <i>O</i> -tetradecyl- <i>sn</i> -glycero-3-phosphocholine
6- <i>O</i> -PC	1,2-Di- <i>O</i> -hexyl- <i>sn</i> -glycero-3-phosphocholine
C <sub>6</sub> -DHPC	1,2-Dihexanoyl- <i>sn</i> -glycero-3-phosphocholine
C <sub>7</sub> -DHPC	1,2-Diheptanoyl- <i>sn</i> -glycero-3-phosphocholine
CF	Cell-free
CHAPS	3-[(3-Cholamidopropyl) dimethylammonio]-1-propanesulfonate
COSY	Correlation spectroscopy
CSA	Chemical shift anisotropy
CTAB	Cetyl-trimethylammonium bromide
DAGK	Diacylglycerol kinase
DDM	Dodecylmaltoside
DHAP	Dihexadecyldimethylammonium bromide
DPC	Dodecylphosphocholine
DTAB	Dodecyl-trimethylammonium bromide
EPR	Electron paramagnetic resonance
gA	Gramicidin A
GpA	Glycophorin A
GPCR	G-protein coupled receptor
HSQC	Heteronuclear single quantum coherence
LDAO	Lauryldimethylamine-oxide
LMPC	Lyso-myristoyl phosphatidylcholine
LMPG	Lyso-myristoyl phosphatidylglycerol
LPPG	Lyso-palmitoyl phosphatidylglycerol
MDD	Multidimensional decomposition
MSP	Membrane scaffold protein
MTSL	1-Oxyl-2,2,5,5-tetramethyl-3-pyrroline-3-methyl-methanethiosulfonate
NMR	Nuclear magnetic resonance

NOE	Nuclear Overhauser enhancement
NOESY	Nuclear Overhauser spectroscopy
NUS	Non-uniform sampling
Omp	Outer membrane protein
PRE	Paramagnetic relaxation enhancement
pSRII	Photosensitive rhodopsin II
RCSA	Residual chemical shift anisotropy
RDC	Residual dipolar coupling
SDS	Sodium dodecyl sulfate
TM	Transmembrane
TROSY	Transverse relaxation spectroscopy
UNC2	Uncoupling protein 2
VDAC	Voltage dependent anion channel

## 1 Introduction

Membrane proteins confer a remarkable array of functionalities to the membranes that define cellular boundaries [1, 2]. They are responsible for the controlled transport of nutrients, electrolytes, signaling agents, and toxins across an otherwise inert lipid bilayer, and also make it possible for a cell to sense and communicate with its environment, a vital process for a wide range of biological events. The fact that alterations in membrane protein function are linked to a number of disease states; e.g., cystic fibrosis, Alzheimer's disease, and long QT syndrome [3–6], and that 50% of known drug targets are membrane proteins [7, 8], has made this class of proteins an attractive target in drug discovery efforts. Consequently there is a high level of interest in understanding how membrane proteins function at the atomic level, and in finding ways in which these functions can be disrupted or enhanced. High-resolution structures greatly facilitate efforts to address these issues, yet at present there are only ~300 unique membrane proteins for which structures have been determined (<http://blanco.biomol.uci.edu/mpstruc/listAll/list>). Although a large number of these have been provided by X-ray crystallography, relatively recent developments in the study of large protein complexes by solution NMR have greatly increased the ability of this approach to provide important insights into membrane protein structure and function.

Solution NMR has unique capabilities to provide structural insights for proteins that are refractory to crystallization [9–11], and to characterize functionally relevant dynamic processes at atomic resolution [12–15]. However, the hydrophobic nature of membrane proteins greatly complicates handling and biophysical analyses in general. This gives rise to significant challenges, particularly for solution NMR of membrane proteins, in (1) development of cost-effective strategies to produce isotopically labeled membrane protein samples, (2) identification of detergent or lipid solutions that can maintain the protein in a folded, soluble state with a complex size that would be compatible with solution NMR, and (3) acquisition of the NMR



data required to determine their structures. We will outline current approaches being used to address these challenges and review the progress being made to increase the range of membrane protein systems that can be studied by solution NMR. It should be noted that solid-state NMR also continues to make impressive advances in the study of membrane protein structure and dynamics, as interested readers can refer to in these authoritative reviews [16–19].

## 2 Production of Membrane Protein Samples for Solution NMR

### 2.1 Factors Affecting Choice of Expression System

Solution NMR of membrane proteins usually requires that uniformly  $^{15}\text{N}$ ,  $^{13}\text{C}$ -labeled samples be produced to facilitate chemical shift and NOE assignments. Samples with uniform or site-specific incorporation of  $^2\text{H}$  atoms are also often needed to reduce the number of unfavorable relaxation pathways that can significantly attenuate the NMR signal in large protein–detergent complexes [20]. Although the price of these isotopes has come down over the last decade, cost-effective production of multiple NMR samples still requires the use of expression systems that can produce a high yield of the target protein from the simplest possible metabolic precursors. *Escherichia coli* is widely regarded as an ideal expression host for this purpose since it is simple to use, can produce high levels of protein, and has available a large variety of expression plasmids and strains for this purpose [21–23]. In addition, condensed phase approaches that allow large volumes of bacterial culture in unlabeled media to be resuspended in reduced volumes of labeled media for protein expression can further reduce associated costs [24–28]. Although yeast and mammalian cell expression systems are also being developed as alternate sources of isotope-labeled protein [29–33], the convenience of the bacterial expression system has preserved its dominance as a host for NMR sample production. Notably, almost all membrane protein structures that have been determined by solution NMR to date were produced using proteins expressed in *E. coli*, or peptides produced through solid-phase chemical synthesis (nicely summarized at <http://www.drorlist.com/nmr/MPNMR.html>).

### 2.2 Membrane Protein Expression in *E. Coli*

Special considerations for membrane protein expression in *E. coli* include issues of targeting; ideally the expressed protein can be incorporated into the bacterial cell membrane, allowing extraction of folded samples from detergent-solubilized cell membranes [34–37]. This has been the case for the small number of polytopic helical membrane protein structures that were successfully determined by solution NMR, namely diacylglycerol kinase (DAGK) [38], the disulfide bond isomerase DsbB [39], sensory rhodopsin II (pSRII) [40], and the mitochondrial uncoupling

protein 2 (UCP2) [41]. However insertion of large amounts of expressed protein in the cell membrane is often not well tolerated by the host [42, 43]. Compromised cell viability may arise from changes in lipid bilayer fluidity [44], or by overwhelming the cytoplasmic protein translocation machinery, fundamentally altering the composition of both the cell envelope and cytoplasmic proteome [45–47]. Special strains have been developed that better tolerate the stresses of toxic protein expression [48] that have proven useful for some membrane proteins [49, 50]. However, even in these systems, many membrane proteins are produced at levels that are too low to facilitate structural studies by solution NMR.

One strategy that has been used to circumvent these toxicity issues is to express the membrane protein as inclusion bodies, thereby avoiding insertion into the membrane [51–53]. These insoluble aggregates of misfolded proteins are usually non-toxic to the host cell [52]. Since they are also typically resistant to proteolytic cleavage, expression levels as high as 25% of the total cell protein has been attained through this approach [53]. While overexpression of some membrane proteins spontaneously gives rise to inclusion body formation, most notably with  $\beta$ -barrels missing their signal sequences, and mitochondrial carrier proteins [54, 55], for other proteins it is possible to use fusion tags that target them to inclusion bodies [56, 57]. For example, expression of a *trp* operon L gene that has been modified to allow translation through its native stop codon generates a polypeptide of 105 residues called *trp* $\Delta$ LE [58, 59]. This polypeptide has a strong tendency to form inclusion bodies either when expressed on its own or when expressed as an N-terminal fusion to smaller membrane proteins (i.e., one to two transmembrane (TM) helices [60–65]).

Many solution NMR structures of polypeptides comprised of a single TM helix have been produced by inclusion body targeting [62, 63, 65–67]. In the case of larger proteins that span the membrane multiple times, the development of high-yielding refolding protocols can present a significant impediment. Nonetheless this strategy has proven to be particularly useful for the production of  $\beta$ -barrel structures, e.g., the bacterial palmitoyltransferase PagP [68], outer membrane proteins OmpA [69, 70] and OmpX [71], the pH-sensitive OmpG porin [72], and the mitochondrial voltage-dependent anion channel VDAC-1 [73]. In addition, there are a number of examples of polytopic helical membrane proteins that have been expressed and refolded from inclusion bodies, e.g., a mammalian G-protein coupled receptor (GPCR) [29, 56, 74], the Y-2 receptor [75] bacteriorhodopsin [76], and a range of mitochondrial carriers [55, 77, 78]. These examples suggest some potential for inclusion body targeted expression of larger proteins, although no membrane protein NMR structure has yet been produced by this approach for any protein containing more than two TM segments.

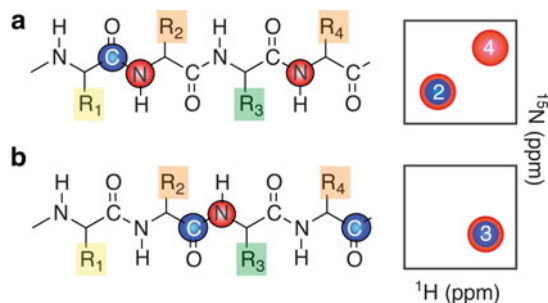
### **2.3 Cell-Free Expression Approaches**

A promising approach that circumvents complications arising from membrane protein toxicity or refolding is the cell free expression system comprised of purified

extracts from bacteria [79, 80] or wheat-germs [81, 82] added to a mixture of tRNAs, amino acids, nucleotides, enzymes, and cofactors [83–87]. The increasing availability of pure isotope-labeled amino acids has allowed a growing number of groups to use this approach to produce specifically labeled samples for NMR [87–97]. In the case of membrane proteins, lipids or detergents can be included in the reaction mix to maintain them in a soluble state during the synthesis [91, 98, 99]. Using these methods it is possible for several milligrams of the target to be synthesized from only a few milliliters of reaction mix.

An exciting development in the cell-free approach to membrane protein production is mounting evidence suggesting that additives normally intended to maintain expressed membrane proteins in solution (e.g., lipids, detergents) might not be necessary. In the absence of added lipids most of the expressed membrane protein tends to precipitate into an insoluble fraction; however, it has been shown that this aggregate can be resolubilized in mild detergents [96, 100–103]. The validity of this method has been supported by functional assays on some of these resolubilized aggregates [100, 101, 103], and in one case by highly similar NMR spectra for samples produced from cell-free precipitates vs conventional *E. coli*-based production systems [88]. Meanwhile the purity of these precipitates tends to be high, potentially eliminating the requirement for subsequent purification steps. Consequently, in some cases it is possible to directly resolubilize the cell-free expression pellet in the desired volume and composition of buffer to allow immediate acquisition of NMR data. Elimination of chromatography, dialysis, fusion protein cleavage, and protein concentration steps provides significant time savings, and also reduces material costs, particularly in the consumption of expensive detergents and lipids [88].

In addition to the advantages in sample purity offered by cell-free membrane protein expression, unique possibilities for amino acid specific labeling are also introduced with this method [88, 92]. In the case of bacterial expression systems, metabolic processes can break down amino acids added to the expression media, leading to label dilution and non-specific incorporation. Use of auxotrophic strains or a limited subset of amino acids can help to get around these issues [104, 105], but neither approach has the flexibility of specific amino acid labeling that is characteristic of cell-free expression systems. This feature has been particularly useful for the development of selective labeling approaches to facilitate backbone assignment of membrane proteins with poor spectral dispersion. Most rely on combinatorial labeling approaches, with samples having specific combinations of  $^{15}\text{N}$ -labeled and/or  $^{13}\text{C}$ -labeled amino acids. This allows inter-residue heteronuclear  $^{15}\text{N}$ – $^{13}\text{C}$  coupling, such as that used in the HNC0 experiment, to rapidly identify adjacent pairs of amino acids [106] (example described in Fig. 1). There is a range of combinatorial strategies that have been developed to maximize the information that can be obtained from a limited number of differently labeled samples [102, 107, 108]. Since membrane proteins have unique sequence and spectral characteristics compared to their water-soluble counterparts, a Monte Carlo method was developed to determine the optimal labeling strategy for a specific protein sequence [88]. Using this approach it was possible to elucidate structures for three different membrane protein structures in an impressively short 8-month period.



**Fig. 1** Example of a combinatorial isotope labeling scheme that identifies adjacent amino acids in the sequence [106]. In this simplified example, backbone structures for the same segment with two differently labeled samples are shown on the *left* in **a** and **b**, with the corresponding  $^1\text{H}$ - $^{15}\text{N}$  2D projection of the HNCO spectrum (*smaller blue peaks*) superimposed on the  $^1\text{H}$ - $^{15}\text{N}$  HSQC spectrum from the same sample (*red peaks*). Peaks are labeled with residue numbers for this tetrapeptide segment. In sample **a**, two of the amino acids with the same side chain (residues 2 and 4, *orange side chains*) are labeled with  $^{15}\text{N}$  (highlighted in *red* in the backbone structure), and one of the amino acids (*yellow*) is labeled with  $^{13}\text{C}$ . The appearance of a peak in the HNCO projection for one of the  $^{15}\text{N}$ -labeled amino acids allows the identification of the preceding amino acid type for that residue (i.e., the sequence is *yellow-orange*). In sample **b** residues 2 and 4 are now labeled with carbonyl  $^{13}\text{C}$  (*blue*), with  $^{15}\text{N}$ -labeling of 'green' amino acids. The single peak that would be seen in the HSQC spectrum for this sample would also be observed in the HNCO projection, allowing the sequence to be extended to *yellow-orange-green* in this example. The number of samples and specific amino acids that would need to be labeled for a full assignment will depend on the method used to find the combination that gives the maximum number of inter-residue correlations with a minimum number of samples [102, 107, 108]

Meanwhile, a combinatorial optimization method has subsequently been developed that allows a wider range of auxiliary conditions to be factored into the design of protein sequence-specific labeling protocols [109]. Overall, these tools should help to make cell-free expression and selective labeling increasingly accessible to a wider range of laboratories in membrane protein structural biology.

### 3 Membrane Mimetic Systems for Solution NMR

#### 3.1 Micelles

One of the challenges for solution NMR of membrane proteins is the identification of conditions that can mimic the native lipid bilayer environment while maintaining the sample in a stable, folded state with a total complex size of  $\sim 100$  kDa or less [110–112]. This has most commonly been achieved through the use of detergents that form smaller micelles ( $\sim 10$ – $30$  kDa) with a relatively high critical micelle concentration (cmc). However, the ability of some detergents to maintain membrane proteins in a water-soluble state sometimes works against structural studies, essentially solubilizing the protein so well that the native contacts are disrupted by

competing interactions with solvent detergent [113–115]. This can be exacerbated by the frequent need to acquire data at elevated temperatures (35 to >40 °C) [112]. Consequently, the identification of a detergent that offers an appropriate balance between solubilization efficiency and preservation of native structure often requires extensive screening and optimization of sample conditions. While a wide range of detergents has proven useful for the study of membrane protein structure and function, a surprisingly narrow selection has so far been used to solve the majority of high-resolution membrane protein solution NMR structures (Table 1). These include micellar systems formed by sodium dodecyl sulfate (SDS), dodecylphosphocholine (DPC), or short-chain phosphatidylcholines (C<sub>6</sub>- or C<sub>7</sub>-DHPC), and small bicelles with C<sub>6</sub>-DHPC/DMPC (Table 1).

### 3.1.1 SDS

SDS was one of the first detergents to be used for solution NMR of membrane proteins [116, 117], and has continued to be used for this purpose, particularly for the study of smaller membrane proteins containing only a single TM or amphipathic helix. SDS micelles have the ability to maintain a high concentration of membrane proteins in solution while forming protein–detergent complexes with size and dynamic properties that are often favorable for NMR spectroscopy. Although SDS is generally classified as a denaturing detergent due to its tendency to unfold water-soluble proteins [114, 118], hydrophobic peptide segments that interact with the micelle usually adopt a helical structure that minimizes exposure of polar backbone atoms to the hydrophobic interior of the micelle in the same way that would occur in lipid bilayers [119–121]. This has made this detergent a convenient solvent for the study of isolated TM segments [122, 123] and amphipathic membrane surface binding helices (e.g., [124–127]).

In spite of the generally accepted classification of SDS as a denaturing detergent, it has been widely used for the study of intermolecular interactions between isolated TM helices [123, 128, 129] or multi-spanning membrane proteins connected by short turn sequences [130, 131]. The suitability of this detergent for the study of TM-helix dimer structures has also been substantiated in a number of systems by comparison of mutational effects on dimer affinity in SDS and bilayer environments [129, 132–134] although small discrepancies do occur in some cases. However, it should be noted that examples of interactions identified in SDS exist that did not appear in the structure of the full-length protein (e.g., DAGK [135]). Moreover, there is a potential for disruptive interactions to occur between SDS and longer inter-helical loops or globular domains (e.g., the Y2 receptor [75]), making it necessary to adopt a cautious approach when interpreting structural information for membrane proteins in SDS. To date only a handful of larger (i.e., more than one TM segment) membrane proteins have been studied by solution NMR in this detergent, although in all cases the structural data was well validated by comparison with complimentary functional or structural data [136–138].

**Table 1** Properties of detergents that have been used for membrane protein structure determination by solution NMR

Charge	Detergent Structure	MW (Da)	CMC (mM)	Aggregation number	Micelle MW (kDa)	Ref
Zwitterionic	DPC $\text{H}_3\text{C}(\text{H}_2\text{C})_{10}\text{H}_2\text{C}-\text{O}-\overset{\ominus}{\text{P}}(\text{O})\text{O}-\overset{\oplus}{\text{N}}(\text{CH}_2\text{CH}_2)_3\text{CH}_3$	352	1.5	50–60	19	[142, 380]
	TDC $\text{H}_3\text{C}(\text{H}_2\text{C})_{12}\text{H}_2\text{C}-\text{O}-\overset{\ominus}{\text{P}}(\text{O})\text{O}-\overset{\oplus}{\text{N}}(\text{CH}_2\text{CH}_2)_3\text{CH}_3$	380	0.12	108	47	[246] <sup>a</sup>
	C <sub>6</sub> -DHPC $\text{H}_3\text{C}(\text{H}_2\text{C})_3\text{H}_2\text{C}-\text{O}-\overset{\oplus}{\text{N}}(\text{CH}_2\text{CH}_2)_3\text{CH}_3-\text{O}-\text{CH}_2-\text{CH}_2-\text{O}-\overset{\ominus}{\text{P}}(\text{O})\text{O}-\text{C}(\text{O})-\text{CH}_2-\text{O}-\text{C}(\text{O})-\text{H}_3\text{C}(\text{H}_2\text{C})_3\text{H}_2\text{C}$	453	14	19–25	16	[180, 388]
	C <sub>7</sub> -DHPC $\text{H}_3\text{C}(\text{H}_2\text{C})_4\text{H}_2\text{C}-\text{O}-\overset{\oplus}{\text{N}}(\text{CH}_2\text{CH}_2)_3\text{CH}_3-\text{O}-\text{CH}_2-\text{CH}_2-\text{O}-\overset{\ominus}{\text{P}}(\text{O})\text{O}-\text{C}(\text{O})-\text{CH}_2-\text{O}-\text{C}(\text{O})-\text{H}_3\text{C}(\text{H}_2\text{C})_4\text{H}_2\text{C}$	482	1.1	42–200	–	[388, 389]
	LDAO $\text{H}_3\text{C}(\text{H}_2\text{C})_{10}\text{H}_2\text{C}-\overset{\oplus}{\text{N}}(\text{O})\text{O}$	229	2	69–73	17	[389]

(continued)

Table 1 (continued)

Charge	Detergent	Structure	MW (Da)	CMC (mM)	Aggregation number	Micelle MW (kDa)	Ref
Anionic	SDS		288	8.2	62	18	[389]
	LMPG		478	0.2-0.3	55	26	[390]
	LPPG		506	0.02-0.6	125	65	[180, 390]

<sup>a</sup>Cmc, aggregation number from Affymetrix/Anatrace

### 3.1.2 Dodecylphosphocholine

Dodecylphosphocholine (DPC) has a zwitterionic headgroup (Table 1) with a reduced tendency to participate in interactions that would disrupt protein structure [139]. It has proven to be particularly useful for NMR structure elucidation of TM helix interactions, as was first demonstrated with the GpA dimer [140]. It has since been used for some of the largest multi-spanning helical membrane protein systems determined to date (e.g., DAGK [38], phospholamban [141], and UCP2 [41]). The ready availability of deuterated DPC, its stable lipid-like headgroup structure, and its relatively small micelle aggregation number and narrow micelle size distribution over a range of conditions [142] are all favorable features contributing to the general utility of this detergent. For these reasons DPC has become one of the most popular detergents for solution NMR of membrane proteins, with more than one-third of the integral membrane protein NMR structures being determined in this detergent (Table 2).

### 3.1.3 Short-Chain Phosphatidylcholines

For those looking for closer mimics of a native phospholipid bilayer, short chain phosphatidylcholines, namely 1,2-diheptanoyl- and 1,2-dihexanoyl-*sn*-glycero-3-phosphocholine (C<sub>7</sub>-DHPC and C<sub>6</sub>-DHPC, respectively), have provided an attractive alternative [143]. These phospholipids differ from those found in biological membranes in the length of the acyl chains, giving rise to a tendency to form micelles instead of bilayer structures. Meanwhile the motionally averaged conformational properties of these lipids resemble those of the long-chain phosphatidylcholine bilayers at higher temperatures [143]. Particularly useful for solution NMR is the low polydispersity and stable size of the C<sub>6</sub>-DHPC micelle over a range of concentrations [144]. However, the best compromise between sample stability and spectral quality was actually found with C<sub>7</sub>-DHPC for the seven TM-helix GPCR pSRII [145]. Although the molecular mass and polydispersity of protein-free C<sub>7</sub>-DHPC micelles is highly dependent on its concentration [144], the complexes formed with pSRII were spectrally homogeneous, illustrating the strong influence of the protein on the properties of the detergent–protein complex as has been previously observed for other proteins [146].

### 3.1.4 Lysolipids

Lysolipids have also been gaining attention as effective, relatively mild solubilization agents [88, 147–150]. The polar glycerol spacer between the headgroup and alkyl chain provides a gentler transition between hydrophobic and charge-containing phases of the micelle, reducing potentially denaturing effects of the headgroup charges. This was illustrated with DAGK, which was shown to retain full activity in lyso-myristoyl phosphatidylglycerol (LMPG) micelles while



**Table 2** Summary of unique integral membrane protein structures determined by solution NMR

Method <sup>a</sup>	Structure class	Protein	Size, kDa <sup>b</sup>	# of TM $\alpha$ ( $\beta$ ) <sup>c</sup>	Detergent	Detergent: protein molar ratio	Temp °C	PDB ID	Ref
(I) Conventional <sup>d</sup>	Single TM helix <sup>e</sup>	Na <sup>+</sup> /H <sup>+</sup> exchanger TM6	3	1	DPC	75	30	2L0E	[373]
		Integrin $\beta$ 3	9	1	DHPC/POPC (3:1) DPC	750	40 (bicelle) 25 (DPC)	2RMZ 2RN0	[162]
	TM helix dimer	Glycophorin A BNip3	4 10	1 $\times$ 2 1 $\times$ 2	DPC DHPC/DMPC (4:1)	60 40	40 40	1AFO 2J5D	[140, 391] [188]
		Erb B2	9	1 $\times$ 2	DHPC/DMPC (4:1)	40	40	2JWA	[287]
		Eph A1	8	1 $\times$ 2	DHPC/DMPC (4:1)	40	40	2K1K 2K1L	[189]
		Eph A2	9	1 $\times$ 2	DHPC/DMPC (4:1)	40	40	2K9Y	[392]
	TM helix trimer	Erb B3 DAPI2-NKG2C	9 11	1 $\times$ 2 1 + 1 $\times$ 2	DPC TDC/SDS (10:1)	45 >450	40 30	2L9U 2L35	[393] [65]
(II) I + TROSY + fractional <sup>2</sup> H-labeling	TM helix oligomer	Phospholamban <sup>f</sup>	30	1 $\times$ 5	DPC	200	30	1ZLL	[338]
		Influenza A M2 tetramer	20	1 $\times$ 4	C <sub>6</sub> -DHPC	400	30	2RLF	[166]
		Influenza B M2 tetramer	14	1 $\times$ 4	C <sub>6</sub> -DHPC	270	32	2KIX	[326]

(III) TROSY, uniform $^2\text{H}$ -labeling, $^1\text{H}^{\text{N}}$ NOEs	$\beta$ -Barrel	OmpA	19 (46)	(8)	DPC	600	50	1G90	[69]
		PagP	20	(8)	DPC $\beta$ OG	570 200	45	1MM4 1MM5	[68]
		OmpG	28	(14)	DPC	<300	40	2JQY	[72]
(IV) III + methyl NOEs		KpOmpA	22	(8)	DHPC	300	40	2K0L	[70]
		OmpX	16	(8)	DHPC	100	30	1Q9F	[71]
		VDAC	32	(19)	LDAO	>300	30	2K4T	[73]
(V) IV + non-methyl NOEs	TM helix heterodimer	Integrin $\beta 3/\alpha \text{II}_b$	9	1 + 1	DHPC/POPC/ POPS (9:2:1)	>400	30	2K9J	[184]
		pSRII	26	7	$\text{C}_7$ -DHPC	-	50	2KSY	[40]
(VI) IV + PREs + RDCs	Multi-TM helix	Ste2p TM 1-2	9	2	LPPG	>500	47	2K9P	[64]
		D8bB	18	4	DPC	-	40	2K73	[39]
		Pf1 coat protein	5	1	DHPC	100	50	2KLV	[323]
(VII) Short-range NOEs + PREs/RDCs	Single TM + amphipathic helix	Stannin	14	1	SDS	>860	60	1ZZA	[394]
		FXYD1	8	1	SDS	>500	40	2JO1	[366]
		FXYD4	7	1	SDS	>500	40	2JP3	[395]
		CD4 TM fragment <sup>e</sup>	8	1	DPC	200	45	2KLU	[396]
		Phospholamban monomer <sup>e</sup>	6	1	DPC	750	50	1N7L	[397]
		KCNE1	16	1	LMPG	80	40	2K21	[67]
		MerF	6	2	SDS	700	60	1WAZ	[137]
ActB	13	2	LMPG	200	45.37	2KSD	[88]		

(continued)

**Table 2** (continued)

Method <sup>a</sup>	Structure class	Protein	Size, kDa <sup>b</sup>	# of TM $\alpha$ ( $\beta$ ) <sup>c</sup>	Detergent	Detergent: protein molar ratio	Temp °C	PDB ID	Ref
		QseC <sup>h</sup>	21	2	LMPG	330	45.37	2KSE	[88]
		KdpD <sup>h</sup>	11	4	LMPG	330	45.37	2KSF	[88]
		Presenilin catalytic fragment	19	3	SDS	>140	40	2KR6	[136]
		DAGK trimer <sup>j</sup>	43	3 × 3	DPC	–	45	2KDC	[38]
		UCP2 <sup>j</sup>	33	6	DPC/DMPC/CL (150:2:1)	190	33	2LCK	[41]

<sup>a</sup>Method used to determine structures are divided into seven different classes, depending on the type of data and experiments that were used for the structure calculation

<sup>b</sup>Total molecular weight for all protein subunits in the complex

<sup>c</sup>Number of  $\beta$ -strands spanning the membrane are denoted in parentheses

<sup>d</sup>Standard NOE-based structure determination protocols typically applied to small water-soluble globular proteins

<sup>e</sup>Selected representatives

<sup>f</sup>Positioning of amphipathic helices in this structure represents ensemble average on the timescale of RDC and NOE restraints

<sup>g</sup>PREs were used to localized solvent-exposed or micelle-bound residues

<sup>h</sup>No NOE restraints directly used, but hydrogen bond restraints were imposed for the TM helices

<sup>i</sup>Data from disulfide bond mapping was used to create intermolecular restraints

<sup>j</sup>No NOE data used

yielding spectra of comparable quality to those obtained in the inactivating DPC conditions used to solve its structure [148]. A potential explanation for this arose from the observation of NOEs between the glycerol spacer of lyso-myristoyl phosphatidylcholine (LMPC) and Trp indole side chains at the peripheral regions of the TM helices; in DPC these same side chains showed NOEs to the phosphocholine headgroup. Based on these observations it was suggested that these indole-phosphocholine interactions were destabilizing, and that the glycerol spacer provided a more appropriate environment for interfacial side chains. The potential utility of detergents with polar spacers to improve the stability of membrane proteins has also been highlighted by synthetic variants of DPC with modifications that mimic the properties of the polar spacer [151]. In this study superior stabilization of OmpX for folding and solution NMR was obtained for DPC with a  $\beta$ -hydroxylated, or ethyl-amide-linked alkyl chain.

Lysolipids have also proven useful for direct resolubilization of membrane proteins from cell-free expression pellets for solution NMR structure determination, as was illustrated for three bacterial membrane proteins [88]. However, these structures were all characterized by loose helical packing, which may reflect a destabilizing influence of LMPG on these proteins. Alternatively, the loose structures could have been a consequence of the type and number of structural restraints used, leaving unanswered the question of possible denaturing effects of this detergent on these structures.

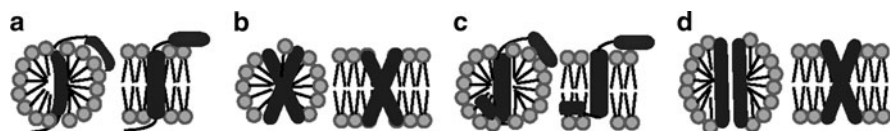
### 3.1.5 Mixed Micelles

For some proteins an appropriate balance between solubilization and stabilization could not be provided by a single detergent system, but was attained using a mixture of detergents. For example, the TM helix  $\zeta\zeta$  dimer could only be solubilized in SDS after its final purification step, but could subsequently be transferred into a less denaturing system by the addition of a fivefold molar excess of DPC over SDS, with samples undergoing aggregation if the ratio of DPC:SDS exceeded  $\sim 10:1$  [152]. Similarly, NMR spectra of the KvAP voltage sensing domain were found to be optimal in 2:1 DPC/LDAO (lauryldimethylamine-oxide), with DPC alone giving rise to exchange-broadened spectra, while LDAO yielded favorable spectral properties but short sample lifetimes [153]. An N-terminal fragment from the Y4 GPCR containing two TM segments also benefited from the use of a detergent mixture (i.e., DPC and lyso-palmitoyl phosphatidylglycerol (LPPG)), since DPC could solubilize the sample but yielded spectra with broad lineshapes, while LPPG was a poor solubilization agent for this sample [154]. The common theme arising from these studies is that detergent mixtures can offer an improved capacity to provide the best compromise between sample solubility and stability. Also, in some cases it may be necessary to include small amounts of phospholipids normally found in the membrane environment, as was required for the mitochondrial uncoupling protein 2 (UCP2) [41].

Aside from mixtures that alter the charge of micelle surfaces, it is also possible to choose detergent combinations that alter the mean size of the micelle hydrophobic core. Analysis of small angle X-ray scattering data showed that the smaller hydrophobic core in short-chain detergent micelles can be increased by the addition of longer-chain detergents [155, 156]. A linear relationship between micelle size and detergent long- to short-chain molar ratios was observed, suggesting that micelle dimensions could be tuned in a straightforward manner to match the size of a protein hydrophobic domain. For a model 2-TM helix system, detergent mixtures that optimized hydrophobic matching between micelle and protein gave rise to the highest quality spectra, and also promoted more compact protein structures [155]. Although the applicability of these trends for other membrane proteins remains to be established, this study has identified an additional parameter that can be explored when optimizing protein–detergent complexes for solution NMR.

### 3.1.6 Potential Drawbacks to the Use of Micellar Membrane-Mimetics

While most solution NMR structures of larger membrane proteins have been elucidated in non-SDS detergent micelles, there are some caveats to keep in mind whenever any detergent is used to study membrane protein structure (Fig. 2). For example, with most NMR-friendly detergents having high cmcs, there is a significant concentration of monomeric detergent in the solvent that can potentially bind to solvent-exposed regions of the protein that do not normally interact with lipids. X-ray crystal structures have provided some examples of this, with detergent being found in the active site cavity of the  $\beta$ -barrel PagP [157], and inside the M2 channel [158]. For smaller membrane proteins this is particularly significant since there are fewer intramolecular interactions that stabilize the protein fold compared to the relatively large number of detergent–protein interactions. This larger proportion of residues that are exposed to solvent detergent increases the potential for



**Fig. 2** Schematic representation of potential changes in integral membrane protein structure that could be imposed by a micellar environment (left hand side of each panel), compared to the native structure in bilayers (*right*). Possible distortions include; (a) micelle-induced curvature in the TM helix or amphipathic helix; (b) monomeric detergent molecules bound to a solvent-exposed region, in this case an aqueous cavity close to the micelle surface; (c) altered relative orientations of amphipathic vs TM helices; (d) loss of tilt relative to other TM segments. In this scenario hydrophobic mismatch between the TM helix and micelle are minimized by distortions in micelle structure that allow hydrophobic protein surfaces to remain in the hydrophobic phase. In the bilayer environment hydrophobic mismatch induces tilt, favoring a non-zero inter-helical crossing angles

detergent–protein interactions to alter the structure. This has been suggested to be a factor in the differences observed between some solid-state structures obtained in lipid membranes and solution-state detergent-solubilized forms [159].

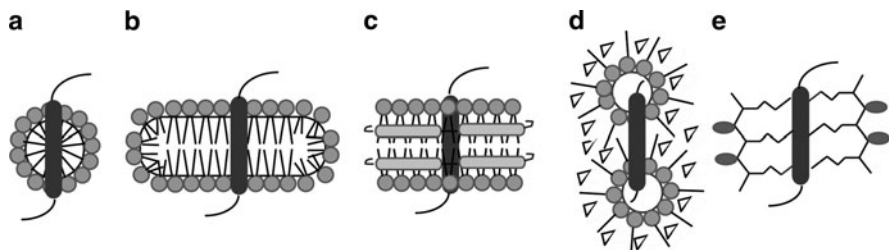
Another potential complication arising from the use of micelles is the significant curvature of the micelle surface, a particularly relevant concern for surface-binding elements. This is capable of inducing curvature in amphipathic helices to maximize burial of the hydrophobic side of the helix in the hydrophobic phase of the micelle [160]. Structural elements with an affinity for the polar regions of the native bilayer could also follow the curved dimension of the headgroup phase, potentially leading to distortions of adjacent parts of the structure (e.g., the voltage dependent potassium channel KvAP [161]). Similarly, mismatch between the size of the hydrophobic core of the micelle and TM segment length can induce curvature in the TM helix to minimize its exposure to the aqueous phase [162]. In addition, the less densely packed, more dynamic state of the detergent monomers makes the hydrophobic phase of the micelle easier to access compared to that of a lipid bilayer [163]. This can promote detergent–protein interactions that may not occur in native lipid membranes, potentially distorting and/or destabilizing other parts of the protein [164, 165].

An illustrative example of how detergents can influence structure is provided by the controversy centered about the influenza A M2 channel. In C<sub>6</sub>-DHPC micelles M2 was found in a tightly packed tetrameric state that bound the inhibitor rimantidine at four allosteric sites in the lipid-facing part of the channel exterior, inhibiting the channel by stabilizing the closed state [166]. In contrast, X-ray crystallographic [158] as well as solid-state NMR data acquired in lipid bilayers [167] on a shorter M2 construct clearly showed a single inhibitor-binding site inside the channel, with no interactions at the exterior allosteric sites. Based on these differences it was suggested that the helix tilt of the M2 TM segment in lipid bilayers could be important in maintaining a more open, active channel conformation [168]. According to this idea the malleability of loosely packed detergent molecules could allow coverage of a larger hydrophobic surface (Fig. 2d), eliminating the environment-induced tilt required to promote larger helix crossing angles. However, functional evidence can be cited that supports the biological relevance of both forms M2 channel [158, 166, 169], suggesting that the two structure types represent different functional states for this protein. This hypothesis is supported by the conformational dynamics observed for M2 in C<sub>6</sub>-DHPC [166] and the structure of a drug-resistant mutant that had a reduced affinity for rimantidine at the allosteric site [170]. Moreover, subsequent experiments have shown that DPC micelles can support a tetrameric state that binds rimantidine at a single site within the channel [171]. This confirms that a micellar environment is capable of capturing the open state of this channel, although the physical properties of the micelle that promote one form of the channel over the other remain one of the interesting questions to be addressed.

### 3.2 *Bicelles*

Avoiding some of the potential problems associated with micellar detergents, small isotropic bicelles have increasingly been used for solution NMR of membrane proteins (reviewed in [172–174]). These are formed by mixtures of up to a fivefold molar excess of short-chain lipid (e.g., C<sub>6</sub>-DHPC) or detergent (e.g., CHAPS) over long-chain lipid (e.g., DMPC, DOPC) [175, 176]. The composition of a bicelle is most accurately described by  $q$ , the molar ratio of lipid and *bicellar* detergent (i.e.,  $[\text{detergent}]_{\text{bicelle}} = [\text{detergent}]_{\text{total}} - \text{cmc}$ ) [112], which for small isotropic bicelles is typically in the range of 0.25–0.5. It has been shown that bicelles with this range of  $q$  values have a disk-shaped morphology, containing a distinct lipid bilayer phase and edges coated by a more mobile detergent phase (Fig. 3) [177–179]. The size of the protein-free small bicelle depends on  $q$ , and can be comparable to the size of commonly used detergent micelles (e.g., ~22 kDa for  $q = 0.15$  [180]). However, spectra obtained for proteins reconstituted in bicelles have generally shown broader peaks than those for the same sample in a micelle [40, 137, 181, 182], with the larger complex size for the bicelle–protein complex contributing to this broadening. Yet in some cases only a bicelle environment could uniquely confer a functionally folded, spectrally homogeneous sample (e.g., the small multidrug resistance pump (Smr) [183]). A bicelle formulation mimicking physiological membrane compositions was also found to be instrumental for structure determination of the weakly interacting integrin  $\alpha\text{IIb}\beta 3$  TM-heterodimer [184], a complex that was not supported by DPC micelles [185].

In contrast with micelles, the introduction of proteins into bicelles may require an additional optimization step, since there are a few different approaches that can be used [186, 187]. However, most membrane proteins that have been reconstituted into small isotropic bicelles for solution NMR could be prepared in solvent-free



**Fig. 3** Schematic diagram showing the general structure of various membrane mimetic systems used for solution NMR studies of membrane proteins. One TM segment is shown embedded in (a) a micelle, (b) a bicelle, (c) a nanodisc, (d) reverse micelles, and (e) amphipols. Polar detergent or lipid headgroups are represented by *spheres*, with hydrophobic acyl chains as *straight lines*. The MSPs surrounding the periphery of the nanodisc are shown as *gray rods*, and co-surfactants or co-solvents that stabilize reverse micelles are shown as *triangles*. Two amphipols are shown surrounding the TM segment in (e), with the polar blocks (*gray*) connected to hydrophobic blocks (*lines*) that interact with the protein

lyophilized (e.g., BNip3 [188], EphA1 [189], integrin  $\alpha$ IIb $\beta$ 3 [184]) or acetone-precipitated (the potassium channel voltage sensing domain [153]) states, allowing direct resolubilization into bicelle solutions. Similarly, proteins that are purified into organic solvents can be directly added to an organic solution of bicelle-forming lipids. As shown for Smr, evaporation of the organic solvent followed by resuspension in aqueous buffer can generate a functional bicelle-solubilized receptor [183].

For proteins that cannot be refolded directly into bicelle solutions, it may be possible to transfer a micelle-solubilized sample into a bicelle solution through on-column exchange, as was done for KCNE3 [190]. Alternatively, for proteins that cannot be folded in bicelle-compatible micelles, it may be necessary first to reconstitute into lipid vesicles [186]. Once in this state, sample concentration and buffer exchange can be achieved by centrifugation, followed by resuspension of the pellet in the bicelle-forming detergent solution. This approach was found to be required for the incorporation of the G protein-coupled receptor CXCR1 into small bicelles [191]. However, only extramembraneous N- and C-terminal regions of the sample could be observed in  $^1\text{H}$ - $^{15}\text{N}$  HSQC spectra, and only for very low  $q$  ratio bicelles (0.1).

When working with bicelles, the stability of the bicelle itself may require special attention since hydrolysis of lipid carboxy-ester bonds changes the composition of the bicelle over time, leading to phase separation [192]. This reaction is accelerated under acidic or basic conditions, with the more water-accessible state of the short-chain phase being particularly susceptible to hydrolysis. For this reason, the hydrolysis-resistant ether-linked analog 6-*O*-PC (1,2-di-*O*-hexyl-*sn*-glycero-3-phosphocholine) is being increasingly used to extend bicelle sample lifetimes [186]. It is also possible to use ether-linked long-chain lipids (e.g., 14-*O*-PC or 16-*O*-PC), although this change in the lipid headgroup can alter the structure of embedded proteins [193]. Bicelle stability can also be improved by spiking the solution with charged amphiphiles to increase charge repulsion between bicelles [194, 195]. Both stabilization strategies were used for Smr, with a 3:1 mixture of short- and long-chain ether lipids (6-*O*-PC and 14-*O*-PC), doped with 10% of a 3:1 mixture of short- and long-chain phosphatidylserines [183]. Even with all these modifications to the basic bicelle mix, the sample half-life was typically on the order of  $\sim$ 1 week, illustrating the inherent difficulties of maintaining some membrane proteins in a folded functional state.

One aspect of any bicelle solution that should not be neglected is the significant concentration of monomeric detergent that exists in equilibrium with the bicelle-bound state. These small amphiphilic molecules may be capable of binding to any exposed hydrophobic patches on the protein that protrude from the bicelle [196], potentially destabilizing the protein [165]. Under these circumstances the slightly larger hydrophobicity of ether lipids may confer an additional advantage since this gives rise to a small decrease in monomeric detergent concentrations [197]. More dramatic reductions in monomeric detergent concentrations can be realized by making bicelles with lower-cmc detergents (e.g., C<sub>7</sub>-DHPC; [198]), although the utility of this bicelle type for solution NMR applications has yet to be demonstrated.



### 3.3 *New Developments in Membrane-Mimetic Solvents for Solution NMR*

While micelles and bicelles have served as the membrane-mimetic solvent for the majority of solution NMR studies on integral membrane proteins, the identification of suitable NMR-compatible solvents remains a formidable challenge. With <30 unique multi-spanning membrane protein structures determined by solution NMR, there is a clear need to expand the current library of solubilization agents for this purpose. In response to this problem, creative approaches are currently being developed that may open the door to solution NMR for a wider range of membrane proteins in the future. Some of the systems that have started to yield promising results include nanodiscs, amphipols, and reverse micelles.

#### 3.3.1 Nanodiscs

Nanodiscs are comprised of a bilayer containing ~130–160 lipids, maintained in a discrete, water-soluble state by the association of two copies of the membrane scaffold protein (MSP) from apolipoprotein A-I wrapped around the hydrophobic rim of the bilayer [199–202] (Fig. 3). Originally developed for the solubilization of functionally active integral membrane proteins, they have since been used for solid-state [203], and more recently, solution NMR applications [153, 204–207].

Incorporation of an integral membrane protein in a nanodisc is usually achieved by incubation of a detergent-solubilized sample with detergent-solubilized lipids, along with MSP. Detergent removal, often with adsorbent beads, can promote the spontaneous formation of nanodisc-reconstituted samples that remain stable for many months [202, 208]. However, a number of variables must be optimized when formulating a reconstitution procedure, with the choice of detergent for membrane protein solubilization being identified as a particularly critical one [200]. Even when solubilized in detergents, a membrane protein may undergo irreversible aggregation that cannot be rescued by reincorporation into a native lipid bilayer. Consequently, high yield reconstitution requires a detergent that can minimize the formation of these aggregates, often with the use of excess nanodisc-forming components.

Once reconstituted into a nanodisc, the membrane protein is embedded in a bilayer phase that shares some similarity to that in the bicelle. However, the absence of detergent in the final nanodisc preparation can prevent the folding/stability problems that may be introduced by the significant concentrations of monomeric detergent in bicelle and micelle solutions. The nanodiscs themselves are also more stable than bicelles, and can be subjected to the same manipulations as soluble proteins (e.g., lyophilization, chromatography, concentration, etc.) [200]. However, there is a significant trade-off in that nanodisc sizes tend to be on the large side for NMR applications, at ~200 kDa. While this remains within the accessible range for relaxation-optimized NMR experiments (see Sect. 4.2), most solution NMR spectra

for nanodisc-embedded membrane proteins have shown broad  $^1\text{H}$  line widths [153, 204, 207]. Rapid axial rotation of smaller membrane proteins within the bilayer can significantly reduce these line widths [205]. Unfortunately this is unlikely to occur for the larger membrane proteins that are more prone to require reconstitution in a detergent-free lipid environment. Consequently, the most common application for nanodisc preparations in solution NMR to date has been to provide a reference 2D correlation spectrum of a bilayer-reconstituted sample to allow comparisons with results obtained in less bilayer-like solvents [153, 205, 209]. The GPCR CC-chemokine receptor 5 (CCR5) has also been reconstituted in nanodiscs, allowing its interaction with its chemokine ligand to be studied by solution NMR [210]. Nanodiscs have additionally proven useful for the study of interactions at membrane surfaces (e.g., peripheral phosphoinositide-binding proteins [206]). Meanwhile further developments in nanodisc formulations and protein labeling may help to expand the utility of these complexes for solution NMR in the future.

### 3.3.2 Amphipols

A distinct class of surfactant that is also being developed for a range of membrane protein applications is the amphipathic polymer, otherwise known as the amphipol [211]. These molecules are typically comprised of an amphipathic “backbone” with hydrophobic branches interspersed with polar or charged groups (reviewed in [212–214]). The result is a polymer that could be thought of as a unimolecular micelle with covalent bonds linking polar headgroups (Fig. 3). The equilibrium between monomeric and micellar states that characterizes unlinked detergent solutions is avoided by the covalent linkages, and only a few amphipol molecules are required to envelop and solubilize the protein target.

Transfer of a membrane protein into amphipols can be achieved by sub-cmc dilution of a detergent-solubilized sample into an amphipol solution. The increase in entropy that results from the release of multiple detergent molecules upon amphipol binding makes the exchange highly favorable [215], allowing functional reconstitution of a wide range of membrane proteins [211, 216–218]. Particularly impressive is the ability of amphipols to refold SDS-solubilized GPCRs to greater yields than could detergent/lipid mixtures that had been identified from extensive screening experiments [219]. This allowed the structure of a GPCR ligand in its receptor-bound state to be determined by solution NMR [220]. Meanwhile, feasibility for solution NMR of the solubilized protein itself has also been demonstrated with the transmembrane domain from OmpA [221]. While a small increase in  $^1\text{H}$  line widths was observed in the amphipol relative to the  $\text{C}_6$ -DHPC spectrum, this could be attributed to the slightly larger size of the amphipol-OmpA complex, along with chemical exchange processes. These results with a first-generation amphipol provide an encouraging indication of the potential that future amphipols could have for membrane protein structural biology. The development of amphipols that can remain soluble under more acidic conditions that decrease rates of amide proton

exchange with solvent, along with methods that improve the homogeneity of prepared samples will be particularly important for this application.

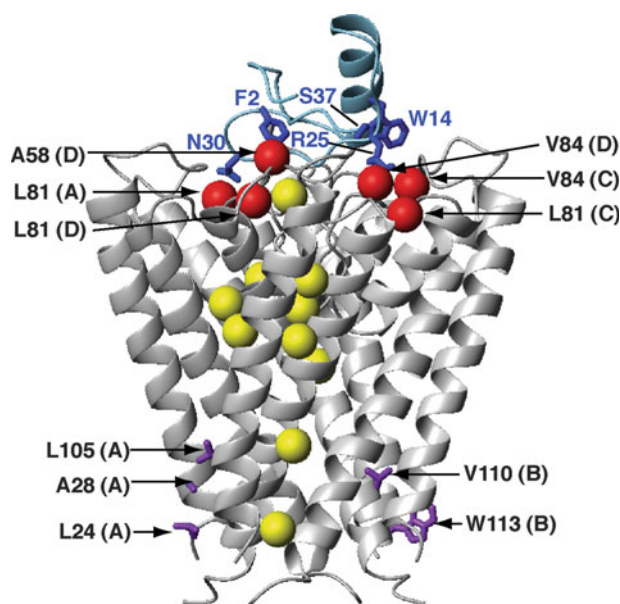
### 3.3.3 Reverse Micelles

A fundamental restriction for solution NMR of solubilized membrane proteins is imposed by the large size of the complexes they make with detergents and lipids, since they usually exhibit slow molecular reorientation rates, giving rise to significant line broadening and decreased coherence transfer efficiencies. To reduce the influence of molecular mass on rotational correlation times, and hence NMR spectral quality, the Wand group has pioneered the use of reverse micelles to encapsulate water-soluble proteins [222–224]. When dissolved in an organic solvent of low viscosity, molecular tumbling rates of these complexes are significantly enhanced compared to those in water. When the increase in molecular reorientation rates exceeds the decrease that comes from the added mass of the reverse micelle, significant improvements in size-sensitive NMR spectral properties can result [225]. For a reverse micelle in pentane this benefit may not appear for a water-soluble protein until its size exceeds ~50 kDa; in contrast, membrane proteins already require detergents for solubilization. Moreover, the benefits of reverse micelle encapsulation can be dramatically improved by the use of solvents with lower viscosities such as butane and propane, which require only slightly elevated pressures to maintain the liquid state in regular NMR tubes [222]. Even larger advantages can be realized by the use of liquid ethane under elevated pressures (~4000 psi) prepared through the use of a special apparatus [226]. In this solvent a >100-kDa protein can be conferred with molecular tumbling properties rivaling those of a 10-kDa protein in aqueous solution.

Although reverse micelle systems were first developed for the study of water-soluble proteins, their application to integral membrane proteins has since been demonstrated with gramicidin A (gA) in a dioctyl sulfosuccinate/pentane system [227]. The pattern of NOEs obtained for the gA peptide dimer in this medium was consistent with its native  $\beta^{6.3}$ -helix. In addition, the intermolecular NOEs showed preservation of native N- to N-terminal intersubunit hydrogen-bonding interactions within the hydrophobic phase. Introduction of gA into this medium was relatively straightforward, since this peptide can be isolated in a lipid and/or detergent-free form and then directly introduced into the reverse micelle system. This was a critical advantage, since the standard reverse micelle systems formulated for water-soluble samples are not compatible with the detergents normally used to purify larger membrane proteins [225].

To extend the utility of reverse micelles to those samples that require detergent or lipid to prevent aggregation and promote folding, reverse micelle formulations have recently been introduced that use detergents capable of forming regular aqueous-phase micelles (e.g., LDAO, dodecyl- or cetyl-trimethylammonium bromide (DTAB or CTAB, respectively)) [228]. These can be used to solubilize membrane proteins in the aqueous micelle conditions usually used during purification.

Dehydration, followed by addition of the organic phase, allows micelles to “flip” out into a reverse micelle configuration while maintaining the hydrophobic domain of the membrane protein in a membrane-mimetic environment. In this complex, each membrane protein is thought to be associated with two reverse micelles that converge around the hydrophobic domain of the protein (Fig. 3) [225, 227, 229]. In some cases co-surfactants (e.g., dihexadecyldimethylammonium bromide (DHAB)) and/or co-solvents (e.g., hexanol) are also required to stabilize this structure. This approach was shown to be successful using the tetrameric KcsA channel as a test system [228]. Reconstitution in reverse micelles was achieved by purification in CTAB, followed by lyophilization and addition of DHAB, hexanol, pentane, and water. The result was a well-resolved  $^1\text{H}$ - $^{15}\text{N}$  HSQC spectrum for KcsA, with significantly enhanced transverse relaxation times relative to those in water (80 ms vs 20 ms). Moreover, potassium ion-dependent chemical shifts in the channel selectivity filter could be observed (Fig. 4), confirming a functional state for this sample. Although protein concentrations in reverse micelles tend to be lower than what can be obtained in



**Fig. 4** Novel insights into KcsA potassium channel structure and function revealed by solution NMR. A ribbon diagram representation of the complex determined for the KcsA tetramer (gray) bound to charybdotoxin (blue) [298]. Highlighted are KcsA methyl groups (red balls) that gave rise to intermolecular NOEs with charybdotoxin side chains (sticks), showing the importance of methyl groups in defining this complex. Selectivity filter residues that were observed to give rise to potassium-dependent chemical shifts are shown in purple [228, 273]. In addition, methyl groups of side chains showing the largest chemical shift differences between active and resting states are shown in yellow for one subunit (residues L24, L40, L59, V70, V76, V91, V95, I100, and L105), indicating widespread changes throughout the TM region between the two states [273]. Residue labels include KcsA subunits A, B, C, and D in parentheses. All ribbon structure figures were made with MOLMOL [398]

aqueous systems (e.g.,  $\sim 0.2$  mM for KcsA), this is offset by the increase in cryogenic probe sensitivity (via an increase in probe quality factor  $Q$ ) that comes from low sample conductance [230]. In general, these results provide an encouraging indication of future benefits that reverse micelles may bring, particularly for the elimination of the requirement for deuterium isotope labeling for some large complexes (described in Sect. 4.1.1).

### 3.4 Approaches for Sample Screening

A consistent theme that has emerged in membrane protein structural biology is that the selection of an appropriate membrane-mimetic solvent is a largely empirical process. This was well illustrated for the pSRII GPCR, which required over 20 different sets of conditions to be systematically explored before C<sub>7</sub>-DHPC was identified as the most appropriate solvent [40]. In this case screening was greatly facilitated by the ability to use the unique absorbance properties of folded pSRII to identify promising conditions rapidly. For those proteins with folding states that cannot be so conveniently monitored it is still possible to screen a wide range of small samples by performing stability measurements. This involves incubation of each sample under conditions that would simulate those of a typical 3D NMR experiment, followed by SDS-PAGE of the soluble fraction separated from aggregated precipitates [37]. Identification of factors that preserve membrane protein solubility can be used to narrow the range of conditions to be examined by more informative but time- and sample-intensive approaches. These include experiments that estimate size (e.g., size exclusion chromatography with or without dynamic light scattering) [40, 138], translational diffusion [146, 231], effective rotational correlation time measurements [70, 138, 207, 232, 233], small angle X-ray scattering [37], analytical ultracentrifugation [234], and spectral quality (e.g., 1D <sup>1</sup>H or 2D <sup>1</sup>H-<sup>15</sup>N HSQC (TROSY) spectra, [37, 149, 159]).

While cryoprobe technology has reduced the sample concentrations required to evaluate the quality of reconstituted samples [235], the need for sample volumes that exceed  $\sim 275$   $\mu$ L limits the number of conditions that can be tested from a single protein preparation. In addition, the detergents themselves can constitute a significant fraction of the expense, particularly since a large excess is required to ensure a 1:1 micelle:protein ratio [111, 236]. However, with the advent of smaller-diameter NMR probes it is now possible to reduce the amount of sample used in each experiment, since the increase in the probe coil length-to-diameter ratio confers an increase in probe mass sensitivity [237–239]. Although acquisition times for biomolecules are lengthened due to the reduction in total sample quantities, it is nonetheless still possible to use sample concentrations similar to those used in conventional probes to obtain 2D spectra within a few hours [240–242]. For example, a 1-mm Bruker TXI microcoil probe with a  $\sim 10$   $\mu$ L sample volume was used to evaluate refolding efficiencies of the *E. coli* outer membrane protein OmpX [151]. In this case it took  $\sim 9$  h to record a <sup>1</sup>H-<sup>15</sup>N TROSY with  $\sim 1.2$  mM detergent

solubilized OmpX. The small quantities involved made it possible to evaluate the NMR spectral properties of refolded samples in a series of novel detergents that were available in a very limited supply. Meanwhile, even faster screening can be achieved if the first set of conditions is evaluated using 1D experiments, with 2D spectra only being acquired for the most promising samples [147]. Also exciting is the recent availability of the cryogenically cooled 1.7-mm microprobe, with an increase in mass sensitivity that compensates for the tenfold reduction in sample volume (30  $\mu\text{L}$  vs 300  $\mu\text{L}$  in a conventional 5-mm probe). With spectral acquisition times and sample concentration requirements being comparable to those of the 5-mm room temperature probe, it is in fact now possible to use these reduced sample volumes for acquisition of all NMR data types required [147, 232].

One aspect of detergent screening that is gaining increasing attention is the concentration of the detergent itself. While it has long been recognized that there is a minimum concentration required for solubilization, samples that exceed this minimum can also exhibit poor spectral quality [232]. Screening experiments with a 1.7-mm microprobe showed that the degradation in OmpX spectral quality associated with higher decylphosphocholine concentrations was caused by an increase in the solution viscosity due to the presence of excess protein-free micelles. Based on these results it could be shown that the optimal detergent concentration for this system approximately corresponded to the number of detergent molecules bound in the protein–detergent complex multiplied by the concentration of OmpX, added to the cmc of the detergent. High detergent-to-protein ratios can also complicate the study of oligomeric states of membrane proteins, with excess micelles reducing the effective concentration of TM segments that can interact [243–245]. However, the practice of concentrating samples in ultracentrifugation devices can increase the concentration of detergent in the final sample, even when the molecular weight cut-off is smaller than the size of the protein-free micelle [234, 246]. This can be particularly problematic for detergents that tend toward greater polydispersity and concentration dependence of micelle size. Consequently, it is important to choose ultrafiltration units with molecular weight cut-offs that can maximize retention of the target protein while minimizing undesired increases in detergent concentrations, particularly for any samples that require concentration or detergent exchange as a final preparation step. It is useful to monitor detergent concentrations during these procedures, with detergent peak intensities in 1D  $^1\text{H}$  NMR spectra providing a convenient read-out of concentration for many of the commonly used detergents for solution NMR [234].

## 4 Methods for Acquisition of Solution NMR Spectra of Membrane Proteins

As discussed above, most of the systems used to reconstitute a membrane protein sample for solution NMR create large, slowly tumbling complexes that suffer from rapid transverse relaxation processes. Size estimations of some of these

protein–detergent complexes show that the detergent can add as much as 60 kDa to the protein molecular weight (e.g., DAGK [247], VDAC-1 [73]). In order to counter the loss of NMR signal sensitivity that is characteristic of these large complexes, it is usually necessary to use approaches originally developed for study of large water-soluble proteins. Those that have found the most widespread utility are deuterium isotope sample labeling and relaxation-optimized pulse sequences.

## ***4.1 Isotope Labeling Schemes for Large Protein Complexes***

### **4.1.1 Deuteration**

To compensate for the unfavorable relaxation properties of protons in large protein complexes, all the non-exchangeable carbon-bound protons can be replaced with deuterium (reviewed in [20, 248, 249]). As the gyromagnetic ratio of  $^2\text{H}$  is 6.7-fold lower than that for  $^1\text{H}$ , many of the relaxation pathways that would otherwise be present in fully protonated samples are greatly attenuated in a deuterated sample. This leads to an increase in  $T_2$  transverse relaxation times, and more effective preservation of signal over the numerous coherence transfer elements that occur in a typical multidimensional NMR pulse sequence. Significant sensitivity gains have been demonstrated for uniformly deuterated proteins in many triple resonance applications with both water-soluble [20] and membrane protein samples [111]. As a result, uniformly deuterated samples are now routinely used to obtain backbone assignments for large membrane proteins.

Since backbone assignment experiments use amide proton magnetization to generate and detect the NMR signal, deuterated backbone amides must first undergo complete exchange with solvent  $^1\text{H}_2\text{O}$  protons [20]. Although the extended exposure to  $^1\text{H}_2\text{O}$  solutions during purification can be sufficient to re-introduce protons at all sites, many of the larger membrane proteins that exhibit good spectral properties are highly stable and therefore resistant to this process [112]. In these cases re-introduction of protons at exchangeable sites can be extremely slow in the core, making it impossible to achieve significant exchange without the exposure to destabilizing conditions. This can be particularly challenging for proteins that have been optimized to maximize stability, potentially necessitating the development of high-yielding unfolding/refolding protocols. This was done for the 4-TM DsbB, which had to be solubilized in a DPC/SDS mixture for solvent exchange before reconstitution back into DPC [39].

One way to get around this problem is to treat the protein as a two-domain system, since non-exchangeable sites tend to be clustered together for several contiguous residues along the protein sequence [250]. As shown for the tetrameric potassium channel protein KcsA in SDS, a complementary pair of samples can be generated; one that retains amide protons at the solvent-exposed sites, and the other with protons at non-exchangeable sites [138]. The former sample is generated using conventional protocols for perdeuteration followed by amide-proton exchange with

solvent  $^1\text{H}_2\text{O}$ . However, for the latter sample, expression is done in a  $^1\text{H}_2\text{O}$ -based medium that is deficient in glucose and instead contains deuterated amino acids from an algal lysate. Subsequent exchange of the purified sample into  $\text{D}_2\text{O}$  allows simplification of the amide spectrum, with only protected site remaining visible. Since the NMR spectral properties of micelle-embedded core regions of membrane protein structure tend to differ from solvent-exposed regions, use of these two labeled samples allows assignment strategies to be specifically tailored to each region.

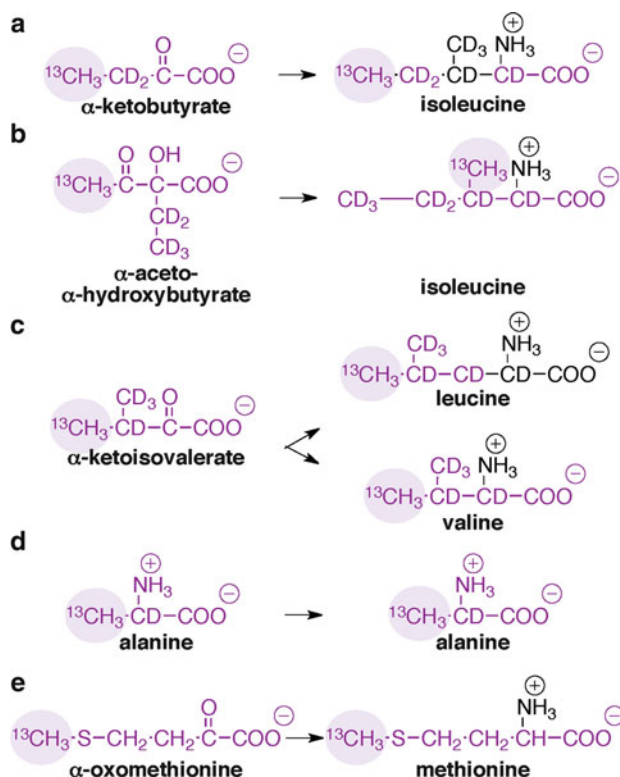
#### 4.1.2 Methyl Protonation

Despite the utility of perdeuteration in the assignment of backbone chemical shifts, the elimination of all but the exchangeable protons impedes structural studies that rely on conventional NOE-based approaches. Although in some cases it is possible to use only amide proton NOEs to obtain a protein global fold, the accuracy of these structures tends to be low due to the small proportion of distance restraints between protons from non-sequential residues (e.g., 5–8 Å backbone pairwise rmsd to target structure [251, 252]). Therefore to increase the number of protons in the protein core while maintaining the benefits of extensive deuteration, a number of methods have been developed to retain protons at specific non-exchangeable sites using a “reverse isotope” labeling approach [20, 248, 249]. Methyl groups have been the principal targets for selective protonation since they are enriched in protein hydrophobic cores [253], making them structurally informative sources of NOE-based restraints [251, 254, 255]. In addition, rapid rotation about the methyl symmetry axis causes its three protons to give rise to a single peak with narrow  $^1\text{H}$  line widths that are additionally narrowed for methyls that terminate flexible amino acid side chains [256]. In the context of a slowly tumbling macromolecule, this rapid methyl rotation also creates ideal conditions for the optimization of relaxation in HMQC-type experiments (methyl-TROSY, described in Sect. 4.2).

A widely used strategy for the selective introduction of methyl protons into deuterated proteins is known as the ILV method, since targeted methyl groups reside in the amino acids Ile, Leu, and Val (reviewed in [256–258]). In this method, deuterated  $\alpha$ -keto acid precursors retaining protons at methyl sites are added to a bacterial expression culture growing in  $\text{D}_2\text{O}$  minimal media approximately 1 h prior to induction of protein expression (Fig. 5) [259, 260]. To maintain a high background of deuterium incorporation, uniformly deuterated glucose is included in the growth as the only other source of carbon. After a relatively short induction period (~4–6 h) designed to maximize incorporation and minimize metabolic scrambling [259], methyl protons from  $\alpha$ -ketobutyrate and  $\alpha$ -ketoisovalerate will be incorporated into the Ile( $\delta$ 1), and Leu/Val methyl groups, respectively. There is a large variety of isotope label combinations available for these  $\alpha$ -keto acids [258]; those shown in Fig. 5 have been the most commonly used for the study of large proteins [257].

The ILV approach has been applied to structure determination of membrane proteins from the  $\beta$ -barrel family, as first shown for OmpX in DHPC [261].





**Fig. 5** Metabolic precursors added to expression media prior to induction of expression for the selective incorporation of  $^{13}\text{C}$  and  $^1\text{H}$  into methyl groups of  $^{12}\text{C}$ ,  $^2\text{H}$ -labeled (a) Ile (C $\delta$ 1-labeled), (b) Ile (C $\gamma$ 2-labeled), (c) Leu and Val, (d) Ala, and (e)  $^{13}\text{C}$ ,  $^1\text{H}$ -Met. The parts of the amino acid that originate from each precursor are shown in purple, with the  $^{13}\text{C}$ ,  $^1\text{H}$ -labeled methyl group highlighted. In all cases the expression medium must include  $^{13}\text{C}$ ,  $^2\text{H}$ -labeled glucose as the carbon source. Selective methyl labeling of Ala also requires media supplementation with deuterated succinate,  $\alpha$ -ketoisovalerate and isoleucine to suppress isotope scrambling. In this procedure glucose and glycerol together are used as the carbon sources prior to addition of methyl labeling agents [263]

NOEs involving methyl protons yielded a fourfold increase in the number of distance restraints over  $^1\text{H}^{\text{N}}\text{-}^1\text{H}^{\text{N}}$  NOEs alone [71], and a twofold increase in NOE restraints for the 210 residue KpOmpA  $\beta$ -barrel [70]. Smaller gains were obtained for the 283 residue VDAC-1, with 324 methyl-associated NOEs adding to the 288 amide proton NOEs [73]. However, in all these cases, the overall fold of the  $\beta$ -barrel is already well-defined by NOEs between amide protons, and hence the impact of methyl NOEs on structure quality tends to be modest, albeit significant, for these folds (e.g., OmpX NOEs from methyls decreased the backbone rmsd to the mean from 2.13 Å to 1.42 Å, [71]).

While the utility of methyl protonation for structure determination of  $\beta$ -barrel folds has been demonstrated, the full benefits of this strategy for structure

determination of helical integral membrane proteins has been more difficult to realize. Even in the case of helical water-soluble proteins it was noted that fewer long-range methyl-based restraints are available relative to other fold types [251]. Similar studies focusing on all-helical membrane protein folds substantiated this observation, with structure accuracies in the 5 Å region even when complete assignment of methyl-associated NOEs was assumed to be possible [262]. In reality the low spectral dispersion that is characteristic of helical membrane proteins can significantly impede chemical shift and NOE assignment for these methyl groups, as was the case for the DAGK trimer [111]. On the other hand, relatively high spectral quality for the pSRII GPCR allowed ~50% of expected inter-helical NOEs involving these methyl groups to be assigned, with the remainder being either absent from the spectrum or buried under strong diagonal signals [40]. Yet this level of NOE assignment was still not sufficient to generate high-resolution structures, making it was necessary to acquire a greater number of long-range restraints involving Ala, Thr, Ile( $\gamma$ 2), and Met methyl group NOEs. In the case of pSRII it was possible to use the ILV methyl assignments to help extend the chemical shift assignments to these methyl groups using a fully protonated sample. This afforded a ~2.5-fold increase in the number of inter-helical NOEs that could be assigned and helped to increase the quality of the resulting structures.

When the spectral quality of a fully protonated sample is not sufficient to help increase the number of methyl proton assignments, other methyl protonation strategies are also available. For example, methyl protonated Ala can be directly incorporated into deuterated proteins so long as a trio of deuterated precursors (Fig. 5) that suppresses metabolic scrambling to undesired sites is also added to the media. [263]. The direct bond between the Ala methyl group and C $\alpha$  backbone atoms make these methyl groups excellent sources of information for backbone structure and dynamics [264]. Importantly, the ILV and Ala-methyl labeling strategies are complementary, allowing protons to be simultaneously introduced to ILV and Ala methyl groups in a single sample. Ile( $\delta$ 1) and Ala methyl groups can also be simultaneously labeled in an alternate strategy that uses deuterated rich media supplemented with appropriately labeled Ala and  $\alpha$ -ketobutyrate [265]. Methionine methyls can similarly be targeted by the inclusion of protonated Met [266] or a selectively protonated  $\alpha$ -keto acid derivative [267] into the D<sub>2</sub>O minimal expression medium (Fig. 5). The C $\gamma$ 2 proton of Ile can also be targeted by including  $\alpha$ -aceto- $\alpha$ -hydroxybutyrate in the minimal media [268]. Meanwhile, selective incorporation of protons beyond the methyl groups has been demonstrated in the stereo-array isotope labeling (SAIL) strategy [269, 270]. This technique uses cell free expression to incorporate a complete suite of synthetically prepared stereo- and regiospecifically <sup>2</sup>H-labeled amino acids, producing a sample with reduced <sup>1</sup>H density that retains a larger number of structurally informative protons. Overall, the general utility of these various strategies will depend on the characteristics of each sample, as well as cost-effective availability of labeled precursors. However, they should prove to be increasingly useful as the number of large membrane proteins being studied by solution NMR continues to increase.

### 4.1.3 Exploring New Frontiers Using Methyl Protonated Samples

Although methyl protons have been widely used to assist structure determination of large protein complexes, their high sensitivity and amenability to TROSY-type sequences (described in Sect. 4.2) has provided the foundation for exciting applications that probe large systems in dynamic equilibrium [256, 257]. This was spectacularly demonstrated for the 670-kDa 20S core-particle proteasome, with the identification of two interchanging structural states for gating residues using  $^{13}\text{CH}_3$ -labeled Met in  $\text{U-}^2\text{H}$ ,  $^{12}\text{C}$ -labeled samples [271], along with interactions inside the proteasome antechamber that maintained ILV-labeled substrate proteins in an unfolded state [272]. Meanwhile, applications with membrane protein samples are also beginning to appear, as shown with the DDM-solubilized KcsA channel [273]. In this study, methyl TROSY spectra from  $\text{U-}^2\text{H}$ , Leu/Val- $[\text{}^{13}\text{CH}_3, \text{}^{12}\text{CD}_3]$ , Tyr- $^1\text{H}$  KcsA showed pH-dependent chemical shifts reflecting interconversion between three different functional states, and coupling between gating residues in two distinct regions of the structure (Fig. 4). Selectively ILV-labeled samples have also been used to characterize the spectroscopically-invisible vesicle-bound states of  $\alpha$ -synuclein by transfer-NOE type experiments involving Leu methyl protons that were detected in the  $^1\text{H-}^{15}\text{N}$  HSQC spectrum of the free state [274].

In an even more ambitious application, methyl group signals allowed ligand-induced changes in the extracellular surface of the 365-residue human  $\beta_2$  adrenergic receptor in DDM to be investigated without the assistance of deuterium labeling [275]. In this study  $^{13}\text{C}$ -methyl groups were introduced through reductive methylation of solvent-exposed Lys  $\epsilon$ - $\text{NH}_2$  groups. The presence of a salt-bridge interaction on the extracellular side could be confirmed, with changes in this interaction being detected in the  $^1\text{H-}^{13}\text{C}$  HMQC spectrum upon agonist binding. This data provided new insights into the structure of the active state, a form that had not yet been captured by X-ray crystallography. The relatively large signal-to-noise ratio provided by methyl groups at the ends of the long flexible Lys side chain was instrumental for the success of this experiment, particularly given the low concentrations of samples that were available (60–200  $\mu\text{M}$ ). These experiments illustrate the great potential for methyl protons to provide a window into previously inaccessible dynamic states of large membrane proteins.

## 4.2 Relaxation Optimization

Given the large size of typical membrane protein samples, selective labeling approaches must usually be accompanied by relaxation-optimized NMR experiments. Just over a decade ago, rapid transverse relaxation rates associated with large, slowly tumbling molecules defined a  $\sim 30$ – $40$  kDa molecular weight ceiling for protein NMR. Beyond this limit, broad peaks and magnetization losses during coherence transfer thwarted attempts to acquire standard spectra required for

structure determination. The introduction of transverse relaxation optimized spectroscopy (TROSY) by Wuthrich and co-workers [276] opened up a wealth of new opportunities for solution NMR on large protein systems, including detergent-solubilized membrane proteins.

First demonstrated for amide  $^1\text{H}$ - $^{15}\text{N}$  [276] followed by aromatic  $^1\text{H}$ - $^{13}\text{C}$  groups [277], the basic idea exploits interference effects between the two dominant mechanisms of relaxation, namely dipole-dipole coupling and chemical shift anisotropy (CSA). The effect of this interference can be seen in an uncoupled COSY spectrum, which for a directly bonded  $^1\text{H}$ - $^{15}\text{N}$  pair gives rise to a doublet in both the  $^1\text{H}$  and  $^{15}\text{N}$  dimensions, each separated by the one-bond coupling constant  $^1J_{\text{HN}}$ . In this quartet, one peak will exhibit a narrower line width relative to the other peaks as a result of partial cancellation of dipole-dipole and CSA relaxation pathways. This effect is most pronounced at higher magnetic field strengths, with maximal cancellation occurring at  $\sim 1$  GHz proton frequency for both  $^1\text{H}$  and  $^{15}\text{N}$  magnetization. In TROSY-type NMR pulse sequences, the coherence that gives rise to the slow-relaxing component of the quartet is preserved while the other components are either eliminated [276] or allowed to decay to negligible levels during the course of the experiment [278, 279]. Although all four coherences would make contributions to the single  $^1\text{H}$ - $^{15}\text{N}$  peak in the decoupled correlation spectrum that is normally run for smaller proteins, interconversion between states that undergo very fast vs slow relaxation give rise to magnetization losses that are dominated by the rapidly decaying coherence. Consequently, for larger proteins ( $\sim 50$  kDa and higher) the signal lost by discarding three parts of the quartet is compensated by the increase in resolution and sensitivity that comes from avoiding the more rapidly relaxing states during acquisition.

Numerous multidimensional triple-resonance NMR experiments have since been designed that isolate this slow-relaxing component during evolution and acquisition periods (reviewed in [280]). Additional strategies that take advantage of the unique relaxation properties of slowly tumbling proteins have since been described that provide additional sensitivity gains for large systems (e.g., polarization transfer schemes that use cross-correlated relaxation [281, 282], or longitudinal relaxation optimization [283]). This has allowed backbone assignments to be obtained for many proteins exceeding 40 kDa, including a number of large membrane protein-detergent complexes [110, 284]. This has usually required the use of deuterium-labeled samples to eliminate  $^1\text{H}$ - $^1\text{H}$  dipole-dipole interactions that also contribute significantly to transverse relaxation rates. In some cases it can be possible to skip deuterium labeling, particularly when only 2D  $^1\text{H}$ - $^{15}\text{N}$  correlation spectra are required, as was the case when screening sample conditions for DAGK [285]. However, the ability to obtain backbone resonance assignments for typical membrane protein samples requires that the TROSY effect be maximized through the use of uniform deuterium incorporation and high spectrometer field strengths (i.e., 700 MHz and greater). Even when these conditions are fulfilled, peak intensities in the TROSY spectrum may be attenuated by microsecond to millisecond timescale exchange processes. This has been commonly observed for membrane proteins in detergent micelles [37, 70, 111, 209, 285-287], making the

optimization of sample conditions that can minimize exchange on this timescale a particularly important consideration for membrane proteins [155].

Relaxation optimization strategies have since been extended to side chain aliphatic groups with the important recognition by Kay and coworkers that the coherence transfer pathways available to a rapidly rotating, isolated  $^{13}\text{C}^1\text{H}_3$  spin system in a slowly tumbling macromolecule undergoes different rates of relaxation [288, 289]. Most significantly, they established that the HMQC is a TROSY type of experiment for isolated  $^{13}\text{C}^1\text{H}_3$  systems since, unlike the HSQC experiment, there is no interconversion of slowly- and rapidly-relaxing coherences. The benefits of this effect are optimal for  $^2\text{H}$ ,  $^{12}\text{C}$ -labeled proteins with specific incorporation of a single  $^{13}\text{C}^1\text{H}_3$  group in a subset of amino acid types (e.g., ILV [290]). Deuterium spin relaxation measurements in  $^{13}\text{CHD}_2$ -labeled methyl groups can also take advantage of this relaxation optimization strategy [291], opening the door to the study of dynamic structural states for extremely large protein complexes, as demonstrated for the 670-kDa 20S proteasome core particle [271].

Although unassigned methyl group resonances can be used as reporters of structural transitions in large proteins [257, 292], any requirement for site-specific information requires that these methyl peaks be assigned. This is straightforward when backbone assignments are available, since TROSY-type sequences can be used to correlate methyl proton or carbon shifts to amide  $^1\text{H}$  and  $^{15}\text{N}$  chemical shifts [293–295]. Alternatively, in some cases it has been possible to transfer assignments made on smaller fragments or isolated subunits to these same components in the intact complex [257, 266, 295]. Meanwhile, efforts are also being made to automate the process of methyl shift assignment, with one protocol using X-ray structure-based chemical shift predictions along with methyl-NOESY data to correctly assign 99% of a 300-kDa ILV-labeled proteasome [296]. This approach may become important for large single-chain membrane protein–detergent complexes, particularly since these samples would not easily be adapted to the fragment-based approach, and  $^1\text{H}$ – $^{15}\text{N}$  spectral quality may not be sufficient for through-bond correlations with methyl groups. However, in the event that backbone assignments are not available or methyl correlation spectra do not allow unambiguous assignment, methyl shift assignments can still be made by using spectra of mutants [266], sometimes paired with NOEs, or paramagnetic relaxation enhancement patterns [297].

## 5 Strategies for Membrane Protein Structure Determination by Solution NMR

### 5.1 NOE-Based Methods of Structure Determination

The approach used for structure determination of a membrane protein varies significantly depending on the size of the complex, the protein fold, and the quality

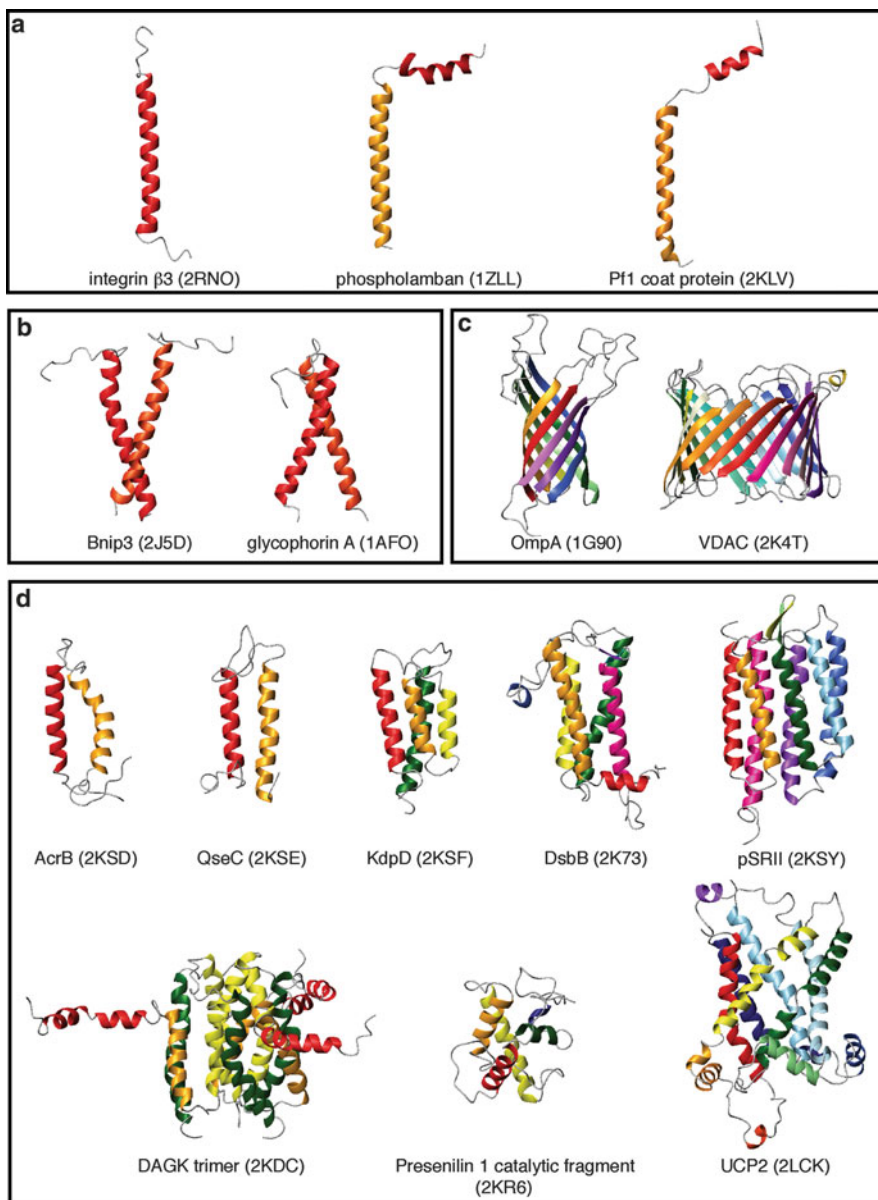
of spectra that can be attained. When high quality spectra of relatively small detergent–protein complexes are available, conventional NOE-based approaches can be applied in a straightforward manner without the need for deuterium labeling or TROSY type experiments. This tends to be the case for structures of single TM helices in monomeric or dimeric forms, examples of which are shown in Table 2, Method I, Fig. 6a. Structures of larger oligomeric states formed by single TM helices have also been determined in this way, although some  $^2\text{H}$ -labeling and TROSY-based backbone assignment experiments were required to deal with the larger size of these complexes (Table 2, Method II, Fig. 6b).

When the length of the polypeptide exceeds that of these single TM-helix constructs, TROSY-based experiments and uniform deuteration are usually used to obtain backbone assignments. In cases where spectral quality or complexity prevents straightforward assignment of backbone resonances from uniformly  $^{15}\text{N}$ ,  $^{13}\text{C}$ ,  $^2\text{H}$ -labeled samples, it is possible to increase the number of assignments with samples selectively  $^{15}\text{N}$ - or  $^{15}\text{N}$ ,  $^{13}\text{C}$ -labeled with a single amino acid. For example, ~20 different selectively labeled samples were used to help assign backbone atoms in the human mitochondrial VDAC  $\beta$ -barrel [73]. Alternatively, if the complex being studied is comprised of more than one polypeptide chain, backbone assignment can be assisted by combining labeled and unlabeled subunits, as was done for the heterotrimeric natural killer cell-activating complex [65]. Additional assignments are also sometimes accessible through the use of a range of temperatures [184] or pH conditions [40] that allow observation of different subsets of resonances. As a last resort it is also possible to use mutagenesis [273, 297, 298], although this approach is not practical for the assignment of a large number of residues.

In cases where spectral quality allows a high level of backbone assignments to be made, it is usually possible to acquire assignable NOEs between amide protons. For  $\beta$ -barrel folds these NOEs can be used with backbone torsion angle input derived from secondary backbone shift data [299], and hydrogen bond restraints inferred from amide solvent exchange to construct an informative global fold (Table 2, Method III, Fig. 6c) [68, 69, 72]. Selectively methyl labeled ILV has also been used for this fold type as an additional source of NOEs in some cases (Table 2, Method IV, Fig. 6c) [70, 71, 73]. For helical membrane proteins, more NOEs are required to obtain a structure of comparable precision (as described in Sect. 4.1.2), usually by extending chemical shift and NOE assignments to other side chain atoms. This requires that 3D  $^1\text{H}$ – $^1\text{H}$  correlation spectra on partially or fully protonated samples are of sufficient quality for assignment, an achievement that that has only been realized in a very small number of cases (see Table 2, Method V).

## 5.2 Non-uniform Sampling

One of the most challenging membrane protein structures determined to a high resolution via an NOE-driven approach is the pSRII GPCR (Fig. 6d), which required additional sensitivity-enhancing approaches to be employed during data



**Fig. 6** Representative samples of integral membrane protein structures solved using solution NMR, comprised of either (a) a single TM helix, (b) a TM helix dimer, (c) a  $\beta$ -barrel, or (d) multiple TM helices from a single chain (with the exception of DAGK, which is a trimer of 3-TM helix subunits). Coloring from the N terminus to C goes from red to orange, yellow, green, blue, purple, and then pink. All PDB accession numbers are indicated in parentheses for each structure

acquisition [40, 145]. Specifically, many of the 3D spectra were acquired using non-uniform sampling in all indirectly detected dimensions [300]. In this method, data is only acquired for a subset of the incremented evolution periods, with more frequent sampling in regions of short evolution times that have larger signal-to-noise ratios [301]. The emphasis on data acquisition in parts of the time domain having the greatest signal intensity leads to increases in sensitivity per unit time over conventional uniformly sampled experiments [302]. Maximum entropy reconstruction was used to process this sparsely sampled dataset, producing multidimensional spectra that retained the resolution of the conventionally sampled experiment in spite of the reduced number of indirect points [303, 304]. In the case of pSRII, the signal-to-noise ratio was further enhanced by adding multiple datasets acquired on one or more samples [40], a useful approach when sample stability limits the time that can be used to acquire a single dataset.

Non-uniform sampling (NUS) also opens the door to experiments of increased dimensionality that might normally require several months of acquisition time to obtain adequate resolution in the indirectly detected domains [302, 305]. This strategy was used to address the problem of chemical shift degeneracy in methyl NOESY data for the 283-residue VDAC  $\beta$ -barrel (Fig. 6d) [73]. Four-dimensional spectra were acquired on this sample to correlate  $^1\text{H}$  shifts with the shift of the directly attached heteroatom on both sides of the NOE interaction [73, 306]. These spectra were processed using multidimensional decomposition (MDD) [307, 308], an approach that is particularly well-suited for the accurate reproduction of signal intensities in spectra having a wide range of peak intensities, as is seen in NOESY spectra [309]. Moreover, significant sensitivity increases for weak peaks were realized by combining multiple datasets during MDD. This technique, called coupled MDD, was used to process a  $^{13}\text{C}$ - $^{13}\text{C}$  separated 4D [ $^1\text{H}$ - $^1\text{H}$ ] NUS-NOESY by co-processing with a  $^1\text{H}$ - $^{13}\text{C}$  HMQC template spectrum [310]. Coupled MDD allowed NOEs involving the 15% of VDAC methyl groups that did not appear in the MDD-processed spectrum to be detected and assigned, many of which were critical for determination of the global fold.

Overall, the application of non-linear sampling and other random sampling techniques to membrane proteins is still in its early days, and does not reflect the diversity of sparse sampling techniques that have been developed to date (reviewed in [302, 305]). Nonetheless these studies demonstrate the significant advantages that can be gained by applying these data acquisition strategies to challenging membrane protein samples, which should motivate an increasing number of research groups to adopt this technology.

### ***5.3 Global Folds in the Absence of Long-Range NOEs***

Even when sample labeling and NMR experiments are optimized to minimize unfavorable relaxation processes and increase sensitivity, it is not unusual for difficulties in side chain and NOE assignment to preclude an NOE-driven approach



to structure determination. In addition, some smaller membrane proteins are comprised of structural elements that only weakly interact with each other, if at all, and would therefore not be expected to give rise to long-range NOEs. Regardless of whether the absence of assignable long-range NOEs is due to structural or spectroscopic causes, in both cases distance restraints are accessible using approaches that do not require assignment of side chain chemical shifts or long-range NOEs. This approach has been used to determine structures for a significant number of membrane proteins (Table 2, Method VII), with residual dipolar couplings (RDCs) and paramagnetic relaxation enhancement (PRE) data being the primary sources of long-range distant restraints.

### 5.3.1 Measurement of RDCs

Although rapid molecular reorientation rates in solution are responsible for the high resolution of liquid-state NMR spectra, this also gives rise to isotropic averaging that prevents the observation of orientation-dependent phenomena. However, it is possible to detect these orientation-dependent effects while retaining the advantages of solution phase spectroscopy by subjecting the protein of interest to a weakly aligning environment [311]. The resulting non-zero averaging allows data such as the RDC and residual chemical shift anisotropy (RCSA) to be acquired and used for structure determination.

From a practical viewpoint, the generation of weakly aligned membrane protein samples poses additional challenges relative to the case for water-soluble proteins. While the first RDC measurements on water-soluble proteins used liquid-crystalline phases of bicelles to induce weak alignment [312], the extension of this measurement to membrane protein samples awaited the development of detergent-resistant alignment media. The first, and one of the most widely used, type of detergent-compatible systems is the aqueous phase of mechanically strained polyacrylamide gels [313]. Asymmetry in gel pores can be straightforwardly created with the assistance of a commercially available gel-stretching apparatus [314], or simply by applying axial compression to gels with diameters smaller than that of the NMR tube interior [313, 315]. A specialized type of NMR tube has also been designed that makes it possible to measure both isotropic and aligned states in the same gel sample, a particularly important advance for accurate measurement of small RCSAs [316].

Variations in the degree of alignment can be obtained by using different gel densities, aspect ratios, and gel cross-linking ratios [313, 314]. However, at the gel concentrations required to maintain mechanical stability of the gel (>4% w/v), diffusion of large protein–detergent complexes into the pores can be hindered, reducing the amount of protein that can be introduced into this medium [68, 247]. Consequently, the first RDC measurements on membrane proteins were done with smaller protein–detergent complexes using higher gel concentrations [137, 160, 314, 317, 318]. More recently it was demonstrated that the ~100-kDa DAGK-DPC complex could be introduced into a 4% hydrated gel by allowing ~2 days for the

system to come to equilibrium [38]. This method differed from typical protocols previously reported since a solution containing the NMR sample is usually used to rehydrate the dried gel, thereby reducing sample losses from dilution. However, this was not possible for the 4% gel, since the presence of detergent in the rehydrating solution did not allow the original gel dimensions to be restored [112]. This reflects the trade-off in mechanical stability that comes with the reduction in polyacrylamide concentrations required to make gels with larger pore sizes.

To address some of the difficulties in generating gel-aligned membrane protein samples, protocols have been developed to perform the gel polymerization in a solution already containing the protein sample, or to introduce an SDS-solubilized sample into a gel via electrophoresis [319]. In the case of copolymerization, a protocol would need to be carefully tailored to the specific features of each protein-detergent system being studied to minimize the potential for undesired covalent modifications during the relatively non-specific polymerization reaction [112]. However, a more general alternative to these neutral gels has also been developed using charged polyacrylamide-based copolymers [320, 321]. Rehydration of charged copolymers is facilitated by a strong electroosmotic effect, providing an increase in mechanical stability over neutral polyacrylamide gels. Consequently, alignment can be reproducibly achieved in the presence of detergents at gel concentrations as low as 2–3%. The larger pore sizes of these lower concentration gels, along with the higher mechanical stability, make it possible for larger protein–detergent complexes to be introduced into the gel by rehydration. This has allowed RDC measurements to be obtained for the 4-TM helix DsbB [39] and the 8-strand OmpA protein [320] both in complex with DPC micelles. In addition, alignment has been demonstrated using copolymers with either positive or negative charge, as well as mixtures of positively and negatively charged copolymers [320, 321]. This permits the electrostatic properties of the alignment medium to be tailored to the charge properties of the sample, and different alignment frame orientations to be generated for the measurement of more than one set of RDCs.

One aligning medium that has been extensively used for soluble proteins is filamentous bacteriophage [322], which was recently been shown to be compatible with modest concentrations (~100 mM) of phosphocholine-based membrane-mimetics so long as the pH of the solution exceeds ~6 [323]. These long negatively charged filamentous structures were also the inspiration for the design of a novel detergent-resistant aligning medium formed by DNA nanotubes [324]. These are self-assembled from concentrated solutions of a 7.3-kb “scaffold” strand mixed with a 170 base “staple” strand, and have been used to induce alignment in the  $\zeta\zeta$  TM dimer [324], and the mitochondrial uncoupling protein 2 [41], both using detergent concentrations in excess of 150 mM. A potentially more convenient alternative has also been developed that generates a similar type of macromolecular structure from the potassium salt of the dinucleotide 2'-deoxyguanylyl-(3',5')-2'-deoxyguanosine [325]. These dinucleotides form G-tetrads that stack into columns with dimensions similar to those of bacteriophage, and have been used for the measurement of RDCs in the LMPG micelle-associated cytoplasmic domain of the

influenza B proton channel [326] and the  $\alpha$ IIB $\beta$ 3 TM helix heterodimer in isotropic bicelles [184].

With the exception of the strained polyacrylamide gel, many of the detergent-resistant media described possess significant charge that can give rise to unfavorable interactions with the alignment media, potentially broadening spectra beyond detection [327]. An alternative that avoids this problem exploits the significant anisotropy of the magnetic susceptibility of lanthanide ions such as Yb<sup>3+</sup> and Dy<sup>3+</sup>. This leads to weak magnetic field-induced alignment that allows RDCs of lanthanide-bound species to be measured [328]. For proteins that do not have a high-affinity metal-binding site, it is possible to introduce a metal chelator via covalent modification of a unique cysteine residue [327, 329, 330] or by engineering a metal-binding EF-hand as was done for the HIV Vpu protein [331] and a modified form of OmpA [332]. However, to maximize the information content from these experiments, high affinity tags with a single susceptibility tensor that is minimally affected by dynamics should be used. For this purpose, sulfhydryl-reactive EDTA based analogs have been designed that, unlike the non-stereospecifically metal-binding precursor, create a single stereoisomer of the metal-bound complex [333].

### 5.3.2 Applications of RDCs in Structure Determination

The magnitude of the RDC depends on the degree of alignment induced in the sample, the inter-nuclear distance separating the two atoms involved, the orientation of the internuclear bond vector with respect to the external magnetic field, and the local dynamics of the atoms involved [311]. When measured for a set of covalently linked backbone atoms with similar dynamic properties, RDCs can be used to determine the relative orientation of the corresponding peptide planes. A large number of NMR experiments are now available to measure these types of RDCs (reviewed in [334]) and inclusion of at least one type of RDC is often performed in standard protein structure determinations (reviewed in [335, 336]), including those for membrane proteins [39, 152, 162, 166, 326, 337, 338]. However, RDCs alone are not enough to determine a unique protein fold when starting from an extended structure, since there is 16-fold degeneracy in the number of peptide plane orientations consistent with the measured couplings [339]. Consequently, RDCs cannot distinguish between structures of different quality at early points in a structure calculation when differences between calculated and actual structures are large (i.e., rmsd > 10 Å) [340, 341]. RDCs are instead used to refine NOE-based structures via a scalable energy function that gradually increases the magnitude of the RDC energy penalty as the NOEs drive the fold closer to its final state [342]. This allows the structure to be oriented into its alignment frame before the couplings refine the structure, an approach that can lead to an increase in the quality and precision of an NOE-based structure [342, 343]. However, these benefits are usually only realized

when multiple RDCs per peptide plane, or RDCs from more than one type of alignment medium, are available [335, 339].

When structures of secondary structure elements or globular domains are known, it is possible to use RDCs to calculate the relative orientation of these structural units. First demonstrated for multi-domain proteins [344, 345], a similar strategy has since been used to determine average protein structures for small membrane proteins in the absence of other long-range restraints, using TM helices as the rigid, or tightly restrained bodies [137, 318, 346]. However, for each structural unit there will be four equivalent orientations that can be defined with respect to a right-handed alignment frame (assuming that the alignment tensor is not axially symmetric) [347]. Hence it is necessary to eliminate this degeneracy through the measurement of two or more sets of RDCs under different aligning conditions [348, 349] or by the recognition of structural restraints that can rule out extraneous solutions [318, 347, 350, 351].

A related application for RDCs has also been described based on the sequence-dependent pattern of RDCs along a helical structure, called a “dipolar wave” by the Opella group [317, 352, 353]. The magnitude and periodicity of the dipolar wave depends on the orientation of the helix, and can allow irregularities in helix structure to be identified. Most commonly, dipolar waves have been used to help determine the location of helices in a protein sequence, allowing these structural elements to be more rigidly restrained over the course of a structure calculation [57, 159, 323, 354]. This is particularly useful for larger helical membrane proteins, since  $\alpha$ -helices are not well defined by the NOEs available in sparsely protonated samples [262].

Aside from helices, RDCs have a more general capacity to recognize related structures. This approach has been used to identify homologous structures in the protein databank [355], which could then be used as a starting point for RDC-based refinement to the final fold [356]. Having relatively few representatives in the PDB, a more useful application for membrane proteins is the identification of structurally homologous protein fragments from the structure database [357] or the structure modeling program Rosetta [358] to generate an initial model for refinement. This approach works well for smaller proteins having around four measured couplings per peptide plane, although ambiguities can arise with multiple structures being equally compatible with the RDC data for some fragments. Hence for larger proteins, it is necessary to supplement RDC data with distance restraints provided by NOEs [358] or paramagnetic relaxation enhancement (PRE) data [41, 351] to resolve these ambiguities.

The power of this fragment-replacement approach was recently demonstrated for the 6-TM helix UCP2 protein, for which an average of 2.2 RDCs per residue could be acquired (Fig. 6d) [41]. These RDCs were used to determine local and secondary structures using a PDB-based fragment database. This resulted in 15 continuous segments of UCP2 with local backbone structures that could be highly-restrained during structure calculation. RDC input at this stage helped to determine relative segment orientations, and PREs provided information on their relative positions. Although no structural information on side chains are obtained in this

approach, the global folds that can be accessed without the requirement for NOE assignment have the exciting potential to increase the number of membrane proteins structures that can be characterized by solution NMR.

### 5.3.3 Paramagnetic Relaxation Enhancements

As shown in the UCP2 study, RDCs on their own are usually not sufficient to determine a unique protein conformation, and long-range distance restraints are often required to resolve structural ambiguities. PREs have frequently been used to provide this type of distance information for membrane proteins both in the absence and presence of RDC data, particularly when long-range NOEs are not available (listed in Table 2, Method VII). PREs arise from dipolar interactions involving the large dipole of an unpaired electron that leads to efficient loss of polarization and coherence in surrounding nuclear spins. There is a relatively straightforward relationship between the paramagnetically enhanced transverse relaxation rate of a proton and the inverse sixth power of its distance from the unpaired electron ( $r^{-6}$ ), allowing this measurement to be converted into distance restraints for structure determination [359, 360]. This is usually achieved by measurement of paramagnetic spin-induced changes in backbone amide  $^1\text{H}$  peak intensities [360], although greater accuracy can be achieved by either full [361] or two-point [362]  $R_2$  relaxation measurements [363]. The range of distances that are accessed by the PRE measurement depends on the type of paramagnetic center introduced into the protein, but typically provides distance information on spins as far away as 25 Å or more.

There are several different types of spin labels that can be introduced into the target protein, usually via covalent modification of a unique cysteine residue and/or inclusion of a metal-chelating peptide in the protein sequence (reviewed in [364, 365]). Many of the metal-chelating tags used to induce alignment for measurement of RDCs (Sect. 5.3.1) can also be used to bind to paramagnetic metals with slow electronic relaxation properties (e.g.,  $\text{Mn}^{2+}$  and  $\text{Cu}^{2+}$ ) that give rise to a significant PRE. However, for membrane proteins, the vast majority of PRE-based restraints have used the sulfhydryl-reactive nitroxide spin label, 1-oxy-2,2,5,5-tetramethyl-3-pyrroline-3-methyl-methanethiosulfonate (MTSL) [38, 39, 41, 67, 88, 136, 354, 366].

To obtain sufficient numbers of PRE-based restraints, it is usually necessary to introduce spin-labels at a diverse range of sites. For example, nine different MTSL-labeled DsbB samples (approximately 1 MTSL-label for every 20 residues), were required to generate an average of 6.5 PREs per residue [39]. Choosing where to place these spin-labels requires identification of sites that can provide informative PRE-based restraints without disrupting the structure. This makes it necessary to have some knowledge of the structure being studied, with secondary structure determination from backbone assignments being the first step in this process. Once labeled, the NMR spectrum is used to confirm that the probe does not alter structure, since only local shift changes should be observed in this case.

For any type of paramagnetic center used, accurate translation of PREs into distance restraints requires that the entire population be spin-labeled, since the extent

of broadening will be underestimated if the labeling reaction is incomplete. In addition, the diamagnetic reference spectrum must not contain any residual paramagnetic contributions. These can be significant when water-soluble reducing agents such as ascorbic acid are used to reduce MTSL, particularly since the membrane-mimetic environment may shield the radical center. This was observed in the case of OmpA, with EPR spectra suggesting that only ~90% of the paramagnetic center was reduced, in spite of attempts to achieve higher levels of completion with a wide range of reducing agents [367]. To avoid this problem it is possible to prepare a separate sample labeled with a diamagnetic analog of MTSL that contains an acetyl group in place of the oxygen radical, allowing a purely diamagnetic reference spectrum to be acquired [367]. Similarly, the reference spectrum for PRE measurements with metal chelating tags is usually obtained with a diamagnetic ion (e.g.,  $\text{Ca}^{2+}$ ,  $\text{Mg}^{2+}$ ) being bound in place of the paramagnetic metal [327, 329, 330].

If care is taken to ensure that samples are completely labeled and that an appropriate diamagnetic reference is used, then it is possible to generate accurate distance restraints from these PRE measurements. For example, residues showing paramagnetic:diamagnetic peak intensity ratios ( $I_{para}/I_{dia}$ ) of 15–85% in MTSL-labeled OmpA had PRE-based distance measurements that were within 2 Å of the corresponding distances determined from the crystal structure [367]. This allows relatively tight restraint bounds to be applied during structure calculations, with  $\pm 2$  Å for the 15–24 Å distance range, and an upper limit of 15 Å being applied for residues showing intensity ratios under 15%. It is also possible to impose lower distance limits of 20–22 Å for residues showing small PREs (>85%), although these did not improve the quality of the OmpA structure [367]. This was the approach used in the structure determination of the 4-TM helix DsbB, which along with RDCs and a handful of long-range NOEs, generated a highly precise ensemble (pairwise backbone rmsd of 0.80 Å) [39]. A comparable level of precision was also achieved for Rv1761, which used slightly looser bounds of  $\pm 3$  Å, since the narrower bounds gave rise to distance violations [159]. More conservative bounds (e.g.  $\pm 4$  Å [67, 351, 360] or greater [41]) or NOE-like assignments into distance bins (i.e., strong, medium weak PREs) have also been used [88, 368], although this does reduce the impact of the PRE on the structure [367].

An important consideration when converting PREs into distances is the influence of spin-label and backbone dynamics. Specifically, the  $r^{-6}$  dependence causes the measured PRE for a particular  $^1\text{H}$  atom to be dominated by conformers having shorter distances to the spin label, even if this state is infrequently sampled [363]. Consequently, PREs involving flexible regions of a protein cannot be accurately converted into a single average distance. Since dynamic regions of the protein can be identified with standard backbone relaxation measurements, PREs involving these residues should be discarded or only used with very loose restraint bounds.

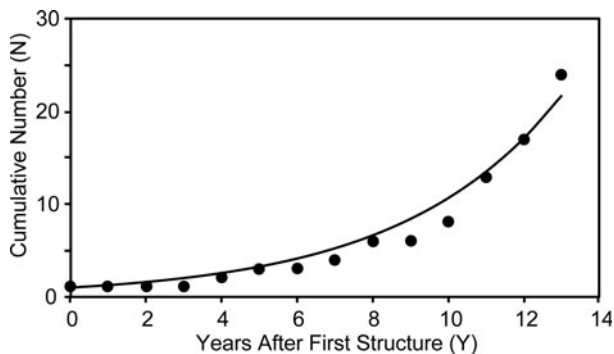
Less straightforward is the treatment of spin labels that are bound to the protein via a flexible linker, where exchange between various rotameric states can occur. To account for the different contributions to the PRE that can arise from this type of exchange, an ensemble approach was developed that refines the structures against PREs in place of PRE-derived distances [363], improving the accuracy of structures

calculated for  $\text{Mn}^{2+}$ -EDTA labeled DNA in complex with SRY. This type of ensemble-based refinement has also been suggested to be necessary for more accurate use of PRE data from the MTLs spin label, particularly since some high-resolution crystal structures of MTLs-modified lysozyme show the presence of more than one spin-label conformer [369, 370]. However, the good agreement that was found between calculated and PRE-measured distances for the relatively rigid barrel region of MTLs-labeled OmpA implies that the presence of multiple conformations may be less significant for this particular spin label [367]. EPR line shapes of MTLs-modified lysozyme also show that the motion of this paramagnetic center is restricted [371]. Hence it appears that carefully measured PREs with MTLs in non-dynamic segments can in most cases be validly converted into long-range distance restraints with tighter bounds (i.e.,  $\pm 2 \text{ \AA}$ ).

Overall, PREs have significantly increased the number of larger membrane protein structures that have been determined by solution NMR in a relatively short period, becoming an established source of structural information. In addition, water-soluble paramagnetic relaxation agents can be used to reveal aqueous-exposed segments [136, 162, 184, 354, 366, 372–374], spin-labeled lipids and detergents can identify lipid-embedded regions [136, 372, 375–381], and intermolecular interactions can be probed with exquisite sensitivity using spin-labeled proteins [243, 382]. This versatility is one of the factors ensuring that paramagnetic effects in NMR will continue to make important contributions to our understanding of membrane protein structure.

## 6 Concluding Remarks

Although membrane proteins continue to present challenges for solution NMR, innovations in sample development, data acquisition, and structure determination strategies have allowed structural insights to be obtained from some of the most demanding systems tackled to date (e.g., GPCRs and large  $\beta$ -barrel channels). In addition to the advances outlined here, these achievements build upon contributions from the handful of veteran laboratories that have worked to apply solution NMR to membrane proteins, even before the introduction of more modern methods for study of large proteins (e.g., [140, 383–386]). Yet in spite of this long history, the field is still considered to be quite young, with only a limited number of groups fully exploiting the potential of solution NMR for membrane protein structure determination. Nonetheless, it is encouraging to note the progress that has been made since the first NMR structure was determined for an integral membrane protein comprised of more than one TM segment back in 1997 (i.e., the glycoporphin A homodimer [140]). As of the end of 2010 the protein structure database holds approximately 24 integral membrane protein structures (counting only those with more than one membrane-spanning segment), putting the field at par with the state of membrane protein crystallography in the late 1990s. However, the rate of new structure accumulation closely follows the exponential rise seen for crystal



**Fig. 7** Progress in integral membrane protein structure determination by solution NMR (structures with more than one membrane-spanning segment only). The cumulative number of integral membrane protein structures determined by solution NMR is shown for each year, starting with the year the first high-resolution structure of a TM-helix dimer was revealed in 1997 [140]. As had been shown for X-ray crystal structures of membrane proteins [399], this accumulation can be fit to the equation  $N = \exp(aY)$ , where  $N$  is the cumulative total of structures for each year following publication of the first structure ( $Y$ ). The scaling parameter  $a = 0.236$  for the best-fit curve (shown as the *solid line*) is almost the same as that previously determined for membrane protein X-ray structures ( $a = 0.242$ ), showing similar rates of progress for the two structure determination techniques. Year 0 for X-ray crystal structures was recognized as 1985 [399], more than 10 years earlier than that for NMR (1997). A useful reference used in the compilation this data was provided by the on-line catalogue of membrane protein NMR structures (<http://www.drorlist.com/nmr/MPNMR.html>)

structures of membrane proteins (Fig. 7) [387], with the main disadvantage being its late start relative to X-ray crystallography. Assuming that solution NMR of membrane proteins continues at this pace, we can expect to see the 100 structure mark surpassed in approximately 6 years. The ability of long-range restraints to be acquired without the need for side chain assignment will be particularly important in meeting these projections, as will the implementation of new sensitivity-enhancing technologies by an increasing number of labs. This analysis shows that, while solution NMR studies of larger membrane proteins may not yet be routine, reaping the rewards of these endeavors will become increasingly feasible for a larger range of membrane proteins than ever before.

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# Protein Structure Determination by Solid-State NMR

Xin Zhao

**Abstract** Membrane proteins are a large, diverse group of proteins, representing about 20–30% of the proteomes of most organisms, serving a multitude of cellular functions and more than 40% of drug targets. Knowledge of a membrane protein structure enables us insight into its function and dynamics, and can be used for further rational drug design. Owing to the intrinsic hydrophobicity, flexibility, and instability of membrane proteins, solid-state NMR may offer an unique opportunity to study membrane protein structure, ligand binding, and activation at atomic resolution in the native membrane environment on a wide ranging time scale. Over the past several years, solid-state NMR has made tremendous progress, showing its capability of determining membrane protein structure, ligand binding, and protein dynamic conformation on a variety of time scales at atomic resolution. In this chapter we will mainly discuss some recent achievements on membrane protein structure determination, ligand conformation and binding, structure changes upon activation, and structure of insoluble fibrous proteins investigated by using magic-angle spinning solid-state NMR from the structural biology point of view. Protein dynamics, sensitivity enhancement, and the possibility of chemical shift-based structure determination in solid-state NMR are also briefly touched upon.

**Keywords** Membrane proteins · MAS solid-state NMR · Structure determination · ligand conformation · structural changes and dynamics

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## 1 Introduction

Membrane proteins are a large, diverse group of proteins, representing about 20–30% of the proteomes of most organisms, serving a multitude of cellular functions, and more than 40% of drug targets [1]. For example, membrane proteins often serve as receptors or provide channels for polar or charged molecules to pass through the cell membrane and maintain a variety of biological processes [2]. Knowledge of a membrane protein structure enables us insight into its function and dynamics, and can be used for further rational drug design. Therefore it is always desirable to have an accurate picture of protein structure in the highest resolution possible. However, owing to their intrinsic hydrophobicity, flexibility, and instability, many fewer structures have been solved by X-ray crystallography for membrane proteins compared to soluble proteins [3, 4]. Moreover, the conformational change of the transmembrane helices upon activation increase the difficulty of capturing the activation state of a membrane protein to a higher resolution by X-ray crystallography [5, 6]. In contrast, solid-state NMR (SSNMR) is a suitable technique to study molecular structure and interactions at atomic level in a variety of sample forms; it can be used to determine a membrane protein structure and probe its conformational dynamics in the native membrane environments.

Over the past several years, SSNMR has made tremendous progress, showing its capability of determining membrane protein structure, ligand binding, and protein dynamic conformation on a variety of time scales at atomic resolution. Many membrane proteins have been investigated by magic-angle spinning (MAS) SSNMR, factors investigated including the following: activation, inhibition, and dynamics of the potassium channel KcsA-Kv1.3 [7–11]; structure, ligand conformation, activation, and dynamics of the G protein coupled receptor – rhodopsin [12–25]; protonation switch mechanism of the human H1 receptor[26]; the

influenza M2 proton channel structure, function and ligand binding [27, 28]; the human prion protein [29]; the structural conversion of neurotoxic amyloid beta oligomers to fibrils [30]; and the structure and dynamics of the retinylidene proteins from bacteria, including bacteriorhodopsin, sensory rhodopsin, halorhodopsin, and proteorhodopsin [31–41], and the structure of the HET-s(218–289) fibril [42, 43]. Some other membrane proteins, such as membrane-embedded enzymes [44–46], histidine kinases [47], ABC transporters [48], and bacterial outer membrane proteins [49, 50], have also been investigated through multidimensional correlation experiments in SSNMR.

In this chapter we will briefly review some of the recent progress in studying membrane proteins by using magic-angle spinning solid-state NMR from biological structure point of view; for a complete overview of the achievements in this field, please refer to the following excellent reviews [19, 51–67].

## 2 Basic Experimental Techniques Used in Solid-State NMR

Unlike solution-state NMR, the resolution and sensitivity of SSNMR are affected heavily by orientation dependent anisotropic spin interactions such as chemical shift anisotropy, homonuclear and heteronuclear dipolar couplings, quadrupole coupling, etc. These interactions generally cannot be averaged out by the molecular tumbling motions in solids, presenting a very broad line shape with poor sensitivity of SSNMR spectra. Combining MAS with cross-polarization (CP), high power proton decoupling, recoupling, and isotopic labeling can achieve high resolution and signal-to-noise ratio in SSNMR for protein structure determination.

### 2.1 Magic-Angle Spinning

MAS is an essential technique in SSNMR for obtaining a high resolution spectrum [68–70]. The basic idea is to spin the sample container (rotor) about an axis, which subtends an angle of  $54.74^\circ$ , the magic-angle, with respect to the static field  $B_0$ . The spatial rotation of the sample introduces time-dependence to anisotropic spin interactions, such as chemical shift anisotropy, homonuclear dipole–dipole couplings, and heteronuclear dipole–dipole couplings, which are averaged out more efficiently as the sample spinning frequency increases. Due to the periodic time-dependence of the CSA and dipolar spin interactions, the broad static line shape breaks up into a center band at the isotropic position and a set of spinning sidebands separated by the spinning frequency. As the spinning frequency increases, the time averaging is more effective, which leads to a decrease in the sideband intensities and an increase in the center band intensity. The advantage of MAS is that both resolution and sensitivity are greatly increased.

The disadvantage is that the spectrum loses all the anisotropic dependent molecular geometry information.

## 2.2 High Power Heteronuclear Decoupling

Hydrogen is a very common element and has a large  $\gamma$  value and almost 100% natural abundance. Its coupling to rare spins with low  $\gamma$  complicates the observation of the rare spins, such as  $^{13}\text{C}$  and  $^{15}\text{N}$ . It is usually necessary to remove heteronuclear dipolar couplings between protons and the observed rare spins by using strong rf irradiation on proton spins, i.e., so-called high power proton decoupling. Continuous-wave (CW) decoupling is a standard scheme for decoupling of heteronuclear dipolar spin interactions. The rf field induces a fast rotation of the proton spin states, averaging out their interaction with the rare spins [71, 72]. For fast MAS, more sophisticated pulse sequences, such as TPPM, C12, R24, XiX, and Spinal64, will achieve a better decoupling efficiency [73–77].

## 2.3 Cross-Polarization

The sensitivity of the rare spin species ( $S$ -spins) with low magnetogyric ratio, such as  $^{13}\text{C}$  and  $^{15}\text{N}$ , can be enhanced by transferring magnetization from abundant spin species ( $I$ -spins) with high magnetogyric ratio, such as  $^1\text{H}$ . Cross-polarization (CP) is the most widely used method in solids to transfer polarization between unlike spin species through heteronuclear dipolar couplings [78, 79]. Magnetization transfer is achieved when the strengths of the two fields match the Hartmann-Hahn (HH) condition  $|\gamma_I B_{rf}^I| = |\gamma_S B_{rf}^S|$ , where  $\gamma_I$  and  $\gamma_S$  are the magnetogyric ratios of the  $I$ -spins and  $S$ -spins, respectively [80]. The enhancement of the  $S$ -spin magnetization is roughly proportional to the ratio of the two magnetogyric ratios  $|\gamma_I/\gamma_S|$ . Ramping or adiabatic-passing one of the  $rf$  fields can improve the reliability and reproducibility of HH-CP, especially in fast MAS experiments [81–84].

## 2.4 Recoupling in MAS Solid-State NMR

Fast MAS leads to high resolution and sensitivity, the basic requirement for sequential assignment of protein structures in SSNMR. However, all anisotropic spin interactions which can be used to extract molecular geometry information are averaged out by MAS. In order to get back the anisotropic spin interactions in the presence of MAS, recoupling techniques were developed to retrieve selectively the anisotropic spin interactions for rotating solids [189–190]. Generally, there are two approaches to reintroduce anisotropic dipolar interactions, either mechanically,



where the recoupling is achieved through rotational resonance, or by rf pulse driven methods, where the recoupling is achieved by applying rf pulse trains. By applying recoupling sequences, certain anisotropic interaction terms can be retrieved in the presence of MAS. Magic-angle spinning NMR in combination with recoupling has become a widely used experimental technique for obtaining molecular structural information in non-crystalline or disordered materials. MAS makes it possible to obtain SSNMR spectra with good resolution and sensitivity, while the recoupling techniques cause a selective restoration of informative anisotropic nuclear spin interactions, which are normally suppressed by magic angle rotation. A wide range of recoupling methods has been developed for homonuclear and heteronuclear spin systems [85–88]. With a combination of recoupling sequences and multidimensional experiment schemes, membrane protein structure and dynamic conformation upon activation can be revealed through the short and long range semi-quantitative distance restraints from multidimensional homonuclear/heteronuclear correlation experiments, such as PDS/DARR, PAR, TEDOR, PAIN-CP, and CHHC/NHHC type experiments [18, 89–92, 127, 191–192], etc.

### 3 Protein Structure Determination by MAS Solid-State NMR

#### 3.1 *Labeling Strategy and Sample Environments*

Isotopic labeling plays a very important role in molecular structure determination in SSNMR. It not only enhances the spectral sensitivity and improves the spectral resolution, but also helps with the resonance assignments of NMR spectra and tackles the specific problem with the structure and dynamics of proteins through the designed labeling scheme.

There are three main approaches to isotopically label proteins: specific, selective, and uniform labeling. All are used extensively in SSNMR studies of protein structure and ligand conformation, depending on the questions that the designed experiment needs to be answered.

Specific labeling normally refers to incorporate a non-uniformly  $^{13}\text{C}$ ,  $^{15}\text{N}$  labeled amino acid into a polypeptide/protein or a ligand at a certain position. It requires solid-phase peptide synthesis or chemical synthesis. This approach has been used extensively to study amyloid peptides, membrane peptides, and GPCR ligand conformation [93–99].

Selective labeling, including forward and reverse labeling, refers to the biosynthetic incorporation of a single type or several types of labeled amino acid(s) – with the rest unlabeled – into the protein expression media. Selective labeling is widely used for the protein resonance assignments and structure calculations in SSNMR. For example, a selectively and extensively labeling scheme using [1, 3- $^{13}\text{C}$ ] and [2- $^{13}\text{C}$ ]-glycerol as the sole carbon source in the bacterial growth medium has been used to determine the  $\alpha$ -spectrin SH3 domain,  $\alpha\text{B}$ -crystalline, and outer membrane

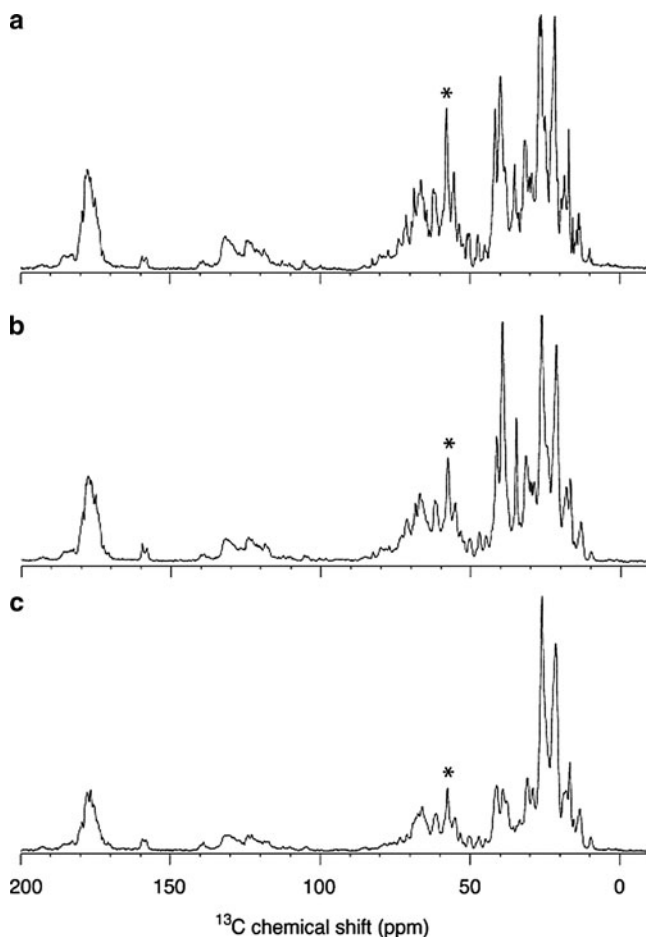
protein G (OmpG) [100, 101]. A selective labeling scheme has also been used in the HEK293 system to express labeled rhodopsin for structure and function study [12, 18, 22, 102–104]. The reverse labeling approach has been designed to study structure and dynamics of the sensory rhodopsin II and the potassium ion channel KcsA-Kv1.3 [105, 106]. This approach has the advantage of decreased numbers of overlapping resonances from the hydrophobic transmembrane helices, making the extensive assignments possible. By taking advantage of the glycolysis pathway, the TEASE  $^{13}\text{C}$  selective labeling scheme has been suggested to probe the transmembrane segments of membrane proteins [107], and this reverse labeling approach may possibly be adaptable for other different applications for its flexibility.

A uniformly labeling scheme is the simplest and most cost-effective biosynthetic labeling method for protein SSNMR. Normally, uniformly  $^{13}\text{C}$ -labeled glucose or glycerol, and  $^{15}\text{N}$ -labeled ammonium chloride or ammonium sulfate are used as the labeled precursors in the growth medium. With a single sample, all the structural constraints can be obtained through a set of correlation experiments and the protein structure can be calculated thereby. This approach has been first demonstrated on microcrystalline proteins of known structure [108–111], and then successfully applied to membrane protein for structure determination [7, 38, 41, 105, 112–114].

Recently, perdeuteration of proteins has been exploited to gain even higher resolution for U- $^{13}\text{C}$  and  $^{15}\text{N}$  labeled proteins in SSNMR by taking advantage of minimizing the  $^1\text{H}$ - $^1\text{H}$  dipolar couplings causing spectral broadening. Perdeuterated soluble microcrystalline proteins have been used to study protein dynamics and interactions [115–117], and structure determination [118]. This approach has also been attempted to express the fully labeled 7-transmembrane protein bacteriorhodopsin for structure investigation, as shown in Fig. 1 [119].

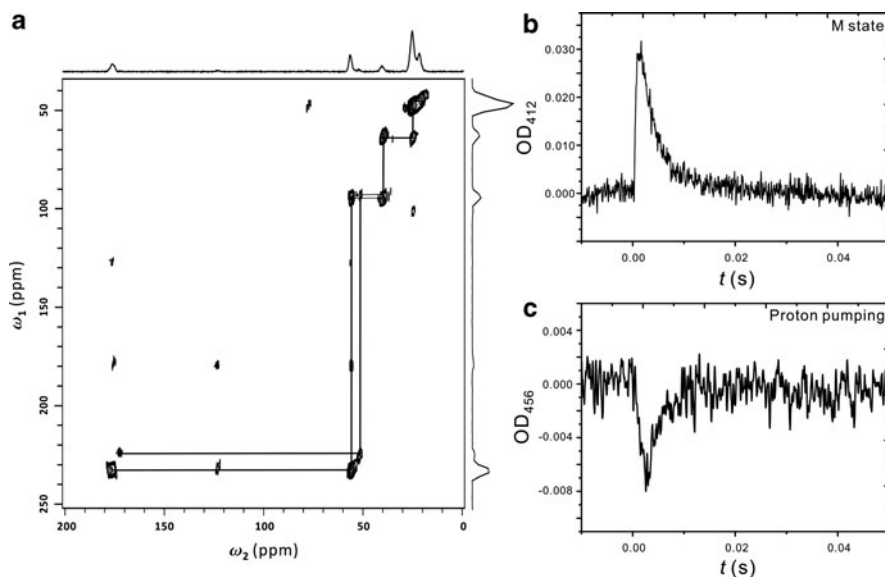
However, with increase of deuteration level, protein expression level may drop and some strains may even be difficult to grow on  $\text{D}_2\text{O}$  [120]. Increasing deuteration level may also influence the resonance frequencies and CP efficiency. Therefore, a good balance needs to be considered from both the NMR and biological sides.

Choosing an appropriate sample environment is not only critical for generating a high resolution NMR spectrum, but also critical for the protein to have correct folding and activity. Obtaining homogeneous sample preparation leads to improved linewidths and therefore spectral resolution, while heterogeneous samples can result in artifacts such as unexpected peaks and peak doubling [120]. Furthermore, it is not enough to show that a protein construct is functional to validate a structure unless the functional assay is performed in the same environment as that used for the structural characterization [121]. For small proteins, nano-/micro-crystalline or nano-disk samples have been proven a good choice for yielding high quality spectra in solids [122–124]. For membrane proteins, samples can be prepared in either detergent micelles, bicelles, or lipid bilayers. Given that membrane proteins function within a bilayer environment, it is more biologically applicable to be able to carry out structural investigations in lipids [120]. Structural data obtained in an appropriate lipid bilayer environment can serve as benchmarks for validating structures determined in other mimetic environments [121]. Figure 2 shows the



**Fig. 1** 1D  $^{13}\text{C}$  NMR spectra (800 MHz for  $^1\text{H}$ ) of CN-bR (a), HDCN-bR (b), and DCN-bR (c), fully hydrated in sodium citrate buffer and collected under very similar conditions. Adapted from [119] with permission from Elsevier B.V

$^{13}\text{C}$ -DQ/SQ correlation experiment on the fully hydrated [ $^{13}\text{C}_6$ ,  $^{15}\text{N}$ ]-Leu-BR sample using a Bruker AV-III 600 MHz wide-bore spectrometer with an MAS rate of 8 kHz (Fig. 2a). The sample activity is carefully assessed through detection of the M state and the proton pumping cycle at 412nm and 456nm by optical dynamic spectroscopy, respectively. It is clear that the resolved narrow spectral linewidths are attributed to the predominant distribution of the Leu residues on the helical segments with similar local environments, and the up-field shifted peak is attributed to the Leu residue located at the loop region. This small shift clearly indicates a homogenous sample condition. The expressed protein is fully functioning as confirmed by capture of the M state signal and the proton pumping cycle signal shown in Fig 2b and c.



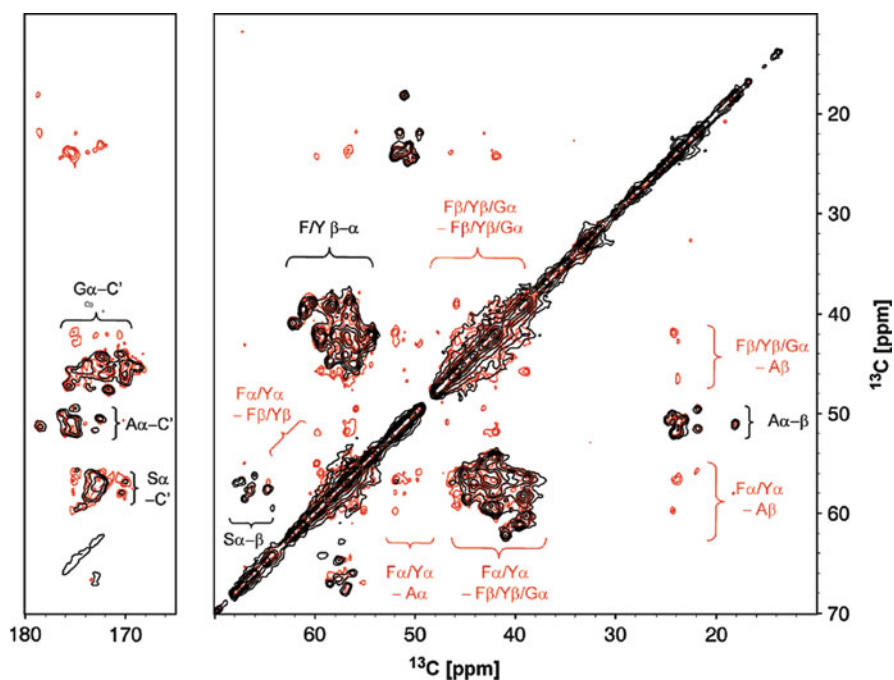
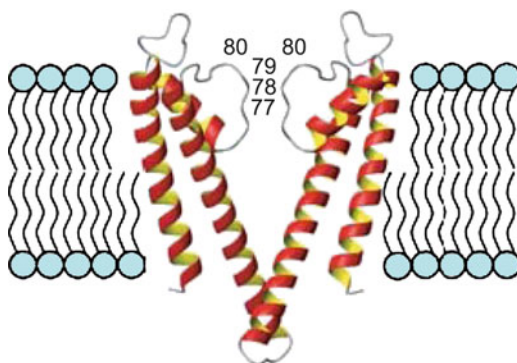
**Fig. 2** 2D  $^{13}\text{C}$  DQ/SQ NMR correlation of the fully hydrated [ $^{13}\text{C}_6$ ,  $^{15}\text{N}$ ]-Leu-BR sample using a Bruker AV-III 600 MHz wide-bore spectrometer with an MAS rate of 8 kHz (a) and the protein function assay through detection of the M state signal and proton pumping cycle signal at 412nm (b) and 456nm (c) by optical dynamic spectroscopy, respectively

### 3.2 Three-Dimensional Structure Determination

Determination of membrane proteins and amyloid fiber structures is still a frontier of structural biology. Over the last few years, much progress has been made in various areas from NMR methodology to sample preparation in order to reduce the dipolar truncation, establish reliable short and long range distance constraints, and improve spectral linewidths and resolution for a full structure determination. The Baldus group has established a method to probe through-space ( $^1\text{H}$ ,  $^1\text{H}$ ) contacts of protonated solid-phase systems in high spectral resolution [125]. The so-called CHHC/NHHC experiment was first demonstrated to solve a micro-crystalline Crh protein [126], and then successfully applied to the membrane bound potassium channel KcsA-Kv1.3 to study its structure activation and inhibition [7]. By combining the  $^1\text{H}$ ... $^1\text{H}$  distance restraints, chemical shift changes, and molecular modeling, they have obtained the first membrane protein structure by SSNMR, as shown in Fig. 3, which provides a deep understanding of channel activation and inhibition [7].

Full or nearly full SSNMR spectral assignments have not been achieved on membrane proteins for quite some time due to the signal overlapping and fast longitudinal and transverse relaxation which causes line broadening and signal intensity loss in many correlation experiments, impeding the pace of SSNMR structure determination of large membrane proteins and protein-complexes. Recently, the Oschkinat group has proposed a novel specific labeling scheme by

**Fig. 3** Solid-state NMR structure of membrane bound potassium channel KcsA-Kv1.3. Adapted from [7] with permission from Nature Publishing Group



**Fig. 4**  $^{13}\text{C}$ - $^{13}\text{C}$  proton-driven spin diffusion spectrum of OmpG-GAFY recorded at 900 MHz with 20 ms (black) and 700 ms (red) mixing time. Intra- (black) and inter-residual (red) cross-peak regions are indicated. Adapted from [49] with permission from the American Chemical Society

using [ $^{15}\text{N}$ , 2, 3- $^{13}\text{C}$ ]-labeled Phe and Tyr residues and fully labeled Gly and Ala residues to restore favorable cross-relaxation properties of the glycerol samples in order to obtain inter-residue cross-peaks [49]. They have expressed the 281-residue OmpG with this labeling strategy (referred to as OmpG-GAFY) and demonstrated its advantage by the  $^{13}\text{C}$ - $^{13}\text{C}$  proton-driven spin diffusion (PDSF) experiment [127, 191–192] as shown in Fig. 4. It is very clear that the Ca-C $\beta$  signals for Phe, Tyr, and Ala, and the Ca-C' cross-peaks for Gly and Ala are all well resolved and the signal intensities are increased significantly [49]. This labeling scheme has several

advantages: (1) it is composed of a low number of small, isolated spin systems, (2) transfer of magnetization into the side chain is thus eliminated and spectral quality enhanced, and (3) the number of inter-residue cross-peaks is significantly increased which is important both for assignment and structure calculation [49]. This labeling scheme or similar ones may be applicable to other large membrane proteins, e.g., the 7-transmembrane family proteins.

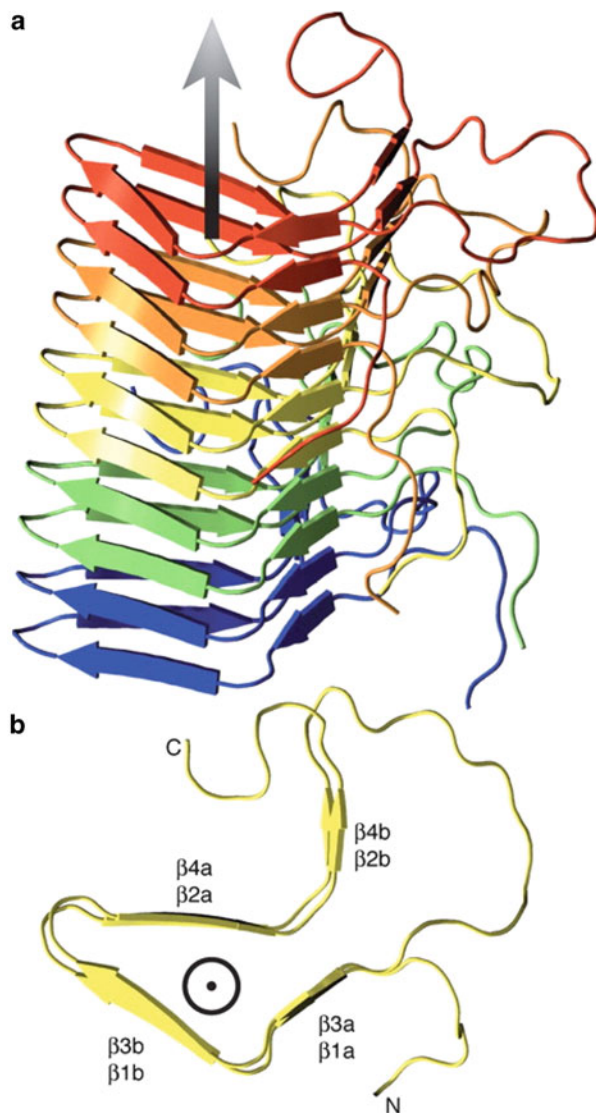
Another great achievement by SSNMR is the structure determination of the HET-s (218–289) fibril by the Meier group [42], as shown in Fig. 5. Total 90  $^{13}\text{C}$ – $^{13}\text{C}$  and 44  $^1\text{H}$ – $^1\text{H}$  distance restraints obtained by the CHHC, NHHC and PDSO experiments, and 74 angle restraints obtained by TALOS [193], were used for structure calculations. The extraordinarily high order in the HET-s prion fibrils can be explained by the well-organized structure obtained by SSNMR.

### 3.3 Ligand Conformation and Binding

The molecular mechanisms of membrane protein activation are at the center of interest in the study of cellular responses to biogenic stimulus and drugs. For example, GPCRs are activated by a wide range of stimuli, including hormones, neurotransmitters, ions, odorants, and photons of light [128]. Knowledge of the three-dimensional structures of several GPCRs, such as rhodopsin,  $\beta_1\text{AR}$ ,  $\beta_2\text{AR}$ ,  $\text{A}_{2\text{a}}\text{R}$ ,  $\text{CXCR4}$ ,  $\text{D}_3\text{R}$ , and  $\text{H}_1\text{R}$  [129–141] have been resolved by X-ray crystallography in either an inactive state or agonists/antagonists bound form at high resolution which open up new possibilities for investigating GPCRs of human therapeutic significance and for rational drug design. However, despite the availability of those crystal structures, a comprehensive understanding of the mechanism of structure activation is still a challenge due to the lack of high resolution structure at the activated state. For example, activation of rhodopsin has still been challenging because of the lack of high resolution structure at atomic level for the activated Meta II state [6]. On the other hand, SSNMR can offer direct measurements at atomic resolution to study protein activation caused by the conformational changes and the ligand binding interactions [65, 120].

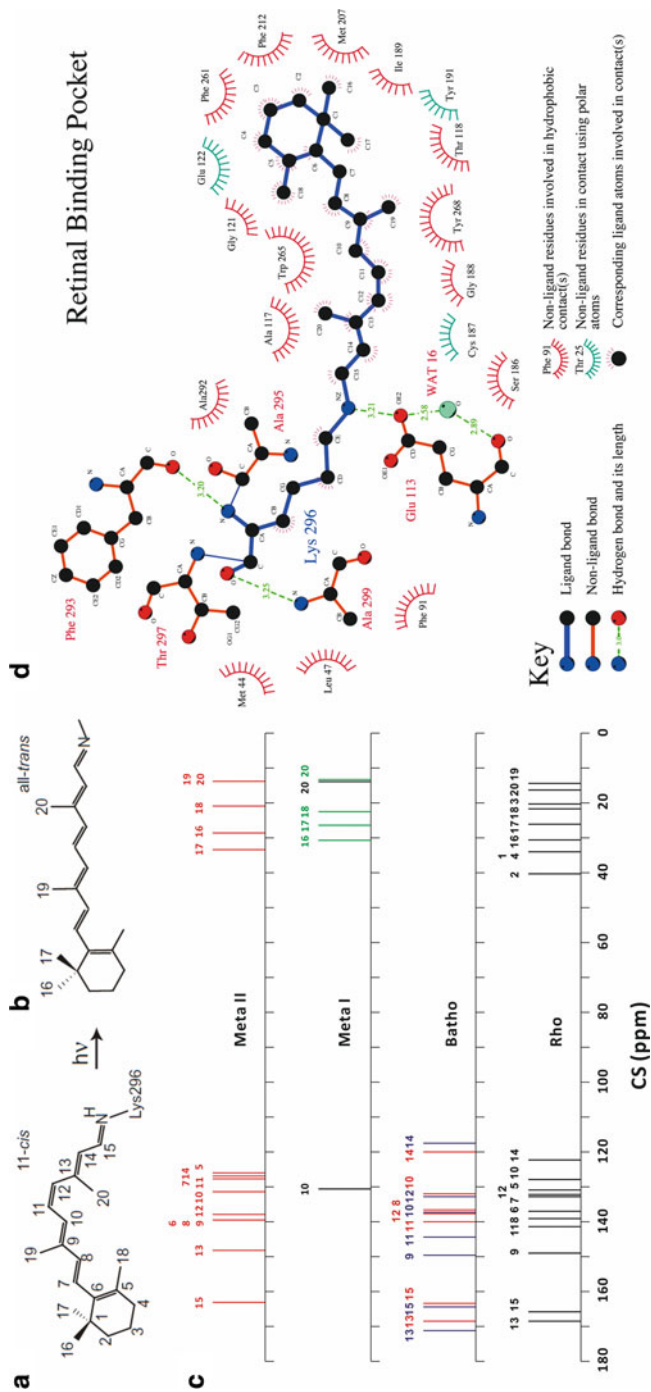
The conformation of retinal chromophore and protein activation upon the 11-*cis* to all-*trans* isomerization of the retinal have been studied extensively by MAS SSNMR through the chemical shift measurements, distance measurements, correlation experiments, and  $^2\text{H}$  NMR [12–18, 20, 22, 23, 25, 95, 103, 104, 142–144]. Complete assignment of the retinal carbons in ground state and partially in the Batho-, Meta I, and Meta II intermediates have been achieved [14, 17, 145–151]. These valuable data allow us insight into the conformational changes of the retinal protonated Schiff base (PSB) complex and the related transmembrane helices through the transition from the ground state to the Batho, Meta I, and Meta II intermediates in the binding pocket, as shown in Fig. 6.

Large down-field or up-field changes have been observed at the C16, C17 of the  $\beta$ -ionone ring, and on the retinal polyene chain around C9–C10–C11–C12 and



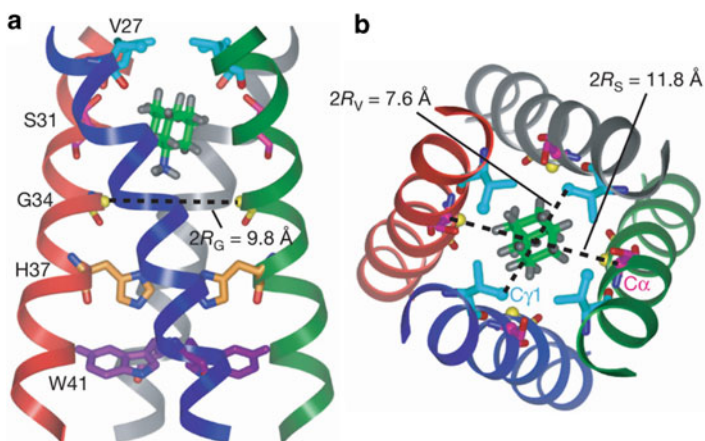
**Fig. 5** Solid-state NMR structure of the HET-s (218–289) fibril. (A) Side view of the five central molecules of the lowest energy structure of the HET-s(218–289) heptamer calculated from the NMR restraints. (B) Top view of the central molecule from (A). beta3 and beta4 lie on top of beta1 and beta2, respectively. Adapted from [42] with permission from the American Association for the Advancement of Science

C13–C14–C15 regions (Fig. 6). The large change at C13, C9, and C17 may be attributed to the break of the Glu113 salt bridge, van der Waals contact, and local hydrophobic environmental changes due to the displacements of the helices upon activation. In addition to the studies on the retinal in rhodopsin, the structure of



**Fig. 6** Comparison of retinal protonated Schiff base chemical shifts in rhodopsin different photocycle states. The chemical structure of retinal chromophore in 11-*cis* (a) and all-*trans* configuration (b). Chemical shifts of the retinal atoms in ground state, and Batho-, Meta I, and Meta II intermediate states (c) and a schematic drawing of the retinal binding pocket containing all residues within 4 Å to the retinal and Lys296 (d). The chemical shifts are adapted from the following references: [146–148] (black); [14] (blue); [149–151] (green); and [17, 145] (red). (d) is adapted from [188] with permission from the Elsevier B. V





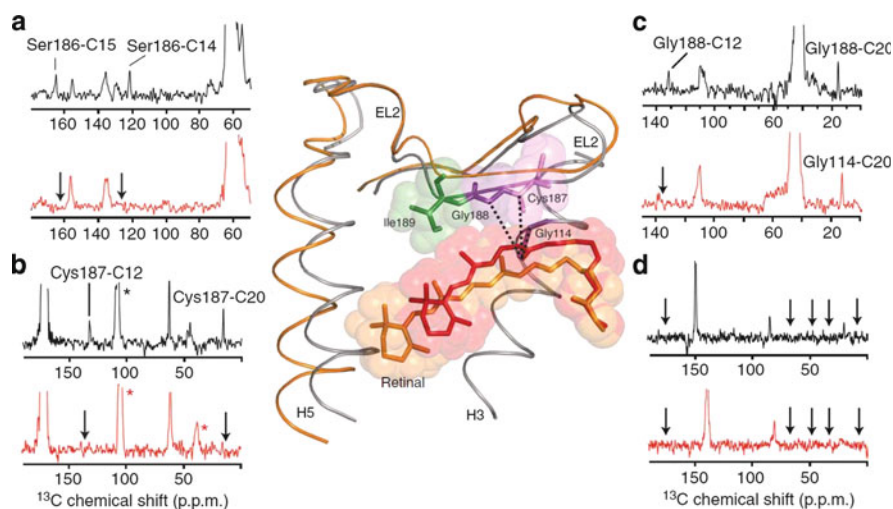
**Fig. 7** The solid-state NMR structure of Amt-bound M2 in lipid bilayers. (A) Side view showing several key residues and Amt in the high-affinity luminal site. The time-averaged Amt orientation is parallel to the channel axis. (B) Top view showing the Ser 31 and Val 27 pore radii. Adapted from [27] with permission from Nature Publishing Group

neurotensin bound or unbound to the NTS1 receptor has also been investigated by SSNMR, revealing a  $\beta$ -strand conformation upon binding [94]. The conformation of the bradykinin (BK) peptide bound to the human bradykinin B2 receptor in DDM, on the other hand, has recently been proposed to have a double S-shape structure [152]. In the case of the human histamine H1 receptor, changes in the protonation state of the ligand histamine binding to the receptor, SSNMR experiments have revealed that the ligand can bind in a different cationic form and a protonation switch might be part of the activation mechanism [26]. Another outstanding example of these types of work is the study of influenza M2 proton channel structure, function, and ligand binding conducted by the group of Hong [27, 28]; through the extensive heteronuclear distance measurements and orientation measurements, they have successfully proposed a structure model of the amantadine binding to M2 in phospholipid bilayers, as shown in Fig. 7. This study has clearly demonstrated the ability of SSNMR to elucidate drug-membrane protein interactions at atomic resolution and this is useful when conducting novel drug design for human therapeutics.

### 3.4 Structural Changes upon Activation

Upon binding of agonists, which typically occurs in proximity to the extracellular opening of the helical bundle, GPCRs undergo a series of structural changes that cascade from the extracellular to the intracellular part of the receptor and ultimately lead to G protein activation [153]. The multiple structural “switches” in rhodopsin

that trigger the conformational changes are involved in activation and formation of the G protein binding site. In the last few years, the Smith group has focused on SSNMR studies of structural changes on the extracellular side of the receptor caused by retinal isomerization. By combined SSNMR chemical shift measurements and 2D-dipolar assisted rotational resonance (DARR) [154] NMR measurements with selective labeling schemes and mutagenesis, they have published several papers addressing the functions of the displacement of EL2 on rhodopsin activation. Figure 8 shows the two-dimensional  $^{13}\text{C}$  DARR NMR spectra of retinal-EL2 interactions – close contact between the retinal  $^{13}\text{C}14$  and  $^{13}\text{C}15$  carbons and  $^{13}\text{C}\beta$ -Ser186 (Fig. 8a), between the retinal  $^{13}\text{C}12$  and  $^{13}\text{C}20$  carbons and  $^{13}\text{C}1$ -Cys187 (Fig. 8b), and between the retinal  $^{13}\text{C}12$  and  $^{13}\text{C}20$  carbons and  $^{13}\text{C}\alpha$ -Gly188 in rhodopsin (Fig. 8c). But the contacts between the retinal  $^{13}\text{C}9$  and  $^{13}\text{C}12$  carbons and U- $^{13}\text{C}_6$ -Ile189 in rhodopsin or Meta II were not observable (Fig. 8d) [18]. All the results indicate that the formation of Meta II is accompanied by the displacement of EL2 away from the retinal binding site and that there is a rearrangement in the hydrogen-bonding networks connecting EL2 with the extracellular ends of transmembrane helices H4, H5, and H6. This displacement is coupled to the rotation of TM5 and breaking of the ionic lock connecting TM3 and TM6 [18]. These comprehensive results may lead to further investigation of the molecular mechanism of the cavity formation between H3, H5, and H6 for G protein binding [155].

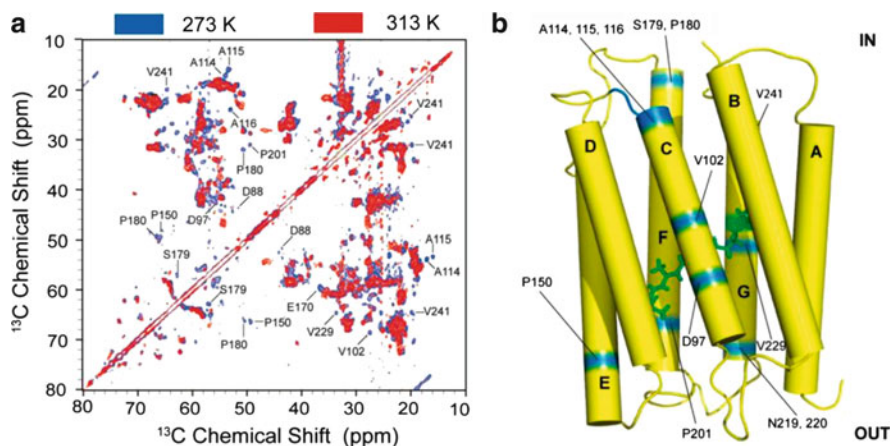


**Fig. 8** Two-dimensional  $^{13}\text{C}$  DARR NMR spectra of retinal-EL2 interactions. Rows from the two-dimensional  $^{13}\text{C}$  DARR NMR spectra of rhodopsin (*black*) and Meta II (*red*) are shown. The rhodopsin crystal structure (*gray*) with the Meta II model (*orange*) obtained from molecular dynamic simulations are shown in the middle of the figure. Adapted from [18] with permission from Nature Publishing Group

### 3.5 Protein Dynamics

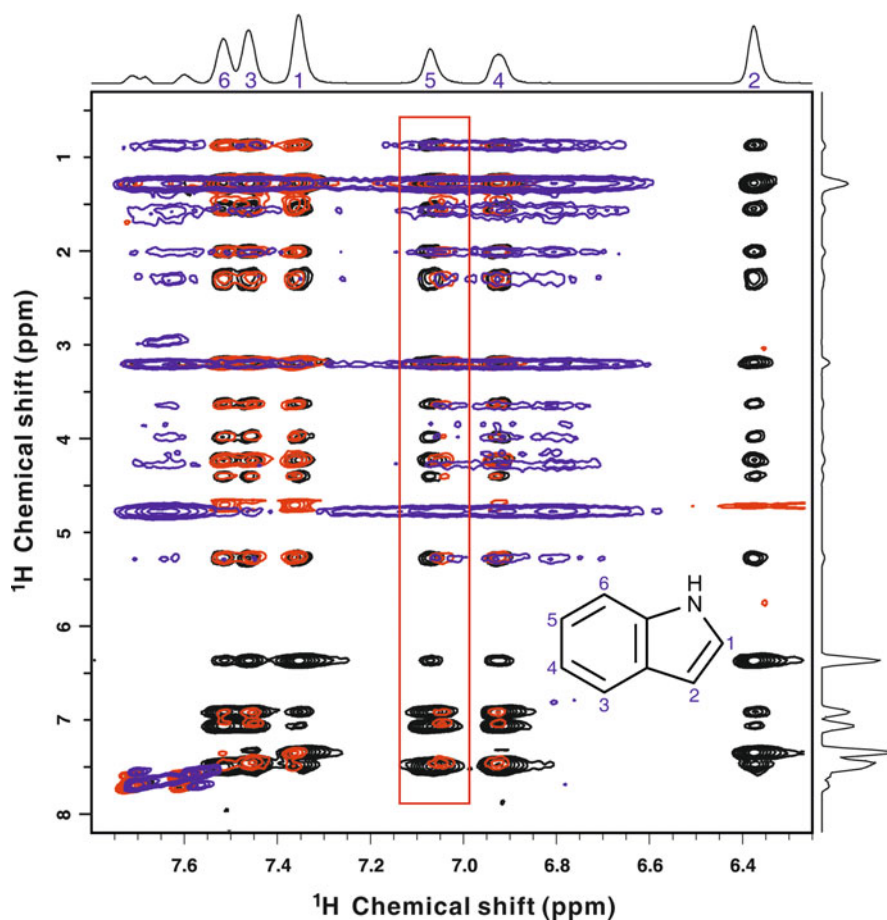
Transmission of signals between cells, within cells, and from the extracellular environment to the cellular interior is essential to life, and the dynamic properties of the signaling proteins are crucial to their functions [156]. Therefore, understanding the dynamic nature of a membrane protein within lipid bilayers is crucial to reveal its function mechanism at the molecular level.

SSNMR may be the best technique to study membrane protein dynamics in the native lipid environment [56, 157, 158], and site specific  $^2\text{H}$  labeling at the methyl groups of Ala, Leu, and Val is most commonly used in SSNMR, and is an excellent probe of the dynamics of membrane proteins. The Glaubitz group has recently reported a dynamic picture of the green proteorhodopsin structure using through-space and through-bond correlation experiments in SSNMR [41]. They have used  $\text{U-}[^{13}\text{C}, ^{15}\text{N}]\text{-PR}$  and reversely labeled  $\text{U-}[^{13}\text{C}, ^{15}\text{N}]_{\text{WHYFI}}\text{-PR}$  samples to establish a clear correlation between hydration water and the mobile J-residues (mainly in flexible loops and tails), as indicated in Fig. 9. Hydration water plays an essential role for enhancing molecular dynamics of residues in tails and interhelical loops, but not in transmembrane domains or rigid, structured loop segments. The result is very important for understanding the dynamic proton pumping mechanism of proteorhodopsin. It also provides an approach to study the site-resolved effects of water and lipid bilayers on the dynamic properties of membrane proteins in general.



**Fig. 9**  $^{13}\text{C}$ -DARR spectra of the aliphatic regions of  $\text{U-}[^{13}\text{C}, ^{15}\text{N}]_{\text{WHYFI}}\text{-PR}$  at 273 (blue, gel phase) and 313 K (red, liquid crystal phase) (a) and the modified homology model of green PR (b). Helical residues influenced by changes in membrane elasticity (labeled in blue) are found in helices C, E, F, and G as well as in loops EC and EF. These residues disappear in the fluid membrane but are visible in the gel phase. This indicates that especially helices C and G but also E and F undergo thermal equilibrium fluctuations in the ground state of PR. Adapted from [41] with permission from the American Chemical Society

Contact and association of hydrophilic peptides and peripheral proteins with cellular membranes are commonly found through the sequence motifs which contain basic and aromatic amino acid residues. Those sequence motifs are not only critical for protein binding but also important for local disruption and penetration of membranes, for recruitment of lipids, and for membrane fusion [159–165]. However, the locations and insertion depths in membrane lipids are quite different as compared with those residues in for transmembrane proteins. It is not straightforward to apply those well established principles generalized from transmembrane proteins to peripheral membrane proteins. For example, we have recently investigated the insertion depth of Trp residue into POPC lipid bilayers in different peptide sequencing, as shown in Fig. 10. Clearly the penetration depth of the side chain of Tryp residue in AWA and VAMP2 peptides are deeper than the



**Fig. 10**  $^1\text{H}$  MAS NOESY spectra of indole (black), AWA (red), and VAMP2 (blue). The indole structure and numbering are for aid of peak assignment

free indole ring in lipid bilayers, as indicated by an up-field shift of the C5 at the indole ring due to the even strong hydrophobic interaction by the  $^1\text{H}$  MAS NOESY experiments. Currently, we are conducting further research on peptide sequencing, composition, and insertion depth with basic and aromatic amino acid residues in different membrane lipids.

$^1\text{H}$  MAS NOESY experiments have been widely used in SSNMR to study peptide–lipid interactions because of the fast axially rotation and segmental motion of membrane lipids in the liquid crystalline phase which average out efficiently the  $^1\text{H}$ - $^1\text{H}$  dipolar couplings, resulting in a high resolution  $^1\text{H}$  spectrum of membrane lipids under the slow MAS frequencies [166], which leads to the rapid applications of the NOESY-type [167] of solution NMR methods to study peptide-membrane interactions in MAS SSNMR [168–170].

The insertion depth of aromatic residues in membrane lipids depends not only on the peptide sequencing and charge but also on the lipid composition, state, hydration level, and peptide/lipid molar ratio, so the insertion presents a complicated dynamic mechanism.

### 3.6 Sensitivity Enhancement

SSNMR sensitivity can be enhanced by manipulating the Boltzmann factor by such means as increasing the field strength, decreasing the system temperature, and transferring proton polarization to a rare nucleus. However, the maximum polarization transferred in the CP experiment is determined by  $\gamma_I/\gamma_S$  and the overall enhancement is not big enough for the free use of multidimensional correlation experiments on low-concentrated membrane proteins, especially for identifying the dynamic switch of GPCRs between an inactive (R) state and an active (R\*) conformation [171]. Dynamic nuclear polarization (DNP) [172, 173], chemically induced nuclear polarization (CIDNP), or photo-chemically induced nuclear polarization (photo-CIDNP) [174–176], and spin-exchange optical pumping (SEOP) [177, 178] have been exploited to deliver much higher spin polarizations. For example, by transferring the polarization of electron spins to nuclei, the MAS-DNP has been successfully applied to study the intermediate states in the photocycle of bacteriorhodopsin, a 7TM light-driven proton pump. The enhanced sensitivity of DNP permitted for the first time the characterization of the retinal conformation in the K, L, and M states [35, 179, 180]. For detailed description of the DNP experiment, please refer to the chapter of “Dynamic Nuclear Polarization: New Methodology and Applications” in this book.

The NMR signal-to-noise ratio is directly proportional to the square root of the number of transitions and the total experimental time is mainly determined by the repetition time between two successive transitions, that is about five times of the spin longitudinal relaxation time  $T_1$ . Therefore, NMR sensitivity enhancement can be achieved by shortening the spin longitudinal relaxation time  $T_1$ . Ishii and colleagues have demonstrated that  $^1\text{H}$   $T_1$  values of the two model proteins,

lysozyme and ubiquitin, in microcrystals can be reduced to 60 ms by Cu–EDTA doping without major degradation in the resolution of their  $^{13}\text{C}$  CP-MAS spectra [181]. They have further presented the paramagnetic relaxation-assisted condensed data collection (PACC) experiment on  $\beta$ -amyloid fibrils and ubiquitin to obtain the two-dimensional  $^{13}\text{C}$ - $^{13}\text{C}$  and  $^{13}\text{C}$ - $^{15}\text{N}$  SSNMR spectra in 1–2 d by reducing the  $^1\text{H}$   $T_1$  down to 50–100 ms with carefully adjusted paramagnetic doping [182]. This sensitivity enhancement through shortening data acquisition time will make those sophisticated and time-consuming multidimensional correlation experiments on macromolecules, such as membrane proteins become realistic.

### 3.7 Structure Determination Based on Chemical Shifts

The chemical shift of a nucleus directly reflects the electron density around it and also reflects the influence of the local environment on it, including the nature of the neighboring atoms and hydrogen-bonding, etc. Therefore, chemical shifts are valuable parameters for protein structure determination. Very recently, three chemical shift-based methods, Cheshire [183], CS-Rosetta [184], and CS23D [185] have been reported to determine the native state structures of proteins in solution up to 130 residues and to a resolution of 2 Å or better. These approaches have also been successfully demonstrated to determine the native structures of GB1, GB3, ubiquitin, and SH3 using SSNMR chemical shift data to a relatively high resolution [186, 187]. These methods may have the potential to determine the properties of even bigger proteins, for example membrane proteins, to a higher resolution with further improvement of the methods.

## 4 Perspective

Magic-angle spinning solid-state NMR has made tremendous progress, showing its capability of determining membrane protein structure, ligand binding, and protein dynamic conformation on a variety of time scales at atomic resolution. Further developments may be threefold:

1. Sample preparation. New attempts might focus on developing new labeling schemes to reduce further the spectral overlap and dipolar truncation, expressing a sufficient amount of functioning membrane proteins using a cell-free expression system, and a new approach to produce high ordered NMR samples in membrane native environment.
2. New NMR methods which should be suitable for fully hydrated membrane samples with better measurement accuracy and less dipolar truncation effects.

3. Integration of all newly developed techniques, including cell-free expression, specific labeling, new pulse sequences, molecular dynamic simulations and DNP.

With all the integrated techniques, we should be in a good position to solve a membrane protein structure based purely on the SSNMR constraints.

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# Dynamic Nuclear Polarization: New Methodology and Applications

Kong Hung Sze, Qinglin Wu, Ho Sum Tse, and Guang Zhu

**Abstract** One way to overcome the intrinsically low sensitivity of Nuclear Magnetic Resonance spectroscopy is to enhance the signal by dynamic nuclear polarization (DNP), where the polarization of high-gyromagnetic ratio ( $\gamma$ ) electrons is transferred to the surrounding nuclei using microwave (MW) irradiation. Recent developments in DNP instrumentations and applications have shown that DNP is one of the most effective methods to increase the nuclear spin polarization in inorganic, organic, and biological materials. It is possible to obtain a solution of molecules containing hyperpolarized nuclei in combination with methods to dissolve rapidly the polarized solid sample. In this chapter, a brief introduction on a theoretical basis and some of new DNP applications in NMR spectroscopy as well as medical applications in Magnetic Resonance Imaging (MRI) are described.

**Keywords** DNP · NMR · MRIh

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## 1 Introduction

Nuclear Magnetic Resonance (NMR) is an important spectroscopic tool for the identification and structural characterization of molecules in chemistry and biochemistry. The most significant limitation of NMR spectroscopy compared to other spectroscopic techniques is its relatively low sensitivity, which thus often requires long measurement times or large amounts of sample, typically half a milliliter (mL) of sample at rather high concentrations. The origin of low sensitivity in NMR is well known to be due to the small magnetic moment of nuclear spins, which yields small Boltzmann polarizations and weak absorption signals. One of the ways to overcome this low signal-to-noise ratio is to enhance the signal by the creation of hyperpolarized transitions. This can be achieved by a process which was named dynamic nuclear polarization (DNP), where the polarization of high-gyromagnetic ratio ( $\gamma$ ) electrons is transferred to the surrounding nuclei using microwave (MW) irradiation.

In 1953, Albert Overhauser [1] first proposed that it was possible to transfer polarization to nuclei from electrons in metals by saturating the electron transition. This idea was not widely accepted until experimentally verified by Carver and Slichter with low field (3 mT) experiments performed on lithium metal and other materials with mobile electrons [2, 3]. This was soon expanded to be applied to solid dielectrics by Abragam and Proctor [4]. Extension of electron-nuclear and other high polarization transfer experiments involving noble gases, parahydrogen, semiconductors, or photosynthetic reaction centers [5–10] to contemporary solid-state and solution experiments is very appealing, since it could significantly enhance the sensitivity in a variety of NMR experiments. In particular, the theoretical enhancement for electronuclear polarization transfers is approximately  $(\gamma_e/\gamma_H)$ , where the ratio is 660, making the gains in sensitivity ideally very large. Accordingly, during the 1960s and 1970s, there were extensive efforts to perform electron nuclear polarization transfer in liquids and solids. In 1980s, work has been carried out to couple DNP to magic-angle-spinning solid-state NMR (MAS-ssNMR).

This concept of nuclear polarization enhancement, originally proposed by Overhauser in 1953, was first experimentally demonstrated in metals and subsequently in liquids, which are two distinct types of systems with mobile electrons.

Thus, DNP is not a new area of scientific endeavor, but rather one undergoing a transition from low to high magnetic fields and frequencies. This chapter outlines the theoretical descriptions of DNP mechanisms followed by recent developments in DNP instrumentations and applications. DNP has proven to be one of the most effective methods to increase the nuclear spin polarization in inorganic, organic, and biological materials. In combination with methods to dissolve rapidly the polarized solid sample it is possible to obtain a solution of molecules containing hyperpolarized nuclei. This has enabled new applications in NMR spectroscopy as well as medical applications in Magnetic Resonance Imaging (MRI).

## 2 Theories of Dynamic Nuclear Polarization

In electron-nuclear based DNP experiments, it is required that the electron paramagnetic resonance (EPR) spectrum be irradiated with microwaves that drive the exchange of polarization between the electrons and the nuclear spins. In the case of liquids, these transitions are based on the Overhauser effect and, in solids, other mechanisms – the solid effect (SE), thermal mixing (TM), or the cross effect (CE) – dominate the polarization transfer process. Since DNP experiments require irradiation of the EPR spectrum, they were confined to relatively low magnetic fields because of the paucity of high frequency microwave sources. In particular, the microwave sources used in both the liquid and solid-state experiments were klystrons that operate at 640 GHz, constraining DNP-MAS experiments to 660 MHz  $^1\text{H}$  frequencies. Thus, for DNP to be applicable to the higher fields employed in contemporary NMR experiments, new instrumental approaches to produce higher frequency microwaves are necessary.

A theoretical analysis of electron and nuclear system requires the quantum mechanical representation. In a DNP experiment, the general static Hamiltonian is written as

$$\begin{aligned}
 H &= H_E + H_N + H_{EN} \\
 &= \omega_{0E}E_Z - \omega_{0N}N_Z + H_{EN}^{\text{is}} + H_{EN}^{\text{di}} \\
 &= \omega_{0E}E_Z - \omega_{0N}N_Z + K_{\text{SE}}(E_ZN_Z + E_YN_Y + E_XN_X) + K_{\text{PSE}}E_XN_Z
 \end{aligned} \tag{1}$$

where  $H_E$  and  $H_N$  are the Hamiltonians for electron and nuclear respectively.  $H_{EN}$  is the hyperfine coupling, which can be separated into the isotropic hyperfine interaction  $H_{EN}^{\text{is}}$  and the anisotropic dipolar coupling  $H_{EN}^{\text{di}}$  between electrons and nucleus. The  $H_{EN}$  can be further expressed in a form where the coefficients  $K_{\text{SE}}$  and  $K_{\text{PSE}}$  denote the secular and pseudosecular hyperfine interactions.  $\omega_{0E}$  and  $\omega_{0N}$  are the electron and nuclear Larmor frequencies.

DNP experiments can be classified based on the polarizing mechanisms. We will discuss the continuous-wave (CW) and time domain polarization mechanisms in the following section.

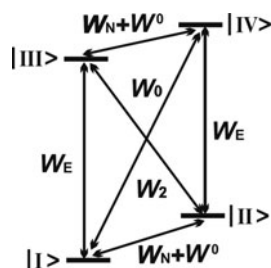
## 2.1 Continuous-Wave DNP Mechanism

Continuous wave DNP polarization-transfer mechanisms can be generally classified into four fields, namely the Overhauser Effect (OE), the Solid Effect (SE), the Cross Effect (CE), and Thermal Mixing (TM). They have all been successfully applied to both solid and liquid samples. It is reported that many DNP applications have been performed at low magnetic fields based on these different polarization effects.

### 2.1.1 Overhauser Effect

The OE DNP mechanism is a relaxation process which relies on the mixing of an electron and a nuclear spin. More specifically, in liquid cases, these relaxation processes are based on time-dependent dipolar and scalar interactions between electrons and nuclei. These interactions are governed by molecular rotational and translational motion in the dipolar case and by chemical exchange and fast relaxation in the scalar case [11]. In solid cases, the mobile electrons are required for the OE relaxation process. These electrons can be offered by the conduction band of metal or one-dimensional organic conductors. In addition, the condition  $\omega_{0E}\tau < 1$  ( $\tau$ , rotational correlation time of paramagnetic species) is required for OE to be established. The latter condition is difficult to satisfy with large value of  $\omega_{0E}$ . In addition,  $\omega_{0E}$  is in proportion with the strength of magnetic field. Therefore, the efficiency of the OE polarizing mechanism sharply decreases at high magnetic fields. However, the OE polarization is extremely useful for liquids, for it is the only practical mechanism for the direct application of liquids. We will give the outlines of the main OE features in a later section. A more detailed description can be found in the recent publication and review by Hofer et al. [12].

Figure 1 shows the energy level diagram of the transitions responsible for the OE.  $W_E$  and  $W_N$  are the rates for the EPR and NMR transitions, respectively.  $W^0$  is the nuclear relaxation rate in the absence of electrons, while  $W_0$  and  $W_2$  are the rates for the zero and double quantum transitions, respectively. In OE polarization, the allowed transitions of EPR are saturated. The observed NMR signals are enhanced by changing the nuclear spin population, which is induced by the zero and double



**Fig. 1** Energy level diagrams and main transition rates for the Overhauser effect (OE) DNP mechanism

quantum transitions. The enhancement factor is defined as  $\xi = N_Z/N_{eq}$ , where  $N_Z$  and  $N_{eq}$  are the nuclei numbers in the polarizing and thermal balance states, respectively. After solving the rate equations for this system, we can obtain the following equation for the signal enhancement:

$$\xi = 1 - \rho\mu\eta\frac{\gamma_E}{\gamma_N} \tag{2}$$

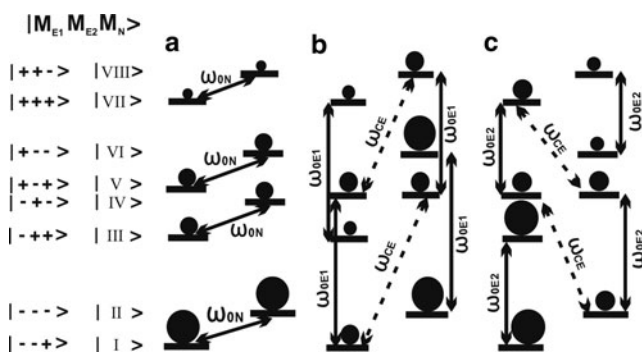
with

$$\rho = \frac{W_2 - W_0}{W_0 + 2W_N + W_2}, \tag{3}$$

$$\mu = \frac{W_0 + 2W_N + W_2}{W_0 + 2W_N + W_2 + W^0}, \tag{4}$$

$$\eta = \frac{\langle E_0 \rangle - \langle E_Z \rangle}{\langle E_0 \rangle}, \tag{5}$$

where  $\eta$  corresponds to the electron transition saturation factor ranging from 0 (zero saturation,  $\langle E_z \rangle = \langle E_0 \rangle$ ) to 1 (complete saturation,  $\langle E_z \rangle = 0$ ). From the above equations, the coupling parameter  $\rho$  can vary from  $-1.0$  to  $0.5$ , corresponding to pure scalar coupling and pure dipolar coupling, respectively. In liquid samples, the dipolar coupling is a major interaction leading to a maximum enhancement of 330. On the other hand,  $\rho$  also depends on the electron Larmor frequency  $\omega_{0E}$  and the degree of molecular motion. This relation is shown in Fig. 2 of the Maly et al. review [13]. It is obvious to find that the OE process is inefficient at high magnetic fields. Many scientists (Armstrong and Han, Grucker et al., and Hofer et al.)



**Fig. 2** (a) Population distribution at thermal equilibrium for a general three-spin system. (b) Saturation of the allowed EPR transitions for one of the dipolar coupled electrons ( $\omega_{0E1}$ ) leads to negative enhancement. (c) Saturation of the transition corresponding to the second electron ( $\omega_{0E2}$ ) leads to positive enhancement.  $M_{E1}$ ,  $M_{E2}$ , and  $M_N$  are the spin states of electron 1, 2, and nucleus. Reproduced with permission from [13]

discussed this phenomenon and show the determination of the coupling parameter  $\rho$  in their articles [12, 14, 15]. The leakage factor  $\mu$  describes the nuclear spin relaxation by electron spins. The value of  $\mu$  can be assumed to lie between 0 and 1. The  $\mu = 0$  represents no relaxation caused by the electron-nuclear coupling, while  $\mu = 1$  means no other relaxation mechanisms ( $W^0 = 0$ ). The saturation factor  $\eta = 1$  when the electron transitions are completely saturated. This condition is important in the application of polarizing agents based on nitroxide radicals, although it depends on the number of hyperfine lines in the spectra. For concurrent time-dependent spin exchange, an additional factor  $\sigma$  can be introduced to describe the attenuate contribution on signal enhancement by scalar coupling. Generally, the value of  $\sigma$  can vary between 0 and 1, corresponding to no contribution and large contribution of the scalar coupling, respectively [13].

### 2.1.2 Solid Effect

The solid effect (SE) is a DNP mechanism which requires states mixing caused by the nonsecular component  $K_{\text{PSE}}$  of the hyperfine coupling [16]. The pseudosecular term  $K_{\text{PSE}}$  contains the form  $E_Z N^+$  and  $E_Z N^-$  (E: Electron spin operator, N: Nucleus spin operator), which leads to a mixing of the states of the system. In SE polarization, the new mix states are generated from the original states with a coefficient  $p$ , which can be calculated by the first order perturbation theory and is given by

$$p = -\frac{3\gamma_E\gamma_N}{4\omega_{0N}R^3} \sin\theta \cos\theta e^{-i\phi} \quad (6)$$

where  $R$ ,  $\theta$ , and  $\phi$  are the polar coordinates describing the electron-nuclear vector. The irradiation probability of zero quantum or double quantum transitions is proportional to  $4p^2$  [16]. Furthermore, this transition probability and sensitivity enhancement will be scaled with  $\omega_{0N}^{-2}$ , since  $p$  is proportional to  $\omega_{0N}$ . Therefore, this correlation has restricted the application of SE in high field DNP experiments.

Practically, the SE requires the use of polarizing agents with a homogeneous EPR line width and an inhomogeneous spectral width smaller than the nuclear Larmor frequency. These agents can ensure that only one of the forbidden transitions is excited at a time. However, the Differential Solid Effect (DSE, [17]) simultaneously exists and leads to partial or complete cancellation of the polarization effect.

### 2.1.3 Cross Effect/Thermal Mixing

The Cross Effect mechanism is based on allowed transitions and involves the interaction of electron spin packets in an inhomogeneously broadened EPR line. A similar effect, found in a homogeneously broadened EPR line, is called thermal mixing. Wollan proposed a method to process the intermediate case of

inhomogeneously and homogeneously broadening at low magnetic field [18]. Since both CE and TM have recently been used at high magnetic fields, the theoretical framework describing these effects needs to be fully established [19–23].

Griffin and co-workers [22] proposed that the CE is defined as a three-spin process, involving the interaction between two dipolar coupled electrons with EPR frequencies  $\omega_{E01}$  and  $\omega_{E02}$  that satisfy the relation

$$\omega_{E02} - \omega_{E01} = \omega_{0N} \quad (7)$$

When the EPR lines are broadened by the inhomogeneous anisotropy, this is the dominant mechanism. The electrons are weakly coupled via electronic cross relaxation. Theoretically, the CE requires that the inhomogeneous breadth  $\Omega$  is larger than the nuclear Larmor frequency  $\omega_{0N}$  to make two effective EPR resonance frequencies. Meanwhile, the homogeneous width must satisfy  $\Delta < \omega_{0N}$ . The biradical chemical polarizing agents can improve the polarization effect of DNP systems to achieve this condition by a dipolar coupling between its two electrons. In contrast, the TM polarization mechanism contains homogeneously broadened EPR lines, where  $\Delta > \omega_{0N}$  is satisfied. This condition requires a high concentration of polarizing agent at high magnetic fields, which will restrict the resolution in an MAS NMR experiment.

Figure 2 shows the thermal equilibrium spin population for a three-spin system. Generally there is no degeneracy present. When an appropriate polarizing agent is used, the energy levels  $|IV \rangle$  and  $|V \rangle$  or  $|VI \rangle$  and  $|III \rangle$  become degenerated (Fig. 2a). Irradiation of EPR transition and CE transitions leads to positive (Fig. 2b) or negative enhancement (Fig. 2c) of the nuclear polarization.

The signal enhancements of CE/TM are scaled with the magnetic field strength because the line width of the EPR spectra decreases along with the increasing strength of the magnetic field. Despite this drawback, the CE/TM polarizing mechanisms have been widely applied to polarize biological solids at high magnetic fields [24, 25].

The TM effect is less efficient compared to the CE. The TM can be described by a series of interacting systems: the electron Zeeman system (Ezs), the electron dipolar system (EDS), and the nuclear Zeeman system (NZS) [19].

## 2.2 Pulsed DNP

S.R. Hartmann and E.L. Hahn reported that pulsed DNP experiments based on coherent polarization transfer such as the Hartmann–Hahn cross polarization (HHCP) are more efficient and do not show a field dependence of the polarization transfer with increasing magnetic field strengths.

HHCP is a fundamental technique widely used in ssNMR spectroscopy. With this mechanism, the polarization of one nucleus can be transferred to another nucleus [26, 27]. This transfer requires the rotating frame Hartmann–Hahn matching condition to be satisfied:

$$\omega_{1N} = \omega_{1E} \quad (8)$$

where  $\omega_{1N}$  and  $\omega_{1E}$  are the field strengths of the applied RF fields. In solid-state NMR experiments, the radio frequency fields are strong so their excitation profiles can cover the entire NMR spectrum and they can efficiently spin-lock both spin species. Thus, the Hartman–Hahn condition can be fulfilled.

For DNP experiments, however, there is large inhomogeneous broadening of high field EPR spectra and the spectral breadth  $\Delta$  can usually exceed several hundred megahertz. In other words, the microwave field magnitude power is one order of magnitude smaller in comparison with the conditions found in ssNMR. The condition of matching electron-nuclear cross polarization (eNCP) cannot be satisfied by modification of the microwave field strength and the RF strength. Therefore, off resonance effects of EPR must be considered in this case.

From thermodynamic theory, the signal enhancement by CP transfer between I (high- $\gamma$ ) and S (low- $\gamma$ ) spins is given by

$$\xi = \left( \frac{\gamma_I}{\gamma_S} \right) \frac{1}{1 + N_S/N_I} \quad (9)$$

The efficient CP transfer between two spins is required so that the concentrations of I is much more than S. In this case, when  $N_S \gg N_I$ ,

$$\xi \approx \frac{\gamma_I}{\gamma_S},$$

and the full ratio  $\gamma_I/\gamma_S$  can be transferred. NMR-CP experiments with the HHCP schemes are typically performed by transferring polarization from abundant high  $\gamma$  spins I to dilute low  $\gamma$  spins S. The condition of efficient CP transfer can be easily satisfied. However, in DNP experiments, the concentration of unpaired electrons from high  $\gamma$  polarizing agents is about four orders of magnitude lower than the concentration of low  $\gamma$   $^1\text{H}$  nucleus. Therefore, the enhancement factor  $\xi < 1$  and only a small amount of electrons can be transferred to  $^1\text{H}$ . The pumping of CP polarization is applied to improve the enhancement because the relaxation time of the electron is much shorter than that of the nucleus.

### 2.2.1 Integrated Solid Effect

Integrated Solid Effect (ISE) was first introduced by Henstra et al. [17]. It can overcome the low efficiency of SE when the homogeneous width is much larger than the nuclear Larmor frequency ( $\Delta \gg \omega_{0N}$ ), in which the polarization effect could be canceled by simultaneous saturation of the forbidden transitions at  $\omega_{0E} \pm \omega_{0N}$ . The ISE can preserve the polarization in the case of  $\Delta \gg \omega_{0N}$  by inverting a forbidden EPR transition prior to saturation of an allowed transition. This effect can be achieved by using a selective inversion pulse after the irradiation on resonance at  $\omega_{0E} \pm \omega_{0N}$ , or applying CW microwave irradiation at a fix frequency

$\omega_{0E}$  while sweeping the magnetic field through the entire EPR line. In the latter case, if the microwave power used for inversion of the electron spin polarization is sufficient, electron spin will go through an adiabatic fast process. This is considered as adiabatic-ISE. It is reported that ISE experiments were performed at low magnetic fields (9 GHz) using either electrons to polarize  $^{29}\text{Si}$  nuclei or transferring the high polarization of a photoexcited triplet to surrounding protons [28–30]. A. Henstra et al. show an ISE build up curve for the adiabatic-ISE. In this experiment, the  $^{29}\text{Si}$  signal in p-type Si at 1.2 K is detected and a factor of about 20 larger enhancements is obtained in an adiabatic-ISE rather than DSE [17].

There are certain limitations in the application of both the ISE and the adiabatic-ISE. The ISE requires high-power microwave pulses for the excitation of a large amount of electrons, while the adiabatic-ISE requires an adiabatic magnetic field sweep. On the one hand, the appropriate microwave sources are not commercially available at high microwave frequencies. On the other, the adiabatic magnetic sweep is difficult to perform at high magnetic fields. Although having these limitations, the ISE has a brilliant future in DNP applications because of the large sensitivity enhancement it can achieve.

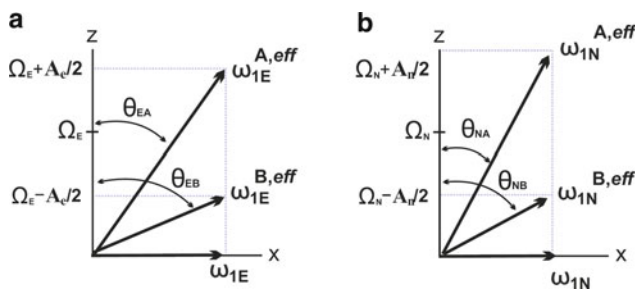
### 2.2.2 NOVEL

Nuclear spin orientation via electron spin locking (NOVEL) is a DNP experiment based on coherent pulses. In NOVEL experiments, the electron magnetizations are locked in the electron rotating frame via electron spin-lock sequences. The Hartmann-Hahn condition (rotating frame or lab-frame) can be satisfied if the field strength of the spin-lock pulse meets the condition  $\omega_{0E} = \omega_{0N}$ ; then the polarization can be transferred from the electrons to the nuclei. van den Heuvel et al. [31] show the nuclear polarization buildup curve for a NOVEL experiment in which the electron polarization from pentacene guest molecules in a photoexcited triplet state is transferred to  $^1\text{H}$  of the naphthalene host crystal. An enhancement of 220 is achieved [31]. To date, NOVEL experiments have only been performed at 9 GHz microwave frequencies [30–33]. Much higher microwave field strengths will be required at higher magnetic fields.

### 2.2.3 Dressed-State Solid Effect

Dressed-state solid effect (DSSE) is a DNP mechanism based on simultaneous near resonant microwave and RF irradiations, which can be established even in the absence of nonsecular hyperfine coupling terms. This experiment was first introduced by Weis et al. [34]. The mechanism of DSSE is illustrated in Fig. 3.  $\omega_{1E}^A$  and  $\omega_{1E}^B$  are two effective fields which have unequal influence at high frequency DNP polarization, because the microwave power is currently limited. It indicates that the EPR transitions cannot be achieved with the same microwave field strength. The matching condition for DSSE polarization is given by





**Fig. 3** Illustration of the electron (a) and nuclear (b) spin effective fields. The effective fields belonging to the EPR and NMR transitions are no longer equal.  $\Omega_E$  and  $\Omega_N$  are the resonant offsets of electrons and nuclei, while  $A_e$  and  $A_n$  are the hyperfine coupling constant of electrons and nuclei.  $\omega_{1E}^{A,eff}$  and  $\omega_{1E}^{B,eff}$  ( $\omega_{1N}^{A,eff}$  and  $\omega_{1N}^{B,eff}$ ) are the unequal effective microwave fields on different resonant offsets of electrons (nuclei).  $\theta_{EA/EB/NA/NB}$  are the angles between the effective microwave fields on electrons/nuclei and the  $z$  dimension

$$\begin{aligned} & \sqrt{(\delta_E + A_e/2)^2 + \omega_{1E}^2} + \sqrt{(\delta_E - A_e/2)^2 + \omega_{1E}^2} \\ & = \pm \left[ \sqrt{(\delta_N + A_n/2)^2 + \omega_{1N}^2} + \sqrt{(\delta_N - A_n/2)^2 + \omega_{1N}^2} \right] \end{aligned} \quad (10)$$

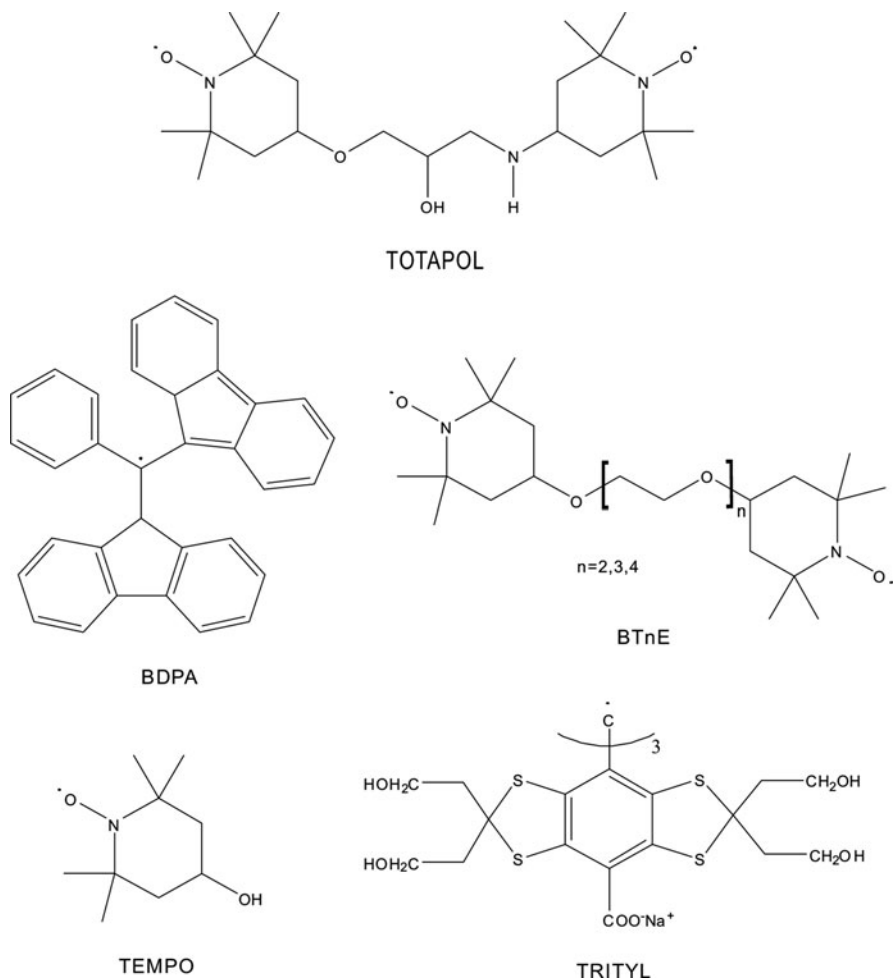
where  $\delta_E$  and  $\delta_N$  are the resonant offsets of electron and nuclear spectra, respectively.  $\omega_{1E}$  and  $\omega_{1N}$  represent the RF field strengths and  $A_e$  and  $A_n$  the hyperfine coupling constants. For NMR-CP experiments, the matching condition is given by the equation

$$\delta_N \approx \pm \sqrt{\omega_{1E}^2 + (A_e/2)^2} - \omega_{1N}^2 \approx \pm \sqrt{\omega_{1E}^2 + (A_e/2)^2} \quad (11)$$

in the case that the microwave irradiation is on resonance and the hyperfine coupling is arbitrary. The expected polarizing region is determined by hyperfine couplings.

### 2.3 Polarizing Agents

Polarizing agents are widely used in DNP experiments. The choices of these agents are strongly dependent on a number of factors, such as the width of the EPR spectrum, the radical solubility and toxicity, the radical reactivity, the temperature dependence of relaxation times, etc. Either exogenous or endogenous radicals can be used as polarizing agents. Figure 4 shows five polarizing agents which



**Fig. 4** Chemical structures of five polarizing agents which have been frequently used in high-field DNP experiments

have been frequently used in high-field DNP experiments. The trityl and BDPA are suitable for SE DNP experiments. The TEMPO is used to allow TM polarizing mechanism, while its biradical derivatives, BTnE and TOTAPOL is used to induce the CE polarizing mechanism. These similar nitroxide based radicals and biradicals are stable and soluble in a variety of solvents, allowing them to be applied in DNP experiments involving both liquid- and solid-states as well as biological samples.

### 2.3.1 Polarizing Agents for the SE

The SE polarizing mechanism requires that the zero and double quantum forbidden transitions are not simultaneously excited and that the enhancement effect is without positive and negative transitions cancellation. To date, two radicals, BDPA and trityl, can satisfy these requirements at high magnetic fields [35, 36]. BDPA is the first polarizing agent applied to investigate the SE mechanism in polystyrene matrix, which reveals an inhomogeneous linewidth of  $\Omega \sim 20$  MHz at a 211-MHz machine [37–39]. However, BDPA is insoluble in aqueous solution, making it very difficult to use in biological science. In contrast, the trityl has good solubility in water and a narrow EPR spectrum width; it has therefore been successfully used in aqueous solution DNP applications with an Oxford HyperSense Oxford DNP polarizer [40].

### 2.3.2 Polarizing Agents for TM and CE

As mention above, under the situation that  $\Omega > \omega_{0N}$ , the TM and CE mechanisms are the main polarizing mechanisms in DNP experiments. The homogeneously broadened EPR leads to TM while the inhomogeneously broadened EPR results in CE. Nitroxide based radicals and biradicals are suitable polarizing agents for TM and CE because their EPR inhomogeneous line widths are about 600MHz at a  $^1\text{H}$  Larmor frequency of 211-MHz. Furthermore, these chemical compounds, which have an EPR spectrum consisting of two narrow sharp lines separated by  $\omega_{0N}$ , are ideal polarizing agents for CE polarization.

## 3 Instrumentation

The instrumentation for DNP experiments has enjoyed great developments in recent years. Generally speaking, on top of an existing NMR system, a microwave source, a waveguide, and a new probe are extra instrumentation required for performing DNP experiments. The microwave source is used to create the microwave and the waveguide will transmit the microwaves from the source to the NMR probe. The new probe is designed with the additional function that it can irradiate the sample with microwaves. In addition, there are some other requirements for special cases, such as performing MAS below 90 K, irradiating the EPR spectrum with maximum enhancement, and so on.

### 3.1 Microwave Sources

The power of microwave sources for DNP depends strongly on the nature of the experiment to be performed. The field strength of the microwave  $\omega_{1E}$  has

a relationship with the microwave power  $P_m$  and the quality factor  $Q_m$  according to the following equation:

$$\omega_{1E} \propto k\sqrt{Q_m \cdot P_m}. \quad (12)$$

Thus a high power microwave source is required for DNP experiments. Currently, the microwave sources can be classified into two different types: solid-state and vacuum electronic devices. Several excellent reviews of the introduction of – and outlook for – microwave source technology are available [41–43].

Among these, gyrotrons and cyclotron resonance masers are high-frequency vacuum electronic devices that have the ability to produce sufficient power in the frequency range of 140–590 GHz for electrons (200–900 MHz for proton nuclei). The electron cyclotron resonance maser can emit the coherent radiation near the relativistic electron cyclotron frequency. The irradiation frequency is given by

$$\omega_g = f \frac{eB_0}{M' m_e c} \quad (13)$$

and

$$M' = \frac{1}{\sqrt{1 - v^2/c^2}}, \quad (14)$$

where  $B_0$  is the strength of the static magnetic field and  $e$  is the charge of the electron.  $f$  is the harmonics of the operation mode,  $m_e$  the mass of the electron, while  $M'$  is a relativistic mass factor given by (14).  $v$  and  $c$  are the speeds of the electron and light, respectively. From (13), the coherent radiation frequency  $\omega_g$  is directly proportional to the strength of the static magnetic field  $B_0$  when the other experimental parameters in (13) remain constant.

### 3.2 Microwave Waveguides

The microwave waveguide is a device which is applied to transmit the microwaves from the source to the NMR probe. The efficient delivery of the microwave irradiation is highly necessary for microwave waveguides. When the microwave frequencies increase, however, the efficiency will decrease. Nowadays, with the help of corrugated waveguides, the loss is almost negligible and the efficiency of the microwave transmission is immensely increased. These devices have already been used in high frequency EPR and DNP applications [24, 44].

### 3.3 *Probes Used for DNP Experiments*

The probes for DNP experiment can be designed with or without a resonant structure. For liquid state DNP experiments, the  $TE_{102}$  rectangular EPR cavity is used for low microwave frequencies polarization, while the cylindrical  $TE_{011}$  cavity is used for high frequencies [45, 46]. In addition, an external RF coil with a slotted cavity is required in a high frequency electron nuclear double resonance (ENDOR) setup to avoid mounting the NMR RF coil inside the resonator. In 1965, Hyde proposed a double resonance probe for ENDOR experiments, which was widely applied later [47]. For MAS experiments, the microwave irradiation is applied vertically to the axis of the rotor in the present design, instead of parallel in the past. Due to space considerations, the microwaves are often placed between the turns of the NMR coil in high frequency situations. The variable tuning circuits are positioned outside the probe and wire transmission is used for the RF power irradiation. This design ensures that the probe also works under low temperature conditions.

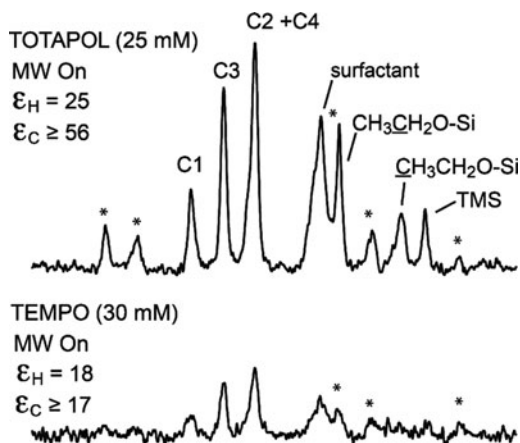
## 4 Applications

NMR is an important spectroscopic tool for the identification and structural characterization of molecules in chemistry and biochemistry. The most significant limitation of NMR compared to other spectroscopic techniques is its intrinsically low sensitivity, which thus often requires long measurement times or large amounts of sample. DNP has been shown to be an effective method to increase the nuclear spin polarization in inorganic, organic, and biological materials; hence DNP has become an attractive technique to boost the sensitivity of NMR signals allowing NMR spectra of small amounts of samples from natural sources or from chemical synthesis to be readily acquired. Perhaps even more interestingly, the availability of the entire hyperpolarized NMR signal in one single scan allows the measurement of transient processes in real time, if applied together with a stopped-flow technique. In combination with the methods to dissolve rapidly the polarized solid sample, it is possible to obtain a solution of molecules containing hyperpolarized nuclei. This has enabled new applications in NMR spectroscopy as well as medical applications in MRI.

### 4.1 *Applications to Small Molecules*

By using incipient wetness impregnation of the solid samples with a solution of organic radical species such as TEMPO or TOTAPOL, surface enhanced NMR spectroscopy by DNP can be carried out to allow fast characterization of functionalized solid surfaces. Polarization is transferred from the radical protons of the solvent to the rare NMR active nuclei at natural abundance on the surface. Lesage et al. [48] have applied surface enhanced NMR spectroscopy by DNP to

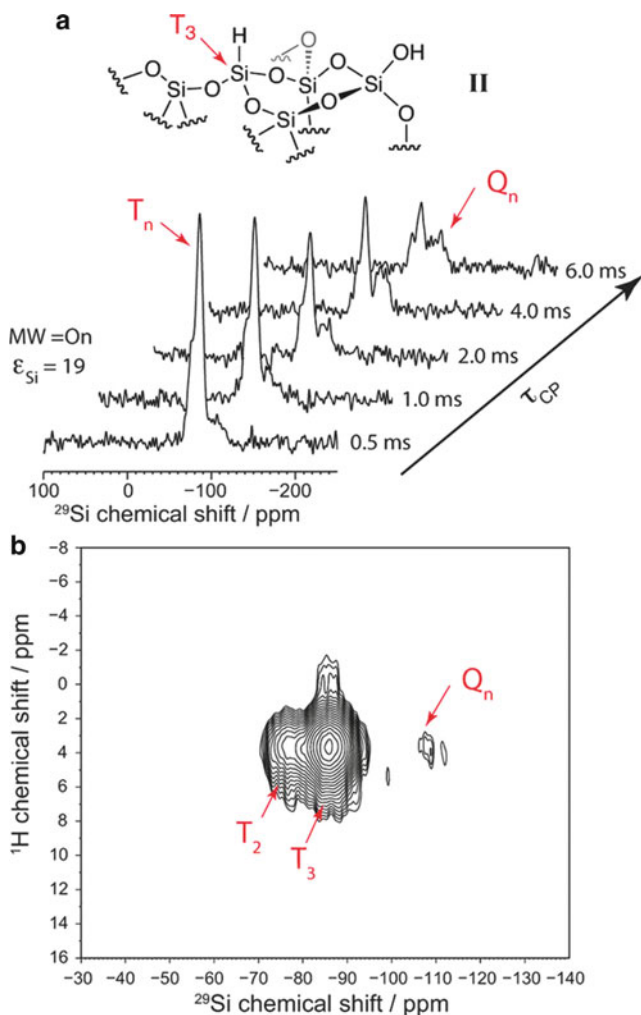
**Fig. 5**  $^{13}\text{C}$  CP MAS spectra of a silica framework with MW irradiation at 263 GHz to induce DNP. The figures compare the best enhancements observed using TEMPO and TOTAPOL radicals. It should be noted that significant DNP enhancement of the alkyl moiety of the surface ethoxy groups was also observed. Reproduced with permission from [48]



yield a 50-fold signal enhancement on  $^{13}\text{C}$  signals for surface species covalently incorporated into a silica framework (Fig. 5). Lelli et al. [49] have applied the technique to study the distribution of surface bonding modes and interaction of functionalized silica materials by observing the  $^{29}\text{Si}$  signals directly (Fig. 6). The remarkable gain in sensitivity and time provided by surface-enhanced silicon-29 DNP NMR spectroscopy is on the order of a factor of 400. This has also allowed the acquisition of the previously inconceivable two-dimensional correlation spectra, enabling more detailed characterizations of these functionalized surfaces [48, 49].

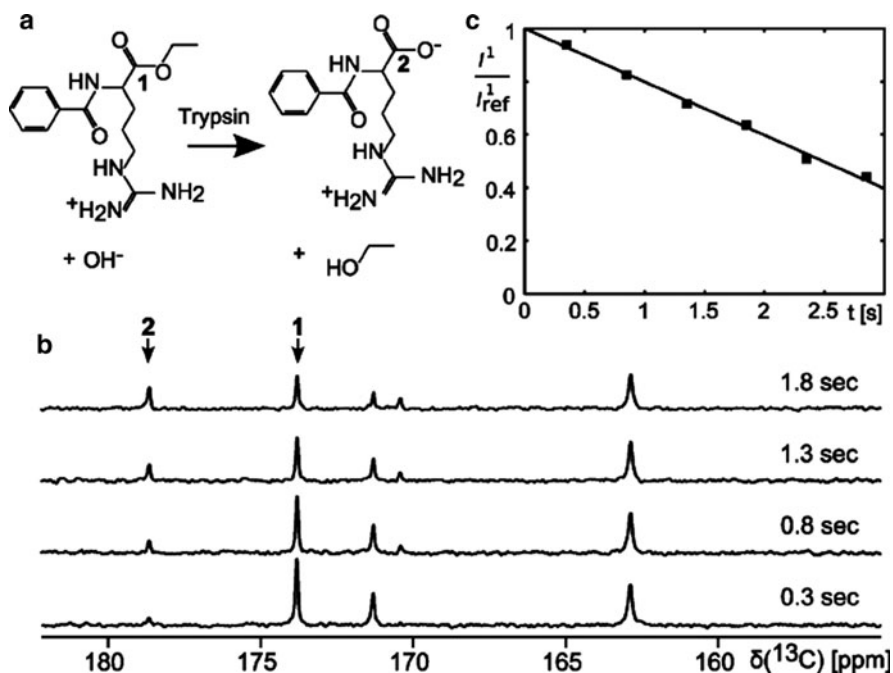
The unprecedented sensitivity observed in DNP-NMR means that sufficient signals can be obtained from a single scan, allowing reactions even far from equilibrium to be studied in real time. Of great importance, the DNP-NMR method is compatible with quantitative rate determination experiments where a single spin in the reactant is labeled in its spin state by a selective radio frequency pulse for subsequent tracking through the reaction, allowing the unambiguous identification of its position in the product molecule. Zeng et al. [50] have demonstrated the application of DNP-NMR to the Diels-Alder reaction of 1,4-diphenylbutadiene with 4-phenyl-1,2,4-triazole-3,5-dione, where reaction rates could be obtained accurately and reproducibly. In particular, the high chemical shift specificity afforded by high-resolution NMR permitted the simultaneous determination of reaction rates and mechanistic information in one experiment. Similarly, time-resolved DNP-enhanced NMR has been used to study the enzymatic reaction of the conversion of *N*-benzoyl-L-arginine-ethyl ester into the product *N*-benzoyl-L-arginine by trypsin enzyme (Fig. 7) [51].

The control of circulating drug concentrations by therapeutic drug monitoring (TDM) is vital for drugs when under- or overdosing may lead to loss of therapeutic efficacy or to adverse effects. Effective TDM depends on effective analytical platforms for the fast detection, identification, and quantification of circulating drugs with a narrow therapeutic range. As a result of the low concentrations of drugs and their metabolites in blood plasma, analytical tools are needed to provide high sensitivity and specificity. NMR spectroscopy is quantitative if care is taken



**Fig. 6** (a) DNP-enhanced silicon-29 CPMAS spectra of compound II as a function of the CP mixing time  $\tau_{CP}$ . (b) Contour plot of a two-dimensional  $^1\text{H}$ - $^{29}\text{Si}$  spectrum of II recorded with DNP. Reproduced with permission from [49]

that signal areas are not affected by different relaxation rates, and the method yields chemical information on the detected molecules in the absence of sample derivatization or sample separation. A principal limitation of conventional  $^1\text{H}$  NMR is, however, the relatively low signal dispersion over a spectral window of approximately 10 ppm, which restricts the capability to resolve chemical compounds in the complex spectral backgrounds of biofluids. The use of  $^{13}\text{C}$  NMR spectroscopy can address poor signal dispersion problems as the  $^{13}\text{C}$  chemical shift dispersion is approximately 20-fold larger than that for  $^1\text{H}$ . Spectral interference of the biofluid



**Fig. 7** Illustration of time-resolved DNP enhanced NMR. (a) Trypsin catalyzed conversion of BAEE into BA. (b) Time-course of the reaction, observed by  $^{13}\text{C}$  NMR at natural abundance, using 3.3 mM DNP enhanced BAEE and 54 mM Trypsin. (c) Linear fit of normalized intensities from **b**, yielding the rate constant  $k_{cat} = 12.1 \text{ s}^{-1}$ . Reproduced with permission from [51]

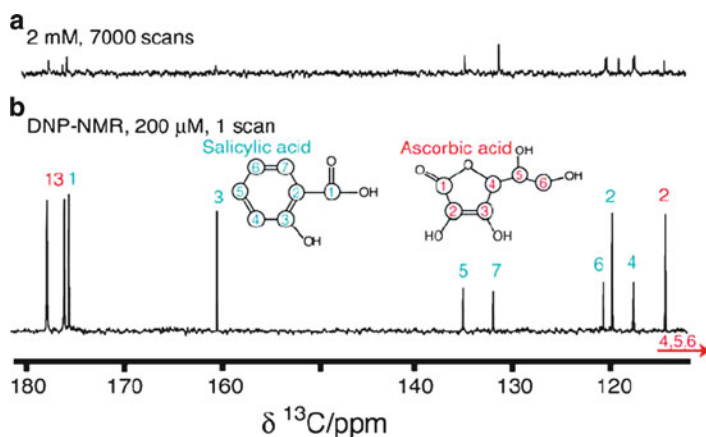
background with  $^{13}\text{C}$ -enriched drugs is further reduced because of the low  $^{13}\text{C}$  abundance (1.1%) in the biofluid background. Therefore  $^{13}\text{C}$  DNP-NMR in the form of the hyperpolarization-dissolution method should afford the sensitivity and spectral resolution for the direct detection and quantification of numerous isotopically labeled circulating drugs and their metabolites in single liquid-state NMR transients. Lerche et al. [52] has applied  $^{13}\text{C}$  DNP-NMR to perform in vitro quantitative assay of drug and its metabolites in blood plasma. The lower limit of detection for the anti-epileptic drug  $^{13}\text{C}$ -carbamazepine and its pharmacologically active metabolite  $^{13}\text{C}$ -carbamazepine-10,11-epoxide is  $0.08 \mu\text{g/mL}$  in rabbit blood plasma analyzed by single-scan  $^{13}\text{C}$  DNP-NMR. Comparison of quantitative DNP-NMR data with an established analytical method (liquid chromatography–mass spectrometry) yields a Pearson correlation coefficient  $r$  of 0.99. These results indicate that  $^{13}\text{C}$  DNP-NMR meets the sensitivity and accuracy requirements for quantitative analysis of the drug and its metabolites in blood plasma. It should be noted that all DNP-NMR determinations were performed without analyte derivatization or sample purification other than plasma protein precipitation. Therefore, quantitative DNP-NMR is an emerging methodology which requires little sample preparation and yields quantitative data with high sensitivity for therapeutic drugs.



Advances with para-hydrogen induced polarization (PHIP) also open up new fields of applications for portable low-field DNP NMR. Glogglar et al. [53] report the possibility of tracing drugs down to the micromolar regime. Selectively polarized nicotine quantities similar to those found in one cigarette and morphine extracted from an opium solution were detectable after polarization with para-hydrogen in single-scan  $^1\text{H}$  NMR experiments. Moreover, they demonstrated the possibility to enhance selectively and detect the  $^1\text{H}$ -signal of drug molecules with PHIP in proton rich standard solutions that would otherwise mask the  $^1\text{H}$  NMR signal of the drug.

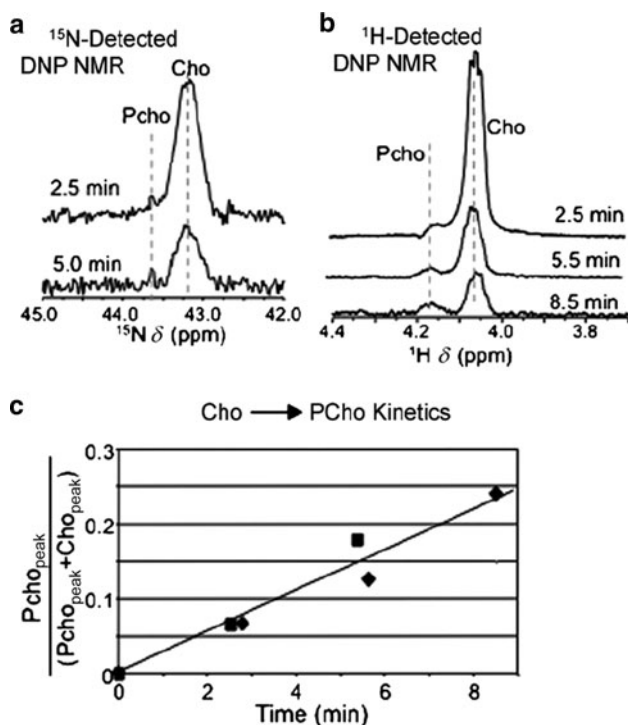
NMR spectroscopy is a well known and versatile technique for the study of molecular interactions, even when these interactions are relatively weak. Signal enhancement by several orders of magnitude through DNP NMR alleviates several practical limitations of NMR-based interaction studies (Fig. 8) [54].  $^{13}\text{C}$  DNP NMR can be applied for ligand binding studies at natural isotopic abundance of  $^{13}\text{C}$  [54]. Resultant screens are easy to interpret and can be performed at  $^{13}\text{C}$  concentrations below micromolar concentration. Of importance, the ligand-detected molecular interaction can be assessed and quantified with enzymatic assays that employ hyperpolarized substrates at varying enzyme inhibitor concentrations. Hence, the physical labeling of nuclear spins by hyperpolarization has offered fast novel in vitro experiments to be performed with low material requirement and without the tedious need for synthetic modifications of target or ligand.

This increase in sensitivity by ex situ DNP-NMR has triggered new research avenues, particularly concerning the in vivo monitoring of metabolism and disease by NMR spectroscopy. So far such gains have mainly materialized for experiments that focus on nonprotonated, low-gamma nuclei, targets favored by their relatively long  $T_1$  relaxation times, which enable them to withstand the transfer from the



**Fig. 8** Comparison of  $^{13}\text{C}$  NMR spectra of a ligand mixture of salicylate and ascorbate recorded on a state-of-the-art NMR spectrometer (18.7 T) with cryogenic probe for signal detection (a) and recorded with DNP-NMR on a 9.4 T spectrometer (b). Ligands were tenfold concentrated for conventional NMR (a) to allow signal detection after 7,000 scans (20 h). Reproduced with permission from [54]

cryogenic hyperpolarizer to the reacting centers of interest. Recent studies have also shown that transferring this hyperpolarization to protons by indirectly detected methods could successfully give rise to  $^1\text{H}$  NMR spectra of hyperpolarized compounds with a high sensitivity. Harris et al. [55] demonstrates that indirectly detected  $^1\text{H}$  NMR spectroscopy can also be exploited as time-resolved hyperpolarized spectroscopy by merging with spatially encoded methods. This method can successfully deliver a series of hyperpolarized  $^1\text{H}$  NMR spectra over a minutes-long timescale. The principles and opportunities presented by this approach were demonstrated by following the *in vitro* phosphorylation of choline by choline kinase, and by tracking acetylcholine's hydrolysis by acetylcholine esterase, an important enzyme partaking in synaptic transmission and neuronal degradation (Fig. 9) [55].



**Fig. 9** NMR spectral changes revealed by a 5 mM solution of hyperpolarized choline upon undergoing phosphorylation by 0.5 units of choline kinase. (a) Emergence of the new phosphocholine resonance shown by directly detected single-pulse  $^{15}\text{N}$  NMR spectroscopy experiments. (b) Emergence of the  $^1\text{H}$  NMR resonance associated with the methylenes in the C2-position of phosphocholine. (c) Comparison between the expected enzyme kinetics of kinase with results afforded by the  $^{15}\text{N}$ - (&) and  $^1\text{H}$ -detected (^) hyperpolarized experiments, as derived from the relative peak ratios of the NMR peaks in (a) and (b). The straight line illustrates the best fit of the combined set of data points, and corresponds to an initial phosphorylation rate of  $0.3 \text{ mM min}^{-1}$  under these conditions. Reproduced with permission from [55]

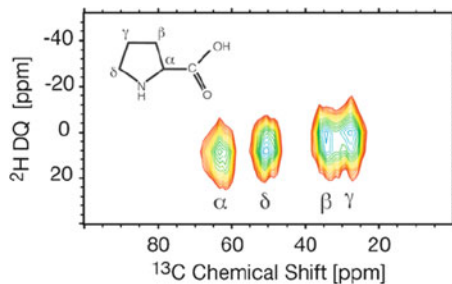
## 4.2 Applications to Biomolecules

ssNMR is a powerful technique for the investigation of membrane-associated peptides and proteins as well as their interactions with lipids, and a variety of conceptually different approaches have been developed for their study. The technique is unique in allowing for the high-resolution investigation of liquid disordered lipid bilayers representing well the characteristics of natural membranes. While magic angle solid-state (MAS) NMR spectroscopy follows approaches that are related to those developed for solution NMR spectroscopy, the use of static uniaxially oriented samples results in angular constraints which also provide information for the detailed analysis of polypeptide structures. Solid-state NMR spectroscopy has already in the past provided valuable structural information for biomolecules and major advancement of the technique can be expected when it is possible to overcome the present lack of sensitivity where recording a two-dimensional spectrum can take many days, even for samples consisting of several milligrams of labeled polypeptide. Although high magnetic field spectrometers have already ameliorated the situation, the novel developments of DNP has promised to boost tremendously the sensitivity of solid-state NMR spectroscopy by about two orders of magnitude [56]. Notably, solid-state DNP-NMR has already been applied on oriented bilayer samples [57], as well as on membrane proteins [56, 58]. With such developments the application of three- and four-dimensional NMR experiments should become possible and much extend the ease of structural investigations on oriented membranes as well as improving the conformational details that can be obtained.

Cyanobacteria are widely used as model organisms of oxygenic photosynthesis due to being the simplest photosynthetic organisms containing both photosystem I and II. Photochemically induced dynamic nuclear polarization (photo-CIDNP)  $^{13}\text{C}$  MAS NMR is a powerful tool in understanding the photosynthesis machinery down to atomic level. Combined with selective isotope enrichment this technique has now opened the door to study primary charge separation in whole living cells. Janssen et al. [59] recently presented the first photo-CIDNP observed in whole cells of the cyanobacterium *Synechocystis*.

In the applications of DNP to MAS spectra of biological systems, including studies of lysozyme [20] and bacteriorhodopsin [60], the enhancements have been smaller, with  $\varepsilon = 40\text{--}50$ , which has limited the application of DNP-NMR to biological samples so far. An exception is the amyloidogenic peptide GNNQQNY7–13 which forms nanocrystals for which the proton  $T_1$  time is long, yielding an  $\varepsilon$  of  $\sim 100$  [25]. Almost a decade ago in studies of model systems it was observed that deuteration of the solvent resulted in significant increases in  $\varepsilon$  and subsequently many DNP experiments have employed  $^2\text{H}$ -labeled solvents such as  $[\text{D}_6]\text{DMSO}$  or  $[\text{D}_8]\text{glycerol}/\text{D}_2\text{O}/\text{H}_2\text{O}$  in an approximately 6:3:1 ratio [21, 22, 61]. Akbey et al. [62] recently showed that perdeuteration of a protein has remarkable effects on the observed DNP enhancements. Superior DNP enhancements are obtained for perdeuterated SH3 samples of up to 3.9 and 18.5 times for  $^{13}\text{C}$  CP-MAS, and  $^{13}\text{C}$

**Fig. 10** Two-dimensional DNP-enhanced  $^2\text{H}$ -DQ- $^{13}\text{C}$  correlation spectrum of U- $[\text{}^2\text{H}_7, \text{}^{13}\text{C}_5]$ -proline recorded at 90 K. Reproduced with permission from [63]



MAS experiments, respectively, compared to the same type of experiments in fully protonated SH3. The optimum perdeuteration is found to be approximately 50% which results in the maximum enhancement of  $\epsilon \sim 148$  in a  $^{13}\text{C}$  MAS NMR spectrum obtained using a  $\text{ZrO}_2$  rotor. Maly et al. [63] also showed that perdeuteration of biological macromolecules for magic angle spinning solid-state NMR spectroscopy can yield high-resolution  $^2\text{H}$ - $^{13}\text{C}$  correlation spectra. They demonstrated that the combination of sample deuteration and DNP yields resolved  $^2\text{H}$ - $^{13}\text{C}$  correlation spectra with a signal enhancement  $\epsilon \geq 700$  compared to a spectrum recorded with microwaves off and otherwise identical conditions. The DNP process is studied using several polarizing agents and the technique is applied to obtain  $^2\text{H}$ - $^{13}\text{C}$  correlation spectra of U- $[\text{}^2\text{H}, \text{}^{13}\text{C}]$  proline (Fig. 10). Therefore the use of perdeuterated proteins in DNP-MAS NMR will open new possibilities in the application of these techniques to difficult biological problems.

### 4.3 Magnetic Resonance Imaging

MRI is one of the most powerful *in vivo* diagnostic tools due to its superb spatial resolution and the amount of information it allows one to obtain, including anatomical and function information as well as flux, perfusion, and diffusion studies. However, compared with other imaging techniques, MRI's major drawback is its inherent low sensitivity, which has precluded its use as a routine tool in the clinic. Recently, great advances have been achieved on the application of MRI with the enormous sensitivity enhancement afforded by DNP ( $>10,000$ -fold) [64]. The strong signal enhancement attainable by DNP has allowed the detection of heteronuclei in MRI, yielding high quality images with very high signal to noise ratios in a few seconds. The applications of hyperpolarized probes to MRI range from vascular imaging, interventional imaging, and perfusion studies to the emerging and challenging field of molecular/metabolic imaging. In fact, the high signal intensities achievable by using hyperpolarized molecules make it possible to detect and image the metabolic products formed upon the administration of the hyperpolarized agent *in vivo* and allow early diagnosis and assessment of diseases in personalized medicine in a non-invasive manner [64].

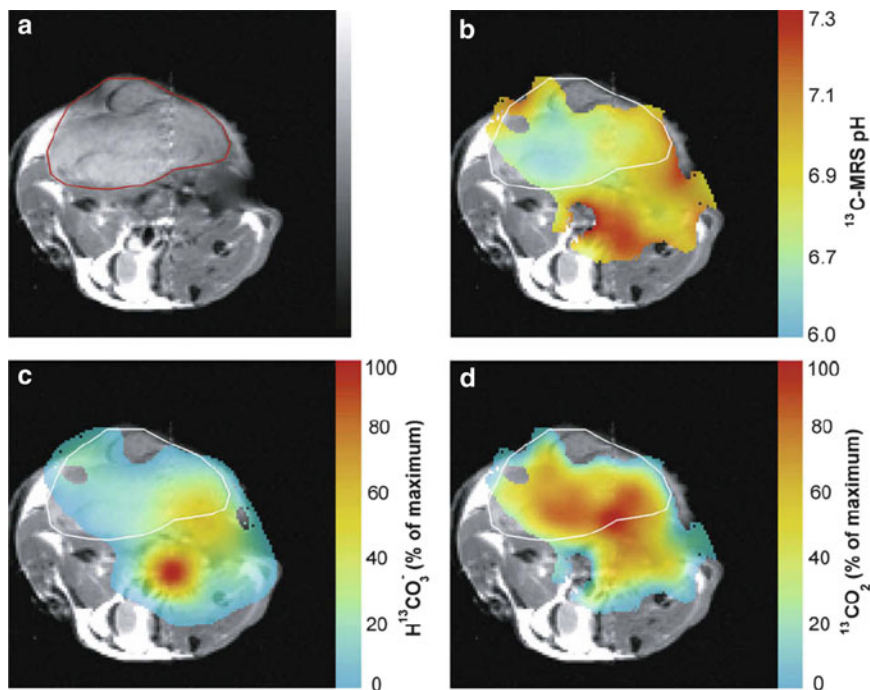
The main limitations of current  $^{13}\text{C}$  based DNP magnetic resonance spectroscopy imaging (MRSI) are, first, the very short lifetime of the hyperpolarization. For example, the  $^{13}\text{C}$ -labeled molecules have a typical life time less than 40 s *in vivo*. And, second, high concentration of the hyperpolarized molecule is required because, following dissolution and intravenous injection, the  $^{13}\text{C}$ -labeled molecule is diluted by about two orders of magnitude before it reaches the tissue of interest [64]. However, there are ways to tackle these limitations. For example, the lifetime can be extended by placing the  $^{13}\text{C}$  label in a position in the molecule where there is no or only very weak proton coupling [65]. Accordingly, highly soluble small metabolites, such as [1- $^{13}\text{C}$ ] pyruvate, [5- $^{13}\text{C}$ ] glutamine, and [1,4- $^{13}\text{C}_2$ ] fumarate, have been designed and they have been effectively applied to obtain greatly improved MRSI results. Furthermore, these specifically  $^{13}\text{C}$  labeled compounds have the advantage of significantly different  $^{13}\text{C}$  chemical shifts from the bulk carbon signal and thus can be readily distinguished in the low resolution spectra measured *in vivo*.

A number of papers dealing with metabolic imaging by  $^{13}\text{C}$ -hyperpolarized pyruvate have appeared in the literature since 2006 [66–76]. Pyruvate is a key molecule in major metabolic and catabolic pathways in mammalian cells, as it is converted to alanine, lactate or carbonate to varying extent depending on the status of the cells. Pyruvic acid naturally forms a glass and has been polarized by up to 40%. Golman et al. [66] reported one of the first examples of tumor imaging *in vivo* by this technique, which showed the maps of distribution of pyruvate, alanine, and lactate after a pyruvate injection into a rat with a P22 tumor. The tumor is clearly revealed by the highest NMR  $^{13}\text{C}$  signal from lactate produced within 30 s [66]. Recently, hyperpolarized pyruvate has been utilized as a therapy response marker [72]. Within 24 h of cytotoxic drug treatment the pyruvate–lactate exchange was substantially reduced in lymphoma tumors [72]. It was shown that monitoring changes in hyperpolarized pyruvate–lactate exchange compares favorably with FEG-PET for detecting treatment response [64]. Thus measurement of pyruvate–lactate exchange may be an alternative to FEG-PET for imaging tumor treatment response in the clinic, in particular for tumors that are not FDG-avid, such as the prostate tumor.

Many pathological states are associated with changes in tissue acid–base balance, including inflammation and ischaemia [77–79]. For instance, most tumors have an acidic extracellular pH compared to normal tissue and this can be correlated with prognosis and response to treatment [80–82]. Despite the importance of pH and its relationship to disease, there is currently no clinical tool available to image the spatial distribution of pH in humans. In principle, tissue pH could be determined from  $^{13}\text{C}$ -MRSI measurements of endogenous  $\text{H}^{13}\text{CO}_3^-$  and  $^{13}\text{CO}_2$  using the Henderson-Hasselbalch equation, if there was sufficient signal-to-noise:

$$\text{pH} = \text{pK}_a + \log\left(\frac{[\text{H}^{13}\text{CO}_3^-]}{[^{13}\text{CO}_2]}\right) \quad (15)$$

One study has shown that the extracellular pH in tumors can be imaged *in vivo* from the ratio of the signal intensities of  $\text{H}^{13}\text{CO}_3^-$  and  $^{13}\text{CO}_2$  following intravenous



**Fig. 11** A pH map of a transverse slice through a mouse implanted with a subcutaneous lymphoma tumor. (a) The *black and white* 1H MR image, obtained at 9.4 T, is shown with the tumor outlined in *red*. (b) The *false-color* <sup>13</sup>C MR pH image is shown separately superimposed over this; the tumor is acidic compared to the surrounding normal tissue. The spatial distribution of H<sup>13</sup>CO<sub>3</sub><sup>-</sup> (c) and <sup>13</sup>CO<sub>2</sub> (d) are also shown in the lower two images. Reproduced with permission from [64]

injection of hyperpolarized H<sup>13</sup>CO<sub>3</sub><sup>-</sup> [83]. This demonstrated that the extracellular pH in a lymphoma tumor was more acidic than the surrounding tissue (Fig. 11).

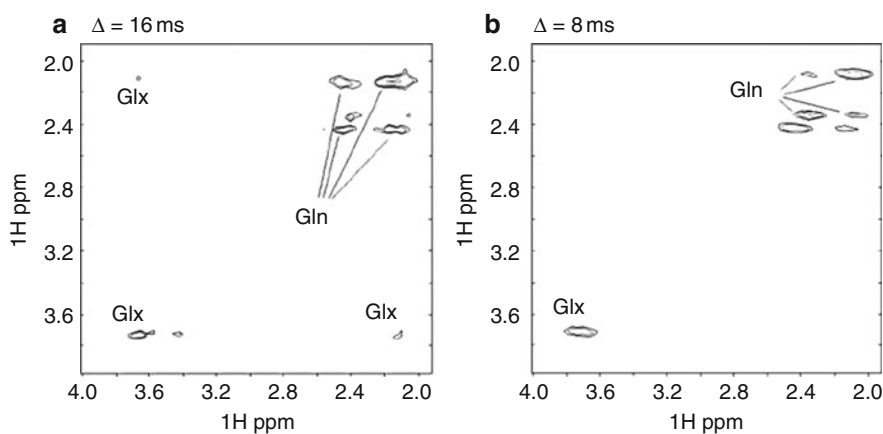
#### 4.4 *Multidimensional Time-Domain Experiment for DNP*

DNP has become an attractive technique to boost the sensitivity of NMR experiments. In the case of *ex situ* polarizations, two-dimensional (2D) spectra are limited by the short lifetime of the polarization after dissolution and sample transfer to a high field NMR magnet. This limitation can be overcome by various approaches. Ludwig et al. [84] showed how the use of <sup>13</sup>C-labeled acetyl tags can help to obtain 2D-HMQC spectra for many small molecules, owing to a nuclear Overhauser enhancement between <sup>13</sup>C spins originating from the long-lived carbonyl carbon, which extends the lifetimes of other <sup>13</sup>C spins with shorter longitudinal relaxation times. For two-dimensional spectroscopy, a pulse sequence using consecutive small flip angle excitations to read out transverse magnetization

while preserving  $z$ -magnetization was employed. This procedure enables one to obtain 2D-HMQC spectra over a time frame of the lifetime of the polarization as exemplified for natural abundance nicotinamide. They have also combined this pulse sequence with non-uniform sampling to make the best use of the limited lifetime of the polarization [84].

Short acquisition time and single scan capability of gradient-assisted ultrafast multidimensional spectroscopy make it possible to record 2D spectra of highly polarized spin systems in the liquid state using DNP in conjunction with fast dissolution. Panek et al. [85] presented a slice selective experiment, suitable for back-to-back acquisition of two independent single-scan 2D experiments from different sample volumes. This scheme maximizes the amount of information obtainable from a sample that is prepolarized with a non-repeatable DNP technique. It is particularly suitable for samples with the short longitudinal relaxation times common to proton NMR spectroscopy. This technique is demonstrated by applying two filtered proton 2D COSY experiments on a DNP-polarized mixture of glutamine and glutamate to amplify selectively the correlation pattern of the protons connected to the beta and gamma carbons of either one of the two amino acids (Fig. 12). Particular emphasis was placed on the reproducibility of the experiments, especially the polarization enhancement. It has been demonstrated that the slice selective approach is a useful extension of the single-scan 2D spectroscopy tool set, enabling dynamic studies of DNP-enhanced samples or, alternatively, the acquisition of different experiments on a single one-time polarized sample.

In addition to the DNP methods described above, laser-polarized xenon can be used for the development of NMR and MRI techniques, and it is possible to enhance the nuclear spin polarization of gaseous xenon by four to five orders of magnitude. More details of this technique can be found in [86, 87].



**Fig. 12** DNP enhanced proton COSY spectra of glutamine and glutamate mixture acquired for two different evolution delays ( $\Delta$ ) in two separate slices. Reproduced with permission from [85]

## 5 Future Perspectives

Recent developments of DNP have shown great potential in solid-state and liquid-state NMR spectroscopy and in MRI. As promising as DNP is, the future development of hyperpolarization transition may require great improvement in the following aspects:

1. Development of new instrumentation in DNP. The new design of instruments for DNP should be able to couple conveniently with the current high field NMR spectrometers.
2. New DNP methods should be able to overcome the limitation of the very short lifetime of the hyperpolarization, and should be integrated into the currently available solid-state and liquid-state NMR experiments.
3. Development of new chemical agents to transfer effectively magnetization of electrons to nuclei for the problem orientated applications in NMR spectroscopy and MRI.

We expect a new wave of DNP research in the next few years, which may result in unprecedented breakthroughs of these techniques, consequently allowing us to detect diseases early in a non-invasive way and to study biomolecules in much lower concentrations.

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