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Toru Tanzawa

# On-chip High-Voltage Generator Design

Design Methodology for Charge Pumps

*Second Edition*

 Springer

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*To the memory of my family*



# Preface

Accordingly, as silicon technology has been advanced, more and more functionalities have been integrated into LSIs. On-chip multiple voltage generation is becoming one of big challenges on circuit and system design. Linear or series regulator is used to convert the external supply voltage into lower and more stable internal voltages. As the number of gates operating simultaneously and the operation frequency increase, AC load current of the regulators also increases. Low-voltage operation and rapid load regulation are becoming design challenges for the voltage down converters. Another type of voltage generator is high-voltage generator or voltage multiplier whose output is higher than the input supply voltage. The voltage multipliers are categorized into two, switching convertor and switched capacitor, with respect to the components used. The former uses an inductor, switch or diode, and AC voltage source whereas the latter uses a capacitor instead of the inductor. Even though the switching convertor has been widely used with discrete chip inductor(s) and capacitor(s), there is little report on implementation of an inductor into ICs because of too low quality factor for large inductance fabricated in current silicon technology. This book aims at discussing thorough high-voltage generator design with the switched-capacitor multiplier technique.

The First Edition has focused on integrated DC-DC voltage multipliers where the DC supply voltage is nominally greater than 1.5 V. In the Second Edition, the design of AC-DC charge pump has been added for those who are interested in the design of RFID and energy harvesting where AC input needs to be transformed into DC. In addition to the new topic, the relationship between output voltage and output current of a charge pump has been updated to be available in wide frequency range with three different conditions of switching devices; (1) diodes, (2) MOSFETs in saturation region, and (3) MOSFETs in triode region. Thus, each I-V equation includes slow to fast switching limit, which would be beneficial to those who worked on research and development of extremely low voltage LSI design. Furthermore, major revisions on DC-DC voltage multipliers have been also made to expand the circuit theories for understanding of power efficiency and for comprehensive design of the system including DC energy transducer and charge pump.

The switched-capacitor multiplier techniques originated with H. Greinacher using a voltage doubler structure for measuring the intensity of ionizing radiation in 1919, E. O. Marx using serial-parallel cells for an impulse voltage generator in 1924, and Cockcroft–Walton using serial capacitor ladders for their experiments on nuclear fission and fusion in 1932. Dickson qualitatively pointed out that the Cockcroft–Walton multiplier had too high sensitivity on parasitic capacitance to realize on-chip multipliers and then theoretically and experimentally showed that the parallel capacitor ladders realized on-chip high-voltage generation for programming Metal–Nitride–Oxide–Semiconductor (MNOS) nonvolatile memory in 1976.

After Dickson’s demonstration, on-chip high-voltage generator has been implemented on Flash memories and LCD drivers and the other semiconductor devices. Accordingly, as the supply voltages of these devices become lower, it gets harder to realize small circuit area, high accuracy, fast ramp rate, and low power at a low supply voltage. This book provides various design techniques for the switched-capacitor on-chip high-voltage generator including charge pump circuits, pump regulators, level shifters, voltage references, and oscillators. The charge pump inputs the supply voltage and a clock, which is generated by the oscillator, and outputs a voltage higher than the supply voltage or a negative voltage. The pump regulator enables the charge pump when the absolute value of the output voltage of the charge pump is lower than the target voltage on the basis of the reference voltage or disables it otherwise. The generated high or negative voltage is transferred to a load through high- or low-level shifters. Chapter 1 surveys system configuration of the on-chip high-voltage generator.

Chapter 2 reviews various topologies of voltage multipliers. Since the charge pump was invented in 1919, various types have been proposed. After several typical types of charge pumps are reviewed, they are compared in terms of the circuit area and the power efficiency. The type that Dickson proposed is found to be the best one as an on-chip generator.

Chapter 3 discusses DC-DC Dickson charge pump. Design equations and equivalent circuit models are derived for the charge pump. Three types of charge transfer gate are considered; switching diode and switching MOSFET in saturation and triode region. Using the model, optimizations are discussed to minimize the circuit area under the condition that the output current or the ramp time is given and to minimize the power dissipation under the condition that the output current is given theoretically. Guideline for comprehensive optimum design is summarized.

Chapter 4 describes AC-DC Dickson charge pump. Two types of AC input are considered; continuous wave with a single frequency and multi-sine wave with multiple frequencies. An analytical, closed-form AC-DC charge pump voltage multiplier model is described to show the dependency of output current and input power on circuit and device parameters for continuous wave AC-DC charge pump. Then, it is expanded for multi-sine wave AC-DC charge pump. Analysis enables circuit designers to estimate circuit parameters, such as the number of stages and capacitance per stages, and device parameters such as saturation current (in the case of diodes) or transconductance (in the case of MOSFETs). In addition, design

optimizations and the impact of AC power source impedance on output power are investigated.

Chapter 5 overviews actual charge pumps composed of capacitors and transfer transistors. Realistic design needs to take parasitic components such as parasitic capacitance at each of both terminals and threshold voltages of the transfer transistors into account. In order to decrease the pump area and to increase the current efficiency, some techniques such as threshold voltage canceling, stage reconfiguration, and faster clocking are presented. Since the supply current has a frequency component as high as the operating clock, noise reduction technique is another concern for pump design. In addition to design technique for individual pump, system level consideration is also important, since there are usually more than one charge pump in a chip. Area reduction can be also done for multiple charge pump system where all the pumps do not work at the same time.

Chapter 6 is devoted to individual circuit block to realize on-chip high-voltage generator. Section 6.1 presents pump regulator. The pump output voltages need to be varied to adjust them to the target voltages. This can be done with the voltage gain of the regulator or the reference voltage changed. The voltage divider which is a main component of the regulator has to have small voltage coefficient and fast transient response enough to make the controlled voltage linear to the trim and stable in time. A regulator for a negative voltage has a circuit configuration different from that for a positive voltage. State of the art is reviewed.

Section 6.2 surveys level shifters. The level shifter shifts the voltage for logic high or low of the input signal to a higher or lower voltage of the output signal. Four types of level shifters are discussed (1) high-level NMOS level shifter, (2) high-level CMOS level shifter, (3) high-voltage depletion NMOS + PMOS level shifter, and (4) low-level CMOS level shifter. The trade-offs between the first three high-voltage shifters are mentioned. The negative voltage can be switched with the low-level shifter. As the supply voltage lowers, operation margins of the level shifters decrease. As the supply voltage lowers, the switching speed becomes slower, eventually infinite, i.e., the level shifter does not work. Some design techniques to lower the minimum supply voltage at which the level shifters are functional are shown.

Section 6.3 deals with oscillators. Without an oscillator, the charge pump never works. In order to make the pump area small, process, voltage, and temperature variations in oscillator frequency need to be minimized. There is the maximum frequency at which the output current is maximized. If the oscillator is designed to have the maximum frequency under the fastest conditions such as fast process corner, high supply voltage, and low temperature, the pump output current is minimum under the slowest condition such as slow process, low supply voltage, and high temperature. It is important to design the oscillator with small variations for squeezing the pump area.

Section 6.4 provides voltage references. Variations in regulated high voltages increase by a factor of the voltage gain of the regulators from those in the reference voltages. Reduction in the variations of the voltage references is a key to make the

high generated voltages well controlled. Some innovated designs for low supply voltage operation are presented as well.

Chapter 7 provides high-voltage generator system design. Multiple pumps are distributed in a die, each of which has sufficiently wide power ground bus lines. Total area including the charge pump circuits and the power bus lines needs to be paid attention for overall area reduction. Design methodology in this regard is shown using an example. Another concern on multiple high voltage generator system design is system level simulation time. Even though the switching pump models are used for the verification, simulation run time is still slow especially for Flash memory where the minimum clock period is 20–50 ns whereas the maximum erase operation period is 1–2 ms. In order to drastically reduce the simulation time, another charge pump model together with a regulator model is presented which makes all the nodes in the regulation feedback loop analogue to eliminate the hard-switching operation. Design and verification flow of integrated high-voltage generator system is summarized.

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# Abbreviations

AC	Alternate current
bjt	Bipolar junction transistor
BL	Bit-line
C	Capacitance of a pump capacitor
$C_B$	Parasitic capacitance at the bottom plate of a pump capacitor
clk	Clock
$C_{OUT}$	Total capacitance of pump capacitors
$C_T$	Parasitic capacitance at the top plate of a pump capacitor
CW	Cockcroft–Walton pump or continuous wave
DC	Direct current
eff	Current efficiency
FET	Field effect transistor
FIB	A type of pump whose $V_{MAX}$ is associated with a Fibonacci number $Fib(N)$ where $N$ is the number of stages
FSL	Fast switching limit
$G_{MAX}$	Maximum voltage gain
$G_V$	Voltage gain
$I_B$	Base current
IC	Integrated circuit
$I_C$	Collector current
$I_{DD}$	Supply current
$I_{DS}$	Drain to source current
$I_{IN}$	Input current
$I_L$	Load current
$I_{LOAD}$	Load current
$I_{OUT}$	Output current
$I_{PP}$	Output current of a positive voltage pump at $V_{OUT}$ of $V_{PP}$
$I_{REG}$	Regulator current
ISM	Industry – Science – Medical
K(N)	4-Port K-matrix of N-stage pump

LCD	Liquid crystal device
LED	Light emitting device
LIN	A type of pump whose $V_{MAX}$ is linear to the number of stages
LSI	Large scale IC
MNOS	Metal nitride oxide semiconductor
MOS	Metal oxide semiconductor
MS	Multi-sine
$N$	Number of stages
$N_{MIN}$	Minimal number of stage
$N_{OPT}$	Optimum number of stages
opamp	Operational amplifier
$P_{IN}$	Input power
$P_{OUT}$	Output power
PV	Photovoltaic
PVT	Process, voltage, and temperature
$Q_{DD}$	Total input charge
qout	Output charge per period
RFID	Radio frequency identification
$R_{LOAD}$	Resistance of a load circuit
$R_{PMP}$	Output impedance of a pump
$R_{PWR}$	Parasitic resistance of power and ground lines
SC	Switched-capacitor
SP	Serial-parallel
SRC	Source
SSL	Slow switching limit
$T$	Clock period of a pump driver clock or temperature
TEG	Thermoelectric generator
$T_{OFF}$	The period when a switch is being turned off
$T_{ON}$	The period when a switch is being turned on
UHF	Ultrahigh frequency
UPS	Utility Power Satellite
$V_{BB}$	Negative output voltage of a charge pump
$V_{BE}$	Base to emitter voltage
$V_{BGR}$	Band-gap reference voltage
$V_{BL}$	Bit-line voltage
$V_{BS}$	Bulk to source voltage
$V_{BV\_CAP}$	Breakdown voltage of a capacitor
$V_{BV\_SW}$	Breakdown voltage of a switch
$V_{CAP}$	Capacitor voltage
$V_D$	Drain voltage
$V_{DD}$	Supply voltage
$V_{DD\_LOCAL}$	Supply voltage at a local interconnection node
$V_{DD\_MIN}$	Minimum operating supply voltage
$V_{DS}$	Drain to source voltage
$V_G$	Gate voltage or voltage gain given by $V_{DD} - V_T$

$V_{GS}$	Gate to source voltage
$V_{IN}$	Input voltage
$V_k$	$k$ -th nodal voltage
$V_{MAX}$	Maximum attainable voltage
$V_{MOD}$	Modulation voltage
$V_{MON}$	Monitored voltage
$V_{OD}$	Overdrive voltage
$V_{OS}$	Offset voltage
$V_{OUT}$	Output voltage
$V_{PP}$	Positive high output voltage of a charge pump
$V_{REF}$	Reference voltage
$V_S$	Source voltage
$V_{SS\_LOCAL}$	Ground voltage at a local interconnection node
$V_{SW}$	Switching voltage
$V_T$	Threshold voltage or thermal voltage $kT/q$
$V_{iD}$	Threshold voltage of a depletion NMOS transistor
$V_{iE}$	Threshold voltage of an enhancement NMOS transistor
$V_{iI}$	Threshold voltage of an intrinsic NMOS transistor
$V_{iP}$	Threshold voltage of a PMOS transistor
WL	Word-line
$\alpha$	Parameter representing a body effect of a MOS transistor
$\alpha_B$	Ratio of $C_B$ to $C$
$\alpha_T$	Ratio of $C_T$ to $C$
$\beta$	Multiplication factor of the collector current to the base current of a bipolar junction transistor
$\eta$	Power efficiency
$\gamma$	Conduction angle
$\Phi_i$	$i$ -th clock phase



# Chapter 1

## System Overview and Key Design Considerations

**Abstract** This chapter describes which categories of voltage converters are covered in this book. Various applications of on-chip high-voltage generators such as memory applications for MNOS, DRAM, NAND Flash, NOR Flash, and phase-change memory, and other electronic devices for motor drivers, white LED drivers, LCD drivers, and energy harvesters are overviewed. System configuration of the on-chip high-voltage generator and key design consideration for the building circuit blocks such as charge pumps, pump regulators, oscillators, level shifters, and voltage references are surveyed.

### 1.1 Applications of On-Chip High-Voltage Generator

Section 1.1 starts with describing which categories of voltage converters are covered in this book. It also overviews various applications of on-chip high-voltage generators such as memory applications for MNOS, DRAM, NAND Flash, NOR Flash, and phase-change memory, and other electronic devices for motor drivers, white LED drivers, LCD drivers, and energy harvesters.

Voltage converters are categorized into two: switching converter (Erickson and Maksimovic 2001) and switched capacitor (Cockcroft and Walton 1932) converter as classified in Table 1.1. Switching converter is composed of one or a few inductors, one or a few capacitors, and one or a few switching devices. Switched capacitor converter is composed of one-to-many capacitors and one-to-many switching devices. The differences are with or without inductor and single or many stages. From the viewpoint of amount of power, the switching converter can be used for applications to generate high power typically larger than 100 mW. On the other hand, switched capacitor converter is used for applications to generate lower power than 100 mW. Presently, degree of integration is all, except for inductors, for switching converter whereas all components for switched capacitor. This is mainly because inductance that integrated inductor can have is much smaller than the value required as well as the input current noise could be much more in switching converter with a single stage. From the viewpoint of voltage gain, that is, the ratio

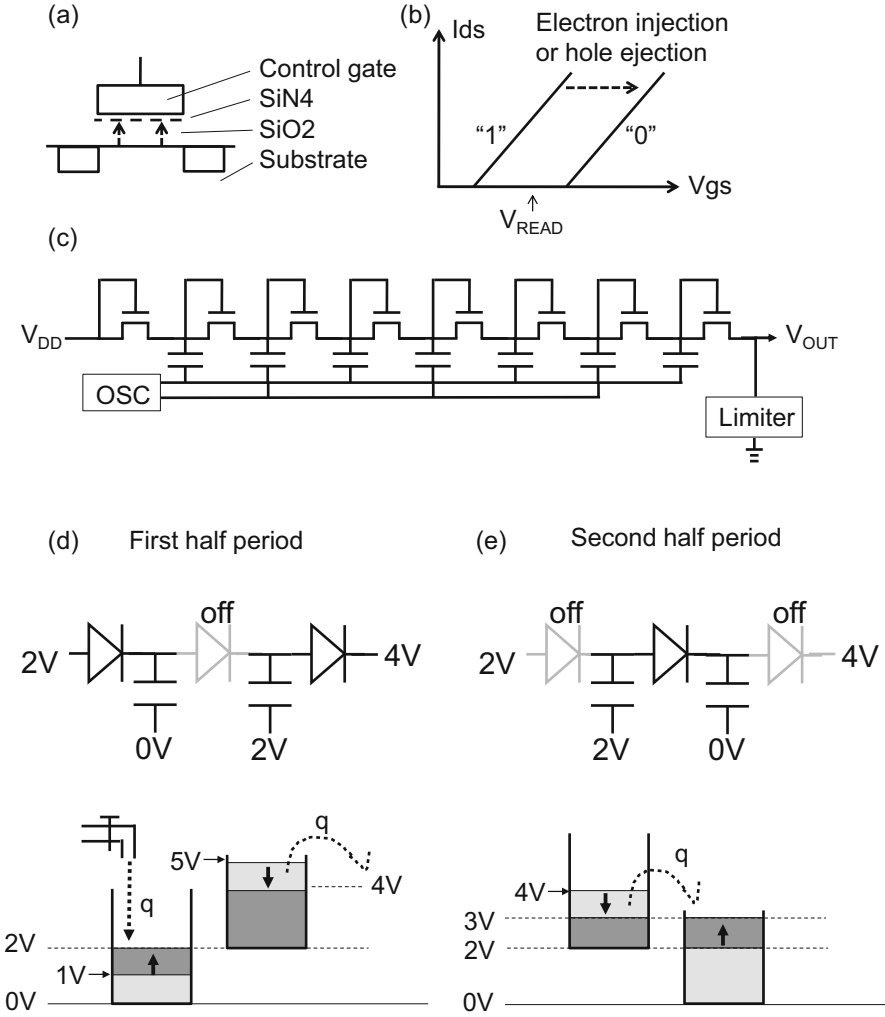
**Table 1.1** Classification of voltage convertors

	Switching converter	Switched capacitor
Components	Inductor	Capacitor
	Capacitor	Switching device
	Switching device	
Feature	High power and low loss	High voltage and low current or low voltage and high current
Integration	Except for inductor	Fully integrated
$G_v \equiv V_{out}/V_{in} > 1$	Boost	Charge pump/voltage multiplier
$1 > G_v > 0$	Buck	Switched capacitor voltage down converter
$G_v < 0$	Buck–boost	Charge pump/voltage multiplier

of the output voltage to the input voltage, there are three categories: greater than one, smaller than one and greater than zero, and smaller than zero. For the switching converter, these are, respectively, called boost converter, buck converter, and buck–boost converter. For the switched capacitor, the first and third are similarly called charge pump or voltage multiplier, and the second is called switched capacitor regulator or voltage down converter. Thus, this book covers these two categories with a voltage gain greater than one or lower than zero for fully integrated high-voltage generation among entire voltage converter system.

Following some figures show applications where on-chip voltage multipliers are used in ICs. A nonvolatile metal–nitride–oxide–semiconductor (MNOS) memory has a nitride film between the control gate and substrate where electrons or holes can trap as shown in Fig. 1.1a. Depending on the charges stored in the film,  $V_{GS}$ – $I_{DS}$  characteristics are varied as described in Fig. 1.1b. The data in memory cells are read with  $V_{READ}$  biased to the control gate. The data is identified as “0” when the memory cell does not flow a sufficient current or as “1” when one flows. To alternate the memory data, the memory needed high voltages of 30–40 V for programming and erasing the data. To significantly reduce the system cost and complexity, an on-chip voltage multiplier was strongly desired. In 1976, Dickson theoretically and experimentally for the first time studied an on-chip high-voltage generator including a charge pump, oscillator, clock drivers, and a limiter, as shown in Fig. 1.1c. The diode is made of a MOSFET whose gate and drain terminals are connected. Dickson used two-phase clock which allowed the clock frequency as fast as possible. Using a seven-stage pump, he successfully generated 40 V from the power supply voltage of 15 V. The capacitors were also implemented using the nitride dielectric available in the MNOS process. Thus, switches and capacitors were integrated in ICs. Design parameters of 2 pF per stage, 7 stages, and 1 MHz realized an output impedance of 3.2 M $\Omega$  and a current supply of an order of 1  $\mu$ A.

Figure 1.1d, e illustrates the image of how the charge pump works. For simplicity, a two-stage pump is shown. As the saying goes, a bucket, water, and the height of the surface of the water are, respectively, used as a capacitor, charge, and the capacitor voltage.  $V_{DD}$  is 2 V and  $V_{OUT}$  is 4 V. In the first half period (Fig. 1.1d), the current to the first capacitor stops when the voltage of the first capacitor reaches 2 V. The current stops flowing from the second capacitor to the output terminal



**Fig. 1.1** MNOS cell structure (a),  $I$ - $V$  curve of memory cells with data 1 and 0 (b), first Si verified on-chip Dickson pump (c), the states of the first (d) and second (e) half periods (Dickson 1976)

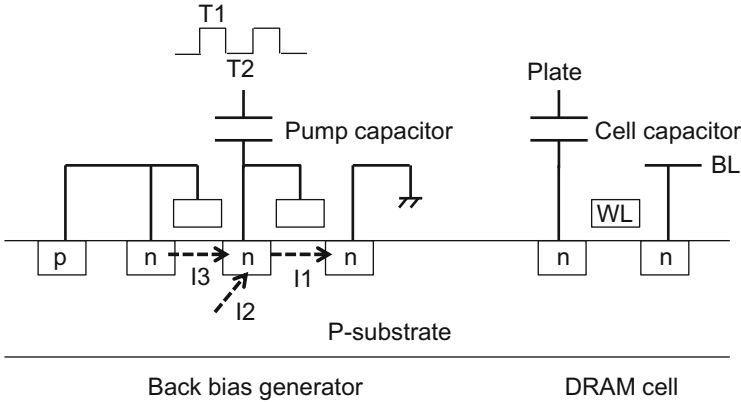
when the capacitor voltage reaches 4 V. At the beginning of the second half of the period (Fig. 1.1e), the capacitor voltage of the first capacitor increases to 4 V, whereas that of the second capacitor decreases to 2 V. This voltage difference between the two capacitors forces to flow the current through the second diode. When the threshold voltage of the diode is ignored, the charge transfer stops when the capacitor voltages are equalized. When the two capacitors are same size, an equilibrium state occurs when the capacitor voltages become 3 V. At the end of the second half of period, the capacitor voltages between the two terminals of the first and second capacitors are, respectively, 1 V and 3 V. At the beginning of the first half of period again, the surface potential at the top terminal becomes 1 V and 5 V,

respectively. The water tap again flows until the surface potential increases to 2 V. Charge transfer from the second capacitor to the output terminal stops when the potential of the second capacitor reaches 4 V. Thus, alternate operations back and forth between the first and second half of periods result in charge transfer from the water tap to the output terminal with the same amount of charge  $q$ .

A dynamic random access memory (DRAM) cell is composed of one transistor and one capacitor as shown at the right-hand side of Fig. 1.2. The data “0” or “1” is stored as amount of charges in the cell capacitor. To read the data, a word-line (WL) is forced high. The amount of charges stored in the cell capacitor modulates the bit-line (BL) voltage, which is sensed and amplified by a sensing circuit. Thus, voltages at WLs and BLs were toggled between 0 V and 5 V during operations when the supply voltage was 5 V. Such a huge voltage swing could make PN junctions of NMOS transistors into forward bias regime locally due to capacitive coupling where it is far from body contacts if the p-type substrate is grounded. If this happens, stored charges could be flown into the substrate, resulting in degradation in data reliability. To avoid it, another negative voltage of  $-5$  V was needed in addition to the power supply voltage of  $+5$  V. The negative voltage was supplied to the substrate to have sufficient operation margin with such a potential localized forward biasing of junctions eliminated.

The  $-5$  V power supply was eliminated by implementing a back bias generator allowing to reduce the system cost and complexity having the negative voltage supply, as shown at the left-hand side of Fig. 1.2. Lee and Breivogel et al. designed the generator to output  $-4.2$  V back bias at zero substrate current and  $-3.5$  V bias at  $5 \mu\text{A}$  substrate current. The output current was needed to be higher than the impact ionization current due to the memory operation. The power dissipation was 1.5 mW. The power efficiency is estimated to be an order of 1 %. Additional advantages are known to be improving the power and speed with smaller junction capacitance at a back bias and steeping the subthreshold slope of transistors. The back bias generator has one stage. The input terminal is connected with the substrate. During  $T_1$  where the clock is high, the capacitor node is made at about  $V_T$  of the switching transistor with the current  $I_1$ . During  $T_2$  where the clock is low, the capacitor node is initially pulled down to about  $V_T - V_{DD}$ . The current  $I_2$  or  $I_3$  flows until the junction or the transistor turns off. Under zero substrate current, the potential of the substrate is made at the lower one of  $2V_T - V_{DD}$  and  $V_T + V_{BE} - V_{DD}$ .

Another application of a charge pump is a motor driver IC, as shown in Fig. 1.3. Because it needs to switch a supply voltage up to 30 V with a peak current of 30 A, a power MOSFET is used. To sufficiently reduce the power dissipation, a channel resistance as low as 40 m $\Omega$  is required. A charge pump of the power IC generates an overdrive voltage for the power MOSFET. The supply voltage for the power IC is ranged in 6–30 V, whereas overdrive voltage is targeted at a voltage higher than 10 V, i.e.,  $V_{PP} > V_{DD} + 10$  V. The clock amplitude is regulated using a Zener diode. The switching diodes are realized by parasitic devices of isolated P-well and N-diffusion, as shown in Fig. 1.3b. The breakdown voltage of the diode is as high as 17 V. The worst-case reverse bias is considered as  $2V_{CLK}$  at the beginning of the

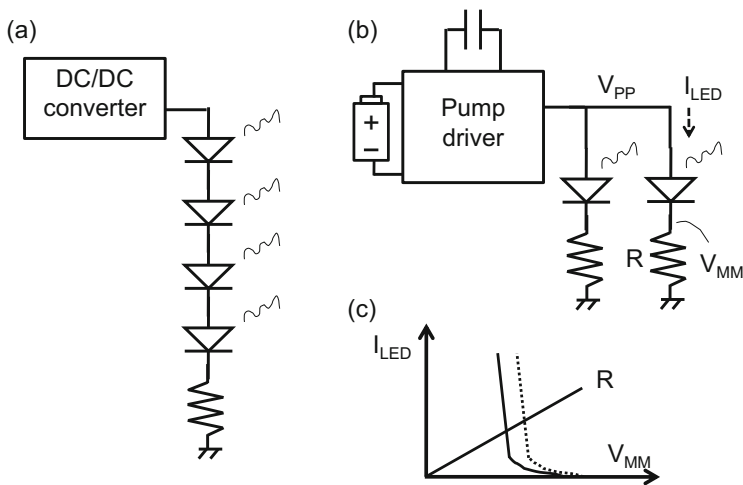
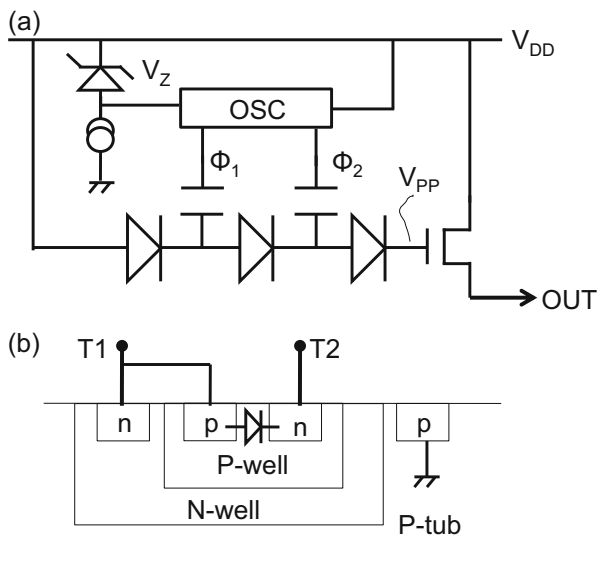


**Fig. 1.2** Back bias generator for DRAM (Lee et al. 1979)

pump operation, where  $V_{CLK}$  is the voltage amplitude of the driving clocks. Thus, the Zener diode with a breakdown voltage of 8 V is used to meet the requirement for  $2V_{CLK} < 17$  V. Considering a sufficient operation margin under an extreme operation temperature range of  $-40$  to  $125$  °C, three-stage structure is used.

Figure 1.4a, b shows two typical configurations of drivers for white light-emitting devices (LEDs). Figure 1.4c describes  $I$ - $V$  characteristics of the structures in Fig. 1.4b, which have similar  $I$ - $V$  curves as forward  $I$ - $V$  curves of diodes. The current increases exponentially as the voltage across the LED increases. Thus, the operating point in the  $I$ - $V$  plane could vary largely if the LED is controlled based on the voltage applied. To make the illumination or the power more stable against variations in the  $I$ - $V$  characteristics per LED, the LED is controlled on a current basis. Simple addition of a resistor to an LED aims at stabilizing the operating point. Red, yellow, or green LED needs about 20 mA at 2 V, whereas white LED does at 3.2–4 V. When DC/DC converter generates 12 V, 5 red or 3 white LEDs can be connected in series in a path as shown in Fig. 1.4a. If 5 paths are needed to have 15 white LEDs in total, the converter with capability to output a current of 100 mA has to be used. For a miniature single white LED, a charge pump IC with a single Li-ion battery with an output voltage of 2.7–3.6 V can be a solution. Whether external capacitors are added or not depends on the total driver size and cost. When adding one discrete capacitor to reduce the cost of the IC with no large pump capacitor is acceptable in terms of its form factor, one could put more numbers of white LED connected in parallel in the system, as shown in Fig. 1.4b. The LED driver IC only includes components of switches and oscillator except for the capacitor. The number of white LED connected in parallel is up to the output current of the charge pump IC. In case that the driver IC outputs 100 mA, for example, one can connect 5 LEDs in parallel. If the system requires only one or a few white LEDs, all the components including the pump capacitor can be integrated.

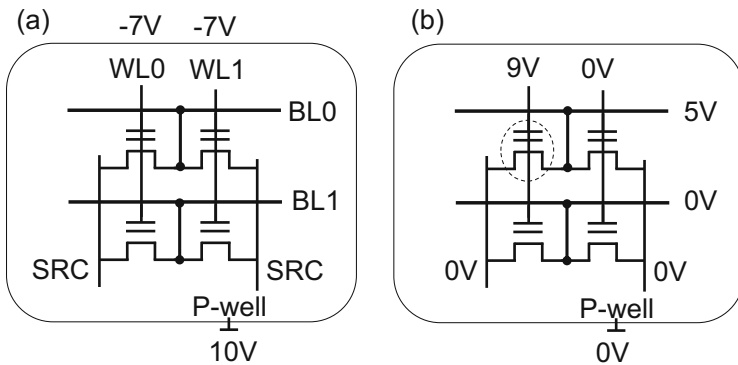
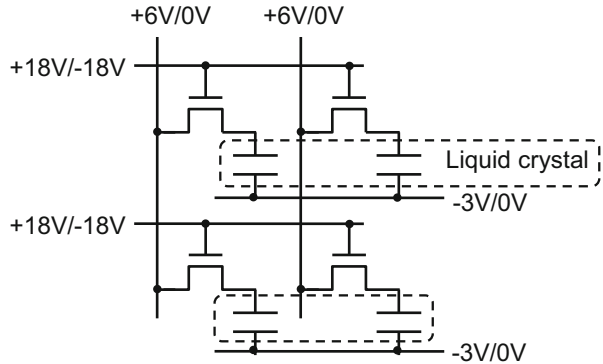
**Fig. 1.3** Pump and load MOS (a) and diode structure (b) of a motor driver IC (Storti et al. 1988)



**Fig. 1.4** White LED driver with (a) DC/DC converter (Chiu and Cheng 2007) and with (b) a charge pump (Wu and Chen 2009), and the operation condition of an LED (c)

A liquid crystal device requires two polarities of two positive voltages and two negative voltages to apply sufficiently high positive and negative voltages to each liquid crystal element aiming at improving the lifetime as shown in Fig. 1.5. Requirement for gate oxide of the transistors is sustaining a voltage of 18 V to fully turn on the pass transistors, which is half in case without generating voltages with two polarities. Otherwise, it would need a high voltage such as 36 V. A single-

**Fig. 1.5** Block diagram of a liquid crystal device (Wu and Chen 2008)

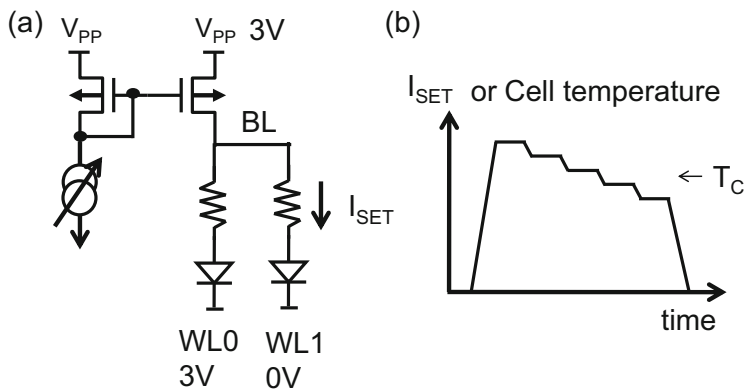


**Fig. 1.6** Channel erase NOR flash memory under an erase bias condition (a) and under a program bias condition (b) (Atsumi et al. 2000)

driver IC generates these four different voltages with a supply current of an order of 10–100  $\mu\text{A}$  because of no direct current to ground.

Another application using dual polarity is a NOR flash memory for erasing the data in a block, as illustrated in Fig. 1.6a. Flash cells are arranged horizontally and vertically, each of which is connected with a common source line (SRC), a bit-line (BL), and a word-line (WL). All cells in a block are placed in a common P-well. A bulk to gate voltage of 17 V needs to generate Fowler–Nordheim tunneling current flowing from the floating gate to the P-well. To allow the switching transistors for SRC and WLs to be scaled for reducing the transistor size, a high erase voltage of 17 V is divided to about half for a positive voltage of 10 V and a negative voltage of  $-7$  V.

Figure 1.6b shows a program bias condition for the NOR flash memory. The cell enclosed by a broken line is under programming with WL and BL supplied by 9 V and 5 V, respectively. Because the scaled flash cell has a relatively low snapback voltage, the bit-line voltage ( $V_{BL}$ ) has to be well controlled. The lower limit is



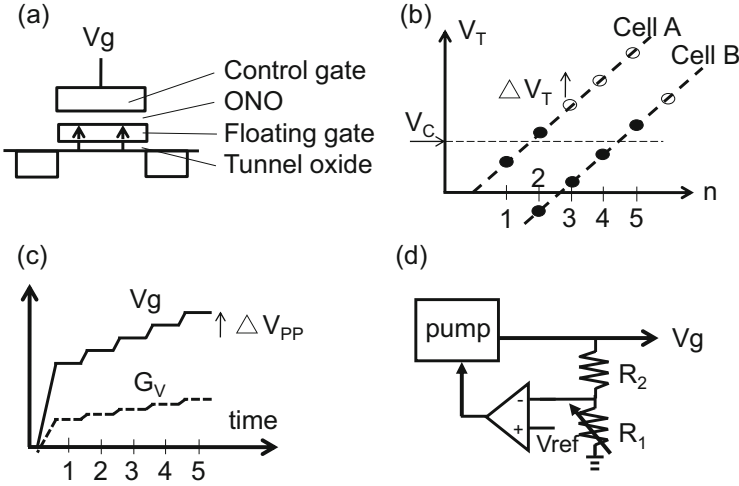
**Fig. 1.7** Set voltage and current generator for phase-change memory (Lee et al. 2008)

determined by the programming speed with hot carrier injection. With too low  $V_{BL}$ , the flash cell could not have sufficient hot electrons to inject to the floating gate. The upper limit is determined by the snapback voltage. When  $V_{BL}$  is directly generated by a pump, a voltage ripple may be so large that the Flash cell can enter the snapback regime. The clamping NMOSFET can control  $V_{BL}$  with much smaller ripple voltage because the load current is determined mainly by the gate voltage as far as the load FET operates in saturation region, resulting in much better stability in programming characteristics.

Figure 1.7a shows phase-change memory elements described as the symbols of resistor, switching diodes, and a set current control circuit. To change into phase crystalline, the memory material needs to be heated up to a critical temperature ( $T_C$ ) and to spend a required time interval at  $T_C$ . Because the memory array has quite large parasitic resistance in bit-lines (BLs) and word-lines (WLs), the input power required to individual memory element should have address dependencies. To program multiple memory cells with a few pulses for fast program operation, the set current as shown in Fig. 1.7b is supplied using a current control circuit with a variable current source. Thus, the boosted voltage  $V_{PP}$  is supplied to the memory elements with various current levels in a single set pulse.

Figure 1.8a illustrates a memory cell structure of NAND Flash memory. Because the floating gate is surrounded by insulator films, charges in the floating gate stay when the voltage difference between the control gate and silicon substrate is low enough. When there are many electrons in the floating gate of a cell, it has the data “0.” When there are few electrons, it has the data “1.” To program the data “0,” the control gate is biased at a high voltage of 20 V while the substrate is grounded. Tunnel phenomenon under a high electric field is known as Fowler–Nordheim tunneling. When the control gate voltage ( $V_g$ ) as shown in Fig. 1.8c is applied, the threshold voltage of the memory cell transistor is shifted as shown in Fig. 1.8b. The incremental step program pulse can reduce entire program time with well-controlled  $V_T$  of programmed cells using the general relation of  $\Delta V_T = \Delta V_{PP}$ . Due to the variation in program characteristics, cell A is programmed with two pulses,



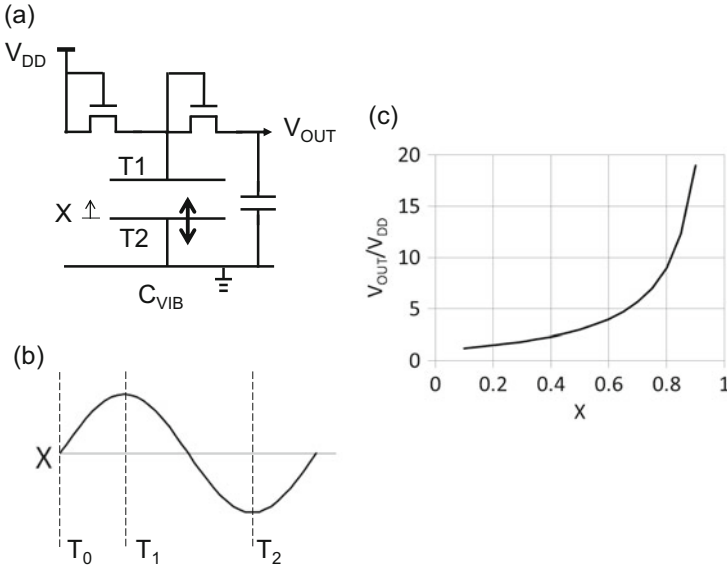


**Fig. 1.8** Incremental step program pulse generation for NAND flash memory (Masuoka et al. 1987; Suh et al. 1995)

whereas cell B is done with five pulses. Once  $V_T$  of a cell becomes greater than a critical value  $V_C$ , the program pulse is no longer applied. Figure 1.8d illustrates the program pulse generator.  $R_1$  of the resistor divider varies to vary the voltage gain  $G_V$  of  $V_G$  to  $V_{REF}$  as shown in Fig. 1.8c. Thus, the generator outputs the incremental step program pulse to control the programmed  $V_{TS}$ .

Energy harvesting has been paid much attention for low-power sensor and wireless applications. Figure 1.9a illustrates vibration energy harvester gathering vibration energy. The second terminal of a capacitor is connected with a mobile plate. The displacement  $X$  is a sine waveform as shown in Fig. 1.9b. Suppose  $X = 0$ ,  $C_{VIB} = C_0$ , and  $V_{CAP} = V_{DD}$  at time  $T_0$ , the charge stored in the pump capacitor is  $Q_0 = C_0 V_{DD}$ . When the displacement is  $+X$  at  $T_1$ ,  $C_{VIB}$  is increased to  $C_0/(1 - X)$ . If there is no transfer transistor connected with the power supply  $V_{DD}$ , the capacitor voltage would be  $Q_0 = C_0 V_{DD} = C_0/(1 - X) V_{CAP}(T_1)$ . Thus,  $V_{CAP}(T_1)$  would be  $(1 - X)V_{DD}$ . With the transfer transistor,  $V_{CAP}(T_1)$  is equalized to  $V_{DD}$ . Thus, the charge stored in the pump capacitor is  $Q_1 = C_0/(1 - X)V_{DD}$ . When the displacement is  $-X$  at  $T_2$ ,  $C_{VIB}$  is reduced to  $C_0/(1 + X)$ . If there is no transfer transistor connected with the output terminal, the capacitor voltage would be  $Q_1 = C_0/(1 - X)V_{DD} = C_0/(1 + X)V_{CAP}(T_2)$ . Thus,  $V_{CAP}(T_2)$  would be  $(1 + X)/(1 - X)V_{DD}$ . Therefore, the maximum attainable output voltage with no current load is  $(1 + X)/(1 - X)V_{DD}$ . Figure 1.9c shows the factor of  $(1 + X)/(1 - X)$  as a function of  $X$ .

Another energy harvester for RFID is illustrated in Fig. 1.10. Unlike that for vibrator energy as shown in Fig. 1.9, the harvester collecting the energy in a radio wave does not require any power supply voltage source. The input power from the antenna varies in a wide range. To protect capacitors and transistors from a high power input, a limiter is required. The capacitor at every even number stage is



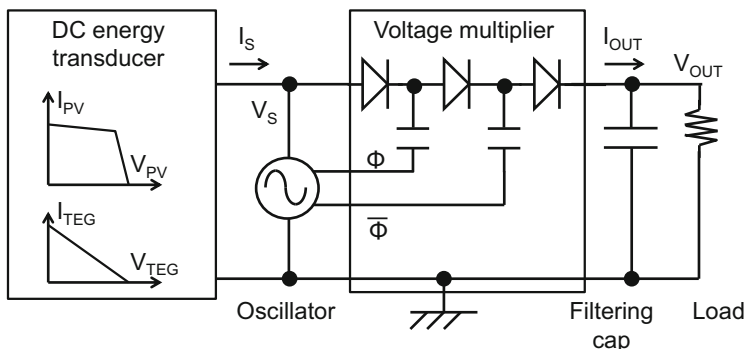
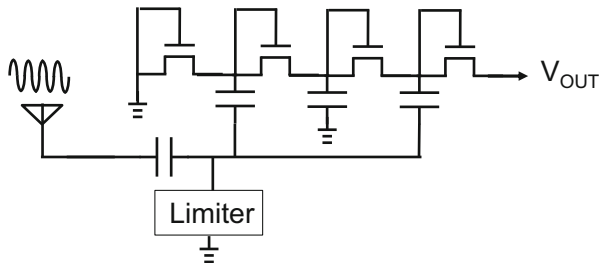
**Fig. 1.9** Energy harvester IC converting from vibration energy (Yen and Lang 2006)

connected with the ground and that at every odd number stage is connected with the common clock line. In comparison with two-phase clock Dickson pump, the single clock pump has the maximum attainable voltage lower by half but the same output impedance, when the same number of stages and same size of capacitors are used. The design of AC-DC voltage multipliers is discussed in Chap. 4 in detail.

Another energy harvester is composed of DC energy transducer like photovoltaic (PV) and thermoelectric generator (TEG) and voltage multiplier as shown in Fig. 1.11. The oscillator has to start operating at a very low voltage because DC energy transducer such as photovoltaic cell and thermoelectric generator can nominally output a voltage as low as a few hundred mV. The output impedance of the DC energy transducer is relatively high so that the charge pump needs to be designed considering the impact of the output impedance of the DC energy transducer. The total output power from the charge pump is significantly affected by the design parameters of the charge pump. This topic is discussed in Chap. 7.

In summary, design parameters of typical integrated high-voltage generator system are as follows. The voltage gain  $G_V$  is required to be 1.5–15. The supply voltage and boosted voltage are, respectively, in a range of 0.5–30 V and 1–40 V. The output current is as low as an order of 1  $\mu$ A especially in case of a high voltage gain and as high as an order of 10 mA especially in case of a low-voltage gain.

**Fig. 1.10** UHF RFID IC  
(Jun et al. 2010)



**Fig. 1.11** Energy harvester with DC energy transducer and voltage multiplier

## 1.2 System and Building Block Design Consideration

Section 1.2 summarizes key design consideration for both systems and circuits, which are discussed in detail in the following chapters.

Figure 1.12 shows on-chip high-voltage generator system and each component circuit block. The charge pump inputs the supply voltage ( $V_{DD}$ ) and the clock which is generated by the oscillator, and outputs a voltage ( $V_{PP}$ ) higher than the supply voltage or a negative voltage. The regulator enables the charge pump when the absolute value of the output voltage of the charge pump is lower than the target voltage on a basis of a reference voltage  $V_{REF}$ , or disables it otherwise.  $V_{PP}$  can vary in time by  $\Delta V_{PP\_DROP}$  due to a finite load current  $I_{LOAD}$  and by  $\Delta V_{PP\_RIPPLE}$  due to a finite response time in feedback loop with the pump regulator. The output voltage of the pump is determined by the reference voltage and the voltage gain of the regulator. To vary the pump output voltage, either reference voltage or voltage gain of the regulator is varied. The generated high or negative voltage is transferred to a load through high- or low-level shifters. The level shifters are controlled by the input supply voltage. The load is capacitive, resistive, or both. Optimization of the charge pump depends on the load characteristics.

According as the supply voltage decreases, the system design becomes more challenging in terms of (1) silicon area, (2) peak and average operation current, (3) ramp-up time, and (4) accuracy in the output voltage in DC and AC. The items

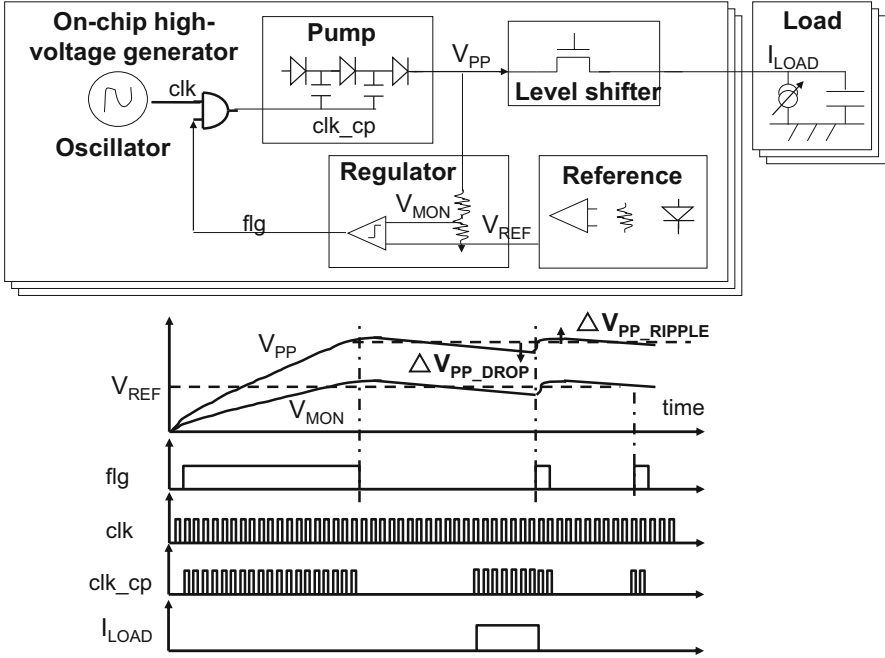


Fig. 1.12 On-chip high-voltage generator system

(1)–(3) are under a trade-off relation. If the ramp-up time needs to be kept constant even with a lower supply voltage, the pump area and operation current would increase. Or, instead, if the pump area needs to remain the same, the output current would decrease, resulting in longer ramp-up time. Therefore, high-voltage generator design requires reconsideration on the entire system due to reduction in the supply voltage.

In addition, reducing the supply voltage while keeping the output voltage level means that the voltage gain is increased. Voltage variations in the reference voltage and the divided voltage of the regulator are amplified with the increased voltage gain, resulting in less accuracy in the output voltage of the generated high voltage. Moreover, IR drop in the power ground lines significantly affects the pump output current especially with lower supply voltage and with multiple high-voltage generators on a chip. Interference between different high-voltage generators occurs via the common impedance of the power ground lines. To take such considerations into design, the parasitic resistance in the power ground lines needs to be included into one of the design parameters.

Because an on-chip high-voltage generator is one of the functional blocks on an LSI, simulation accuracy and run time have to be reasonable when all the blocks are simulated together with the generator. However, the high-frequency clock for driving the charge pump and the charge-transfer operation in the pump tends to make the simulations very slow. Thus, it is important to model the generator properly so that both the accuracy and simulation time are reasonable.

**Table 1.2** Design considerations for each block

Circuit block	Design considerations
Charge pump	Circuit topology choice to minimize the silicon area
	Devices available as capacitors and switches in technology given
	Input/output voltage/current characteristics as a function of device and circuit parameters
	Equivalent circuit model
	Power efficiency
	Design optimization for the clock to maximize the output current
	Design optimizations for the number of stages to minimize the total pump area, the rise time, or the input power
	Switching diode design with $V_T$ canceling techniques
	Capacitor design
	Wide $V_{DD}$ operation
	Area efficient multiple pump system design with reconfiguration technique
	Noise and ripple reduction design
	Standby and active pump design
Circuit model with a reasonable simulation time	
Pump regulator	Resistor design
	Reduction of variations in regulated voltages
	Trimming capability
	Response time reduction
	Negative voltage detection
Oscillator	Reduction in process, voltage, and temperature variations
	Bi-stable oscillator with a high and low duty of 50 %
	Four-phase clock generation
Level shifter	Circuit topology choice according to availability of high-voltage transistors
	Switching speed
	Energy per switching
	Minimum operating voltage
	High-voltage relaxation design
Voltage reference	Circuit topology choice according to availability of bipolar junction transistors
	Reduction in process, voltage, and temperature variations
	Minimum operating voltage

Table 1.2 summarizes design considerations for each block when circuit blocks composing an on-chip high-voltage generator are designed. Once the required voltage gain which is defined by the ratio of the high generated voltage to the supply voltage and the ratio of the parasitic capacitance of the pumping capacitor to the capacitance of the pumping capacitor are given for one's design, one can choose the best topology to minimize the charge pump circuit area. In case where those ratios are, respectively, higher than 5 and 0.03 typically, one should use the topology which Dickson experimented not only for the smallest area but also for the least power. For given transistors as switching devices in charge pumps, one can draw a graph showing the transistors can operate at how fast clock frequency. Then,

using the equivalent pump model, one can determine the design parameters such as the number of stages and capacitance per stage.

Pump regulators need to be designed with a potential variation in the output voltage of the pump considered. If it is larger than the required one, trimming capability needs to be implemented. Even if the output voltage is far from the target, trimming can adjust the output voltage closely to the target. Because the current flowing through the resistor divider needs to be small enough not to affect the net output current of the charge pump, resistance of the voltage divider tends to be relatively large. Adding switching devices for trimming can also increase RC time constant of the divider, which results in slow response from the time when the output voltage of the pump reaches the target to the time when the opamp detects it. According to the response delay, the pump operation continues to increase the output voltage, which creates the ripple in the output voltage. Therefore, the response time improvement is required to stabilize the output voltage.

Oscillators driving the charge pumps directly affect the pump output current. Higher frequency results in larger output current under a nominal condition. Thus, PVT (process, voltage, and temperature) variations in the frequency lead those both in the output and input current. If the pump is designed so that the output current at the slow condition meets the required one, the peak power is seen at the fast condition. Reduction in PVT variations in the oscillator is a key to make the pump performance stable.

Circuit topology choice due to the device availability and minimum operation voltage are common design concerns for level shifters and voltage references. In addition, level shifters need fast switching speed and robustness on high-voltage stress. One has to make sure of long-term operation under high-voltage stress.

In the following chapters, each design consideration is discussed.

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# Chapter 2

## Basics of Charge Pump Circuit

**Abstract** This chapter discusses circuit theory of the charge pump circuit. Since it was invented, various types have been proposed. After several typical types of charge pumps are reviewed, they are compared in terms of the circuit area and the power efficiency. The type that Dickson proposed is found to be the best one as an on-chip generator where the parasitic capacitance is 1–10 % of the pump capacitor.

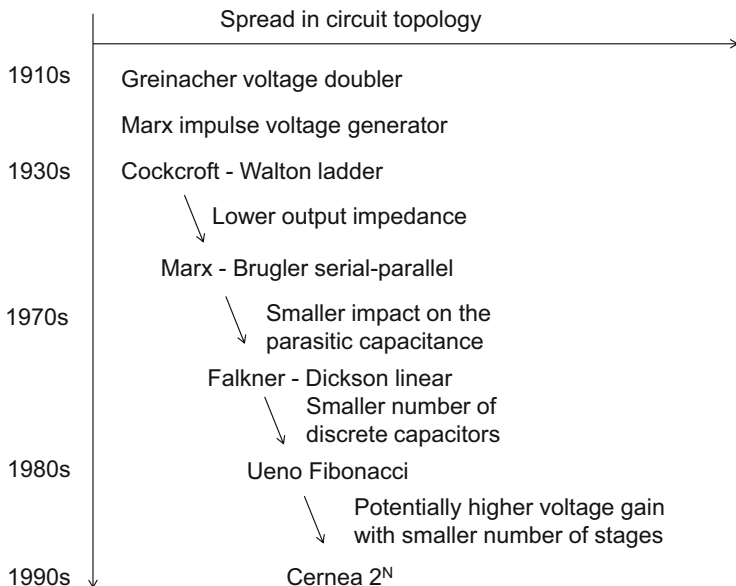
This chapter is composed of the following. Section 2.1 reviews several pump topologies and qualitative comparison among them. Section 2.2 shows matrix expression of charge pump cell to commonly analyze the circuit behavior of each topology. The following Sects. 2.3–2.7, respectively, present operation analysis of each pump cell, i.e., Greinacher and Cockcroft–Walton cell, Marx–Brugler serial–parallel cell, Falkner–Dickson cell, Ueno–Fibonacci cell, and Cernea–2 N cell. Section 2.8 compares them quantitatively. The results suggest that Dickson cell is the best topology of integrated high-voltage multiplier because of the largest voltage gain and smallest circuit area.

### 2.1 Pump Topologies and Qualitative Comparison

This section begins with a brief history of several topologies of charge pump and their background on the critical characteristic parameters, i.e., the output impedance and the maximum attainable voltage. Operation of the initial topology as known as Cockcroft–Walton multiplier is discussed and the characteristic parameters are shown. Optimum design for maximizing the output power is, respectively, given under the conditions of resistive load and current load. After that, several topologies of pump are described which aim at having lower output impedance for higher output current at a given output voltage. Qualitative sensitivity analysis on the parasitic capacitance of pump capacitors suggests that larger number of serially connected capacitors results in larger impact of the parasitic capacitance on the output current.

The switched-capacitor (SC) multiplier originated with Greinacher and Cockcroft–Walton (CW) using serial capacitor ladders independently. Because the CW multiplier had a relatively large output impedance with an order of  $N^3$ ,

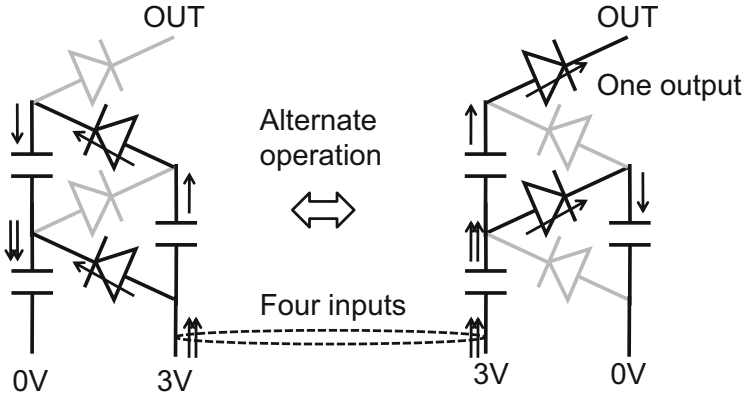




**Fig. 2.1** History of two-phase clock charge pump voltage multipliers

where  $N$  is the number of stages, various types of multipliers with different topologies have been proposed to reduce the output impedance. By alternately switching the state from in-serial to in-parallel and vice versa, Brugler theoretically showed that the serial-parallel (SP) multiplier, which is the same topology as Mark impulse voltage generator, had lower output impedance with an order of  $N^1$  than that of CW. Falkner suggested that parallel capacitor ladders reduced the output impedance as well. Dickson theoretically and experimentally showed that the output impedance of the parallel capacitor ladders was proportional to  $N^1$ . Another direction for improving SC performance is to increase the maximum attainable voltage gain  $G_{MAX}$ . Ueno et al. proposed the Fibonacci SC multiplier whose  $G_{MAX}$  is given by the  $N$ th Fibonacci number of approximately  $1.16 \exp(0.483N)$ . The multipliers whose  $G_{MAX}$  is given by  $2^N$  were proposed by Ueno et al. with multiphase switching clocks and by Cernea with two-phase clocks. Figure 2.1 briefly summarizes the history of two-phase clock charge pump voltage multiplier. Note that recent integrated high-voltage generators are mainly based on the Dickson linear pump topology.

Performance analysis and design methodologies have also been done for the multipliers as described above. To determine an optimum multiplier topology under specific conditions, comparisons among those multipliers on circuit performance have also been made. Before advancing quantitative analysis, this chapter starts with qualitative analysis on which multiplier is optimum with respect to circuit area under the condition that a given current is output at a given output voltage with a given parasitic capacitance.



**Fig. 2.2** Three-stage CW in steady state (Cockcroft and Walton 1932)

Figure 2.2 shows how the CW circuit works. The number of stage is defined by the number of capacitors, i.e., three in this example. The number of diodes is four, larger by one than the number of stages. Because the CW works with two-phase clock, one only needs to take care of these two half of periods. The diodes in gray are not under conduction state. One arrow indicates amount of charge  $Q$  which flows into the output terminal in a period. In the left hand side figure, a same amount of charge  $Q$  flows through each diode under a steady state. The top most two capacitors flow the same amount of  $Q$  to meet the condition that the current is continuous. Thus, the bottom most capacitor in the left branch and the power supply in the right branch flow amount of charges of  $2Q$  to meet Kirchhoff's law. In the right-hand side figure, both diodes and topmost two capacitors, respectively, flow  $Q$ . The bottom capacitor and the power supply in the left branch flow  $2Q$ . As a result, one has to input  $4Q$  to output  $1Q$  per period. As one can easily guess, the current efficiency defined by the output current  $I_{OUT}$  over the input current  $I_{IN}$  is as shown by  $I_{OUT}/I_{IN} = 1/(\#diodes) = 1/(\#stages + 1)$ .

Based on Brugar's approach for theoretical steady state equation, one can calculate the relation between  $I_{OUT}$  and  $V_{OUT}$  using Figs. 2.3, 2.4, 2.5, and 2.6.  $V^-$  indicates the first half period and  $V^+$  does the second half period. Each of  $V^-_i$  and  $V^+_i$  ( $i = 1, 2, 3$ ) indicates the voltage difference between two terminals of each capacitor. The assumptions used here are that the period is too long to be able to neglect any RC time delay and that the threshold voltage of each diode is zero. Starting with Fig. 2.3,  $V_3^-$  is equalized to  $V_{DD}$ . Because  $V_3^+$  is the voltage after amount of charges  $2Q$  is transferred through the diode, it should be given by  $V_3^+ = V_3^- - 2Q/C$ . As a result, it is solved as  $V_3^+ = V_{DD} - 2Q/C$ :

$$V_3^- = V_{DD} \tag{2.1}$$

$$V_3^+ = V_3^- - 2Q/C \tag{2.2}$$

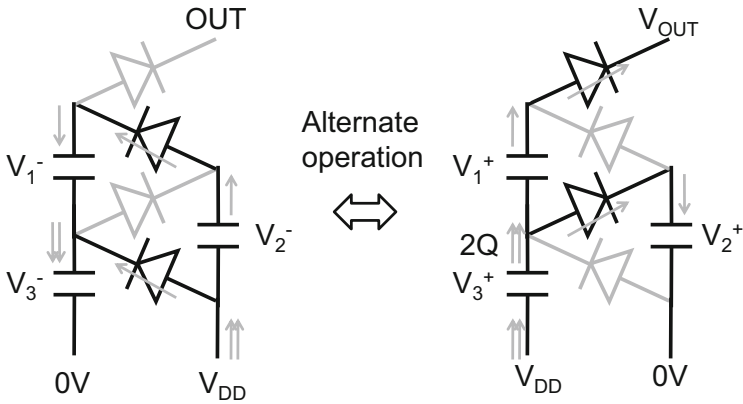


Fig. 2.3 Relation of  $V_3^+$  and  $V_3^-$  to the other nodal voltages in steady state

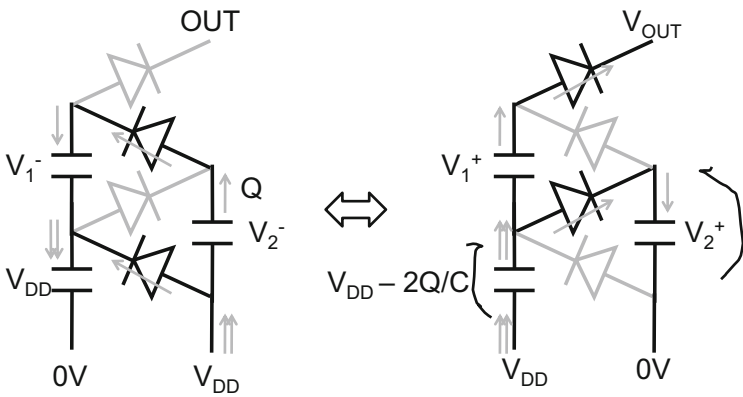


Fig. 2.4 Relation of  $V_2^+$  and  $V_2^-$  to the other nodal voltages in steady state

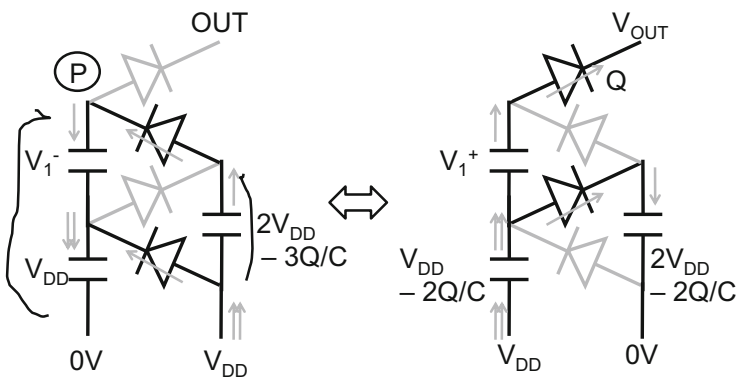
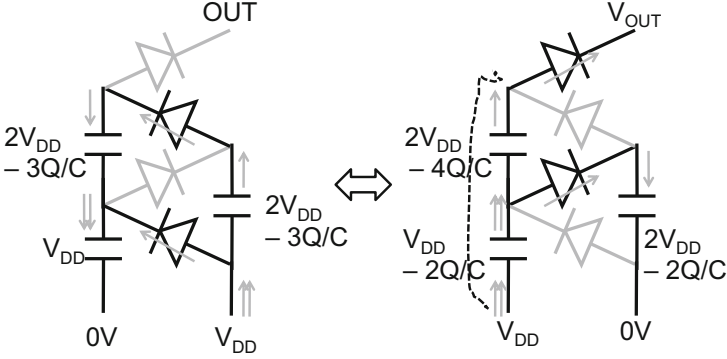


Fig. 2.5 Relation of  $V_1^+$  and  $V_1^-$  to the other nodal voltages in steady state



**Fig. 2.6** Capacitor voltages in steady state

Next, Fig. 2.4 focuses on the relation of  $V_2^+$  and  $V_2^-$  to the other nodal voltages. As shown in the right-hand side,  $V_2^+$  is equalized to  $V_3^+$ , resulting in  $2V_{DD} - 2Q/C = V_2^+$ . As shown in the left-hand side,  $V_2^-$  is lower by  $Q/C$  than  $V_2^+$ . From these two equations,  $V_2^-$  is given by  $V_2^- = 2V_{DD} - 3Q/C$ :

$$V_2^+ = 2V_{DD} - 2Q/C \quad (2.3)$$

$$V_2^- = V_2^+ - Q/C = 2V_{DD} - 3Q/C \quad (2.4)$$

Similarly, Fig. 2.5 focuses on the relation of  $V_1^+$  and  $V_1^-$  to the other nodal voltages. As shown in the left-hand side, the potential at the point  $P$  is calculated by two ways. The first one is  $V_1^- + V_{DD}$  in the left path. The right path results in  $V_2^- + V_{DD} = (2V_{DD} - 3Q/C) + V_{DD} = (3V_{DD} - 3Q/C)$ . By equating these two, one has Eq. (2.5). The right-hand side figure simply indicates that  $V_1^+$  is lower by  $Q/C$  than  $V_1^-$ , thereby Eq. (2.6):

$$V_1^- = 2V_{DD} - 3Q/C \quad (2.5)$$

$$V_1^+ = 2V_{DD} - 4Q/C \quad (2.6)$$

Finally, Fig. 2.6 shows capacitor voltages.  $V_{OUT}$  is calculated with the sum of the capacitor voltages in the left path plus  $V_{DD}$  in the right-hand side figure:

$$V_{OUT} = (2V_{DD} - 4Q/C) + (V_{DD} - 2Q/C) + V_{DD} = 4V_{DD} - 6Q/C \quad (2.7)$$

Thus,  $V_{OUT}$  has two terms. The first term is proportional to  $V_{DD}$ . The multiplication factor of 4 is resulted from the number of capacitors that is the number of stages plus one from  $V_{DD}$  of the clock amplitude. The second term is proportional to  $Q$ . The multiplication factor is larger than the number of stages. This fact is resulted from the fact that amount of charges transferred to the next stage increases as the capacitor position gets closer to  $V_{OUT}$ . Thus, the sum of the multiplication factors tends to be higher as the number of stages increases. This means that the effective

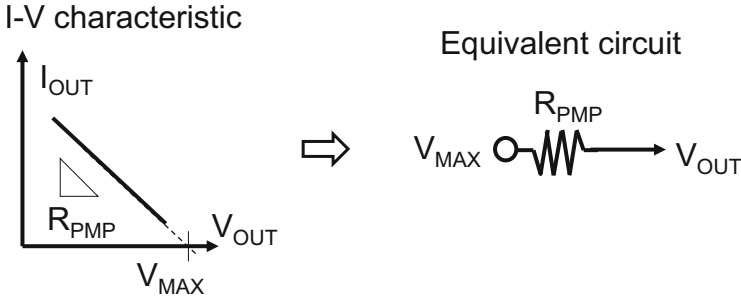


Fig. 2.7 Relation of  $I_{OUT}$  to  $V_{OUT}$  and an equivalent circuit in steady state

impedance of the CW multiplier rapidly increases as the number of stages increases.

What does Eq. (2.7) suggest? Introducing the cycle time  $T$  of the clock, the average output current  $I_{OUT}$  is expressed by Eq. (2.8), where  $V_{MAX}$  is the maximum attainable output voltage when  $I_{OUT}$  is zero as shown by Eq. (2.9) and  $R_{PMP}$  is the effective impedance of the pump as shown by Eq. (2.10), which will be derived in the next section:

$$I_{OUT} \equiv Q/T = (V_{MAX} - V_{OUT})/R_{PMP} \quad (2.8)$$

$$V_{MAX} = 4V_{DD} \rightarrow (N + 1)V_{DD} \quad (2.9)$$

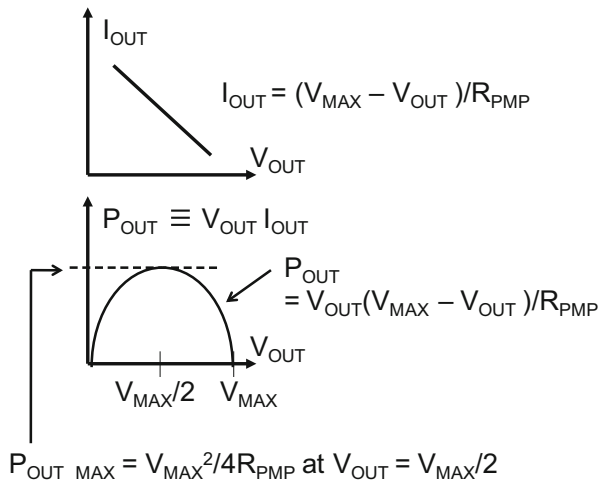
$$R_{PMP} = 6T/C \rightarrow \sim (N + 1)^3/12T/C \quad (2.10)$$

Every topology of charge pumps has a similar  $I-V$  curve with these two characteristic parameters. The equivalent circuit is a simple voltage source and a linear resistor as illustrated in Fig. 2.7. In this example of three-stage CW pump,  $V_{MAX}$  is  $4V_{DD}$  and  $R_{PMP}$  is  $6T/C$ . One can qualitatively consider the power of three in Eq. (2.10) as follows. One comes from the amount of charges proportional to the number of stages, another one comes from  $k$ th capacitor from the bottom transferring the amount of charges proportional to  $k$ , and the last one comes from the amount of charges summed in all the capacitors in the left path. Each of those three factors is proportional to  $N$ , resulting in the power of three. More general and comprehensive discussions are done in the next section.

What else is resulted from the  $I-V$  equation is the optimum operating point where the output power is maximized as shown in Fig. 2.8. The above graph is the  $I_{OUT}-V_{OUT}$  curve. The output power is a multiple of  $I_{OUT}$  with  $V_{OUT}$ , resulting in a quadratic function. The maximum is given at a half of  $V_{MAX}$  because the  $X$ -interceptions occur at zero and  $V_{MAX}$ . The maximum power is then given by Eq. (2.11):

$$P_{OUT\_MAX} = V_{MAX}^2/4R_{PMP} \text{ at } V_{OUT} = V_{MAX}/2 \quad (2.11)$$

**Fig. 2.8** Conditions for maximizing the output power



One may have different load conditions such a resistive load and a current load. No matter what the load is, the optimum operating point in terms of maximizing the output power is at a half of  $V_{MAX}$ , as shown in Fig. 2.9. In case of a resistive load, one can maximize the output power with designing  $R_{PMP}$  matched with  $R_L$ , which is the so-called impedance match:

$$R_{PMP} = R_L \quad (2.12)$$

In case of a current load, one can maximize the output power when the following relation between  $R_{PMP}$  and  $V_{MAX}$  is met:

$$V_{MAX}/R_{PMP} = 2I_L \quad (2.13)$$

Note that maximizing the output power under a given voltage of  $V_{PP}$  is equivalent to maximizing the output current at  $V_{PP}$ . Equations (2.11) to (2.13) are independent of a type of charge pump topology as far as the  $I_{OUT} - V_{OUT}$  characteristic is the same form.

Because of quite high impedance with the CW pump with a relatively large voltage gain, there has not been lots of practice to implement the CW pump. One example implementation of CW in ICs is shown in Fig. 2.10. The key feature of the CW over the other types of pump is that every diode and capacitor sees a voltage difference of  $V_{DD}$  or less. This means that one can construct the pump with low-voltage devices, resulting in smaller circuit area with scaled devices. The circuit designers need to make sure that any device wouldn't be broken down under any emergent case such as a sudden power shutdown and a sudden short of the output node to the ground. Under such circumstance, a high voltage may appear in any low-voltage device.

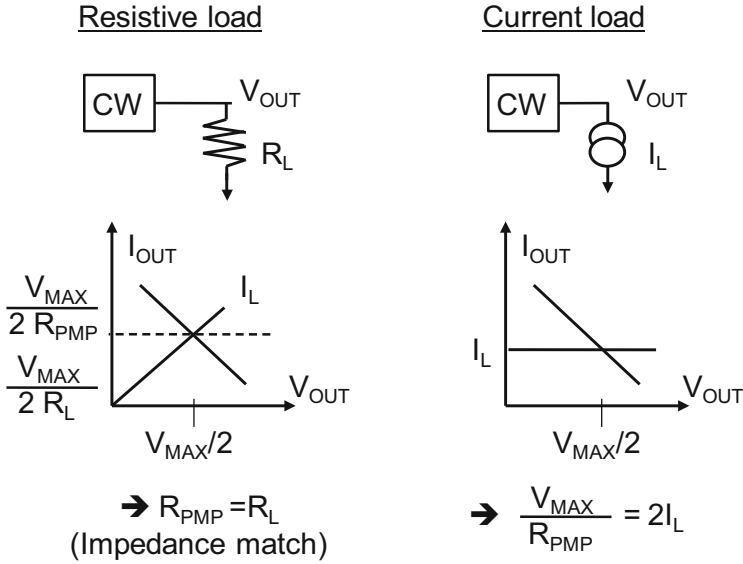


Fig. 2.9 Conditions for maximizing the output power

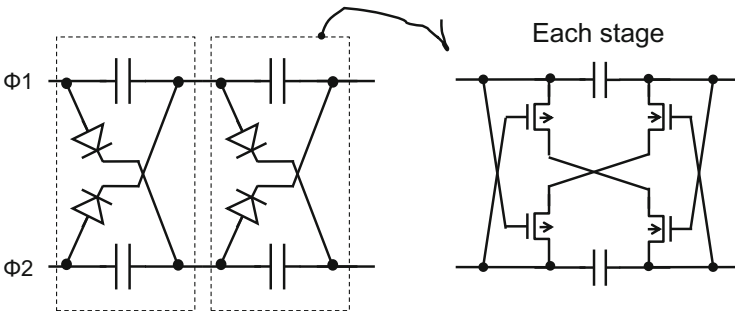
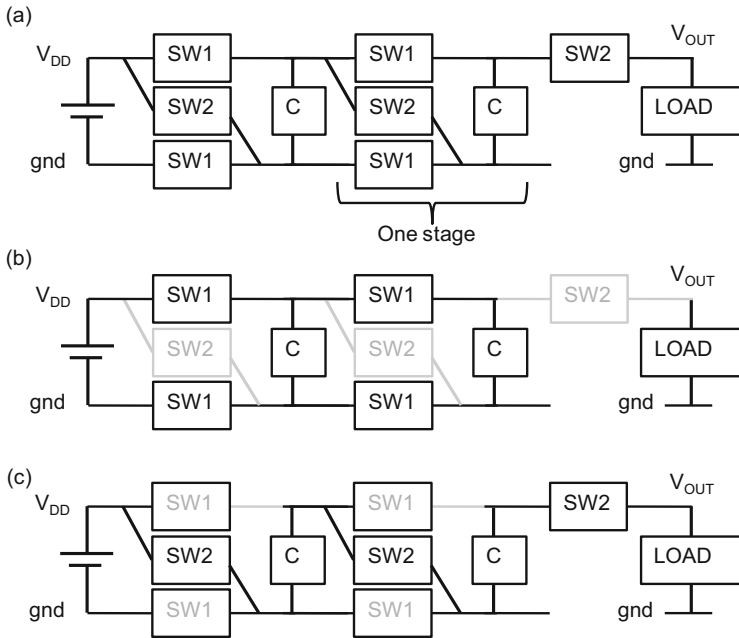
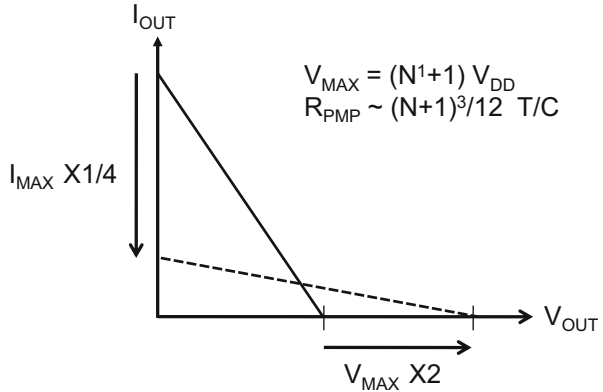


Fig. 2.10 Implementation of CW in ICs (Zhang et al. 2009)

How significant is the power of three with respect to the number of stages in the output impedance of the CW pump? If one needs to double the number of stages to increase  $V_{MAX}$  twice, the output impedance decreases by a factor of eight. Then, the maximum output current where the output voltage is zero decreases by a factor of four, as shown in Fig. 2.11. When one designs the operating point at a half of  $V_{MAX}$ , the output current can decrease by a factor of four as well. Thus, the reduction rate in  $I_{OUT}$  over  $V_{OUT}$  is proportional to the squared number of stages. Thus, the CW multiplier may not be good for the cases where a large voltage gain is needed.

Brugler theoretically showed that there was another topology where the output impedance could be reduced as illustrated in Fig. 2.12a. Adding two more switches per stage, the capacitors can be switched from in-parallel (b) to in-series

**Fig. 2.11**  $I_{OUT}-V_{OUT}$  characteristic when the number of stages is doubled in CW

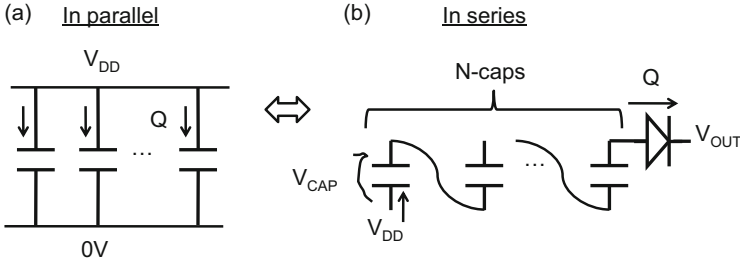


**Fig. 2.12** Serial-parallel switched capacitor with lower  $R_{PMP}$  (Marx 1928; Brugler 1971)

(c) alternately, which is the same topology as Marx impulse voltage generator. All the capacitors are charged to  $V_{DD}$  in a parallel period and are connected in-series between  $V_{DD}$  and the output terminal. Hereinafter, one can call this type of pump serial-parallel or SP.

The procedure to extract the  $I_{OUT} - V_{OUT}$  equation is much easier than the case of CW using Fig. 2.13. Assuming  $Q$  is the amount of charge to be transferred to the output terminal in in-series period. Each capacitor loses the same amount of  $Q$  in





**Fig. 2.13** Two phases of SP

this period. Thus, each capacitor needs to be charged by  $Q$  in in-parallel period. Before charging  $Q$ , each capacitor voltage should be  $V_{DD} - Q/C \equiv V_{CAP}$ . Thus,  $V_{OUT}$  can be related to  $V_{CAP}$  as Eq. (2.14):

$$V_{OUT} = V_{DD} + NV_{CAP} = (N + 1)V_{DD} - NQ/C \quad (2.14)$$

As a result,

$$I_{OUT} \equiv Q/T = (V_{MAX} - V_{OUT})/R_{PMP} \quad (2.15a)$$

$$V_{MAX} = (N + 1)V_{DD} \quad (2.15b)$$

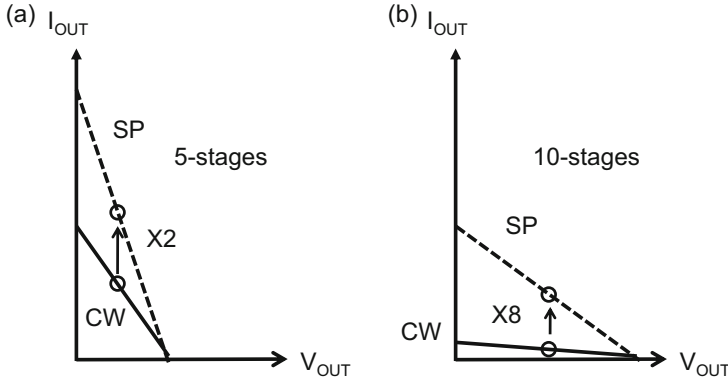
$$R_{PMP} = N^1 T/C \quad (2.15c)$$

Thus, the output impedance is proportional to  $N^1$ . To output  $Q$  in a period, each capacitor doesn't need to do extra work than getting  $Q$  from the power supply. The current efficiency defined by the total output current over the total input current is  $1/(\text{the number of capacitors} + 1)$  as same as that of the CW. There is no advantage in the current efficiency with the serial-parallel pump.

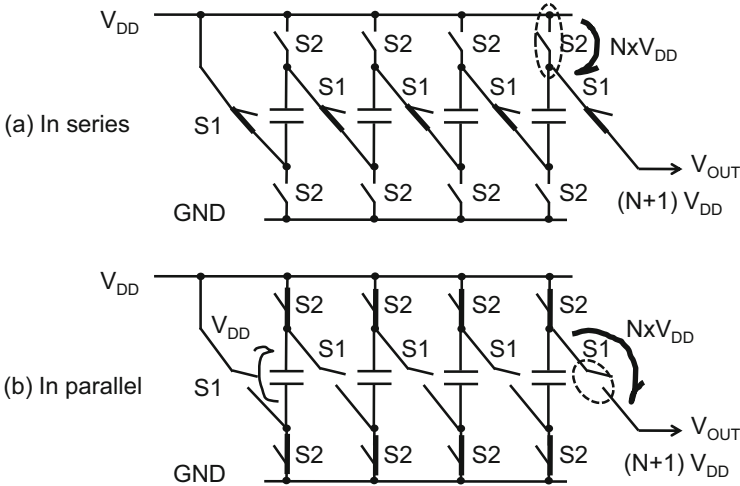
A question here is how significant lower impedance is with the serial-parallel pump. Figure 2.14a, b shows, respectively, 5- and 10-stage pumps'  $I$ - $V$  characteristics. The broken lines show the SP and the solid lines show CW. All the capacitors are assumed to be same. Under the condition, the SP has larger output current than the CW does especially when the number of stage is larger.

Figure 2.15 shows the requirement for breakdown voltages for the capacitors and switches used in the SP. The capacitor voltage in parallel state is equal to  $V_{DD}$  and that in serial state is lower than  $V_{DD}$  by  $Q/C$ . Therefore, the capacitor could be made of a low-voltage device, which enables to reduce the capacitor area with higher capacitance density. On the other hand, the switches used closely to the output terminal see  $N$  times higher than  $V_{DD}$  for both states, resulting in requirement for high-voltage switching devices.

Table 2.1 summarizes comparison of SP with CW in terms of the pump characteristics' parameters and the voltage requirements for capacitors and switches. The maximum attainable voltage  $V_{MAX}$  is no difference. The output impedance of SP is



**Fig. 2.14** Comparisons of 5- and 10-stage pumps'  $I$ - $V$  characteristics between CW in solid lines and SP in broken lines



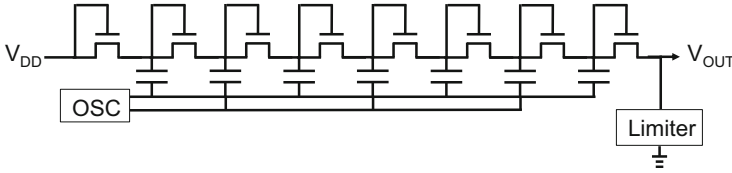
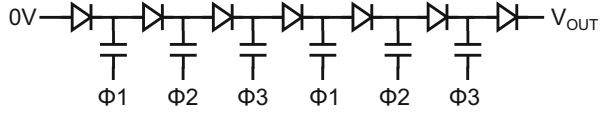
**Fig. 2.15** Requirement for breakdown voltage in SP

**Table 2.1** Comparison in characteristic parameters between CW and SP

	$V_{MAX}$	$R_{PMP}$	$V_{BV\_CAP}$	$V_{BV\_SW}$
CW	$(N + 1) V_{DD}$	$(N + 1)^3 / 12 T / C$	$V_{DD}$	$1V_{DD}$
SP	$(N + 1) V_{DD}$	$N^1 T / C$	$V_{DD}$	$NV_{DD}$

proportional to  $N^1$ , whereas that of CW is to  $N^3$ . The maximum voltage applied to a capacitor is same to be  $V_{DD}$ . The maximum voltage applied to a switch of the CW is  $1V_{DD}$ , whereas that of the SP is  $NV_{DD}$ . From the system view point, one needs to have high-voltage switches to connect the output terminal to a load. Thus, high-voltage devices should be available in designing the LSIs. So, requirement

**Fig. 2.16** Another proposal for lower  $R_{PMP}$  (Falkner 1973)



**Fig. 2.17** First Si verified on-chip Dickson pump (Dickson 1976a, b)

of high-voltage device for a switch in the SP itself shouldn't be considered as a drawback. But, the maximum operating clock frequency could be affected by the high-voltage device, because a high-voltage device is typically slower than a low-voltage device. Relation between scaling of device and operating frequency will be discussed in details in Chap. 3.

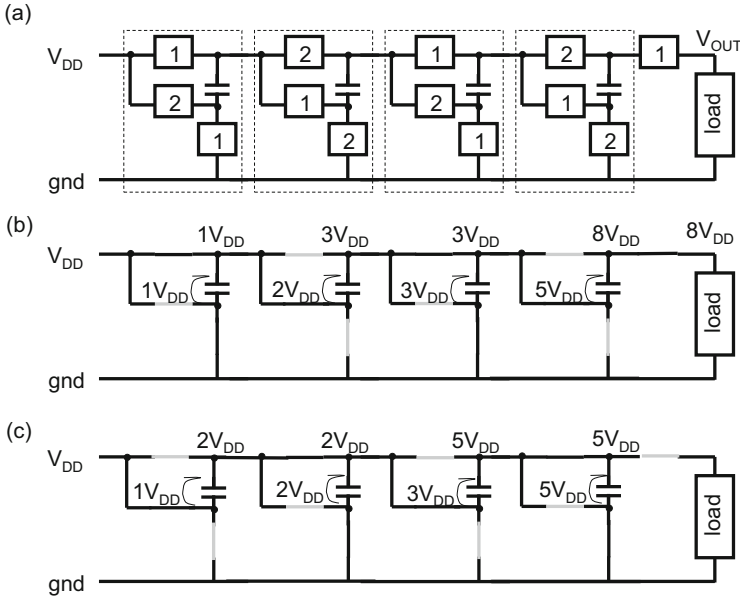
Falkner schematically showed another pump topology with a lower  $R_{PMP}$  than the CW, as shown in Fig. 2.16. The circuit has three phase clock, but it is not the essence. Key point is that each capacitor is connected with next one or two stages in parallel at a time. Unlike the CW has the state with half of stages connected in-series and the SP has that with all stages connected in-series. The numbers of switches or diodes are that of capacitors plus one, which is the same condition as the CW.

In 1976, Dickson theoretically and experimentally for the first time studied an on-chip high-voltage generator including a charge pump, oscillator, clock drivers, and a limiter, as shown in Fig. 2.17. The diode was made of a MOSFET whose gate and drain terminals are connected to play the same role as a rectifying diode. Dickson used two-phase clock which allowed the clock frequency faster than the three-phase clock of Fig. 2.16. Using a seven stage pump, a high voltage of 40 V was successfully generated from the power supply voltage of 15 V. The operation principle will be discussed in Chap. 3 in details.

Another type of two-phase pump was proposed by Ueno et al. aiming at reducing the number of capacitors for low cost small form factor discrete applications, as shown in Fig. 2.18. The interesting characteristic is the maximum attainable voltage is given by Fibonacci number, i.e., 2, 3, 5, 8, 13, and 21:

$$\text{Fib}(N) = \text{Fib}(N - 1) + \text{Fib}(N - 2) \tag{2.16}$$

where  $\text{Fib}(1) = 2$ ,  $\text{Fib}(2) = 3$ . As the number of stages increases, the voltage gain increases more rapidly than the number of stages. For example, when one needs to have  $V_{MAX}$  of 13, one only needs five stages with Fibonacci pump, whereas 12 stages with CW, SP, or Dickson pump. Each stage has one capacitor and three

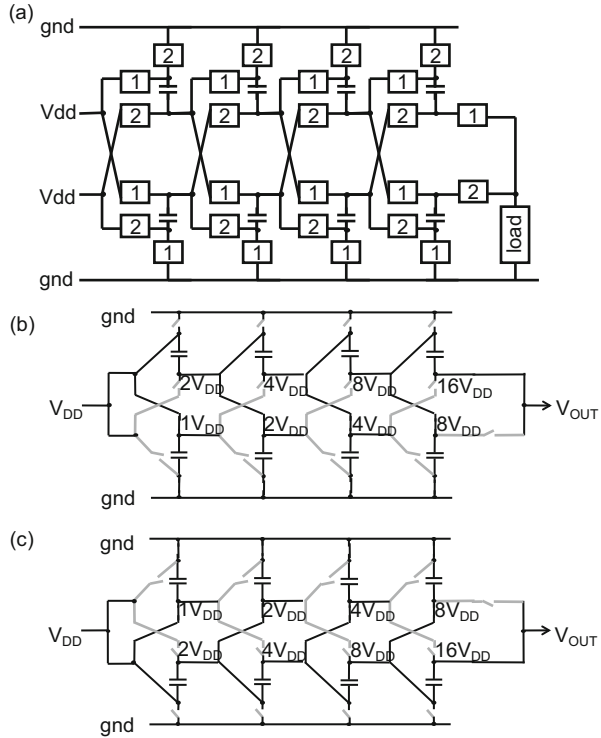


**Fig. 2.18** Fibonacci-type multiplier (Ueno et al. 1991)

switches, as shown in Fig. 2.18a. The number 1 and 2 in the boxes indicate that the switch marked as 1 turns on in a first half period and turns off in a second half of period and the switch marked as 2 turns on in the second half period and turns off in the first half of period. Figure 2.18b shows the connection states in the first half period. Even number of stages is connected in series with the output terminal and odd number of stages is connected in parallel to the serial ones; in other words,  $(2k - 1)$ th stage is connected with  $2k$ th stage in parallel. The nodal voltages shown are valid only when the output current is zero. Figure 2.18c shows the connection states in the second half period. The situations are complementary to the first half period. Thus, a half of stages are in series and the other half of stages are in parallel, alternately.

The last one is  $2^N$  multiplier as shown in Fig. 2.19. When the number of stages connected between the input and the output is  $N$ , the required number of capacitors is  $2N$  because two arrays are required to complete the multiplier unlike the other types of pump. Figure 2.19b shows a first half period. The upper stages are connected in series with the output terminal, whereas the lower stages are connected in parallel with the upper stages, or in other words,  $k$ th lower stage is connected in parallel with  $k$ th upper stage. The voltage values shown in Fig. 2.19b are those in case of no load current. As the number of stages increases by one, the maximum attainable voltage increases by a factor of two in an ideal case where no parasitic capacitance is considered. One may consider the number of stages of  $2^N$  multiplier is smaller than that of the Fibonacci pump. But, the number of capacitors of  $2N$  multiplier is larger than that of the Fibonacci one because the  $2^N$  pump needs

**Fig. 2.19** 2-Phase  $2^N$  multiplier (Cerna 1995)

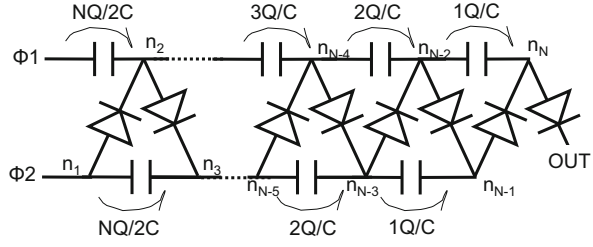


two arrays. For example, when a maximum attainable voltage gain of 16 is required,  $2^N$  pump needs at least eight capacitors as shown in Fig. 2.19, whereas Fibonacci pump does six capacitors.

Several topologies of two phase pump are overviewed. Now one should have a question about which topology should be selected for ICs as on-chip high-voltage generator. To answer the question, one has to take the two factors in terms of parasitic elements into consideration. The first one is a finite threshold voltage  $V_T$  of a real switching device. But, it simply reduces the voltage amplitude at each capacitor node from  $V_{DD}$  and doesn't affect the comparison between different topologies. Once can replace  $V_{DD}$  with  $V_{DD} - V_T$ . Besides, it can be mitigated with several design techniques to effectively eliminate  $V_T$ . State of the art will be overviewed in Chap. 5. The second one is a finite parasitic capacitance ( $C_P$ ) of a real capacitor and switch. Unfortunately, there is no design technique to eliminate the parasitic capacitance. Therefore, sensitivity of  $C_P$  on the pump performance could determine the best topology for integration because the ratio of  $C_P$  to the integrated capacitor can be much larger than that of the discrete capacitor.

Figure 2.20 illustrates  $N$ -stage CW pump. The values shown are the voltage amplitude of each capacitor between two half periods of cycle, which are suggested

**Fig. 2.20** Voltage amplitude of each capacitor in  $N$ -stage CW



by Eqs. (2.1)–(2.6) in the case of three stages. Therefore, the voltage amplitude at the  $(N - k)$ th node,  $(V_{N-k}^+ - V_{N-k}^-)$ , is calculated as  $f(N - k)Q/C$ , where

$$f(N - k) = \sum_{i=1}^{k/2} i = k(k + 2)/8 \quad (2.17)$$

Suppose that each node has the parasitic capacitance  $C_P$ . The power supply driving the clocks  $\Phi 1, 2$  charges  $C_P f(N - k)Q/C$  for  $(N - k)$ th node. Hence, the total amount of charge to  $C_P$  of all nodes,  $Q_P$ , is

$$\begin{aligned} Q_P &= C_P \sum_{k=1}^N k(k + 2)Q/8C \\ &= [N(N + 1)(2N + 7)/48](C_P/C)Q \end{aligned} \quad (2.18)$$

According as the number of stages increases,  $Q_P$  increases with the cube of the number of stages. When  $Q_P$  becomes compatible to  $Q$ , the voltage at each node would decrease from the ideal cases such as Eqs. (2.1)–(2.6) because the voltage amplitude reduces accordingly, resulting in invalidity of Eq. (2.18). For now, one uses Eq. (2.18) as the first-order estimate.

Its worth of taking a look at the impact of parasitic capacitance on  $I$ – $V$  of the SP. Figure 2.21 illustrates an in-series state with no  $C_P$  in ideal case (a), that with  $C_P$  in real case (b), and an in-parallel state (c). The SP works changing the state between (b) and (c), alternately. When all the capacitors are connected in parallel with the power supply, there is no impact of the parasitic capacitance on stored amount of charge in the capacitors. When the capacitors are connected in series, if the parasitic capacitance is negligibly small, each capacitor transfers a same amount of charge  $Q$  to the next capacitor, resulting in outputting  $Q$ , as shown in (a). However, if the parasitic capacitance is not negligible, the transferred charge is reduced at every node. To be worse, the charge loss at an upper node is larger than that at a lower node. Simply assuming  $k$ th capacitor reduces the charge  $q_k$  proportional to  $kV_{DD}$ , which is the voltage amplitude from in-parallel state (c) to in-series state (b), the sum of charge loss from the bottom to the top,  $\Sigma kV_{DD}$ , would be proportional to  $N^2$ . This means that the charge loss increases as the square of the voltage gain. The output charge  $Q$  could be eventually down to zero when the parasitic capacitance and the number of stages are large. For example, when  $C_P/C = 0.1$ , the charge loss

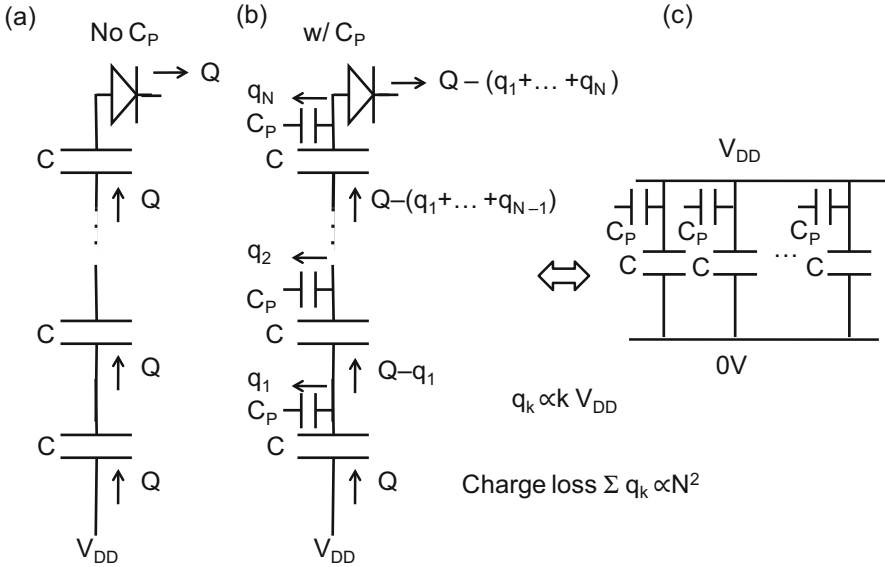


Fig. 2.21 Impact of parasitic capacitance on  $I$ - $V$ s: SP

of  $C_P/CN^2$  becomes greater than 1 with  $N$  of 4. This means that one never have a voltage gain of 5 or larger. Therefore, the impact of the parasitic capacitance is very large in the SP as well as the CW.

What about the Dickson pump? Figure 2.22 illustrates three stages of a Dickson pump. The charge supplied from the power supply is  $Q$  independent of the capacitor location. Assuming each capacitor loses the charge  $q$  due to the parasitic capacitance, every capacitor can transfer  $Q - q$  independent of the capacitor location unlike CW and SP. The difference from the SP is that each capacitor of the Dickson gains the input charge of  $Q$  and loses  $q$ . On the other hand, all the stages of the SP have only one input terminal as shown in Fig. 2.21b.

To simplify the estimates of the impact of the parasitic capacitance in Fibonacci and  $2^N$  pumps, the special case where the output voltage is at the maximum attainable voltage is considered here. According to Fig. 2.18b, c, the voltage amplitude of  $k$ th stage between the first and second half periods in the Fibonacci pump can be expressed by  $Fib(k) - Fib(k-1) = Fib(k-2)$ . Similarly, based on Fig. 2.19b, c, the  $2^N$  pump has  $(V_k^+ - V_k^-)$  of  $2^N - 2^{N-1} = 2^{N-1}$ .

Table 2.2 summarizes the comparison table among the five types of pump. Charge loss due to a parasitic capacitance is proportional to the voltage amplitude at each node in one period. This means that the larger voltage amplitude the larger charge loss. In this regard, the pumps except for the Dickson have more significant impact on the parasitic capacitance than the Dickson. The charge loss is accumulated when the number of input is small. The pumps except for Dickson have one or two inputs only. Thus, the accumulated charge loss is much larger than the Dickson.

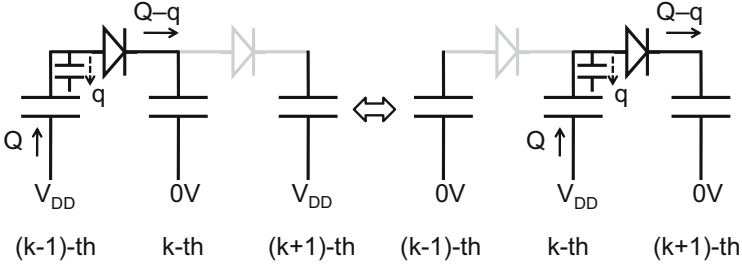


Fig. 2.22 Impact of parasitic capacitance on  $I$ - $V$ s: Dickson

Table 2.2 Qualitative comparison between five topologies of pump

	CW	SP	FIB	$2^N$	Dickson
$(V_k^+ - V_k^-)/V_{DD}$	$\sim k^2$	$\sim k^1$	Fib( $k-2$ )	$2^{k-1}$	1
# of input terminal	2	1	1	2	$N + 1$

From these qualitative view points, the Dickson seems to have the least sensitivity of the parasitic capacitance. But, the next question is if its valid quantitatively too.

## 2.2 Matrix Expression of Charge Pump Cell

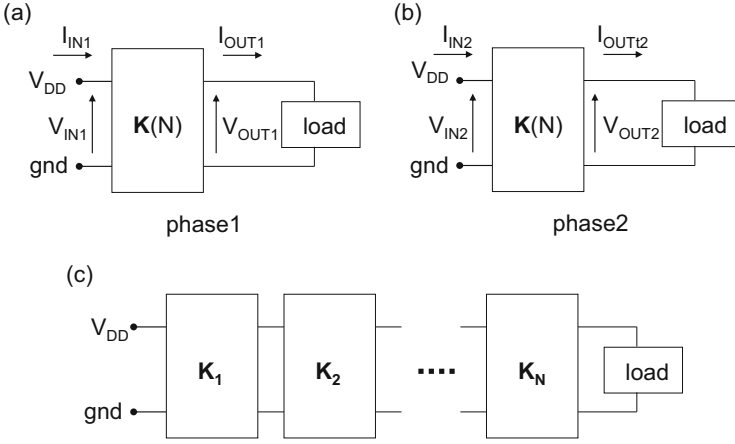
All the two-phase charge pump multipliers discussed in this section have the same symbolical structure as shown in Fig. 2.23, using the two-port transfer matrix  $\mathbf{K}(N)$  that was introduced by Harada et al., where  $N$  is the number of stages.  $\mathbf{K}(N)$  connects the input and output voltages and currents as shown by Eq. (2.19), where a subscript number 1 or 2 indicates phase 1 or 2 as shown in Fig. 2.23a, b. Each stage has a similar four-port structure, as shown in Fig. 2.23c, where  $\mathbf{K}_j$  is the matrix representing  $j$ th stage. Considering the fact that the output of  $j$ th stage is the input of  $(j + 1)$ th stage,  $N$  matrices are simply combined into  $\mathbf{K}(N)$  as shown by Eq. (2.20):

$$[V_{IN1} \ V_{IN2} \ I_{IN1} \ I_{IN2}]^T = \mathbf{K}(N)[V_{OUT1} \ V_{OUT2} \ I_{OUT1} \ I_{OUT2}]^T \quad (2.19)$$

$$\mathbf{K}(N) = \mathbf{K}_1 \mathbf{K}_2 \cdots \mathbf{K}_N \quad (2.20)$$

In the following section, these various types of switched capacitor multiplier are reviewed under the ideal condition where the parasitic capacitance is small enough to be ignored in the analysis, the operation frequency is so slow that internal capacitor nodes are fully charged and discharged in each half of period, and the clock amplitude is high enough to eliminate the effect of the threshold voltages of diodes or switching transistors. Then, the optimum multiplier is identified among serial-parallel, linear, Fibonacci, and  $2^N$  multipliers where the impact of the





**Fig. 2.23**  $K$ -matrix expression of a charge pump multiplier (Harada et al. 1992a, b)

parasitic capacitance is considered. Two-port transfer matrix for calculating an output and input voltage and current of SP, FIB, and  $2^N$  cells with parasitic capacitance at capacitor nodes, which greatly affects the pump performance, is introduced. Numerical results on circuit area and current efficiency as a function of output voltage and parasitic capacitance are shown by using the transfer matrix. The optimum on-chip multiplier with minimum circuit area is then identified to be a Dickson charge pump.

### 2.3 Greinacher–Cockcroft–Walton (CW) Multiplier

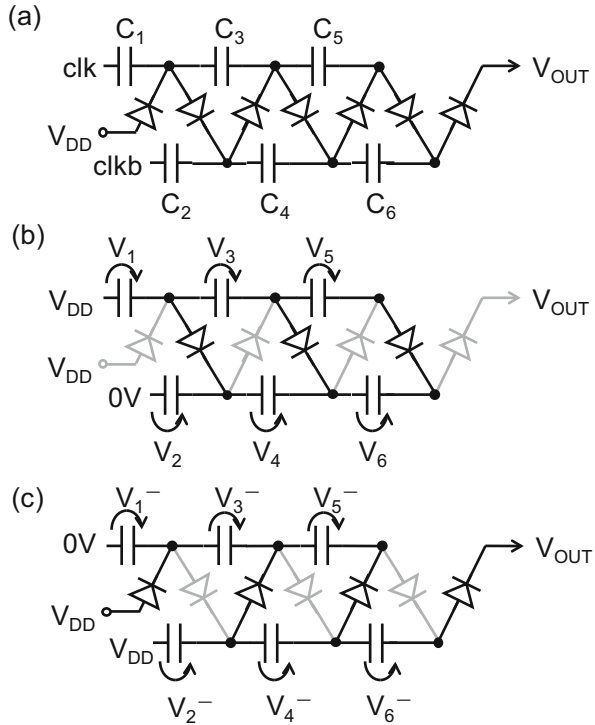
Figure 2.24 illustrates a serial ladder multiplier proposed by Greinacher and Cockcroft–Walton. The number of stage ( $N$ ) is defined by the number of capacitor. The number of diodes is  $N + 1$ . In Fig. 2.24,  $N$  is 6. Each half of them is serially connected and driven by complementary clocks  $\text{clk}$  or  $\text{clkb}$ . Figure 2.24, respectively, shows the first and second half periods. The clock has two voltage states with  $V_{DD}$  and 0 V. Capacitor voltages  $V_k$  and  $V_k^-$  ( $1 \leq i \leq 6$ ) are defined at the end of each half period. The following equations hold:

$$V_{DD} + V_1 = V_2 \quad (2.21)$$

$$V_{DD} + V_1 + V_3 = V_2 + V_4 \quad (2.22)$$

$$V_{DD} + \sum_{i=1}^j V_{2i-1} = \sum_{i=1}^j V_{2i} \quad (2.23)$$

**Fig. 2.24** Six-stage Greinacher–Cockcroft–Walton multiplier (Cockcroft and Walton 1932)



$$V_1^- = V_{DD} \tag{2.24}$$

$$V_{DD} + V_2^- = V_1^- + V_3^- \tag{2.25}$$

$$V_{DD} + \sum_{i=1}^j V_{2i}^- = \sum_{i=1}^j V_{2i-1}^- \tag{2.26}$$

Charge transferred through each diode in half period is same in steady state. The charge in  $C_1$  is transferred to  $C_2$ ,  $C_4$ , and  $C_6$  in the first half period in Fig. 2.24b, in total  $3q$ , when each transferred charge is written as  $q$ . Similarly, the charge in  $C_3$  is transferred to  $C_4$ , and  $C_6$ , and the charge in  $C_5$  is transferred to  $C_6$ . Charges of  $2q$  and  $1q$  are, respectively, discharged from  $C_3$  and  $C_5$ . In case where the number of stage is even  $N$ , the similar consideration results in Eqs. (2.27)–(2.29):

$$C_1 V_1 + \frac{Nq}{2} = C_1 V_1^- \tag{2.27}$$

$$C_3 V_3 + \left(\frac{N}{2} - 1\right)q = C_3 V_3^- \tag{2.28}$$

$$C_{2k-1}V_{2k-1} + \left(\frac{N}{2} - k + 1\right)q = C_{2k-1}V_{2k-1}^- \quad (2.29)$$

Similarly, the charge in  $C_2$  is transferred to  $C_3, C_5$ , and the output terminal in the second half period in Fig. 2.24c, in total  $3q$ . Thus,

$$C_2V_2 - \frac{Nq}{2} = C_2V_2^- \quad (2.30)$$

$$C_4V_4 - \left(\frac{N}{2} - 1\right)q = C_4V_4^- \quad (2.31)$$

$$C_{2k}V_{2k} - \left(\frac{N}{2} - k + 1\right)q = C_{2k}V_{2k}^- \quad (2.32)$$

From Eqs. (2.21)–(2.32),  $V_2^-, V_4^-, V_{2k}^-$  are calculated as Eqs. (2.33)–(2.35):

$$V_2^- = 2V_{DD} - \frac{Nq}{2C_1} - \frac{Nq}{2C_2} \quad (2.33)$$

$$V_4^- = 2V_{DD} - \frac{Nq}{2C_1} - \frac{Nq}{2C_2} - \left(\frac{N}{2} - 1\right)\frac{q}{C_3} - \left(\frac{N}{2} - 1\right)\frac{q}{C_4} \quad (2.34)$$

$$\begin{aligned} V_{2k}^- &= 2V_{DD} - \sum_{i=1}^k \left(\frac{N}{2} - i + 1\right)\frac{q}{C_{2i-1}} \\ &\quad - \sum_{i=1}^k \left(\frac{N}{2} - i + 1\right)\frac{q}{C_{2i}} \end{aligned} \quad (2.35)$$

The output voltage  $V_{OUT}$  is the sum of  $V_{DD}, V_2^-, V_4^-, \dots, V_N^-$  based on Fig. 2.24c resulting in Eq. (2.36):

$$\begin{aligned} V_{OUT} &= V_{DD} + \sum_{k=1}^{N/2} V_{2k}^- \\ &= (N+1)V_{DD} - \sum_{k=1}^{N/2} \left(\frac{N}{2} - k + 1\right)^2 \frac{q}{C_{2k-1}} \\ &\quad - \sum_{k=1}^{N/2} \left(\frac{N}{2} - k + 1\right)^2 \frac{q}{C_{2k}} \end{aligned} \quad (2.36)$$

The relation between  $I_{OUT}$  and  $V_{OUT}$  is calculated as Eqs. (2.37)–(2.39), where the cycle time is assume to be one:

$$I_{\text{OUT}} \equiv q = \frac{V_{\text{MAX}} - V_{\text{OUT}}}{R_{\text{PMP}}} \quad (2.37)$$

$$V_{\text{MAX}} = (N + 1)V_{\text{DD}} \quad (2.38)$$

$$R_{\text{PMP}}(C_1, C_2, \dots, C_N) = \sum_{j=1}^{N/2} \left(\frac{N}{2} - j + 1\right)^2 \frac{1}{C_{2j-1}} + \sum_{j=1}^{N/2} \left(\frac{N}{2} - j + 1\right)^2 \frac{1}{C_{2j}} \quad (2.39)$$

In case where the capacitance of all the capacitors is same as  $C_0 = C_{\text{TOT}}/N$ , where  $C_{\text{TOT}}$  is the total capacitance, Eq. (2.39) is rewritten as Eq. (2.40):

$$\begin{aligned} R_{\text{PMP}} &= \frac{N(N+1)(N+2)}{12} \frac{1}{C_0} \\ &= \frac{N^2(N+1)(N+2)}{12} \frac{1}{C_{\text{TOT}}} \end{aligned} \quad (2.40)$$

In case where the capacitance is weighted so that  $R_{\text{OUT}}$  is minimized under the condition that the total capacitance is constant, one can use Lagrange multiplier introducing functions  $f$  and  $g$ , and a parameter  $\lambda$  as follows:

$$f(C_1, C_2, \dots, C_N) \equiv \sum_{j=1}^N C_j - C_{\text{TOT}} = 0 \quad (2.41a)$$

$$g(C_1, C_2, \dots, C_N, \lambda) \equiv R_{\text{PMP}}(C_1, C_2, \dots, C_N) - \lambda f(C_1, C_2, \dots, C_N) \quad (2.41b)$$

$$\frac{\partial}{\partial C_{2j-1}} g(C_1, C_2, \dots, C_N) = -\left(\frac{N}{2} - j + 1\right)^2 \frac{1}{C_{2j-1}^2} - \lambda = 0 \quad (2.41c)$$

$$\frac{\partial}{\partial C_{2j}} g(C_1, C_2, \dots, C_N) = -\left(\frac{N}{2} - j + 1\right)^2 \frac{1}{C_{2j}^2} - \lambda = 0 \quad (2.41d)$$

Equations (2.41c) and (2.41d) hold when Eq. (2.42) holds:

$$C_1 : C_2 : C_3 : C_4 : \dots : C_{N-1} : C_N = \frac{N}{2} : \frac{N}{2} : \left(\frac{N}{2} - 1\right) : \left(\frac{N}{2} - 1\right) : \dots : 1 : 1 \quad (2.42)$$

Equation (2.39) results in Eq. (2.43):

$$R_{\text{PMP}} = \frac{N^2(N+2)^2}{16} \frac{1}{C_{\text{TOT}}} \quad (2.43)$$

With an effort optimizing each capacitor as Eq. (2.42), one can reduce the output resistance by a factor of about 25 %:

$$\frac{R_{\text{PMP\_unifrom\_C}}}{R_{\text{PMP\_weighted\_C}}} = \frac{4N + 1}{3N + 2} \quad (2.44)$$

Similarly, in case of odd  $N$ , Eqs. (2.37) and (2.38) hold, but Eqs. (2.39), (2.40), (2.43), and (2.44) are, respectively, replaced with Eqs. (2.45), (2.46), (2.47), and (2.48):

$$R_{\text{PMP}}(C_1, C_2, \dots, C_N) = \sum_{j=1}^{(N+1)/2} \left( \frac{N+1}{2} - j + 1 \right)^2 \frac{1}{C_{2j-1}} + \sum_{j=1}^{(N-1)/2} \left( \frac{N-1}{2} - j + 1 \right)^2 \frac{1}{C_{2j}} \quad (2.45)$$

$$R_{\text{PMP}} = \frac{(N+1)(N^2 + 2N + 3)}{12} \frac{1}{C_0} = \frac{N(N+1)(N^2 + 2N + 3)}{12} \frac{1}{C_{\text{TOT}}} \quad (2.46)$$

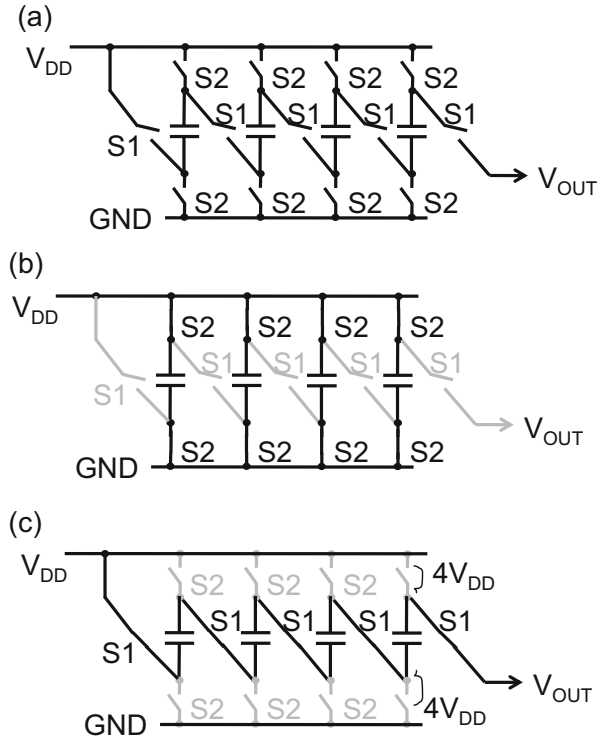
$$R_{\text{PMP}} = \frac{(N+1)^4}{16} \frac{1}{C_{\text{TOT}}} \quad (2.47)$$

$$\frac{R_{\text{PMP\_unifrom\_C}}}{R_{\text{PMP\_weighted\_C}}} = \frac{4N(N^2 + 2N + 3)}{3(N+1)^3} \quad (2.48)$$

## 2.4 Serial–Parallel (SP) Multiplier

Figure 2.25a shows a serial–parallel multiplier. Figure 2.25b, c shows how each capacitor is connected one another in each half period. All the capacitors are connected in parallel between the supply voltage  $V_{\text{DD}}$  and the ground in the first half period (b) and in series between  $V_{\text{DD}}$  and the output voltage  $V_{\text{OUT}}$  in the second half period (c). The capacitor voltage between two terminals of each capacitor is  $V_{\text{DD}}$  in the first half period and  $(V_{\text{OUT}} - V_{\text{DD}})/N$  in the second half period. When the charge transferred to the output terminal in a period is  $q$ ,  $T$  is the period, and  $C$  is capacitance of each capacitor, the output current  $I_{\text{OUT}}$  is given by the following:

**Fig. 2.25** (a) Serial-parallel multiplier, (b) in-parallel state, and (c) in-series state (Marx 1928; Brugler 1971; Mihara et al. 1995)



$$I_{OUT} = q/T = \frac{NC}{T} [(N + 1)V_{DD} - V_{OUT}] \tag{2.49}$$

$I_{OUT}$  is rewritten by

$$I_{OUT} = \frac{V_{MAX} - V_{OUT}}{R_{PMP}} \tag{2.50}$$

where

$$V_{MAX} = (N + 1)V_{DD} \tag{2.51}$$

$$R_{PMP} = \frac{TN}{C} \tag{2.52}$$

It is noted that the output resistance  $R_{PMP}$  is proportional to  $N^1$  in SP which is much less dependency on  $N$  than CW with a dependency of  $N^3$ . Since each capacitor needs to be charged by the same amount of  $q$  in the first period, the input current supplied by  $V_{DD}$  is given by

$$I_{DD} = (N + 1)I_{OUT} \tag{2.53}$$

When the maximum attainable voltage gain is defined by

$$G_V \equiv V_{MAX}/V_{DD} = N + 1 \tag{2.54}$$

The current efficiency is given by

$$\text{eff} \equiv I_{OUT}/I_{DD} = 1/G_V \tag{2.55}$$

Next, let us take parasitic capacitance into analysis. Four-stage SP pump is also expressed by Fig. 2.26a. Every stage is identical, thereby represented as  $\mathbf{K}(1)_{12}$ . Each stage has two operation states as shown in Fig. 2.26b, c, where  $C$  is the multiplier capacitor,  $C_T$  is the parasitic capacitance at one of the terminals of  $C$ , and  $C_B$  is the parasitic capacitance at the other terminal of  $C$ . In steady states, the following equations hold with the assumption that any parasitic resistance can be ignored:

$$V_{IN1} = V_{OUT1} \tag{2.56}$$

$$(I_{IN1} - I_{OUT1})T/2 = q_1 + C_T(V_{OUT1} - V_{OUT2}) \tag{2.57}$$

$$q_1 = C(V_{OUT1} - V_{OUT2} + V_{IN2}) \tag{2.58}$$

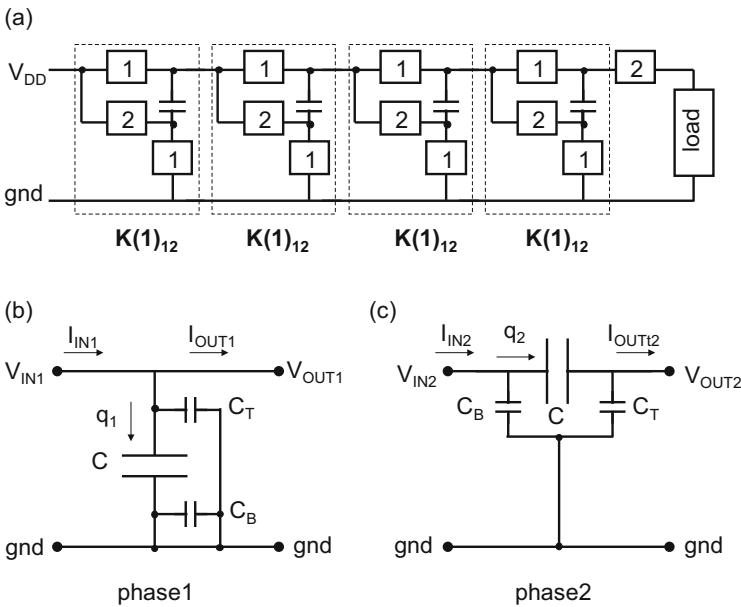


Fig. 2.26 Four-stage SP (a) and two alternate states (b), (c) of each stage

$$q_2 = I_{IN2}T/2 - C_B V_{IN2} = I_{OUT2}T/2 + C_T(V_{OUT2} - V_{OUT1}) \quad (2.59)$$

where  $q_1$  and  $q_2$  are the charge flowing into  $C$  in phases 1 and 2, respectively, and  $T$  is a cycle time. From the steady-state condition of  $q_1 = q_2$ , the  $K$ -matrix in case of 1 stage,  $\mathbf{K}(1)_{12}$ , is calculated as Eq. (2.60) based on Eqs. (2.56)–(2.59):

$$\mathbf{K}(1)_{12} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ -(1 + \alpha_T) & (1 + \alpha_T) & 0 & R \\ 0 & 0 & 1 & 1 \\ -1/r_T - (1 + \alpha_T)/r_B & 1/r_T + (1 + \alpha_T)/r_B & 0 & (1 + \alpha_B) \end{bmatrix} \quad (2.60)$$

where  $\alpha_i = C_i/C$ ,  $r_i = T/2C_i$  ( $i = T, B$ ), and  $R = T/2C$ .

As shown by Eq. (2.20), the entire  $K$ -matrix of  $N$ -stage SP multiplier is calculated by multiplying  $\mathbf{K}(1)_{12}$  by  $N$ -times, resulting in Eq. (2.61):

$$\mathbf{K}_{SP}(N) = (\mathbf{K}(1)_{12})^N \quad (2.61)$$

From Fig. 2.23 and Eq. (2.19), the output current of SP multiplier  $I_{OUT}$  is  $I_{OUT2}/2$  since the averaged time of period for  $I_{OUT}$  is twice as long as that for  $I_{OUT2}$ , the output voltage  $V_{OUT}$  is  $V_{OUT2}$ ,  $I_{OUT1} = 0$ , and  $V_{IN1} = V_{IN2} = V_{DD}$ . Thus, the following equation holds:

$$[V_{DD} \ V_{DD} \ I_{IN1} \ I_{IN2}]^T = \mathbf{K}_{SP}(N)[V_{OUT1} \ V_{OUT} \ 0 \ 2I_{OUT}]^T \quad (2.62)$$

The relation between  $V_{OUT}$  and  $I_{OUT}$  is calculated by the first and second rows of Eq. (2.62) by eliminating  $V_{OUT1}$ . The total current consumption  $I_{IN}$  is calculated by  $(I_{IN1} + I_{IN2})/2$  with certain values of  $V_{OUT}$  and  $I_{OUT}$ . One can easily calculate the output voltage–current characteristics and the current consumption or efficiency with the circuit parameters, such as  $C$ ,  $C_T$ ,  $C_B$ ,  $T$ ,  $N$ , given by using a simple matrix calculator Eq. (2.62).

## 2.5 Falkner-Dickson Linear (LIN) Multiplier

Figure 2.27 illustrates the Dickson charge pump circuit. A charge pump with an even number of stages is considered in this subsection, but a similar analysis in the case of an odd number stage charge pump can be carried out.  $q$  is defined as the charge transferred from one capacitor to the next one during one cycle, and  $Q_i$  ( $1 \leq i \leq N$ ) are defined as the charges stored in the capacitors  $C_i$  at time  $j$ . Figure 2.28 illustrates connection of the first stage with the input terminal and connection between the second and third stages at time  $j$  (a), connection between the first and second stages at time  $j + 1/2$  (b), and connection of the last stage with  $V_{OUT}$  at time  $j$  (c).

From Fig. 2.28a, the following relations hold under the condition that the diode  $D_1$  is cut off at time  $j$ :



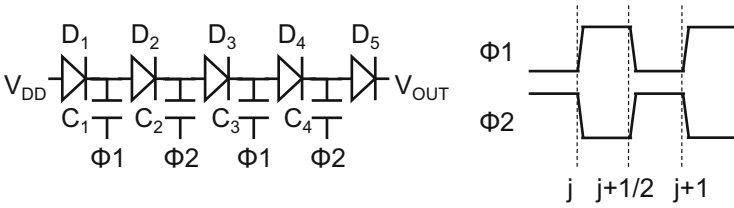
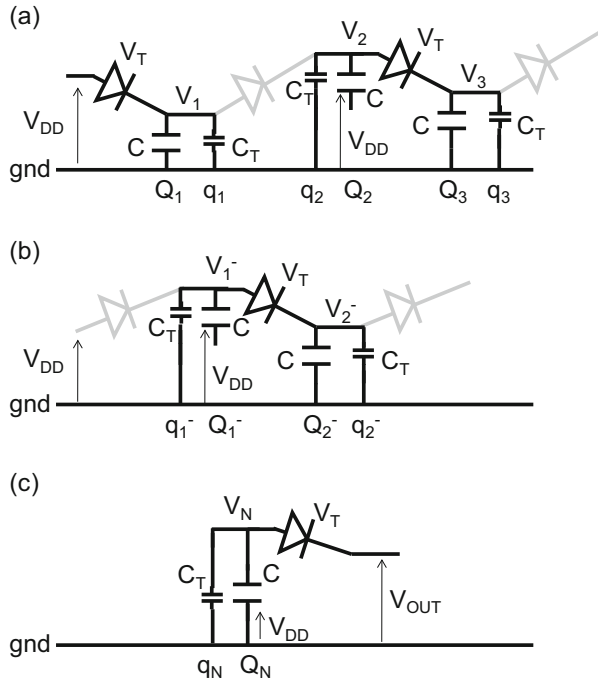


Fig. 2.27 Four-stage Dickson pump (Falkner 1973, Dickson 1976a, b)

Fig. 2.28 Relations between next neighbors



$$V_1 = V_{DD} - V_T \tag{2.63a}$$

$$Q_1 = C(V_{DD} - V_T) \tag{2.63b}$$

$$q_1 = C_T(V_{DD} - V_T) \tag{2.63c}$$

where  $V_{DD}$  is the supply voltage and  $V_T$  the subthreshold voltage. The difference between the total amount of charge stored in  $C$  and  $C_T$  at the first stage at time  $j$  and that at time  $j + 1/2$  is  $q$  under the steady-state condition:

$$(Q_1 + q_1) - (Q_1^- + q_1^-) = q \quad (2.64a)$$

$$Q_1^- = C(V_1^- - V_{DD}) \quad (2.64b)$$

$$q_1^- = C_T V_1^- \quad (2.64c)$$

From Eqs. (2.63b), (2.63c), (2.64a), (2.64b), and (2.64c),

$$V_1^- = (V_{DD} - V_T) + \frac{V_{DD}}{1 + \alpha_T} - \frac{q}{C + C_T} \quad (2.65)$$

Similarly,

$$(Q_2^- + q_2^-) - (Q_2 + q_2) = q \quad (2.66a)$$

$$V_2 = Q_2/C + V_{DD} = q_2/C_T \quad (2.66b)$$

$$V_2^- = Q_2^-/C = q_2^-/C_T \quad (2.66c)$$

From Eqs. (2.66a), (2.66b), and (2.66c),

$$V_2^- = \frac{Q_2}{C} + \frac{\alpha_T V_{DD}}{1 + \alpha_T} + \frac{q}{C + C_T} \quad (2.67)$$

From the condition Eq. (2.68) that the diode  $D_2$  is cut off at time  $j + 1/2$ ,

$$V_1^- - V_2^- = V_T \quad (2.68)$$

and Eqs. (2.65) and (2.67),

$$Q_2 = C \left( \frac{2V_{DD}}{1 + \alpha_T} - 2V_T \right) - \frac{2q}{1 + \alpha_T} \quad (2.69)$$

Similarly, from the condition that the diode  $D_3$  is cut off at time  $j$ , and Eqs. (2.66b) and (2.69),

$$Q_3 = C \left( \frac{2}{1 + \alpha_T} + 1 \right) V_{DD} - 3CV_T - \frac{2q}{1 + \alpha_T} \quad (2.70)$$

Repeating the similar procedure, general  $Q(2k - 1)$  and  $Q(2k)$  are calculated to be as follows:

$$Q_{2k-1} = \left( \frac{2(k-1)}{1 + \alpha_T} + 1 \right) CV_{DD} - (2k-1)CV_T - \frac{2(k-1)q}{1 + \alpha_T} \quad (2.71a)$$

$$Q_{2k} = 2kC \left( \frac{V_{DD}}{1 + \alpha_T} - V_T - \frac{1}{1 + \alpha_T} \frac{q}{C} \right) \quad (2.71b)$$

From Fig. 2.28c, the following relation holds under the condition that the diode  $D_{N+1}$  is cut off at time  $j$ :

$$V_{DD} + \frac{Q_N}{C} - V_T = V_{OUT} \quad (2.72)$$

From Eqs. (2.71b) and (2.72), the output voltage–current characteristic, using  $2k = N$ ,

$$q = \frac{C + C_T}{N} \left[ \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1)V_T - V_{OUT} \right] \quad (2.73)$$

From Eqs. (2.71a), (2.71b), and (2.73), the charge stored in each charge pump capacitor is represented by

$$Q_{2k-1} = \frac{2(k-1)}{N} C(V_{OUT} - V_{DD} + V_T) + C(V_{DD} - V_T) \quad (2.74a)$$

$$Q_{2k} = \frac{2k}{N} C(V_{OUT} - V_{DD} + V_T) \quad (2.74b)$$

The output current  $I_{OUT}$  is given by the following:

$$I_{OUT} \equiv \frac{q}{T} = \frac{C + C_T}{NT} \left[ \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1)V_T - V_{OUT} \right] \quad (2.75)$$

where  $T$  is the clock period. Equation (2.75) was originally derived by Dickson in 1978.  $I_{OUT}$  is rewritten by

$$I_{OUT} = \frac{V_{MAX} - V_{OUT}}{R_{PMP}} \quad (2.76)$$

where

$$V_{MAX} = V_{DD} + N \left( \frac{V_{DD}}{1 + \alpha_T} - V_T \right) - V_T \quad (2.77)$$

$$R_{PMP} = \frac{NT}{C(1 + \alpha_T)} \quad (2.78)$$

$V_{MAX}$  is considered as the sum of the initial voltage input  $V_{DD}$ ,  $N$  stages' voltage gain, each of which is  $V_{DD}/(1 + \alpha_T) - V_T$ , and the voltage drop in  $D_{N+1}$ . It is noted that the output resistance  $R_{PMP}$  and the maximum attainable voltage  $V_{MAX}$  are same

as those of the serial–parallel pump as shown in Sect. 2.4 in the case of  $\alpha_T = V_T = 0$ . Similar to SP, since each capacitor needs to be charged by the same amount of  $q$  in the first period, the input current supplied by  $V_{DD}$  is given by

$$I_{DD} = (N + 1)I_{OUT} \quad (2.79)$$

in the ideal case where  $\alpha_T = \alpha_B = 0$ . When the maximum attainable voltage gain is defined by

$$G_V \equiv V_{MAX}/V_{DD} = N + 1 \quad (2.80)$$

the current efficiency is given by

$$\text{eff} = \frac{1}{G_V} \quad (2.81)$$

Next, another procedure using  $K$ -matrix is discussed.

Figure 2.29 shows another expression of the Dickson linear multiplier with four stages. Because each stage has one input and one output in a period, as shown in Fig. 2.30, one can simply express the input and output voltages (currents) as  $V_{IN}$  and  $V_{OUT}$  ( $I_{IN}$  and  $I_{OUT}$ ) without a suffix of 1 or 2. Then, the following equations hold in case of  $V_T = 0$  in phase 2:

$$I_{IN}T = q + C_T(V_{IN} - V_{OUT}) \quad (2.82)$$

$$q = C(V_{IN} - V_{OUT} + V_{DD}) \quad (2.83)$$

and in phase 1,

$$q = I_{OUT}T + C_T(V_{OUT} - V_{IN}) \quad (2.84)$$

$$I_{DD}(1)T = I_{OUT}T + C_T(V_{OUT} - V_{IN}) + C_B V_{DD} \quad (2.85)$$

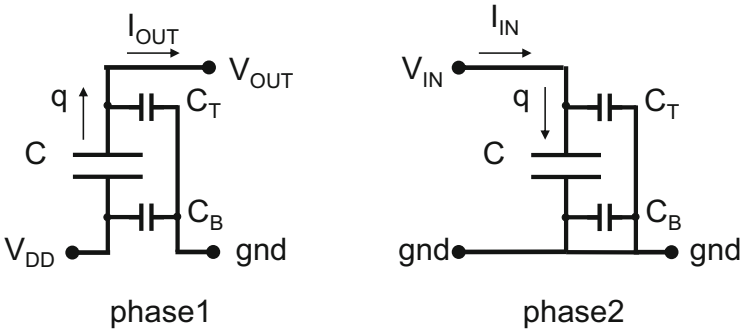
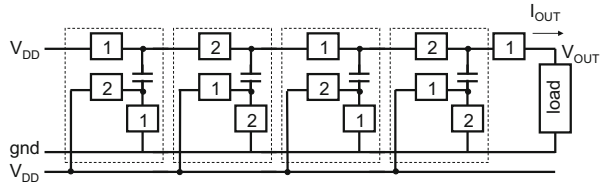
where  $I_{DD}(1)$  is the current supplied by  $V_{DD}$  per stage. From Eqs. (2.82)–(2.84),

$$V_{OUT} - V_{IN} = \frac{V_{DD}}{1 + \alpha_T} - \frac{TI_{OUT}}{C(1 + \alpha_T)} \quad (2.86)$$

Since Eq. (2.86) represents the voltage gain per stage, the total voltage gain of  $N$ -stage multiplier is given by Eq. (2.87), which is the same as Dickson's result Eq. (2.75) in the case of  $V_T = 0$ :

$$V_{OUT}(N) = \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - \frac{TN I_{OUT}}{C(1 + \alpha_T)} \quad (2.87)$$

**Fig. 2.29** Linear multiplier with four stages



**Fig. 2.30** Two alternate states of the linear multiplier

The total input current from the voltage supply into  $N$ -stage multiplier,  $I_{DD}(N)$ , is calculated with Eq. (2.85) by multiplying  $N$  and by adding one  $I_{IN}$  from the input of the first stage:

$$I_{DD}(N) = NI_{DD}(1) + I_{IN} \tag{2.88}$$

From Eqs. (2.82), (2.84), (2.85), and (2.88),

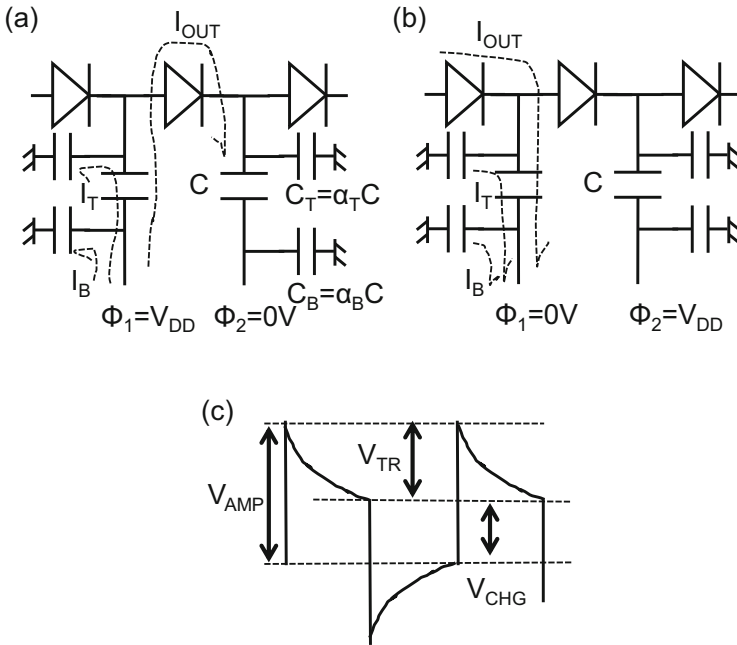
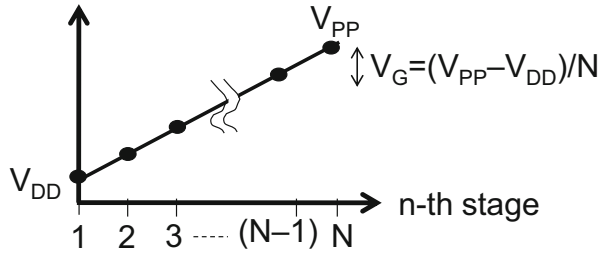
$$I_{DD}(N) = (N + 1)I_{OUT} + C_T(V_{OUT}(N) - V_{DD})/T + NC_B V_{DD}/T \tag{2.89}$$

Thus, the relationship between the output voltage and the output current and between the input and output currents of linear multiplier don't require matrix calculations, but are analytically resolved as Eqs. (2.87) and (2.89), respectively.

The meaning of  $R_{PMP}$  in Eq. (2.78) is considered. Figure 2.31 illustrates the averaged voltage at each stage of  $N$ -stage Dickson pump. The difference voltage  $V_G$  between the next neighbor stages is  $(V_{PP} - V_{DD})/N$ . When  $V_{PP}$  is increased by  $\Delta V_{PP}$ ,  $V_G$  is increased by  $\Delta V_G = \Delta V_{PP}/N$  and the output charge  $Q$  is decreased by  $\Delta Q$ . These two are related via  $\Delta Q = C\Delta V_G$ . The output resistance is defined by  $\Delta V_{PP}/(\Delta Q/T)$ , resulting in  $NT/C$ .

Three components in  $I_{DD}$  given by Eq. (2.89) can be identified as follows. Figure 2.32 shows the input current components of (1) the current from a pump capacitor to the next one which is same as the output current in steady-state  $I_{OUT}$ , (2) the charging current to the parasitic capacitance at the top place ( $C_T = \alpha_T C$ ) of each pump capacitor  $I_T$ , and (3) the charging current to the parasitic capacitance at

**Fig. 2.31** Averaged voltage at each stage of  $N$ -stage Dickson pump



**Fig. 2.32** Two neighbor stages of a charge pump in a first (a) and second (b) half period, and a voltage waveform at the top plate of a pump capacitor in a steady state (c)

the bottom place ( $C_B = \alpha_B C$ ) of each pump capacitor  $I_B$ , where  $C$  is the capacitance of the pump capacitor. These current components flow from the power supply  $V_{DD}$  in a half period (a) and flow to the ground in another half period (b). At the clock edge, the top plate node of each capacitor has the amplitude given by Eq. (2.90):

$$V_{AMP} = V_{DD} / (1 + \alpha_T) \tag{2.90}$$

The capacitor voltage is reduced by

$$V_{TR} = T I_{OUT} / (C + C_T) \tag{2.91}$$

due to a charge transfer of  $TI_{\text{OUT}}$ . Thus, the voltage amplitude  $V_{\text{CHG}}$  between the beginning and end of the clock high and the charging current  $I_T$  to the top-plate parasitic capacitance are given by

$$V_{\text{CHG}} = V_{\text{AMP}} - V_{\text{TR}} = (V_{\text{PP}} - V_{\text{DD}} + (N + 1)V_T)/N \quad (2.92a)$$

$$I_T = C_T V_{\text{CHG}}/T \quad (2.92b)$$

where the  $V_{\text{OUT}} - I_{\text{OUT}}$  relation Eq. (2.75) is used. The current charging the bottom plate parasitic capacitance is given by Eq. (2.93):

$$I_B = C_B V_{\text{DD}}/T \quad (2.93)$$

As a result, the total input current of  $N$ -stage pump is

$$\begin{aligned} I_{\text{DD}} &= (N + 1)I_{\text{PP}} + NI_T + NI_B \\ &= (N + 1)I_{\text{PP}} + \alpha_T C(V_{\text{PP}} - V_{\text{DD}} + (N + 1)V_T)/T \\ &\quad + N\alpha_B C_B V_{\text{DD}}/T \end{aligned} \quad (2.94)$$

Equation (2.94) in case of  $V_T = 0$  V is equivalent to Eq. (2.89).

Figure 2.33a illustrates a three stage linear pump operating with a single phase clock. Two of three stages contribute to charge pumping, resulting in a lower maximum attainable voltage than a two phase clock pump with the same number of stages, as shown by Eq. (2.95a) where  $V_{\text{MAX}_{1-\phi}}$  is the voltage amplitude of the single clock and  $V_{\text{MAX}_{2-\phi}}$  is the voltage amplitude of the two-phase clock. However, the output impedance is the same because Fig. 2.31 is valid regardless of the number of phases, as shown by Eq. (2.95b). As a result, the single-phase clock pump has the  $V_{\text{OUT}} - I_{\text{OUT}}$  line as shown in Fig. 2.33b in comparison with that of the two-phase clock pump:

$$V_{\text{MAX}_{1-\phi}} = V_{\text{DD}}(N + 1)/2 = V_{\text{MAX}_{2-\phi}}/2 \quad (2.95a)$$

$$R_{\text{PMP}} = \frac{T}{C} N \quad (2.95b)$$

## 2.6 Fibonacci (FIB) Multiplier

This section starts with zero parasitic capacitance and then analyzes the Fibonacci multiplier with a finite parasitic capacitance.

Figure 2.34a illustrates a Fibonacci multiplier with four stages, which work with a two-phase clock. The squares show switches, and the numbers 1 and 2 inside indicate turning on in phases 1 and 2 of the clock, respectively. The two-port transfer matrix  $\mathbf{K}(N)$  is again defined by Eq. (2.96):

$$[V_{\text{IN1}} \ V_{\text{IN2}} \ I_{\text{IN1}} \ I_{\text{IN2}}]^T = \mathbf{K}(N)[V_{\text{OUT1}} \ V_{\text{OUT2}} \ I_{\text{OUT1}} \ I_{\text{OUT2}}]^T \quad (2.96)$$

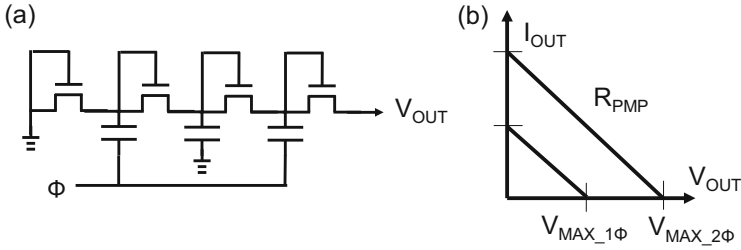


Fig. 2.33 Linear pump operating with a single-phase clock (a) and its  $V_{OUT} - I_{OUT}$  line (b)

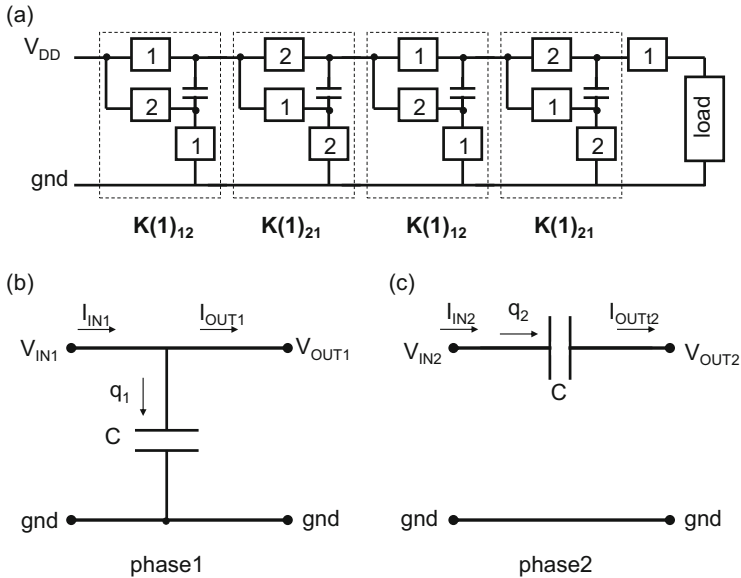


Fig. 2.34 Four-stage Fibonacci multiplier (Harada et al. 1992a, b)

Each stage of the multiplier has two operation states as shown in Fig. 2.34b, c, where  $C$  is the multiplier capacitor. In steady states, the following equations hold with the assumption that any parasitic resistance and capacitance can be ignored:

$$V_{IN1} = V_{OUT1} \tag{2.97}$$

$$(I_{IN1} - I_{OUT1})T/2 = q_1 \tag{2.98}$$

$$q_1 = C(V_{OUT1} - V_{OUT2} + V_{IN2}) \tag{2.99}$$

$$q_2 = I_{IN2}T/2 = I_{OUT2}T/2 \tag{2.100}$$



From the steady-state condition of  $q_1 = q_2$ , the  $K$ -matrix in case of 1 stage,  $\mathbf{K}(1)_{12}$  is calculated as Eq. (2.101):

$$\mathbf{K}(1)_{12} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ -1 & 1 & 0 & R \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad (2.101)$$

where  $R = T/2C$ . The  $K$ -matrix of Ueno Fibonacci multiplier for even stages as shown in Fig. 2.34a,  $\mathbf{K}(1)_{21}$ , is given by simply exchanging the suffix 1 with 2 for  $V_{IN}$  and  $V_{OUT}$  of Eqs. (2.97)–(2.100):

$$\mathbf{K}(1)_{21} = \begin{bmatrix} 1 & -1 & R & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 1 \end{bmatrix} \quad (2.102)$$

As shown in Fig. 2.34a, the entire  $K$ -matrix of  $N$ -stage UF multiplier is calculated by Eq. (2.103) in the case of even number of stages and by Eq. (2.104) in the case of odd number of stages, respectively:

$$\mathbf{K}_{UF}(2n) = (\mathbf{K}(1)_{12}\mathbf{K}(1)_{21})^n \quad (2.103)$$

$$\mathbf{K}_{UF}(2n+1) = (\mathbf{K}(1)_{12}\mathbf{K}(1)_{21})^n \mathbf{K}(1)_{12} \quad (2.104)$$

For Ueno Fibonacci multiplier with odd number of stages, the output current  $I_{OUT}$  is  $I_{OUT2}/2$  since the averaged time of period for  $I_{OUT}$  is twice as long as that for  $I_{OUT2}$ , the output voltage  $V_{OUT}$  is  $V_{OUT2}$ ,  $I_{OUT1} = 0$ , and  $V_{IN1} = V_{IN2} = V_{DD}$ . Thus, the following equation holds for Ueno Fibonacci multiplier with odd number of stages:

$$\begin{bmatrix} V_{DD} & V_{DD} & I_{IN1} & I_{IN2} \end{bmatrix}^T = \mathbf{K}_{UF}(2n+1) \begin{bmatrix} V_{OUT1} & V_{OUT} & 0 & 2I_{OUT} \end{bmatrix}^T \quad (2.105)$$

The relation between  $V_{OUT}$  and  $I_{OUT}$  is calculated by the first and second rows of Eq. (2.105) by eliminating  $V_{OUT1}$ . The total current consumption  $I_{IN}$  is calculated by  $(I_{IN1} + I_{IN2})/2$  with certain values of  $V_{OUT}$  and  $I_{OUT}$ . For the UF multiplier with even number of stages, it is valid when the conditions of  $I_{OUT} = I_{OUT1}/2$ ,  $V_{OUT} = V_{OUT1}$ , and  $I_{OUT2} = 0$  are used instead. Thus,

$$\begin{bmatrix} V_{DD} & V_{DD} & I_{IN1} & I_{IN2} \end{bmatrix}^T = \mathbf{K}_{UF}(2n) \begin{bmatrix} V_{OUT} & V_{OUT2} & 2I_{OUT} & 0 \end{bmatrix}^T \quad (2.106)$$

One can easily calculate the output voltage–current characteristics and the current consumption or efficiency with the circuit parameters, such as  $C$ ,  $T$ , and  $N$ , given by using a simple matrix calculator Eqs. (2.105) or (2.106):

**Table 2.3**  $V_{\text{MAX}}$  and  $R_{\text{PMP}}$  as a function of the number of stage  $N$

$N$	$V_{\text{MAX}}(N)V_{\text{DD}}$	$R_{\text{PMP}}(N)$
1	2	$1^2/C_1$
2	3	$1^2/C_2 + 1^2/C_1$
3	5	$1^2/C_3 + 1^2/C_2 + 2^2/C_1$
4	8	$1^2/C_4 + 1^2/C_3 + 2^2/C_2 + 3^2/C_1$
$N$	$F(n+1)$	$\sum_{j=0}^{n-1} F(j)^2 / C_{N-j}$

$$I_{\text{OUT}} = \frac{V_{\text{MAX}} - V_{\text{OUT}}}{R_{\text{PMP}}} \quad (2.107)$$

where  $V_{\text{IN}1,2}$  is  $V_{\text{DD}}$  and  $F(j)$  is  $j$ th Fibonacci number and  $F(0) = F(1) = 1$ ,  $F(j+2) = F(j+1) + F(j)$ . In case where the capacitance of each capacitor is not same among  $N$  capacitors,  $R_{\text{PMP}}$  is generally written as

$$R_{\text{PMP}}(N) = \sum_{j=0}^{N-1} \frac{F(j)^2}{C_{N-j}} \quad (2.108)$$

Table 2.3 summarizes the characteristic parameters of UF multipliers.

Under the condition that the total capacitor area is given, optimum distribution exists so that  $R_{\text{PMP}}$  is minimized. One can use Lagrange multiplier introducing functions  $f$  and  $g$ , and a parameter  $\lambda$  as follows, where  $C_{\text{TOT}}$  is the total capacitance of  $N$  capacitors:

$$f(C_1, C_2, \dots, C_N) \equiv \sum_{j=1}^N C_j - C_{\text{TOT}} = 0 \quad (2.109a)$$

$$g(C_1, C_2, \dots, C_N, \lambda) \equiv R_{\text{PMP}}(C_1, C_2, \dots, C_N) - \lambda f(C_1, C_2, \dots, C_N) \quad (2.109b)$$

$$\frac{\partial}{\partial C_j} g(C_1, C_2, \dots, C_N) = - \left( \frac{F(N-j)}{C_j} \right)^2 - \lambda = 0 \quad (2.109c)$$

Equation (2.109c) holds when Eq. (2.110) holds:

$$C_1 : C_2 : C_3 : C_4 : \dots : C_{N-1} : C_N = F(N-1) : F(N-2) : \dots : F(0) \quad (2.110)$$

$$R_{\text{PMP}}(N) = \frac{N}{C_{\text{TOT}}} \left( \sum_{j=0}^{N-1} F(j) \right)^2 \quad (2.111)$$

When the maximum attainable voltage gain is defined by

$$G_V \equiv V_{\text{MAX}}/V_{\text{DD}} = F(N + 1) \quad (2.112)$$

the current efficiency is given by

$$\text{eff} = \frac{1}{G_V} \quad (2.113)$$

$$I_{\text{DD}} = G_V I_{\text{OUT}} \quad (2.114)$$

Next, the circuit analysis is done in case where the parasitic capacitance is considered.

Each stage has two operation states as shown in Fig. 2.35b, c, where  $C$  is the multiplier capacitor,  $C_T$  is the parasitic capacitance at one of the terminals of  $C$ , and  $C_B$  is the parasitic capacitance at the other terminal of  $C$ . In steady states, the following equations hold with the assumption that any parasitic resistance can be ignored:

$$V_{\text{IN1}} = V_{\text{OUT1}} \quad (2.115)$$

$$(I_{\text{IN1}} - I_{\text{OUT1}})T/2 = q_1 + C_1(V_{\text{OUT1}} - V_{\text{OUT2}}) \quad (2.116)$$

$$q_1 = C(V_{\text{OUT1}} - V_{\text{OUT2}} + V_{\text{IN2}}) \quad (2.117)$$

$$q_2 = I_{\text{IN2}}T/2 - C_2V_{\text{IN2}} = I_{\text{OUT2}}T/2 + C_1(V_{\text{OUT2}} - V_{\text{OUT1}}) \quad (2.118)$$

where  $q_1$  and  $q_2$  are the charge flowing into  $C$  in phase 1 and 2, respectively. From the steady-state condition of  $q_1 = q_2$ , the  $K$ -matrix in case of one stage,  $\mathbf{K}_P(1)_{12}$  is calculated as Eq. (2.119):

$$\mathbf{K}_P(1)_{12} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ -(1 + \alpha_T) & (1 + \alpha_T) & 0 & R \\ 0 & 0 & 1 & 1 \\ -1/r_T - (1 + \alpha_T)/r_B & 1/r_T + (1 + \alpha_T)/r_B & 0 & (1 + \alpha_B) \end{bmatrix} \quad (2.119)$$

where  $\alpha_i = C_i/C$ ,  $r_i = T/2C_i$  ( $i = T, B$ ), and  $R = T/2C$ . The  $K$ -matrix of Fibonacci multiplier with even stages as shown in Fig. 2.35a,  $\mathbf{K}_P(1)_{21}$ , is given by simply exchanging the suffix 1 with 2 for  $V_{\text{IN}}$  and  $V_{\text{OUT}}$  of Eqs. (2.115)–(2.118):

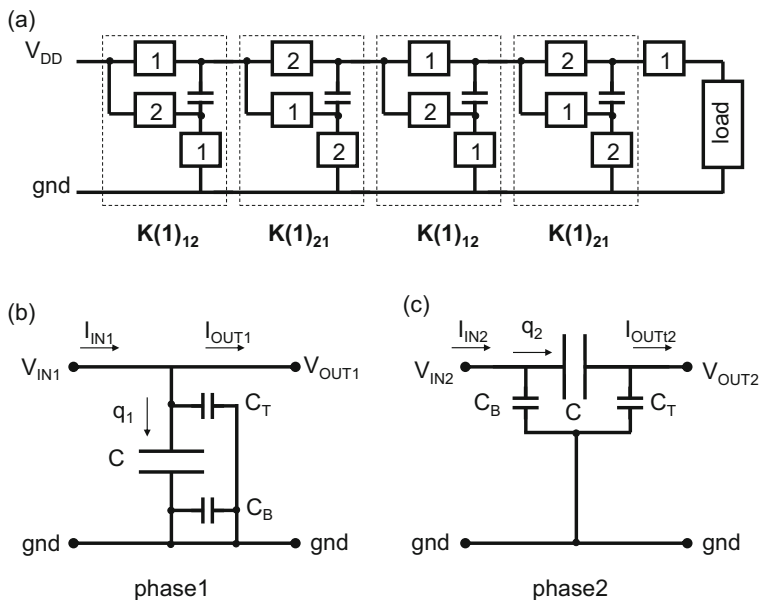


Fig. 2.35 Four-stage Fibonacci multiplier with a finite parasitic capacitance considered

$$\mathbf{K}_P(1)_{21} = \begin{bmatrix} (1 + \alpha_T) & -(1 + \alpha_T) & R & 0 \\ 0 & 1 & 0 & 0 \\ 1/r_T + (1 + \alpha_T)/r_B & -1/r_T - (1 + \alpha_T)/r_B & (1 + \alpha_B) & 0 \\ 0 & 0 & 1 & 1 \end{bmatrix} \quad (2.120)$$

As shown in Fig. 2.35a, entire  $K$ -matrix of  $N$ -stage FIB multiplier is calculated by Eq. (2.121) in the case of even number of stages and by Eq. (2.122) in the case of odd number of stages, respectively:

$$\mathbf{K}_P(2n) = (\mathbf{K}_P(1)_{12} \mathbf{K}_P(1)_{21})^n \quad (2.121)$$

$$\mathbf{K}_P(2n + 1) = (\mathbf{K}_P(1)_{12} \mathbf{K}_P(1)_{21})^n \mathbf{K}_P(1)_{12} \quad (2.122)$$

For the UF with odd number of stages, the output current  $I_{OUT}$  is  $I_{OUT2}/2$  since the averaged time of period for  $I_{OUT}$  is twice as long as that for  $I_{OUT2}$ , the output voltage  $V_{OUT}$  is  $V_{OUT2}$ ,  $I_{OUT1} = 0$ , and  $V_{IN1} = V_{IN2} = V_{DD}$ . Thus, the following equation holds for the UF with odd number of stages:

$$[V_{DD} \quad V_{DD} \quad I_{IN1} \quad I_{IN2}]^T = \mathbf{K}_P(2n + 1)[V_{OUT1} \quad V_{OUT} \quad 0 \quad 2I_{OUT}]^T \quad (2.123)$$

The relation between  $V_{OUT}$  and  $I_{OUT}$  is calculated by the first and second row of Eq. (2.123) by eliminating  $V_{OUT1}$ . The total current consumption  $I_{IN}$  is calculated

by  $(I_{IN1} + I_{IN2})/2$  with certain values of  $V_{OUT}$  and  $I_{OUT}$ . For the UF multiplier with even number of stages, it is valid when the conditions of  $I_{OUT} = I_{OUT1}/2$ ,  $V_{OUT} = V_{OUT1}$ , and  $I_{OUT2} = 0$  are used instead. Thus,

$$[V_{DD} \ V_{DD} \ I_{IN1} \ I_{IN2}]^T = \mathbf{K}_P(2n)[V_{OUT} \ V_{OUT2} \ 2I_{OUT} \ 0]^T \quad (2.124)$$

## 2.7 $2^N$ Multiplier

This section starts with zero parasitic capacitance and then analyzes the  $2^N$  multiplier with a finite parasitic capacitance.

Figure 2.36 shows the one with four stages. Figure 2.37 illustrates two alternate states. Equations (2.125) and (2.126) hold in phase 1:

$$I_{IN1}T/2 = C_a(V_{IN1} - V_{OUT2} + V_{IN2}) + C_b(V_{IN1} - V_{OUT1} + V_{IN2}) \quad (2.125)$$

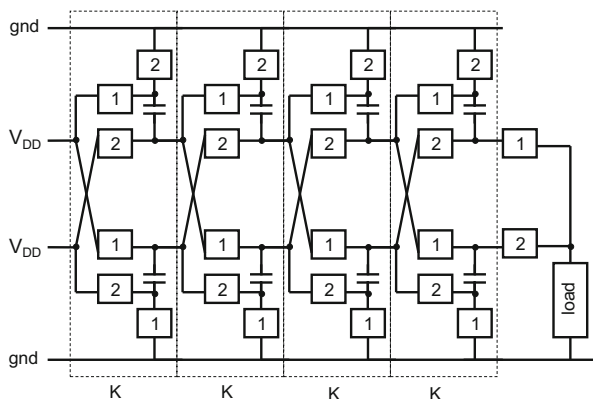
$$I_{OUT1}T/2 = C_b(V_{IN2} - V_{OUT1} + V_{IN1}) \quad (2.126)$$

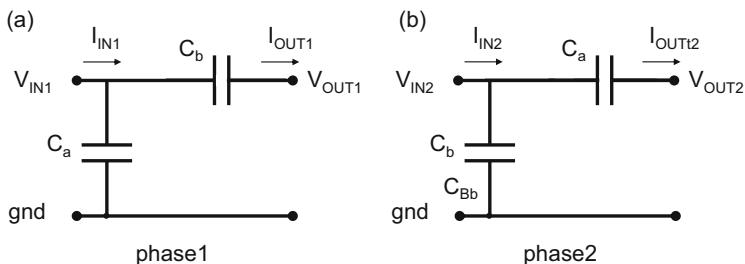
When  $C_a = C_b = C/2$ , phases 1 and 2 are identical. In this case, the input and output voltage (current) can be written by  $V_{IN}$  and  $V_{OUT}$  ( $I_{IN}$  and  $I_{OUT}$ ) without differentiating the two states. Using this symmetry between phases 1 and 2,  $K$ -matrix can be reduced to  $2 \times 2$ . The matrix for  $j$ th stage is given by

$$\tilde{\mathbf{K}}(j) = \begin{bmatrix} \frac{1}{2} & R(j) \\ 0 & 2 \end{bmatrix} \quad (2.127)$$

$$R(j) = \frac{1}{2C(j)} \quad (2.128)$$

**Fig. 2.36** Four-stage  $2^N$  multiplier (Cernea 1995)





**Fig. 2.37** Two phases of  $2^N$  multiplier with no parasitic capacitance

where  $T$  is assumed to be one. The matrix for  $n$ -stage pump  $\mathbf{K}(n)$  is written by

$$\mathbf{K}(n) = \mathbf{K}(n-1)\tilde{\mathbf{K}}(n) \quad (2.129)$$

$$\mathbf{K}(n) = \begin{bmatrix} a(n) & e(n) \\ b(n) & d(n) \end{bmatrix} \quad (2.130)$$

From Eqs. (2.127), (2.129), and (2.130),

$$\begin{aligned} \begin{bmatrix} a(n) & e(n) \\ b(n) & d(n) \end{bmatrix} &= \begin{bmatrix} a(n-1) & e(n-1) \\ b(n-1) & d(n-1) \end{bmatrix} \begin{bmatrix} \frac{1}{2} & R(n) \\ 0 & 2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{1}{2}a(n-1) & R(n)a(n-1) + 2e(n-1) \\ \frac{1}{2}b(n-1) & R(n)b(n-1) + 2d(n-1) \end{bmatrix} \end{aligned} \quad (2.131)$$

Using the initial condition,

$$\begin{bmatrix} a(1) & e(1) \\ b(1) & d(1) \end{bmatrix} = \begin{bmatrix} \frac{1}{2} & R(1) \\ 0 & 2 \end{bmatrix} \quad (2.132)$$

The components are solved as

$$a(n) = 2^{-n} \quad (2.133)$$

$$b(n) = 0 \quad (2.134)$$

$$e(n) = 2^{-n+1}R(n) + 2e(n-1) \quad (2.135)$$

$$d(n) = 2^n \quad (2.136)$$

From Eqs. (2.132) and (2.135),  $e(n)$  is solved as

$$e(n) = \sum_{j=1}^n 2^{n-2j+1}R(j) \quad (2.137)$$

The output resistance of the pump is calculated as

$$R_{\text{PMP}}(N) = \frac{e(N)}{a(N)} = \sum_{j=1}^N 2^{2N-2j} R(j) = \sum_{j=1}^N \frac{2^{2N-2j}}{C(j)} \quad (2.138)$$

In case where all capacitors are equivalently divided, i.e.,

$$C(j) = \frac{C_{\text{TOT}}}{N} \quad (2.139a)$$

$R_{\text{PMP}}$  is given by

$$R_{\text{PMP}}(N) = \frac{N}{C_{\text{TOT}}} \sum_{j=1}^N 2^{2N-2j} = \frac{N}{C_{\text{TOT}}} \left( \frac{2}{3} N(N-1)(2N-1) + 1 \right) \quad (2.139b)$$

In case where each capacitor is weighted so that the output resistance is minimized, one can use Lagrange multiplier introducing functions  $f$  and  $g$ , and a parameter  $\lambda$  as follows:

$$f(C_1, C_2, \dots, C_N) \equiv \sum_{j=1}^N C_j - C_{\text{TOT}} = 0 \quad (2.140a)$$

$$g(C_1, C_2, \dots, C_N, \lambda) \equiv R_{\text{PMP}}(C_1, C_2, \dots, C_N) - \lambda f(C_1, C_2, \dots, C_N) \quad (2.140b)$$

$$\frac{\partial}{\partial C_j} g(C_1, C_2, \dots, C_N) = -\frac{2^{2N-2j}}{C_j^2} - \lambda = 0 \quad (2.140c)$$

From Eq. (2.140c),

$$\frac{2^{2N-2}}{C_1^2} = \frac{2^{2N-4}}{C_2^2} = \dots = \frac{2^0}{C_N^2} \quad (2.140d)$$

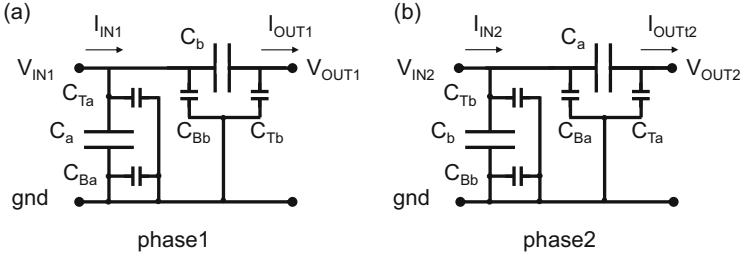
Therefore,

$$C(j) = C_{\text{TOT}} \frac{2^{N-j}}{2^N - 1} \quad (2.140e)$$

From Eqs. (2.138) and (2.140e),  $R_{\text{PMP}}$  is given by

$$R_{\text{PMP}}(N) = \frac{(2^N - 1)^2}{C_{\text{TOT}}} \quad (2.140f)$$

Next, the  $2^N$  pump with a finite parasitic capacitance is considered using Fig. 2.38 instead of Fig. 2.37.



**Fig. 2.38** Two phases of  $2^N$  multiplier with a finite parasitic capacitance  $C_T, C_B$

Equations (2.141) and (2.142) hold in phase 1:

$$I_{IN1}T/2 = C_a(V_{IN1} - V_{OUT2} + V_{IN2}) + C_{Ta}(V_{IN1} - V_{OUT2}) + C_{Bb}V_{IN1} + C_b(V_{IN1} - V_{OUT1} + V_{IN2}) \quad (2.141)$$

$$I_{OUT1}T/2 = C_b(V_{IN2} - V_{OUT1} + V_{IN1}) + C_{Tb}(V_{IN2} - V_{OUT1}) \quad (2.142)$$

When  $C_a = C_b = C/2$ ,  $C_{Ta} = C_{Tb} = C_T/2$ , and  $C_{Ba} = C_{Bb} = C_B/2$ , where a factor of 2 is included for two array structure, phases 1 and 2 are identical. In this case, the input and output voltage (current) can be written by  $V_{IN}$  and  $V_{OUT}$  ( $I_{IN}$  and  $I_{OUT}$ ) without differentiating the two states. Using this symmetry between phases 1 and 2,  $K$ -matrix can be reduced to  $2 \times 2$ :

$$\begin{bmatrix} V_{IN} \\ I_{IN} \end{bmatrix} = \mathbf{K}_{2N}(N) \begin{bmatrix} V_{OUT} \\ I_{OUT} \end{bmatrix} \quad (2.143)$$

$$\mathbf{K}_{2N}(N) = \mathbf{K}(1)^N \quad (2.144)$$

$$\mathbf{K}(1) = \frac{1}{2 + \alpha_T} \begin{bmatrix} 1 + \alpha_T & 2R \\ \alpha_T + \alpha_B + \alpha_T\alpha_B & 4 + \alpha_T + \alpha_B \end{bmatrix} \quad (2.145)$$

where  $\alpha_i = C_i/C$ ,  $r_i = T/2C_i$  ( $i = T, B$ ), and  $R = T/2C$ . The output voltage and current relation are calculated in the first row of Eq. (2.143) with  $V_{IN} = V_{DD}$  and  $I_{IN}$  is calculated in the second row.

Another  $2^N$  charge pump using multi-phase clock is known (Ueno et al. 1986), but its discussion is omitted in this chapter because the area would increase at a rate of the number of phases.

## 2.8 Comparison of Five Topologies

### 2.8.1 Ideal Case Where the Parasitic Capacitance Is Negligibly Small

Table 2.4 summarizes pump characteristic parameters such as the maximum attainable output voltage ( $V_{MAX}$ ) in case of  $V_{DD} = 1$  or equivalently the voltage gain



**Table 2.4** Characteristic parameters in case of negligibly small parasitic capacitance

		$V_{\text{MAX}}$ $G_V$	$R_{\text{PMP}}$	Uniform cap	Weighted cap	Each cap $C(j)$	Current efficiency	$V_{\text{CAP}}$ (Max)	$V_{\text{SW}}$ (Max)
CW	(N: even)	$N+1$	$\frac{N^2(N+1)(N+2)}{12C_{\text{TOT}}}$	$\frac{N^2(N+2)^2}{16C_{\text{TOT}}}$	$C_{2j-1}$	$C_{2j}$	$\frac{1}{G_V}$	1	1
	(N: odd)		$\frac{N(N+1)(N^2+2N+3)}{12C_{\text{TOT}}}$	$\frac{(N+1)^4}{16C_{\text{TOT}}}$	$\propto \left(\frac{N+1}{2} - j + 1\right)$	$\propto \left(\frac{N-1}{2} - j + 1\right)$			
SP			$\frac{N^2}{C_{\text{TOT}}}$	$\frac{N^2}{C_{\text{TOT}}}$	$C_{\text{TOT}}/N$			1	$N$
Dickson								$N$	1
Fibonacci		$F(N+1)$	$\frac{N}{C_{\text{TOT}}} \sum_{j=0}^{N-1} F(j)^2$	$\frac{N}{C_{\text{TOT}}} \left( \sum_{j=0}^{N-1} F(j) \right)^2$	$\propto F(N-j)$			$F(N-1)$	
2 N		$2^N$	$\frac{N}{C_{\text{TOT}}} \left( \frac{2}{3} N(N-1)(2N-1) + 1 \right)$	$\frac{(2^N - 1)^2}{C_{\text{TOT}}}$	$C_{\text{TOT}} \frac{2^{N-j}}{2^N - 1}$			$2^{N-1}$	

( $G_V$ ), the output impedance ( $R_{PMP}$ ), the current efficiency, the maximum voltage applied to two terminals of a pumping capacitor ( $V_{CAP}$ ), and the maximum voltage applied to a switching device ( $V_{SW}$ ). Except for the serial–parallel and the liner pumps, there is an optimum weight for the pumping capacitors to make  $R_{PMP}$  minimized under the condition of a given area. The table includes both cases of non-weighted, i.e., equal sized capacitor and weighted capacitor.

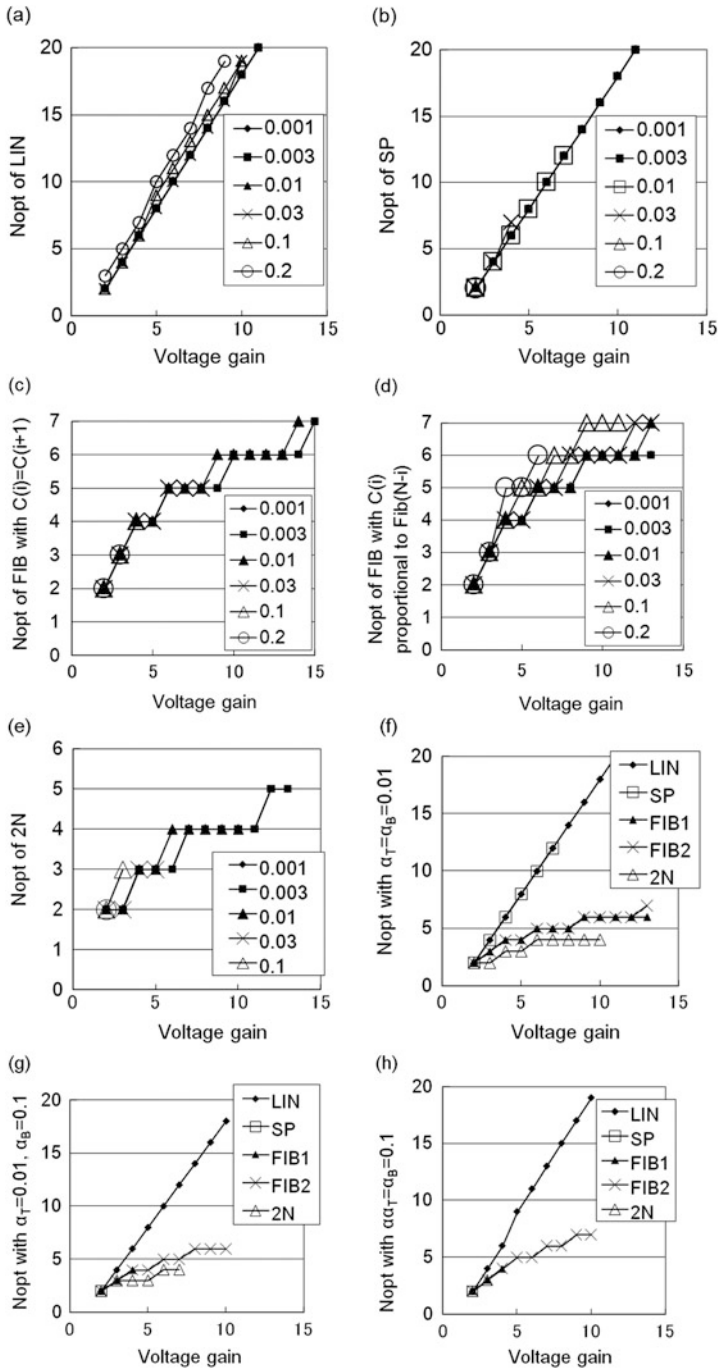
### 2.8.2 Area and Current Efficiency Comparison

Next, let us compare those five topologies in more realistic case where the parasitic capacitance is not negligibly small, which is valid for on-chip high-voltage generation.

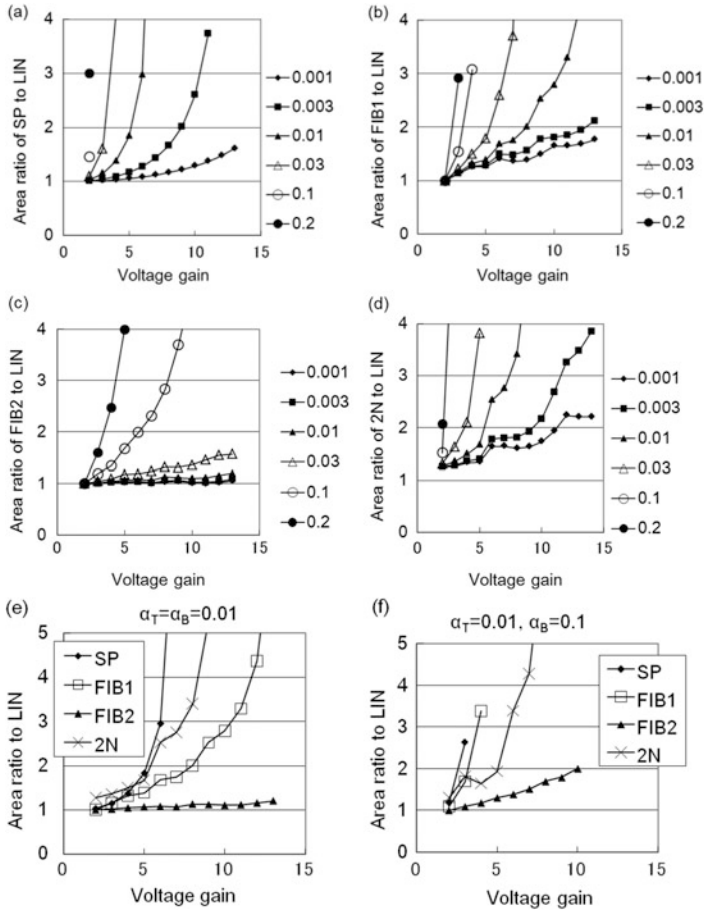
The optimum number of stages ( $N_{OPT}$ ) is determined under the condition that the output current maximizes with a constant entire capacitor area  $\Sigma C(i)$  using the  $K$ -matrix. This procedure is done for various output voltages and parasitic capacitance conditions for each of the multipliers. Then, the capacitor is calculated to output a certain current at a given output voltage with a given parasitic capacitance. Thus, the multipliers are designed and compared with respect to the total capacitor area. In the following figures, for simplicity, it is assumed that (1)  $C_i$  is proportional to  $C$  so that  $\alpha_i$  doesn't depend on  $C$ , (2)  $\alpha_T = \alpha_B$  except for a special case with  $\alpha_T = 0.01$  and  $\alpha_B = 0.1$ , (3) every stage has a same value for  $C$  except for Fibonacci case2 (FIB2) where the values for  $C$ s are varied per stage as described later.

#### (a) Optimum number of stages

Figure 2.39a–e shows the optimum stages ( $N_{OPT}$ ) as a function of voltage gain with a constant current load. The Fibonacci multiplier has the smallest output resistance when  $C(i)$  is proportional to  $\text{Fib}(N - i)$ , where  $i$  indicates  $i$ th stage and  $\text{Fib}(j)$  is the  $j$ th Fibonacci number. SP has a linear dependency on the voltage gain ( $G_V$ ) as the Dickson (LIN). On the other hand,  $N_{OPT}$  of the other multipliers has dependencies as  $\log(G_V)$ . Figure 2.39f, g, h compares five multipliers with  $\alpha_T = \alpha_B = 0.01$  (f),  $\alpha_T = 0.01$  and  $\alpha_B = 0.1$  (g), and  $\alpha_T = \alpha_B = 0.1$  (h). Even with  $\alpha_T = \alpha_B = 0.01$ , SP cannot generate a voltage gain higher than 7, whereas the others can generate voltage gains higher than 10 or more, as shown in Fig. 2.39f. This result shows that the parasitic capacitance decreases the output current as the number of stages in the series increases. With  $\alpha_T$  of 0.1, which is a typical value in cases of integrated multipliers; however, only LIN and FIB2 can generate a voltage gain of 10 or more. Neither FIB1 nor  $2^N$  can generate a voltage gain of 8 or more, as shown in Fig. 2.39g, h. With  $\alpha_T = \alpha_B = 0.01$ ,  $2^N$  only needs the number of stages of 4, which is less by 2 than FIB2 does in case of a voltage gain of 10 as shown in Fig. 2.39f.



**Fig. 2.39** Optimum number of stages as a function of voltage gain with various  $\alpha_T = \alpha_B$  of 0.001 to 0.2 for (a) linear (LIN), (b) serial–parallel (SP), (c) Fibonacci with  $C(i + 1) = C(i)$  ( $i = 1, \dots, N - 1$ ) (FIB1), (d) Fibonacci with weighted  $C$  (FIB2), and (e)  $2^N$  multipliers. Comparisons of the optimum number of stages as a function of voltage gain with  $\alpha_T = \alpha_B = 0.01$  (f), with  $\alpha_T = 0.01, \alpha_B = 0.1$  (g), and with  $\alpha_T = \alpha_B = 0.1$  (h). Comparison of the multiplication factors as a function of  $\alpha_T(j)$  and  $\alpha_B(k)$  for LIN with different optimization methods (Tanzawa 2010)



**Fig. 2.40** Area ratio of (a) SP, (b) FIB1, (c) FIB2, (d) 2N with optimum number of stages to LIN as a function of voltage gain with various  $\alpha_T = \alpha_B$  of 0.001 to 0.2. Comparisons of the area ratio as a function of voltage gain with  $\alpha_T = \alpha_B = 0.01$  (e), with  $\alpha_T = 0.01, \alpha_B = 0.1$  (f), and with  $\alpha_T = \alpha_B = 0.1$  (g). Comparisons of the area ratio as a function of  $\alpha_T (= \alpha_B)$  with a voltage gain of 3 (h), 6 (i), and 10 (j) (Tanzawa 2010)

(b) Circuit area

Using the values for  $N_{OPT}$  calculated, as shown in Fig. 2.39, the values for  $C$  per stage are calculated to output a specific current for each voltage gain, each  $\alpha$ , and each multiplier. Figure 2.40 shows the circuit area, which is defined by  $\Sigma C(i)$  as a measure, compared to that of LIN. As shown in Fig. 2.40a, SP can have an equivalent area as LIN does as far as both  $\alpha$  and voltage gains are small, but SP becomes very sensitive to  $\alpha$  higher than 0.01. FIB1, FIB2, and  $2^N$  have smaller sensitivity than SP, as shown in Fig. 2.40b–d, but they also become very sensitive to  $\alpha$  higher than 0.03. Figure 2.40e–g

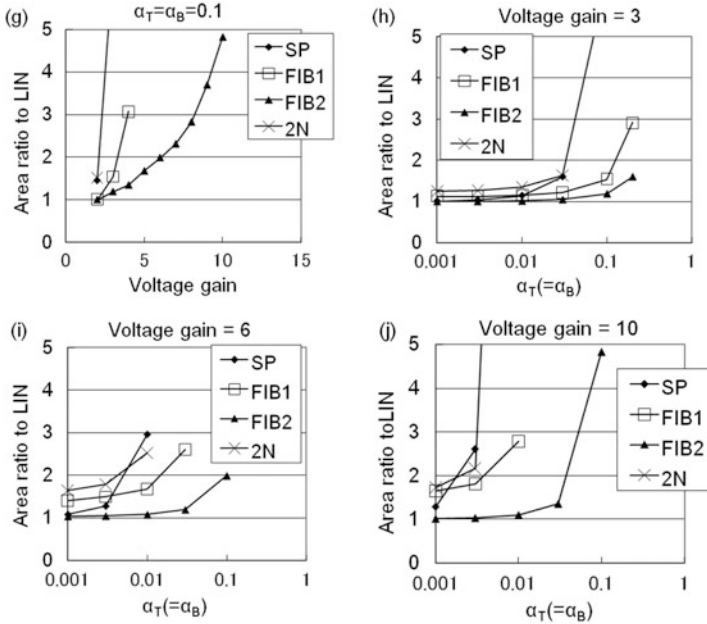


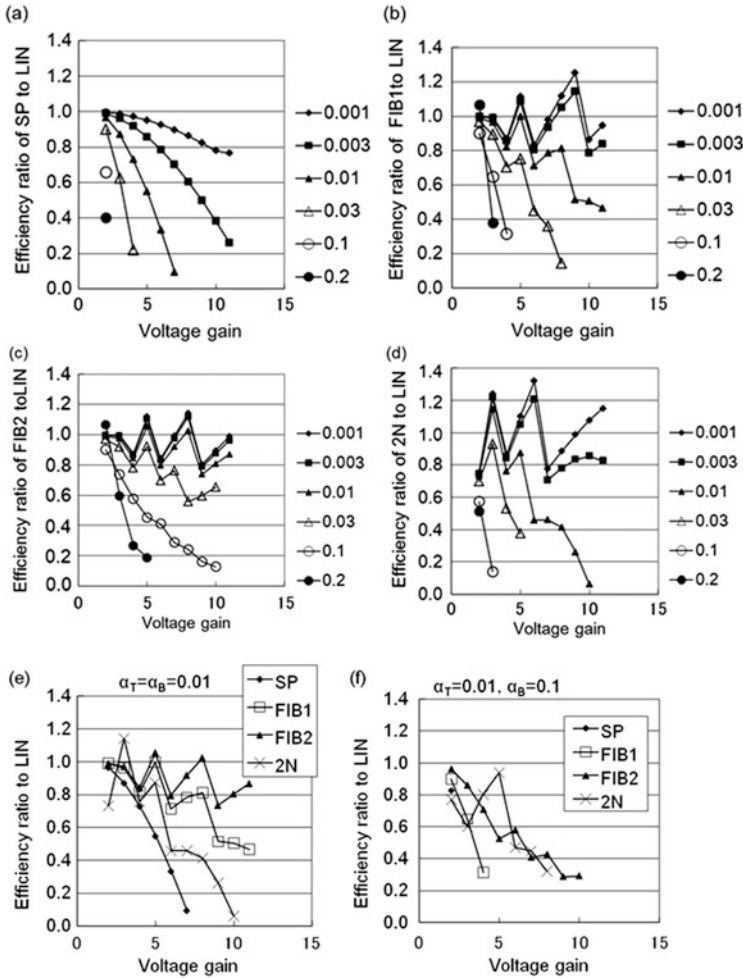
Fig. 2.40 (continued)

compares the area ratios of four multipliers to LIN with  $\alpha_T = \alpha_B = 0.01$  (e),  $\alpha_T = 0.01$  and  $\alpha_B = 0.1$  (f), and  $\alpha_T = \alpha_B = 0.1$  (g). Among those four, only FIB2 has similar area as LIN in case of  $\alpha_B$  of 0.01, but each needs much more area than LIN in case of  $\alpha_B$  of 0.1. For example, FIB2 with a voltage gain of 10 needs an area that is five times larger than LIN in case of  $\alpha_T = \alpha_B = 0.1$ , as shown in Fig. 2.40g.

Thus, LIN has minimum total capacitor area among the multipliers in case of  $\alpha_T$  and  $\alpha_B$  of 0.1 or higher which are typical numbers in integrated circuits. In case of  $\alpha_T$  and  $\alpha_B$  of 0.01 or smaller, LIN and FIB2 have smaller area than the others do under the condition of a voltage gain of 10 or smaller. Such a small parasitic capacitance is realized in discrete application.

### (c) Current efficiency

Current efficiency is defined by  $I_{OUT}/I_{IN}$ . Figure 2.41 shows ratios of the efficiency of the other four multipliers to that of LIN. Because the efficiency strongly depends on the number of stages, nonmonotonic dependencies on the voltage gain are observed in FIB1, FIB2, and  $2^N$  due to different dependencies of  $N_{OPT}$  on the voltage gain against LIN. In case of  $\alpha_T$  and  $\alpha_B$  of 0.01, FIB2 and  $2^N$  have similar efficiencies as LIN within  $\pm 30\%$  up to a voltage gain of 10, as shown in Fig. 2.41e. However, in case of  $\alpha$  of 0.1, any multiplier decreases the efficiency monotonically, as shown in Fig. 2.41f, g. The efficiency of FIB2 is degraded to about 0.3 of that of LIN at a voltage gain of 10 in case of  $\alpha_T = 0.01$  and  $\alpha_B = 0.1$ , as shown in Fig. 2.41f.



**Fig. 2.41** Current efficiency ratio of (a) SP, (b) FIB1, (c) FIB2, (d) 2N with optimum number of stages to LIN as a function of voltage gain with various  $\alpha_T = \alpha_B$  of 0.001 to 0.2. Comparison of the efficiency ratio as a function of voltage gain with  $\alpha_T = \alpha_B = 0.01$  (e), with  $\alpha_T = 0.01, \alpha_B = 0.1$  (f), and with  $\alpha_T = \alpha_B = 0.1$  (g). Comparisons of the efficiency ratio as a function of  $\alpha_T (= \alpha_B)$  with a voltage gain of 3 (h), 6 (i), and 10 (j) (Tanzawa 2010)

(d) Number of discrete capacitors

In order to compare the number of discrete capacitor components among the multipliers for discrete applications,  $2^N$  multiplier needs to use  $N_{CAP}$  defined by  $2N_{OPT}$  rather than  $N_{OPT}$  itself as shown in Fig. 2.39e.  $N_{CAP}$  of the rest of the multipliers is same as  $N_{OPT}$ . Figure 2.42a, b is, respectively, identical to Fig. 2.39f, h except for the vertical axis. Figure 2.42 shows FIB has the least number of capacitors among the multipliers. For example,  $N_{CAP}$  of FIB is about one-third of that of LIN in case of a voltage gain of 10 and  $\alpha_T = \alpha_B = 0.01$ .

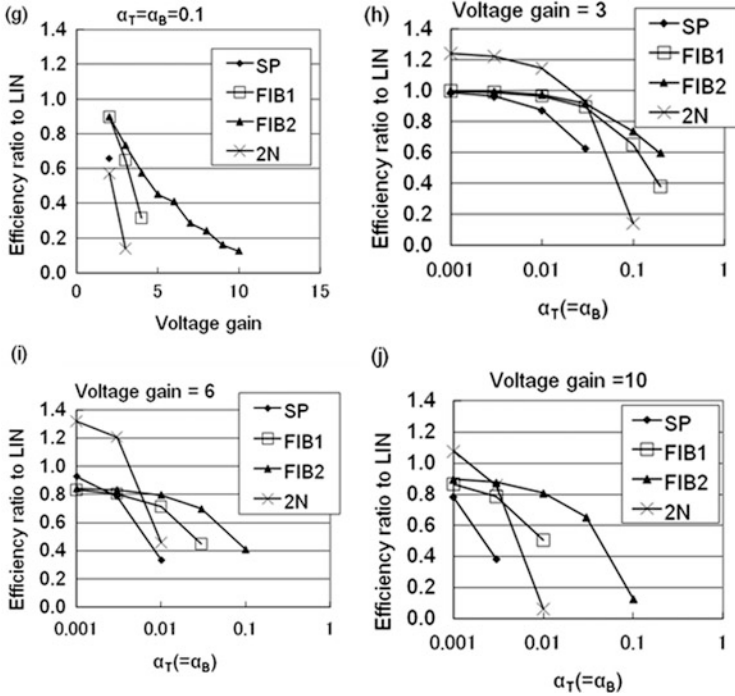


Fig. 2.41 (continued)

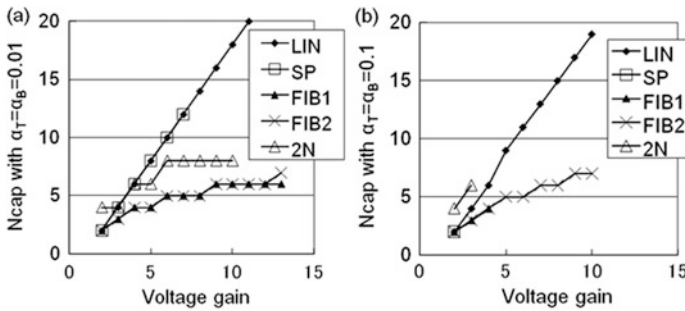


Fig. 2.42 Comparison of the number of discrete capacitors (Tanzawa 2010)

As a result, the linear Dickson cell is the best for integration because of the smallest total capacitor area and the highest current or power efficiency under the assumption that the parasitic capacitance is not smaller than 10 % of the multiplier capacitance, and Fibonacci cell is the best for discrete application (Makowski and Maksimovic 1995) because of the minimum number of capacitor components with moderate current or power efficiency under the assumption that the parasitic capacitance is not larger than 1 % of the multiplier capacitance.

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# Chapter 3

## Design of DC-DC Dickson Charge Pump

**Abstract** This chapter expands upon circuit theory of DC-DC Dickson charge pump and discusses how the charge pump voltage multiplier is optimally designed. Design equations and equivalent circuit models are derived for the charge pump. Power efficiency of the charge pump is described as a function of device and design parameters. Using the model, optimizations are discussed to minimize the circuit area under various conditions that the output current, the ramp time, and the power dissipation are given theoretically. Guideline for comprehensive optimum design is also described.

This chapter is composed of the following. Section 3.1 describes a dynamic behavior of charge pump and the equivalent circuit under the condition that the operation frequency is low enough. Section 3.2 expands upon circuit equations at a moderate to high operation frequency where the switching device is a diode, MOSFET in saturation region, and MOSFET in triode region, respectively. Section 3.3 shows how power efficiency of charge pumps is determined taking the dependence of the output voltage, the threshold voltage of switching device, and parasitic capacitance into consideration. Section 3.4 discusses several optimizations of the circuit with respect to circuit area, rise time, and power, as well as guideline for comprehensive optimum design. Section 3.5 summarizes key design equations.

### 3.1 Circuit Analysis Under Low-Frequency Operation

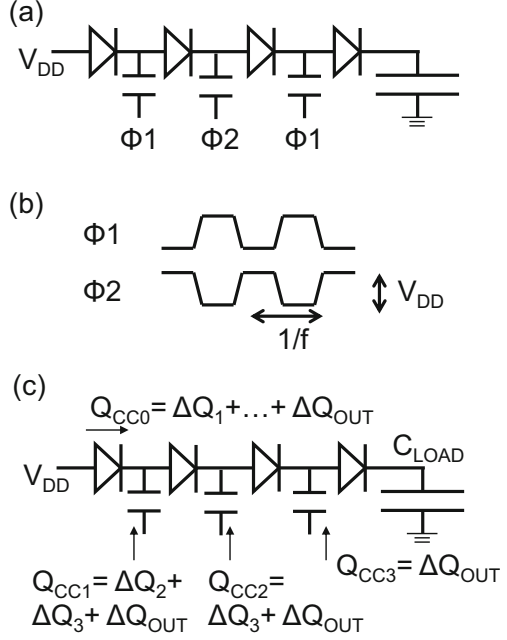
#### 3.1.1 Dynamic Behavior

This subsection discusses dynamic behavior of the Dickson charge pump and extracts the equivalent circuit model.

Figure 3.1 illustrates three-stage Dickson pump.  $\Delta Q_i$  ( $i = 1-3$ , OUT) indicates increased amount of charges in each capacitor in the time period  $T_R$ .

In order to derive the recurrence formula of the output voltage, the total charge consumed by the charge pump during boosting is obtained by two different methods; by using the charge stored in each capacitor as shown in Fig. 3.2a, b and by using the sum of the charge consumed by the charge pump in one cycle time

**Fig. 3.1** Charges supplied by  $V_{DD}$  in  $T_R$



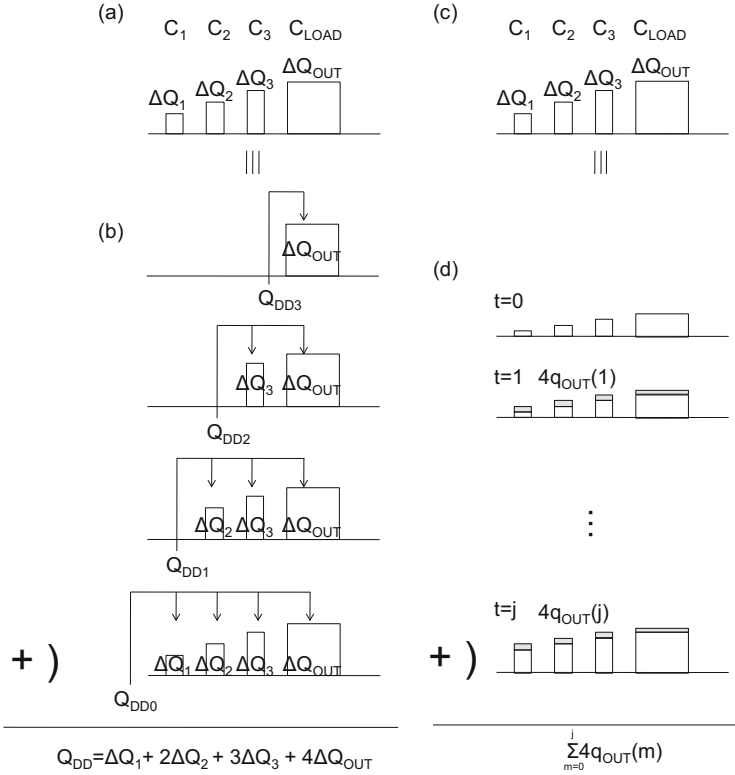
as shown in Fig. 3.2c, d. Firstly, the total charge  $Q_{DD}^d(j)$  consumed by the charge pump during the arbitrary time  $j$  is calculated using the charges stored in the charge pump capacitors, based on Figs. 3.1 and 3.2a, b. The total charge  $Q_{DD}(k, j)$  ( $1 \leq k \leq N$ ) consumed by the driver  $I(k)$  driving the capacitor  $C(k)$  during  $j$  equals the total charge transferred from the capacitor  $C(k)$  to the next one  $C(k+1)$  through the diode  $D(k+1)$  during  $j$ , as illustrated in Fig. 3.1. Therefore,  $Q_{DD}(k, j)$  equals the total charge increase in the capacitors  $C(k+1)$ ,  $C(k+2)$ ,  $\dots$ ,  $C(N)$  and  $C_{LOAD}$  during  $j$ , where  $C_{LOAD}$  is the load capacitance of the charge pump circuit. Similarly, the total charge  $Q_{DD}(0, j)$  supplied by the input voltage  $V_{DD}$  at the left-hand side of Fig. 3.1 equals the total charge increase in all capacitors including  $C_{LOAD}$ . Therefore, if  $Q(i, j)$  and  $Q_{LOAD}(i)$  are the charges stored in the capacitors  $C(i)$  ( $1 \leq i \leq N$ ) and  $C_{LOAD}$  at  $j$ , respectively, for  $1 \leq k \leq N-1$ :

$$Q_{DD}(k, j) = \sum_{i=k+1}^N [Q(i, j) - Q(i, 0)] + [Q_{LOAD}(j) - Q_{LOAD}(0)] \quad (3.1)$$

and

$$Q_{DD}(N, j) = Q_{LOAD}(j) - Q_{LOAD}(0) \quad (3.2)$$

The total consumed charge  $Q_{DD}^d(j)$  is the sum of all charges  $Q_{DD}(k, j)$   $1 \leq k \leq N$ , so that



**Fig. 3.2** Two methods for calculating the total input charges: method 1 with (a) and (b) and method 2 with (c) and (d)

$$\begin{aligned}
 Q_{DD}^d(j) &= \sum_{k=0}^N Q_{DD}(k, j) \\
 &= \sum_{k=1}^N k [Q(k, j) - Q(k, 0)] + (N + 1)[Q_{LOAD}(j) - Q_{LOAD}(0)]
 \end{aligned}
 \tag{3.3}$$

The following initial conditions can be assumed:

$$Q(2k, 0) = 0 \tag{3.4}$$

$$Q(2k - 1, 0) = C(V_{DD} - V_T) \tag{3.5}$$

and

$$Q_{\text{LOAD}}(0) = C_{\text{LOAD}}(V_{\text{DD}} - V_{\text{T}}) \quad (3.6)$$

which satisfy Eqs. (2.74a) and (2.74b). Under the assumption that Eqs. (2.74a) and (2.74b) hold during boosting, Eq. (3.3) results in

$$Q_{\text{DD}}^{\text{d}}(j) = (N + 1)C_{\text{OUT}}(V_{\text{OUT}}(j) - V_{\text{G}}) \quad (3.7)$$

$$C_{\text{OUT}} \equiv C_{\text{LOAD}} + C_{\text{PMP}} \quad (3.8)$$

$$C_{\text{PMP}} = (1 + \alpha_{\text{T}})A(N)C \quad (3.9)$$

where  $A(N)$  is a function of  $N$ ,

$$A(N) = \frac{4N^2 + 3N + 2}{12(N + 1)} \quad (3.10a)$$

$$A(N) = \frac{4N^2 - N - 3}{12N} \quad (3.10b)$$

for even and odd  $N$ , respectively.  $C_{\text{PMP}}$  is about one-third of the total charge pump capacitance,  $NC/3$ , and its error is less than 3 % for even  $N \geq 4$  and less than 7 % for odd  $N \geq 5$ :

$$C_{\text{PMP}} \approx (1 + \alpha_{\text{T}})CN/3 \quad (3.11)$$

Another expression for  $Q_{\text{DD}}^{\text{d}}(j)$  is derived below, based on Fig. 3.2c, d. Since the charge  $q_{\text{DD}}^{\text{s}}$  supplied by the power supply in a cycle time in steady state is equal to the charge  $q$  transferred to the capacitor  $C(1)$  through the diode  $D(1)$  plus the charge  $Nq$  transferred from  $N$  capacitors  $C(k)$  ( $1 \leq k \leq N$ ) to the next ones:

$$q_{\text{DD}}^{\text{s}} = (N + 1)q \quad (3.12)$$

Like the above equation in steady state, the relation between the supplied charge  $q_{\text{DD}}^{\text{d}}(j)$  and output charge increase  $q_{\text{OUT}}(j)$  in a cycle time from  $j$  to  $j + 1$  during boosting:

$$q_{\text{DD}}^{\text{d}}(j) = (N + 1)q_{\text{OUT}}(j) \quad (3.13)$$

Under the assumption that Eq. (3.13) holds even during boosting, the total supplied charge during  $j$ ,  $Q_{\text{DD}}^{\text{d}}(j)$ , is given by

$$\begin{aligned}
Q_{DD}^d(j) &= \sum_{m=0}^j q_{DD}^d(m) \\
&= (N+1) \sum_{m=0}^j \frac{(1+\alpha_T)C}{N} \left[ N \left( \frac{V_{DD}}{1+\alpha_T} - V_T \right) + V_{DD} - V_T - V_{OUT}(m) \right]
\end{aligned} \tag{3.14}$$

where Eq. (2.73) is used. Combining Eqs. (3.7) with (3.14),

$$\begin{aligned}
C_{OUT}(V_{OUT}(j) - V_{DD} - V_T) \\
= \sum_{m=0}^j \frac{(1+\alpha_T)C}{N} \left[ N \left( \frac{V_{DD}}{1+\alpha_T} - V_T \right) + V_{DD} - V_T - V_{OUT}(m) \right]
\end{aligned} \tag{3.15}$$

Since Eq. (3.15) holds for arbitrary  $j$ , the recurrence formula for  $V_{OUT}$  holds as follows:

$$\begin{aligned}
C_{OUT}(V_{OUT}(j+1) - V_{OUT}(j)) \\
= \frac{(1+\alpha_T)C}{N} \left[ N \left( \frac{V_{DD}}{1+\alpha_T} - V_T \right) + V_{DD} - V_T - V_{OUT}(j+1) \right]
\end{aligned} \tag{3.16}$$

Using the initial condition of  $V_{OUT}(0) = V_{DD} - V_T$  from Eqs. (3.6) to (3.16) is solved as

$$V_{OUT}(j) = N \left( \frac{V_{DD}}{1+\alpha_T} - V_T \right) + V_{DD} - V_T - N \left( \frac{V_{DD}}{1+\alpha_T} - V_T \right) \beta^j \tag{3.17}$$

$$\beta = \left( 1 + \frac{(1+\alpha_T)C}{NC_{OUT}} \right)^{-1} \tag{3.18}$$

As a result, the rise time  $T_R$  that the output voltage  $V_{OUT}(j)$  rises from  $V_{DD} - V_T$  to  $V_{PP}$ , which satisfies  $V_{OUT}(T_R) = V_{PP}$ , is solved as

$$T_R = \ln \left( 1 - \frac{V_{PP} - V_{DD} + V_T}{N(V_{DD}/(1+\alpha_T) - V_T)} \right) / \ln \beta \tag{3.19}$$

It is noted that this term should be multiplied by the cycle time of the driving clocks in practice, because Eq. (3.19) is expressed by the number of clock cycles. From Eq. (3.19), the mean current consumption  $I_{DD}^d$  during  $T_R$  can be obtained as

$$\begin{aligned}
I_{DD}^d &\equiv Q_{DD}^d(T_R)/T_R \\
&= (N+1)C_{OUT}(V_{PP} - V_{DD} + V_T)/T_R
\end{aligned} \tag{3.20a}$$

$C_{OUT}$  can be regarded as the total load capacitance during boosting. Therefore, it is considered that  $C_{PMP}$  represents the self-load capacitance of the charge pump

itself.  $C_{PMP}$  is about one-third of the total charge pump capacitance,  $NC/3$ , and its error is less than 3 % for even  $N \geq 4$  and less than 7 % for odd  $N \geq 5$ . When the parasitic capacitance at the bottom nodes of the pumping capacitors ( $\alpha_B C$ ) is taken into account for the current consumption, Eq. (3.20a) needs to be replaced with Eq. (3.20b):

$$I_{DD}^d = (N + 1)C_{OUT}(V_{PP} - V_{DD} + V_T)/T_R + \alpha_B N C V_{DD}/T \quad (3.20b)$$

### 3.1.2 Equivalent Circuit Model

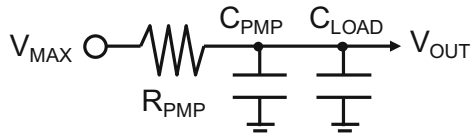
Although the output voltage  $V_{OUT}(j)$  is actually a staircase waveform, it can be regarded as a smooth function in case the rise time is sufficiently large compared with the cycle time of the driving clocks. In this case, Eq. (3.16) indicates the equivalent circuit of the charge pump as shown in Fig. 3.3.  $R_{PMP}$  represents the output series resistance of the charge pump and is given by  $N/C(1 + \alpha_T)$  (as mentioned above, this is multiplied by the cycle time of the driving clocks and has the same dimension as resistor).  $V_{MAX}$  is the maximum output voltage of the charge pump,  $N(V_{DD}/(1 + \alpha_T) - V_T) + V_{DD} - V_T$ .  $C_{PMP}$  expressed by Eqs. (3.9) and (3.10a)–(3.10b) indicates the self-load capacitance of the charge pump and is connected in parallel with the output load capacitance  $C_{LOAD}$ .

In order to compute the rise time and the current consumption accurately, only the cutoff condition of the transfer diodes and the charge conservation rule are used. A charge pump circuit with an even number of stages is considered. Since the charges  $Q(2k - 1, j)$  stored in the capacitors  $C(2k - 1)$  at time  $j$  are transferred to the next ones  $C(2k)$  by time  $j + 1/2$ , the following relations hold if the charge conservation rule is assumed:

$$\begin{aligned} Q(2k - 1, j) + Q(2k, j) \\ = Q(2k - 1, j + 1/2) + Q(2k, j + 1/2) \end{aligned} \quad (3.21)$$

Note that the charges stored in the parasitic capacitors are canceled out each other. From the condition that the diode  $D(2k)$  is cut off at time  $j + 1/2$ :

**Fig. 3.3** Equivalent pump model



$$V_{MAX} = \left(\frac{N}{1 + \alpha_T} + 1\right)V_{DD} - (N + 1)V_{TH}$$

$$R_{PMP} = \frac{N}{(1 + \alpha_T)Cf} \quad C_{PMP} \approx (1 + \alpha_T)CN / 3$$

$$\frac{Q(2k, j + 1/2)}{C} - \frac{Q(2k - 1, j + 1/2)}{C} = \frac{V_{DD}}{1 + \alpha_T} - V_T \quad (3.22)$$

Similarly, the charges  $Q(2k, j + 1/2)$  stored in the capacitors  $C(2k)$  at time  $j + 1/2$  are transferred to the capacitors  $C(2k + 1)$  by time  $j + 1$ :

$$\begin{aligned} & Q(2k, j + 1) + Q(2k + 1, j + 1) \\ &= Q(2k, j + 1/2) + Q(2k + 1, j + 1/2) \end{aligned} \quad (3.23)$$

$$\frac{Q(2k + 1, j + 1)}{C} - \frac{Q(2k, j + 1)}{C} = \frac{V_{DD}}{1 + \alpha_T} - V_T \quad (3.24)$$

And also,

$$\begin{aligned} & Q_{LOAD}(j + 1) + Q(N, j + 1) \\ &= Q_{LOAD}(j + 1/2) + Q(N, j + 1/2) \end{aligned} \quad (3.25)$$

$$\frac{Q_{LOAD}(j + 1)}{C_{LOAD}} - \frac{Q(N, j + 1)}{C} = \frac{V_{DD}}{1 + \alpha_T} - V_T \quad (3.26)$$

Furthermore,

$$Q(1, j) = Q(1, j + 1) = C(V_{DD} - V_T) \quad (3.27)$$

Eliminating the intermediate states at time  $j + 1/2$ ,  $Q(k, j + 1/2)$ , and  $Q_{LOAD}(j + 1/2)$ , from the above equations, for more than three stages,

$$\begin{aligned} Q(2, j + 1) &= \frac{1}{4} \left[ Q(2, j) + Q(3, j) + Q(4, j) \right. \\ &\quad \left. - C \left( \frac{1 - \alpha_T}{1 + \alpha_T} V_{DD} - V_T \right) \right] \end{aligned} \quad (3.28)$$

$$\begin{aligned} Q(3, j + 1) &= \frac{1}{4} \left[ Q(2, j) + Q(3, j) + Q(4, j) \right. \\ &\quad \left. + C \left( \frac{3 + \alpha_T}{1 + \alpha_T} V_{DD} - 3V_T \right) \right] \end{aligned} \quad (3.29)$$

$$\begin{aligned} Q(2k, j + 1) &= \frac{1}{4} \left[ Q(2k - 1, j) + Q(2k, j) + Q(2k + 1, j) \right. \\ &\quad \left. + Q(2k + 2, j) - 2C \left( \frac{V_{DD}}{1 + \alpha_T} - V_T \right) \right] \end{aligned} \quad (3.30)$$

$$\begin{aligned} Q(2k + 1, j + 1) &= \frac{1}{4} \left[ Q(2k - 1, j) + Q(2k, j) + Q(2k + 1, j) \right. \\ &\quad \left. + Q(2k + 2, j) + 2C \left( \frac{V_{DD}}{1 + \alpha_T} - V_T \right) \right] \end{aligned} \quad (3.31)$$



$$Q(N, j+1) = \frac{C}{2(C_{\text{LOAD}} + C)} [Q(N-1, j) + Q(N, j) + 2Q_{\text{LOAD}}(j) - (2C_{\text{LOAD}} - C) \left( \frac{V_{\text{DD}}}{1 + \alpha_{\text{T}}} - V_{\text{T}} \right)] \quad (3.32)$$

$$Q_{\text{LOAD}}(j+1) = \frac{C_{\text{LOAD}}}{2(C_{\text{LOAD}} + C)} [Q(N-1, j) + Q(N, j) + 2Q_{\text{LOAD}}(j) + 3C \left( \frac{V_{\text{DD}}}{1 + \alpha_{\text{T}}} - V_{\text{T}} \right)] \quad (3.33)$$

Equations (3.30) and (3.31) hold for more than five stages and  $2 \leq k \leq N/2 - 1$ . The stored charges  $Q(k, j)$  ( $1 \leq k \leq N$ ) and  $Q_{\text{LOAD}}(j)$  can be iteratively computed using the initial condition of Eqs. (3.4)–(3.6). The rise time can be obtained by the time that the output voltage  $Q_{\text{LOAD}}/C_{\text{LOAD}}$  rises from the initial voltage  $V_{\text{DD}} - V_{\text{T}}$  to the final voltage  $V_{\text{pp}}$ . In case of a charge pump with one or two stages the equations like the above can be analytically solved for each  $Q(k, j)$  and  $Q_{\text{LOAD}}(j)$ , so that the rise time can be solved exactly. On the other hand, the solution of the rise time for a charge pump with three stages can be obtained approximately rather than exactly because of the nonlinear equation for the rise time. This approximation introduces an error of only a few percent to the solution.

The charge supplied to the charge pump during one cycle from  $j$  to  $j+1$ ,  $q_{\text{DD}}^{\text{d}}(j)$ , is the sum of the charges supplied by the drivers  $I(2k-1)$ , which are the charge increases in the capacitors  $C(2k)$ ,  $(1 + \alpha_{\text{T}}) (Q(2k, j+1/2) - Q(2k, j))$ , the charges supplied by the drivers  $I(2k)$  and  $I(N)$ , which are the charge increases in the capacitors  $C(2k+1)$  and  $C_{\text{LOAD}}$ ,  $(1 + \alpha_{\text{T}}) (Q(2k+1, j+1) - Q(2k+1, j+1/2))$  and  $Q_{\text{LOAD}}(j+1) - Q_{\text{LOAD}}(j+1/2)$ , respectively, the charge supplied to the capacitor  $C(1)$  by the input voltage,  $(1 + \alpha_{\text{T}}) (Q(1, j+1) - Q(1, j+1/2))$ , and the charge supplied to the parasitic capacitance at the bottom nodes of pumping capacitors,  $N\alpha_{\text{B}}CV_{\text{DD}}$ . By using Eqs. (3.21)–(3.27),

$$q_{\text{DD}}^{\text{d}}(j) = (1 + \alpha_{\text{T}}) \left( 2C \left( \frac{V_{\text{DD}}}{1 + \alpha_{\text{T}}} - V_{\text{T}} \right) - Q(2, j) \right) - \sum_{k=2}^{N/2} (1 + \alpha_{\text{T}}) (Q(2k, j) - CV_{\text{G}}) + \sum_{k=1}^{N/2-1} (1 + \alpha_{\text{T}}) Q(2k+1, j+1) + (Q_{\text{LOAD}}(j+1) - Q_{\text{LOAD}}(j)) + \alpha_{\text{B}} N C V_{\text{DD}} \quad (3.34)$$

Therefore, the total charge supplied to the charge pump during boosting,  $Q_{DD}^d(T_R)$ , can be iteratively computed by

$$Q_{DD}^d(T_R) = \sum_{j=0}^{T_R} q_{DD}^d(j) \quad (3.35)$$

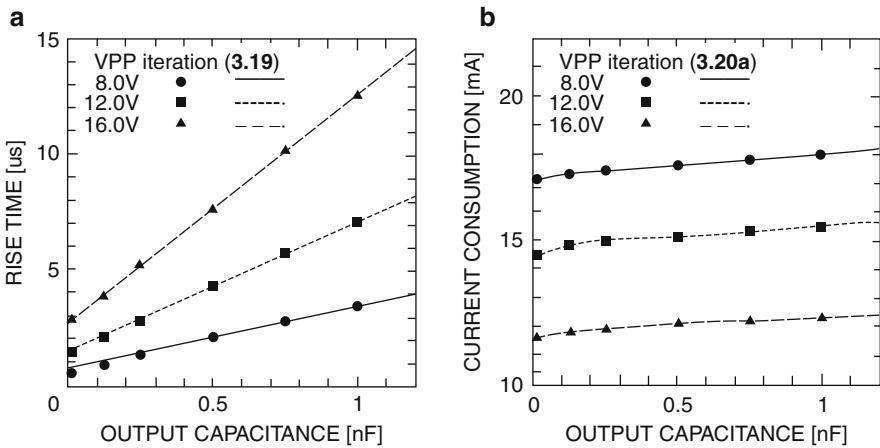
As a result, the average current consumption during boosting,  $I_{DD}^d$ , can be calculated by

$$I_{DD}^d = Q_{DD}^d(T_R)/T_R \quad (3.36)$$

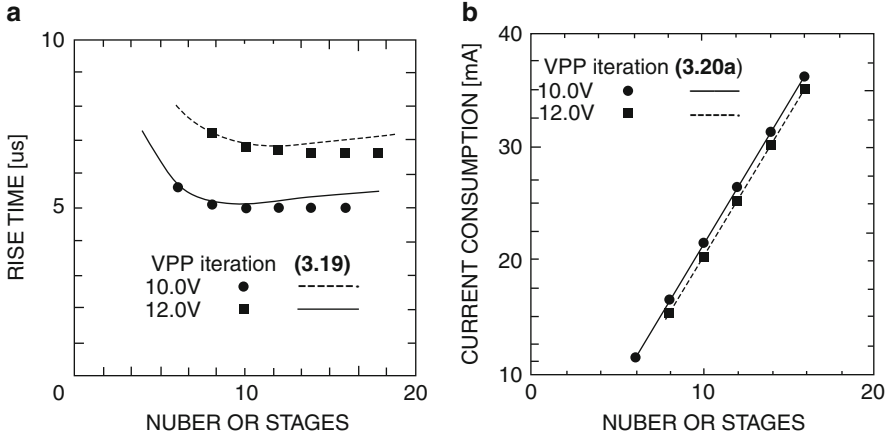
The rise time and the current consumption computed have been in good agreement with the SPICE simulation results within 5 %. Therefore, the verification of the analytical results is made by the comparison with the iteration method as below.

Figures 3.4, 3.5, 3.6, and 3.7, respectively, show the dependence of the rise time and the current consumption on the output load capacitance (Fig. 3.4), the number of stages (Fig. 3.5), the boosted voltage (Fig. 3.6), and the supply voltage (Fig. 3.7). As shown in Fig. 3.4a, the rise time increases proportionally to the output capacitance. The y-intersection in Fig. 3.4a indicates the rise time in case of no output load capacitance ( $C_{LOAD} = 0$ ), and the self-load capacitance of the charge pump, which has been estimated by the analysis as about one-third of the total charge pump capacitance, is in good agreement with the iteration method. The current consumption during boosting has small dependence on the output load capacitance, as shown in Fig. 3.4b.

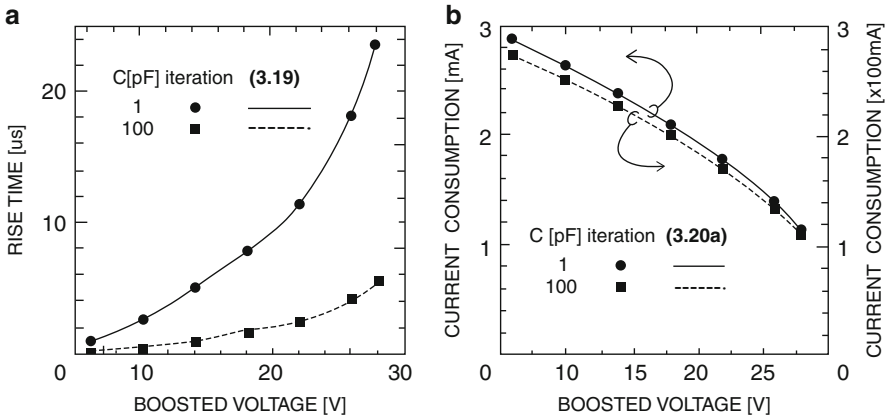
As shown in Fig. 3.5a, the rise time iteratively computed by Eqs. (3.28)–(3.33) is constant for a large number of stages, while the rise time calculated by the analytical expression slightly increases with the number of stages because of the



**Fig. 3.4** Dependence of the rise time (a) and the current consumption (b) on the output load capacitance under the condition of  $N = 8$ ,  $V_{DD} = 3.0$  V,  $V_T = 0.6$  V,  $C = 100$  pF,  $\alpha_T = \alpha_B = 0$ , and the cycle time of driving clocks,  $T = 100$  ns

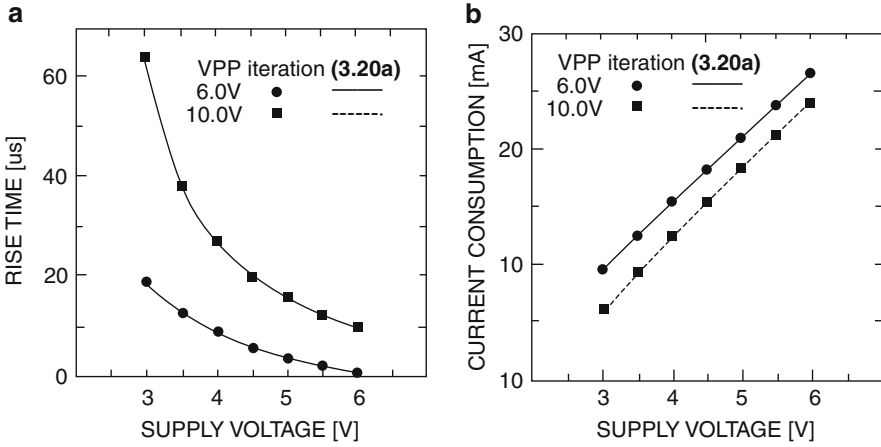


**Fig. 3.5** Dependence of the rise time (a) and the current consumption (b) on the number of stages under the condition of  $V_{DD}=3.0$  V,  $V_T=0.6$  V,  $C=100$  pF,  $\alpha_T=\alpha_B=0$ ,  $C_{LOAD}=1$  nF and  $T=100$  ns



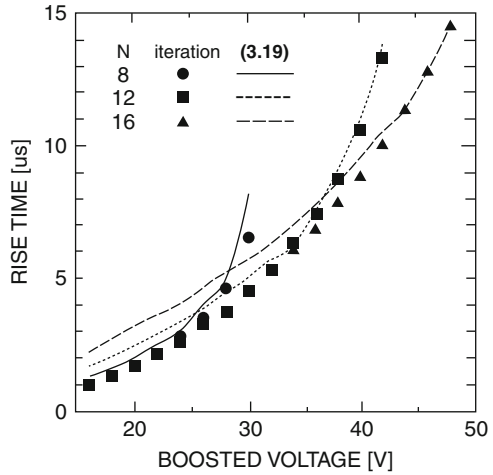
**Fig. 3.6** Dependence of the rise time (a) and the current consumption (b) on the boosted voltage under the condition of  $N=8$ ,  $V_{DD}=4.0$  V,  $V_T=0.6$  V,  $\alpha_T=\alpha_B=0$ ,  $C_{LOAD}=10$  pF, and  $T=100$  ns

increasing self-load capacitance  $C_{PMP}$ . Figure 3.5a indicates the rise time doesn't depend on the excess number of stages in actual and the error of the analytical expression increases as the boosted voltage becomes much smaller than the maximum output voltage  $N(V_{DD}/(1+\alpha_T) - V_T) + V_{DD} - V_T$ . This suggests that the assumption that the charge pump is kept at steady state even during boosting doesn't hold in such case. The constant rise time and the total supplied charge proportional to the number of stages result in a current consumption that is increasing with the number of stages (Fig. 3.5b). The discrepancy between analytical and iterative results in Fig. 3.5a is attributed to the inaccuracy in the self-load



**Fig. 3.7** Dependence of the rise time (a) and the current consumption (b) on the supply voltage under the condition of  $N = 4$ ,  $V_T = 0.6$  V,  $\alpha_T = \alpha_B = 0$ ,  $C = 100$  pF,  $C_{LOAD} = 10$  nF, and  $T = 100$  ns

**Fig. 3.8** Dependence of the rise time on the boosted voltage under the condition of  $V_{DD} = 4.0$  V,  $V_T = 0.6$  V,  $\alpha_T = \alpha_B = 0$ ,  $C_{LOAD} = 0$  pF, and  $T = 100$  ns



capacitance  $C_{PMP}$ , while this discrepancy doesn't appear in Fig. 3.5b. This is because the discrepancy of the rise time  $T_R$  is canceled by that of the total supplied charge  $Q_{DD}^d(T_R)$  in Eq. (3.20a), which is also increasing with the number of stages. The rise time and the current consumption show a large dependence on the boosted voltage (Fig. 3.6a, b) and the supply voltage (Fig. 3.7a, b). Even in case that the charge pump capacitance is ten times larger than the output capacitance as shown in Fig. 3.6 (in case of  $C = 100$  pF and  $C_{LOAD} = 10$  pF), the analytical expression agrees with the iteration method.

Figure 3.8 shows the dependence of the rise time on the output voltage under the condition of no output load capacitance. In this case, the charge pump circuit has

only a self-load capacitance. The analytical expression Eq. (3.19) in which the load capacitance  $C_{\text{LOAD}}$  is set to 0 agrees with the iteration method in case that the boosted voltage  $V_{\text{PP}}$  is not much smaller than the maximum output voltage of  $N$  ( $V_{\text{DD}}/(1 + \alpha_{\text{T}}) - V_{\text{T}}) + V_{\text{DD}} - V_{\text{T}}$ . However, in case of a small boosted voltage, the rise time given by the iteration method is independent of the number of stages. On the other hand, the rise time given by the analytical expression increases with the number of stages.

As mentioned above, the difference between the analytical expression and the iteration method increases as the boosted voltage becomes much smaller than the maximum output voltage, or in other words, the number of stages becomes excessively large compared with the number of stages necessary for the boosted voltage. In such case, the analytical results of Eqs. (3.19) and (3.20a) cannot use. In a typical case that the boosted voltage is not smaller than one-fourth of the maximum output voltage, the analytical results agree with the simulation results computed by the iteration method within 10 % for the rise time and within 2 % for the current consumption.

### 3.1.3 Input and Output Power in Dynamic State

The power consumption  $P_{\text{IN}}$ , the output power  $P_{\text{OUT}}$ , and the power efficiency  $R_{\text{PWR}}$  during boosting are defined as

$$P_{\text{IN}} \equiv \sum_{j=0}^{T_{\text{R}}} q_{\text{DD}}^{\text{d}}(j)V_{\text{DD}}/T_{\text{R}} \quad (3.37)$$

$$P_{\text{OUT}} \equiv \sum_{j=0}^{T_{\text{R}}} q_{\text{OUT}}(j)V_{\text{OUT}}(j)/T_{\text{R}} \quad (3.38)$$

$$R_{\text{PWR}} \equiv P_{\text{OUT}}/P_{\text{IN}} \quad (3.39)$$

By using Eq. (3.17) for  $V_{\text{OUT}}(j)$ , Eq. (2.73) for  $q_{\text{OUT}}$ , and Eq. (3.13) for  $q_{\text{DD}}^{\text{d}}(j)$ , these values can be calculated as

$$P_{\text{IN}} = (N + 1)C_{\text{OUT}}(V_{\text{PP}} - V_{\text{G}})V_{\text{DD}}/T_{\text{R}} \quad (3.40)$$

$$P_{\text{OUT}} = \frac{1}{2}C_{\text{OUT}}(V_{\text{PP}}^2 - V_{\text{G}}^2)/T_{\text{R}} \quad (3.41)$$

$$R_{\text{PWR}} = \frac{V_{\text{PP}} + V_{\text{G}}}{2(N + 1)V_{\text{DD}}} \quad (3.42)$$

where  $V_{\text{G}}$  is  $V_{\text{DD}} - V_{\text{T}}$ .

### 3.1.4 Body Effect of Transfer Transistors

$C_{\text{PMP}}$ ,  $V_{\text{OUT}}(j)$ , and  $T_{\text{R}}$  can be expanded in case where the body effect of transfer transistors should be taken into account in the cut-off condition. Following the approach that Witters et al. made, the body effect of transfer transistors is expressed by a parameter  $\alpha$  as

$$V_{\text{S}} = \alpha(V_{\text{D}} - V_{\text{T}}) \quad (3.43)$$

where  $V_{\text{S}}$  is the source follower voltage,  $V_{\text{D}}$  is the voltage applied on the drain terminal which is shorted to the gate terminal, and  $V_{\text{T}}$  is the threshold voltage at no back bias. Thus, the following equations hold:

$$Q(1) = \alpha C V_{\text{G}} \quad (3.44)$$

$$Q(2k-1) = \sum_{i=1}^{2k-1} \alpha^i (C V_{\text{G}} - q_{\text{OUT}}) - \alpha^{2(k-1)} q_{\text{OUT}} \quad (3.45)$$

$$Q(2k) = \sum_{i=1}^{2k} \alpha^i \left( C V_{\text{G}} - \frac{q_{\text{OUT}}}{\alpha} \right) \quad (3.46)$$

$$Q(N) = C \left( \frac{V_{\text{OUT}}}{\alpha} - V_{\text{G}} \right) \quad (3.47)$$

Therefore, the output voltage–current characteristic with the body effect of transfer transistors is derived by

$$q_{\text{OUT}} = C \left( \sum_{i=1}^{N+1} \alpha^i V_{\text{G}} - V_{\text{OUT}} \right) / \sum_{i=1}^N \alpha^i \quad (3.48)$$

In this case, the recurrence formula for the output voltage holds as follows:

$$\begin{aligned} & C_{\text{OUT}}(V_{\text{OUT}}(j+1) - V_{\text{OUT}}(j)) \\ &= C \left( \sum_{i=1}^{N+1} \alpha^i V_{\text{G}} - V_{\text{OUT}}(j+1) \right) / \sum_{i=1}^N \alpha^i \end{aligned} \quad (3.49)$$

where the self-load capacitance  $C_{\text{PMP}}$  included in the total load capacitance  $C_{\text{OUT}}$  is, respectively, expressed for even and odd  $N$  by

$$C_{\text{PMP}} = \frac{1}{(N+1)(1-\alpha^N)} \left[ \frac{\alpha N^2 + (N+1)^2 - 1}{4\alpha} - \frac{1 - (N+1)\alpha^N + N\alpha^{N+1}}{(1-\alpha)^2} \right] C \quad (3.50a)$$

$$C_{\text{PMP}} = \frac{1}{(N+1)(1-\alpha^N)} \left[ \frac{\alpha(N+1)^2 + N^2 - 1}{4\alpha} - \frac{1 - (N+1)\alpha^N + N\alpha^{N+1}}{(1-\alpha)^2} \right] C \quad (3.50b)$$

Using the initial condition of  $V_{\text{OUT}}(0) = V_G$ , Eq. (3.49) is solved as

$$V_{\text{OUT}}(j) = \sum_{i=1}^{N+1} \alpha^i V_G - \left( \sum_{i=1}^{N+1} \alpha^i - 1 \right) V_G \beta^j \quad (3.51)$$

Therefore, the rise time that the output voltage rises from  $V_G$  to  $V_{\text{PP}}$  is

$$T_R = \ln \left[ 1 - (V_{\text{PP}} - V_G) / \left( \sum_{i=1}^{N+1} \alpha^i - 1 \right) V_G \right] / \ln \beta \quad (3.52)$$

If the transistors do not suffer from the body effect, i.e.,  $\alpha = 1$ , Eqs. (3.51) and (3.52) reduce to Eqs. (3.17) and (3.19) in case of  $\alpha_T = \alpha_B = 0$ , respectively.

## 3.2 Circuit Analysis Under Medium- to High-Frequency Operation

It has been assumed so far that all the diodes turn off at the end of a half period and all the amount of charge are transferred to the next capacitor. This subsection discusses more realistic case where this assumption is not valid to figure out an optimum clock frequency.

When the clock frequency is low enough, the output current is proportional to the clock frequency. Figure 3.9c indicates that the output charge per a half period is never affected when the clock frequency is slower than 5 MHz because there is no current at 100 ns. When the clock frequency increases to 25 MHz, the current at 20 ns is finite, as shown in Fig. 3.9d. One can guess that you may never gain the output current with a faster clock than 25 MHz. At 100 MHz, there may be little chance to transfer any charge, as shown in Fig. 3.9e. As a result, the frequency vs. output current curve could be a quadratic function in a real pump as shown in Fig. 3.9b.

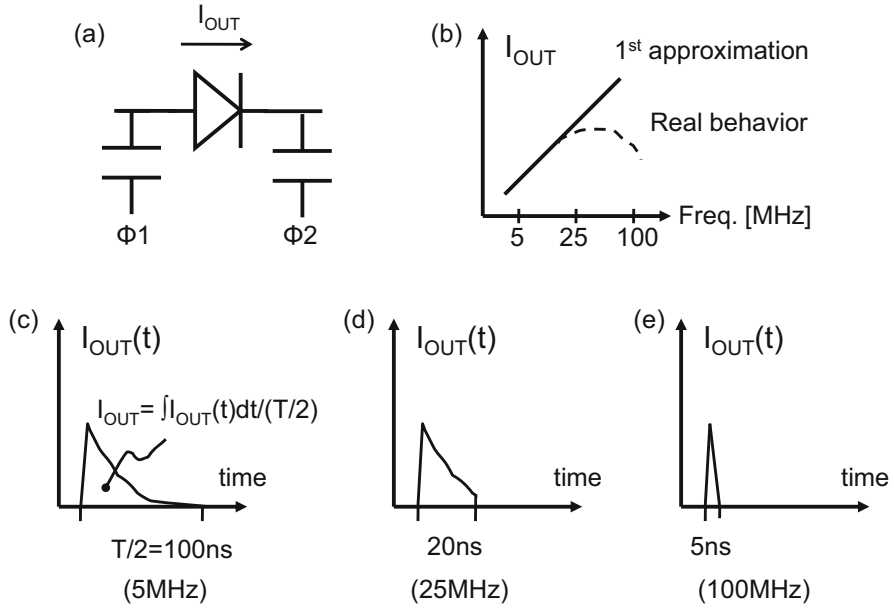


Fig. 3.9 Ideal and real frequency response to the output current

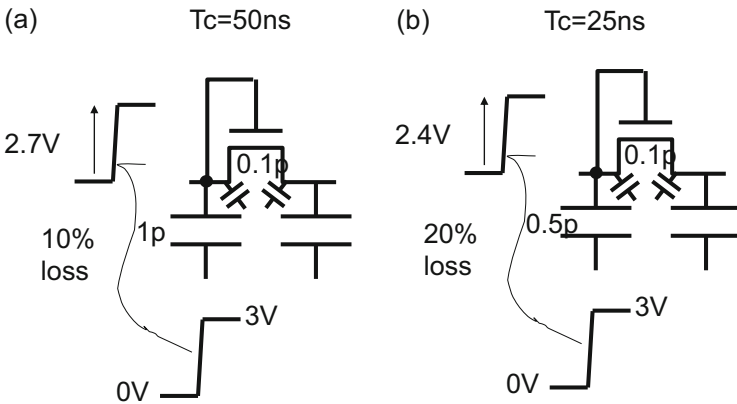


Fig. 3.10 Scalability in frequency

Circuit designers have to reduce the circuit area as much as possible for a small die size. It is possible to cut the capacitor partially if one uses a fast clock. The question is how much capacitor area can cut to maintain the output current with a faster clock. Let  $C$  and  $C_T$  be the capacitance of the pump capacitor and parasitic capacitance, respectively.  $\alpha_T$  is defined by  $C_T/C$ . Assume  $C$  is 1 pF and  $C_T$  is 0.1 pF for 50 ns clock, as shown in Fig. 3.10a. The effective voltage amplitude of the capacitor would be 10 % lower than the supply voltage due to  $\alpha_T$  of 10 %. When one



considers that the clock frequency increases twice, whereas the capacitor decreases half to reduce the total pump area with a faster clock and without changing the size of transfer gate, the effective clock amplitude could be reduced by 0.3 V due to an increased  $\alpha_T$  of 20 %, as shown in Fig. 3.10b. This means that this scenario is broken.

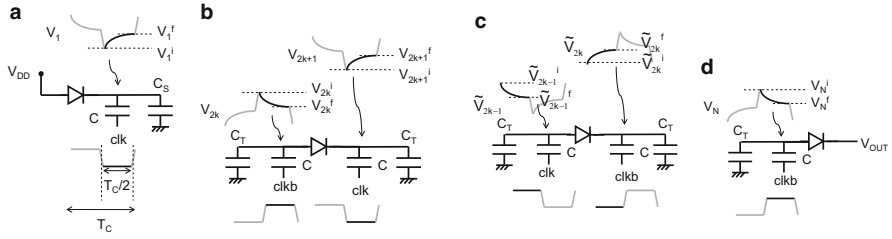
Circuit analysis and modeling has been addressed in previous subsection in cases where the charges are fully transferred in every half period. To minimize the silicon area for the pump that outputs a required current at a given output voltage, a faster clock frequency is preferred. However, the output current could decrease while the clock is running faster than a critical point at which the timing margin for four nonoverlapped clocks of the clocks is no longer negligible to the period while the switches turn on. This part discusses a switch-resistance-aware Dickson charge pump model. Equations between  $V_{OUT}$  and  $I_{OUT}$  and between the input current  $I_{IN}$  and  $I_{OUT}$  are determined in cases where the charges are not fully transferred due to the resistance of the switches ( $R$ ). In Sect. 3.2, the impact of  $R$  on  $I_{OUT}$  is investigated and optimization of the clock frequency and the transistor size are presented to maximize  $I_{OUT}$  under a given circuit area in a given technology.

Recently, energy harvesting is becoming increasingly important for autonomous sensor networks and implantable electronic devices. Photovoltaic (PV) cells and thermoelectric generators (TEGs) are used as DC power sources. In some applications where a small form factor is a prime concern, a Dickson charge pump DC-DC switched-capacitor multiplier can be integrated into a power management circuit to (1) eliminate the need for an inductor and (2) enable DC voltage of energy transducers as low as 0.5 V to supply a higher voltage such 1.5 V to digital and analog circuits, as shown in Fig. 1.11 of Chap. 1. To minimize the operation current for control circuits and the oscillator, the switches are realized by simple diodes, e.g., p-n diodes, Schottky barrier diodes, or metal–oxide–semiconductor field-effect transistors (MOSFETs), whose gate terminals are connected with their drain terminals.

Section 3.2.1 aims at providing another output voltage–current equation to design a charge pump that inputs a low DC voltage and uses diodes as switching devices for energy harvesting. One can estimate the impact of the electrical characteristics of the switching diodes or MOSFETs on the charge pump performance.

### 3.2.1 DC-DC Charge Pump Using Switching Diodes

Figure 3.11a–d show several combinations among the pump in steady state, where  $V_{DD}$  is the supply voltage,  $C$  the main capacitor per stage,  $C_T$  the stray capacitance at the top plate of the main capacitor,  $T_C$  the period of the clock,  $V_k = V_k(t)$  the voltage at node  $k$  ( $1 \leq k \leq N$ ),  $N$  the number of stages,  $V_k^i$  the initial voltage of  $V_k$  in the first half of the period,  $V_k^f$  the final voltage of  $V_k$  in the first half of the period,  $\tilde{V}_k^{i(f)}$  the initial (final) voltage of  $V_k$  in the second half of the period, and  $V_{OUT}$  the



**Fig. 3.11** Four representative combinations between neighboring capacitors. (a) first, (b) 2  $k$ -th to (2  $k$  + 1)-th stages, and (d) last stages, in the first half of period, followed by (c) (2  $k$  – 1)-th to 2  $k$ -th stages in the second half of the period

output voltage. Each diode has a voltage ( $V_{\text{DIO}}$ ) – current ( $I_{\text{DIO}}$ ) characteristic as given by Eq. (3.53), where  $I_S$  is the saturation current and  $V_T$  is the parameter determining the slope in the  $V_{\text{DIO}} - \log(I_{\text{DIO}})$  plot, which is proportional to the thermal voltage ( $kT/q$ ):

$$I_{\text{DIO}} = I_S \exp\left(\frac{V_{\text{DIO}}}{V_T}\right) \quad (3.53)$$

Here, the parasitic series resistance of the diode is assumed to be small enough. This is valid as far as the operating points are limited in the linear region of the diode current.

The parasitic resistance of the capacitor and the impedance of clocks  $clk$  and  $clkb$  are also assumed to be negligibly small in comparison with the switching device. From Fig. 3.11a–d, Eqs. (3.54a)–(3.54d) hold:

$$C(1 + \alpha_T) \frac{dV_1}{dt} = I_S \exp\left(\frac{V_{\text{DD}} - V_1}{V_T}\right) \quad (3.54a)$$

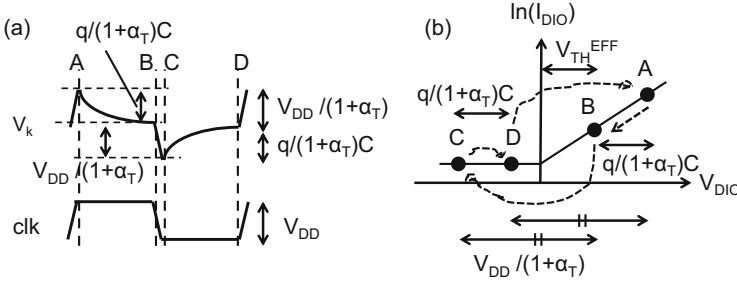
$$C(1 + \alpha_T) \frac{dV_{2k+1}}{dt} = -C(1 + \alpha_T) \frac{dV_{2k}}{dt} = I_S \exp\left(\frac{V_{2k} - V_{2k+1}}{V_T}\right) \quad (3.54b)$$

$$C(1 + \alpha_T) \frac{d\tilde{V}_{2k}}{dt} = -C(1 + \alpha_T) \frac{d\tilde{V}_{2k-1}}{dt} = I_S \exp\left(\frac{\tilde{V}_{2k-1} - \tilde{V}_{2k}}{V_T}\right) \quad (3.54c)$$

$$C(1 + \alpha_T) \frac{dV_N}{dt} = -I_S \exp\left(\frac{V_N - V_{\text{OUT}}}{V_T}\right) \quad (3.54d)$$

where  $\alpha_T$  is  $C_T/C$  which represents the gate overdrive loss.

Figure 3.12 shows a trajectory of a capacitor node in the I–V plane. The diode flows a reverse current when  $V_{\text{DIO}}$  is negative, which is assumed to be negligibly small in this paper. The diode starts conducting at the clock rise edge as shown by point A. According to Eq. (3.53), an amount of charges is transferred to a next stage. The capacitor voltage decreases by  $q/(1 + \alpha_T)C$ , where  $q$  is the charge transferred



**Fig. 3.12** Trajectory of a capacitor node in the I-V plane

to the output terminal in a period. At the falling edge B, one can have a finite residual potential as shown by  $V_{TH}^{EFF}$  which represents the effective threshold voltage. A voltage swing of  $V_{DD}$  in clk translates into that of the nodal voltage of  $V_{DD}/(1 + \alpha_T)$  in the transition from point B to C. Charge transfer from the previous stage increases the capacitor voltage the same as  $q/(1 + \alpha_T)C$  in the transition from point C to D. Another voltage swing of  $V_{DD}$  in clk translates into that of the nodal voltage of  $V_{DD}/(1 + \alpha_T)$  in the transition from point D to the original A. Thus, the initial and final voltages at each capacitor node in each half period are connected to one another as follows:

$$\tilde{V}_{2k-1}^f - \frac{V_{DD}}{1 + \alpha_T} = V_{2k-1}^i \quad (3.55a)$$

$$\tilde{V}_{2k-1}^i = V_{2k-1}^f + \frac{V_{DD}}{1 + \alpha_T} \quad (3.55b)$$

$$\tilde{V}_{2k}^f + \frac{V_{DD}}{1 + \alpha_T} = V_{2k}^i \quad (3.55c)$$

$$\tilde{V}_{2k}^i = V_{2k}^f - \frac{V_{DD}}{1 + \alpha_T} \quad (3.55d)$$

Steady state indicates the following relations at each node:

$$V_1^i = V_1^f - \frac{q}{C(1 + \alpha_T)} \quad (3.56a)$$

$$V_{2k}^i = V_{2k}^f + \frac{q}{C(1 + \alpha_T)} \quad (3.56b)$$

$$V_{2k+1}^i = V_{2k+1}^f - \frac{q}{C(1 + \alpha_T)} \quad (3.56c)$$

$$V_N^i = V_N^f + \frac{q}{C(1 + \alpha_T)} \quad (3.56d)$$

where  $N$  is assumed to be an even number. However, the final results do not depend on whether  $N$  is even or odd.

From the conditions where  $V_1(0) = V_1^i$  and  $V_1(T_C/2) = V_1^f$ , Eq. (3.54a) is solved to be Eq. (3.57a) using Eq. (3.56a):

$$\exp\left(\frac{V_1^f}{V_T}\right) = \frac{T_C I_S}{2(1 + \alpha_T) C V_T} \exp\left(\frac{V_{DD}}{V_T}\right) / \left(1 - \exp\left(-\frac{q}{(1 + \alpha_T) C V_T}\right)\right) \quad (3.57a)$$

Similarly, Eqs. (3.54b)–(3.54d), respectively, result in Eqs. (3.57b)–(3.57d) using Eqs. (3.55a)–(3.55d) and (3.56b)–(3.56d):

$$\exp\left(\frac{V_{2k+1}^f - V_{2k}^f}{V_T}\right) = \frac{T_C I_S}{(1 + \alpha_T) C V_T} / \left(1 - \exp\left(-\frac{2q}{(1 + \alpha_T) C V_T}\right)\right) \quad (3.57b)$$

$$\exp\left(\frac{V_{2k}^f - V_{2k-1}^f}{V_T}\right) = \frac{T_C I_S}{(1 + \alpha_T) C V_T} \exp\left(\frac{2C V_{DD} - 2q}{(1 + \alpha_T) C V_T}\right) / \left(1 - \exp\left(-\frac{2q}{(1 + \alpha_T) C V_T}\right)\right) \quad (3.57c)$$

$$\exp\left(\frac{V_{OUT} - V_N^f}{V_T}\right) = \frac{T_C I_S}{2(1 + \alpha_T) C V_T} / \left(1 - \exp\left(-\frac{q}{(1 + \alpha_T) C V_T}\right)\right) \quad (3.57d)$$

From the multiplication of the following  $(N + 1)$  equations,

$$\begin{aligned} \exp\left(\frac{V_1^f}{V_T}\right) &= \frac{T_C I_S}{2(1 + \alpha_T) C V_T} \exp\left(\frac{V_{DD}}{V_T}\right) / \left(1 - \exp\left(-\frac{q}{(1 + \alpha_T) C V_T}\right)\right) \\ \exp\left(\frac{V_2^f - V_1^f}{V_T}\right) &= \frac{T_C I_S}{(1 + \alpha_T) C V_T} \exp\left(\frac{2C V_{DD} - 2q}{(1 + \alpha_T) C V_T}\right) / \left(1 - \exp\left(-\frac{2q}{(1 + \alpha_T) C V_T}\right)\right) \\ \exp\left(\frac{V_3^f - V_2^f}{V_T}\right) &= \frac{T_C I_S}{(1 + \alpha_T) C V_T} / \left(1 - \exp\left(-\frac{2q}{(1 + \alpha_T) C V_T}\right)\right) \\ &\dots \\ \exp\left(\frac{V_{2k}^f - V_{2k-1}^f}{V_T}\right) &= \frac{T_C I_S}{(1 + \alpha_T) C V_T} \exp\left(\frac{2C V_{DD} - 2q}{(1 + \alpha_T) C V_T}\right) / \left(1 - \exp\left(-\frac{2q}{(1 + \alpha_T) C V_T}\right)\right) \\ \exp\left(\frac{V_{2k+1}^f - V_{2k}^f}{V_T}\right) &= \frac{T_C I_S}{(1 + \alpha_T) C V_T} / \left(1 - \exp\left(-\frac{2q}{(1 + \alpha_T) C V_T}\right)\right) \\ &\dots \\ \exp\left(\frac{V_{OUT} - V_N^f}{V_T}\right) &= \frac{T_C I_S}{2(1 + \alpha_T) C V_T} / \left(1 - \exp\left(-\frac{q}{(1 + \alpha_T) C V_T}\right)\right) \end{aligned} \quad (3.58)$$

One can have  $V_{OUT} - q$  characteristic as shown in Eq. (3.59):

$$\exp\left(\frac{V_{OUT}}{V_T}\right) = \frac{1}{4} \left( \frac{T_C I_S}{(1 + \alpha_T) C V_T} \right)^{N+1} \exp\left(\frac{(N + 1 + \alpha_T) C V_{DD} - Nq}{(1 + \alpha_T) C V_T}\right) \times \left(1 - \exp\left(-\frac{q}{(1 + \alpha_T) C V_T}\right)\right)^{-2} \left(1 - \exp\left(-\frac{2q}{(1 + \alpha_T) C V_T}\right)\right)^{-N+1} \quad (3.59)$$

The amount of charges  $q$  is related to an averaged output current  $I_{OUT}$ , where  $f$  is the clock frequency:

$$q = I_{OUT} T_C = I_{OUT} / f \quad (3.60)$$

Equations (3.59) and (3.60) provide the  $V_{OUT} - I_{OUT}$  characteristic. When  $q \gg (1 + \beta) C V_T$ , Eqs. (3.59) and (3.60) are simplified to be Eq. (3.61):

$$I_{OUT} = (V_{MAX} - V_{OUT}) / R_{PMP} \quad (3.61)$$

where the effective output impedance of the charge pump  $R_{PMP}$  and the maximum attainable voltage  $V_{MAX}$  are given by Eqs. (3.62) and (3.63a), respectively:

$$R_{PMP} = \frac{N}{(1 + \alpha_T) C f} \quad (3.62)$$

$$V_{MAX} = \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1) V_{TH}^{EFF} \quad (3.63a)$$

The effective threshold voltage  $V_{TH}^{EFF}$  is given by Eq. (3.63b):

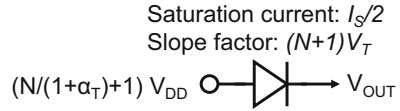
$$V_{TH}^{EFF} = V_T \ln \left( 4^{\frac{1}{N+1}} \frac{(1 + \alpha_T) f C V_T}{I_S} \right) \quad (3.63b)$$

Equations (3.61)–(3.63) show the same circuit equation as the original one Eq. (2.75) under a low operation frequency as discussed in Sect. 2.5 or in slow switching limit (SSL) condition. Thus, Eq. (3.59) includes the original model as an approximation when  $q \gg (1 + \alpha_T) C V_T$ . Using Eq. (3.63b), one can estimate the effective threshold voltage when characteristics of the diodes and design parameters are given. A linearized equivalent circuit model is shown in Fig. 3.3.

Conversely, when the operating frequency is too high to transfer charges from one stage to the next completely, Eq. (3.59) is approximated to Eq. (3.64) in a fast switching limit (FSL) condition:

$$I_{OUT} = \frac{I_S}{2} \exp \left[ \frac{(N / (1 + \alpha_T) + 1) V_{DD} - V_{OUT}}{V_T (N + 1)} \right] \quad (3.64)$$

**Fig. 3.13** Equivalent circuit of a charge pump with switching diodes in fast switching limit condition



**Table 3.1** Device and circuit parameters used for verification

Parameter	Diode	MOSFET
$T_C$	1 ns to 1 $\mu$ s	
$V_{DD}$	0.5 V	
$V_{PP}$	0 V to 5 V	
$C$	1 pF to 10 pF	
$I_S$	200 pA to 20 nA	40 nA
$V_T$	26 mV	50 mV
$k$	–	0.9 mA/V <sup>2</sup>
$V_{TH}$	–	0.2 V
$\alpha_T$	0.05	
$N$	12	

**Table 3.2** SPICE diode model parameters used for verification

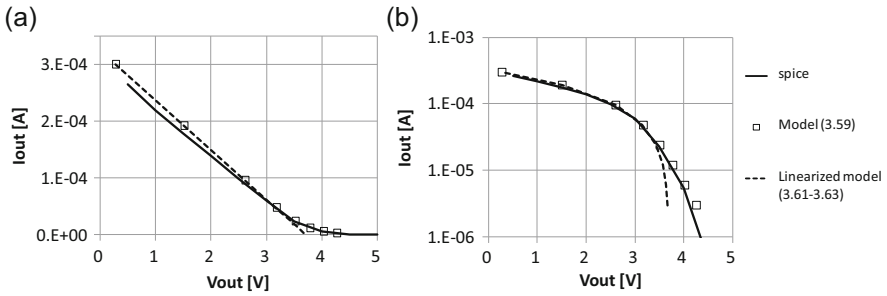
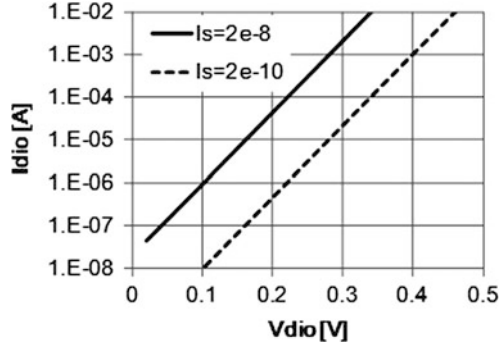
Name	Parameter	Value	Unit
Level	Junction diode model	1	No unit.
EG	Band-gap energy	0.69	eV
$N_P$	Emission coefficient	1	No unit
RS	Parasitic resistance	0.1	Ohm
IS	Saturation current	1e-8, 1e-10	A

Equation (3.64) indicates an equivalent circuit in FSL as shown in Fig. 3.13. Equation (3.64) is the same equation as the diode current–voltage equation Eq. (3.53) when one replace the original parameters  $I_S$ ,  $V_{DIO}$ , and  $V_T$  with  $I_S/2$ ,  $V_{MAX}-V_{OUT}$ , and  $(N+1)V_T$ , respectively. A factor of 2 comes from the fact that the output current flows in half cycle. A factor of  $(N+1)$  comes from the fact that there are  $(N+1)$  diodes connected in series between the input and output terminals. Equation (3.64) has no frequency term. Thus, in FSL, the output current is determined by the characteristics of switching diodes.

To verify the model equation, SPICE simulations were run using the parameters shown in the column “Diode” of Table 3.1. The diode and MOSFET models are, respectively, based on a Schottky barrier diode as shown in Table 3.2. The rise and fall time of 0.1 ns was used. Figure 3.14 shows  $V_{DIO} - \log(I_{DIO})$  used in the simulation.

Figure 3.15 compares the model Eq. (3.59) and the linearized model Eqs. (3.61–3.63) with SPICE. The slopes are well matched with one another except for low current regime. On the other hand, the maximum attainable voltage which is defined by the X intercept for the linearized model Eqs. (3.61)–(3.63) is smaller than those of the original model Eq. (3.59) and SPICE result. This is because the effective threshold voltage decreases as the operation current decreases. However,

**Fig. 3.14**  $V_{DIO}-\log(I_{DIO})$  used in the SPICE simulation



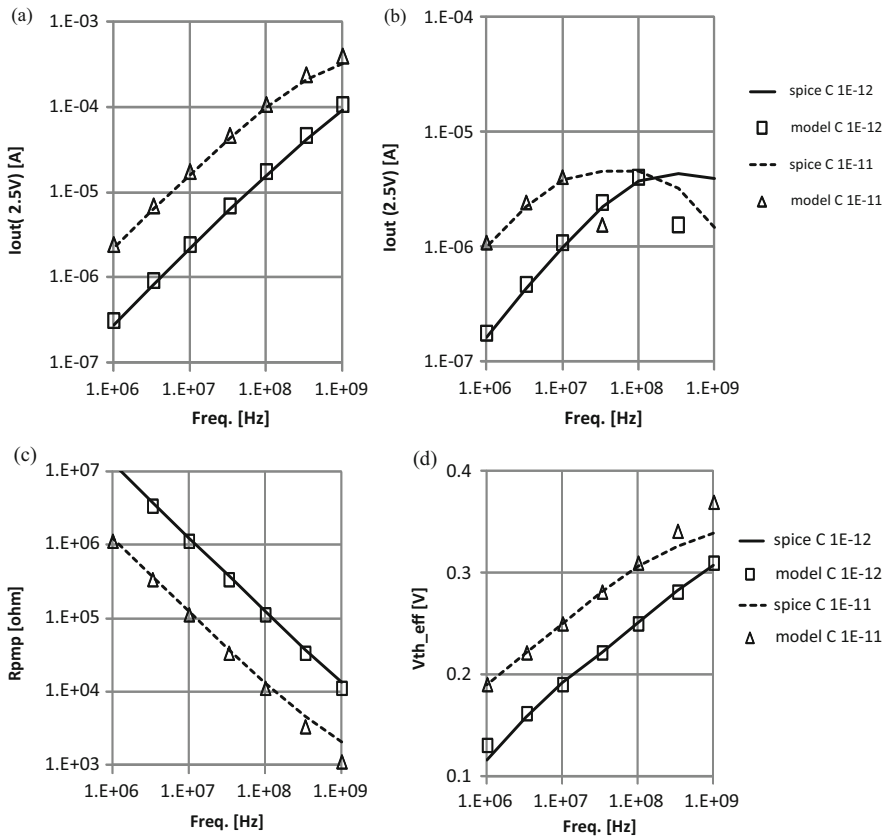
**Fig. 3.15**  $V_{OUT}$  vs.  $I_{OUT}$  when  $I_S = 20$  nA,  $C = 10$  pF, Freq. = 100 MHz in linear (a) and log (b) plots

in a nominal regime where the designed output voltage is approximately one half of the maximum attainable output voltage so that the output power is maximized as shown in Sect. 3.4 later, one can consider that the linearized model is sufficiently accurate.

Figure 3.16a, b, respectively, compares the linearized model Eqs. (3.61)–(3.63) with simulations in terms of  $V_{OUT} - I_{OUT}$  when  $I_S = 20$  nA and 200 pA. All data points for  $I_S = 20$  nA are matched well whereas some data points at higher frequencies for  $I_S = 200$  pA in the model are unmatched. Because they became negative, the data points such as  $C = 10$  pF and Freq.  $> 100$  MHz were omitted from Fig. 3.16b.

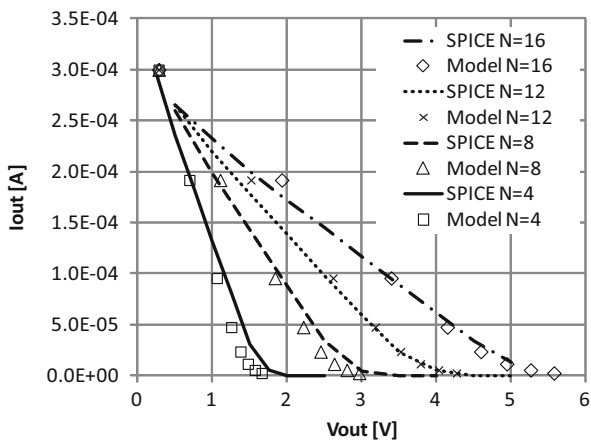
To see which of  $R_{PMP}$  or  $V_{TH}^{EFF}$  starts deviating from those of simulated results, the parameters were extracted from  $V_{OUT} - I_{OUT}$  lines, as shown in Fig. 3.16c, d.  $V_{TH}^{EFF}$  approaches  $V_{DD}$  of 0.5 V as the frequency increases in case of  $C = 10$  pF. Because the forward current through the switching diode decreases, the reverse leakage becomes relatively significant in the SPICE simulation. Conversely,  $V_{TH}^{EFF}$  continuously increases in the proposed model. As a result,  $I_{OUT}$  becomes negative in an extremely high-frequency range.

As shown in Fig. 3.17, the model Eq. (3.59) is also in good agreement with SPICE results for various numbers of stages.



**Fig. 3.16** Frequency vs.  $I_{OUT}$  when  $I_S = 20$  nA (a), 200pA (b) at  $V_{OUT} = 2.5$  V. Frequency vs.  $R_{PMP}$  (c),  $V_{TH}^{EFF}$  (d) when  $I_S = 200$ pA,  $V_{OUT} = 2.5$  V

**Fig. 3.17**  $V_{OUT}$  vs.  $I_{OUT}$  comparison of model Eq. (3.59) with SPICE when  $I_S = 20$  nA,  $C = 10$  pF, Freq. = 100 MHz,  $N = 4, 8, 12, 16$





### 3.2.2 DC-DC Charge Pump Using Switching MOSFET in Saturation Region

The switching device can be MOSFET whose gate is connected with the drain. This subsection discusses the case where one needs to use Eq. (3.65) instead of Eq. (3.53):

$$I_{\text{MOS}} = k(V_{\text{MOS}} - V_{\text{TH}})^2 \quad (3.65)$$

where  $V_{\text{MOS}}$  is the drain to source voltage,  $I_{\text{MOS}}$  is the drain to source current,  $k$  is transconductance, and  $V_{\text{TH}}$  is the threshold voltage of the MOSFET. One needs to use the differential equations Eqs. (3.66a–3.66d) instead of Eqs. (3.54a–3.54d):

$$C(1 + \alpha_{\text{T}}) \frac{dV_1}{dt} = k(V_{\text{DD}} - V_1 - V_{\text{TH}_{1}})^2 \quad (3.66a)$$

$$C(1 + \alpha_{\text{T}}) \frac{dV_{2k+1}}{dt} = -C(1 + \alpha_{\text{T}}) \frac{dV_{2k}}{dt} = k(V_{2k} - V_{2k+1} - V_{\text{TH}_{2k+1}})^2 \quad (3.66b)$$

$$C(1 + \beta) \frac{d\tilde{V}_{2k}}{dt} = -C(1 + \beta) \frac{d\tilde{V}_{2k-1}}{dt} = k(\tilde{V}_{2k-1} - \tilde{V}_{2k} - V_{\text{TH}_{2k}})^2 \quad (3.66c)$$

$$C(1 + \alpha_{\text{T}}) \frac{dV_{\text{N}}}{dt} = -k(V_{\text{N}} - V_{\text{OUT}} - V_{\text{TH}_{\text{N}+1}})^2 \quad (3.66d)$$

Using Eqs. (3.55a–3.55d) and Eqs. (3.56a–3.56d), Eqs. (3.65a–3.65d) are solved to be Eqs. (3.67a–3.67d), respectively:

$$V_{\text{DD}} - V_1^{\text{f}} - V_{\text{TH}_{1}} = \frac{V_{\text{OD1}}}{2} \quad (3.67a)$$

$$V_{2k-1}^{\text{f}} - V_{2k}^{\text{f}} = V_{\text{OD2}} + \frac{2q}{(1 + \alpha_{\text{T}})C} - \frac{2V_{\text{DD}}}{1 + \alpha_{\text{T}}} + V_{\text{TH}_{2k}} \quad (3.67b)$$

$$V_{2k}^{\text{f}} - V_{2k+1}^{\text{f}} = V_{\text{OD2}} + V_{\text{TH}_{2k+1}} \quad (3.67c)$$

$$V_{\text{N}}^{\text{f}} - V_{\text{OUT}} - V_{\text{TH}_{\text{N}+1}} = \frac{V_{\text{OD1}}}{2} \quad (3.67d)$$

where the effective overdrive voltages  $V_{\text{OD1}}$  and  $V_{\text{OD2}}$  are given by Eqs. (3.68a) and (3.68b), respectively:

$$V_{\text{OD1}} = \sqrt{\frac{q^2}{(1 + \alpha_{\text{T}})^2 C^2} + \frac{8q}{kT_{\text{C}}}} - \frac{q}{(1 + \alpha_{\text{T}})C} \quad (3.68a)$$

$$V_{\text{OD2}} = \sqrt{\frac{q^2}{(1 + \alpha_{\text{T}})^2 C^2} + \frac{2q}{kT_{\text{C}}}} - \frac{q}{(1 + \alpha_{\text{T}})C} \quad (3.68b)$$

Adding Eqs. (3.67a–3.67d) with respect to  $k = 1 \dots N/2$  results in Eq. (3.69):

$$\begin{aligned} & \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1) V_{TH}^{AVG} - V_{OUT} \\ &= (N - 1) \sqrt{\frac{q^2}{(1 + \alpha_T)^2 C^2} + \frac{2fq}{k}} + \sqrt{\frac{q^2}{(1 + \alpha_T)^2 C^2} + \frac{8fq}{k}} \end{aligned} \quad (3.69)$$

where the average  $V_{TH}^{AVG}$  is

$$V_{TH}^{AVG} = \sum_{k=1}^{N+1} V_{TH,k} / (N + 1) \quad (3.70)$$

Equations (3.69) and (3.60) provide the  $V_{OUT} - I_{OUT}$  characteristic of the square model. When the clock frequency is very low, Eq. (3.69) is approximated to be Eq. (3.71), where  $R_{PMP}$  and  $V_{MAX}$  are given by Eq. (3.72) and Eq. (3.73), respectively. Equations (3.71)–(3.73) are the same circuit equations as the original ones under a low operation frequency as discussed in Sect. 3.1, where the effective threshold voltage is given by Eq. (3.70):

$$I_{OUT} = (V_{MAX} - V_{OUT}) / R_{PMP} \quad (3.71)$$

$$R_{PMP} = \frac{N}{(1 + \alpha_T) C f} \quad (3.72)$$

$$V_{MAX} = \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1) V_{TH}^{AVG} \quad (3.73)$$

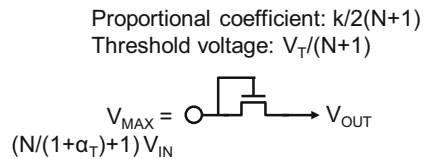
Conversely, when the clock frequency is very high, Eq. (3.69) is approximated to be Eq. (3.74), where the effective voltage gain per stage  $V_G^{EFF}$  is given by Eq. (3.75):

$$I_{OUT} = k (V_G^{EFF} - V_{TH}^{AVG})^2 / 2 \quad (3.74)$$

$$V_G^{EFF} = \left( \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - V_{OUT} \right) / (N + 1) \quad (3.75)$$

Equation (3.74) indicates an equivalent circuit in FSL as shown in Fig. 3.18. Equation (3.74) is the same equation as the current–voltage equation Eq. (3.65)

**Fig. 3.18** Equivalent circuit of a charge pump with switching FETs in fast switching limit condition



when one replace the original parameters  $k$ ,  $V_{\text{MOS}}$ , and  $V_{\text{TH}}$  with  $k/2(N+1)$ ,  $V_{\text{MAX}}-V_{\text{OUT}}$ , and  $(N+1)V_{\text{TH}}^{\text{AVG}}$ , respectively. A factor of 2 comes from the fact that the output current flows in half cycle. A factor of  $(N+1)$  comes from the fact that there are  $(N+1)$  FETs connected in series between the input and output terminals. Equation (3.74) has no frequency term. Thus, in FSL, the output current is determined by the characteristics of switching FETs.

The square model Eq. (3.69) is compared with SPICE simulation using the parameters shown in the column ‘‘MOSFET’’ of Table 3.1. Figure 3.19a, b, respectively, compares  $V_{\text{GS}} - I_{\text{DS}}$  in linear and log scale. A parameter set of  $k = 0.9 \text{ mA/V}^2$  and  $V_{\text{TH}} = 0.2 \text{ V}$  fits the SPICE model well at  $V_{\text{GS}}$  of 0.2 V and above.

At  $V_{\text{GS}}$  of 0.3 V or below, an exponential curve with a parameter set of  $I_{\text{S}} = 40 \text{ nA}$  and  $V_{\text{T}} = 50 \text{ mV}$  fits well as shown in Fig. 3.19b. Figure 3.19c shows the dependency of  $V_{\text{TH}}$  on  $V_{\text{S}}$ . Because the body of the MOSFETs is grounded in this experiment, the threshold voltage of  $k$ -th stage,  $V_{\text{TH}_k}$ , is approximately given by

$$V_{\text{TH}_k} = V_{\text{TH}} \left( V_{\text{S}} = V_{\text{DD}} + \frac{k}{N}(V_{\text{PP}} - V_{\text{DD}}) \right) \quad (3.76)$$

Equation (3.76) can be used for Eq. (3.69). When  $V_{\text{TH}} - V_{\text{S}}$  slope is linear as shown in Fig. 3.19c, one can simply calculate Eq. (3.76) using the averaged  $V_{\text{TH}}$  of the first and last MOSFETs Eq. (3.77):

$$V_{\text{TH}}^{\text{AVG}} = (V_{\text{TH}}(V_{\text{S}} = V_{\text{DD}}) + V_{\text{TH}}(V_{\text{S}} = V_{\text{PP}}))/2 \quad (3.77)$$

Figure 3.20a compares  $I_{\text{OUT}}-V_{\text{OUT}}$  curves. The exponential model Eq. (3.59) fits the SPICE result much better than the square model. Figure 3.20b, c, respectively, shows the frequency dependencies of  $R_{\text{PMP}}$  and  $V_{\text{TH}}^{\text{EFF}}$ . The exponential model Eq. (3.59) fits the SPICE result in terms of  $R_{\text{PMP}}$  when the clock frequency is not very high. The slow clock approximation of the square model Eq. (3.69) cannot have the effective threshold voltage lower than  $V_{\text{TH}}^{\text{EFF}}$  whereas the exponential model Eq. (3.59) can. Therefore, the exponential model Eq. (3.59) is closer to the SPICE result in terms of  $V_{\text{TH}}^{\text{EFF}}$  when the frequency is not very high, where the MOSFETs can enter subthreshold regime. Conversely, when the frequency is very high, the SPICE result shows that  $V_{\text{TH}}^{\text{EFF}}$  becomes a weak function of the frequency. In total, one should use the exponential model rather than the square model for the initial design stage prior to SPICE simulation when the clock frequency is not very high.

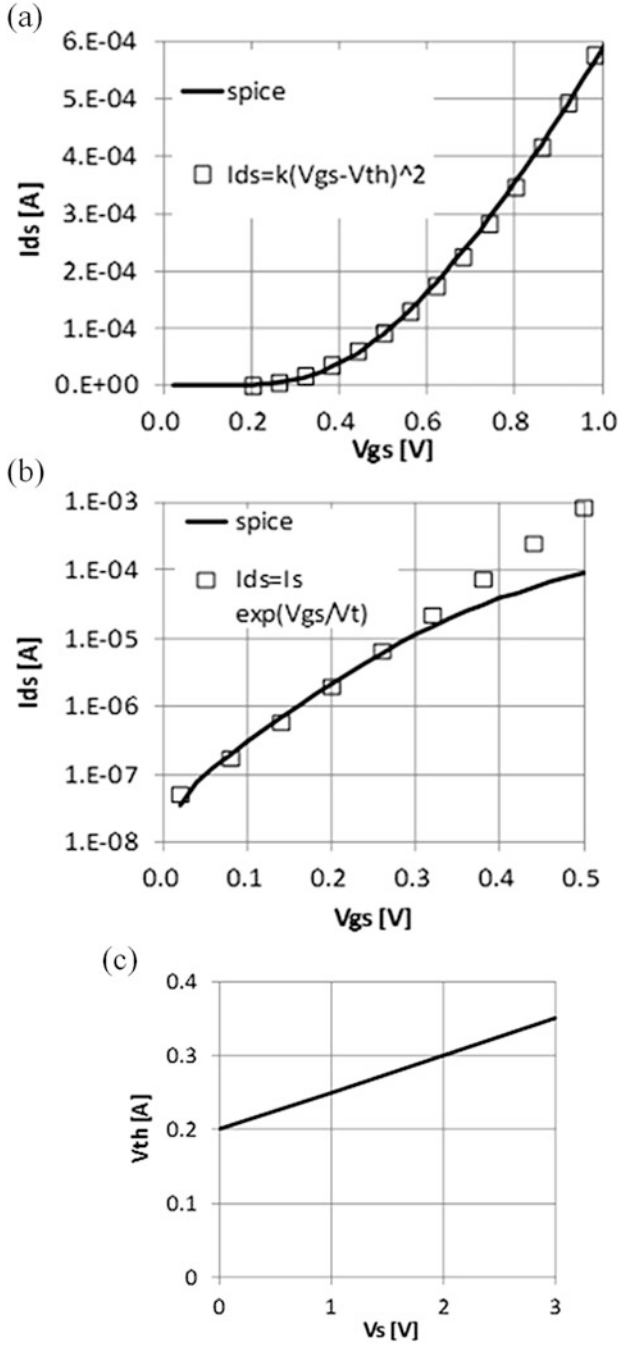
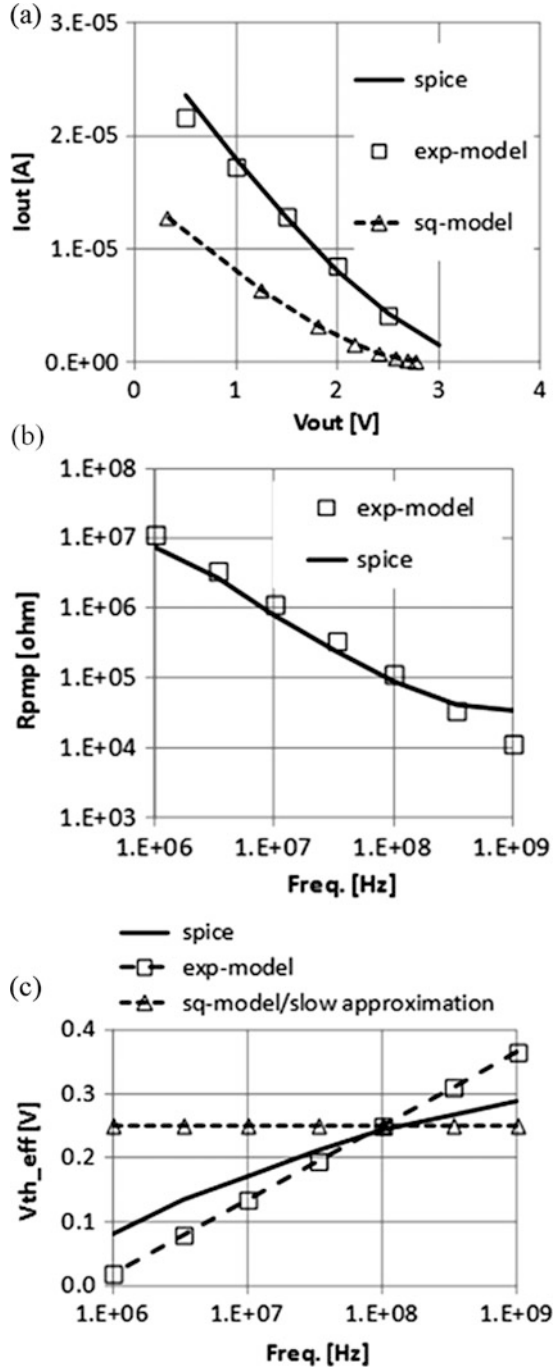


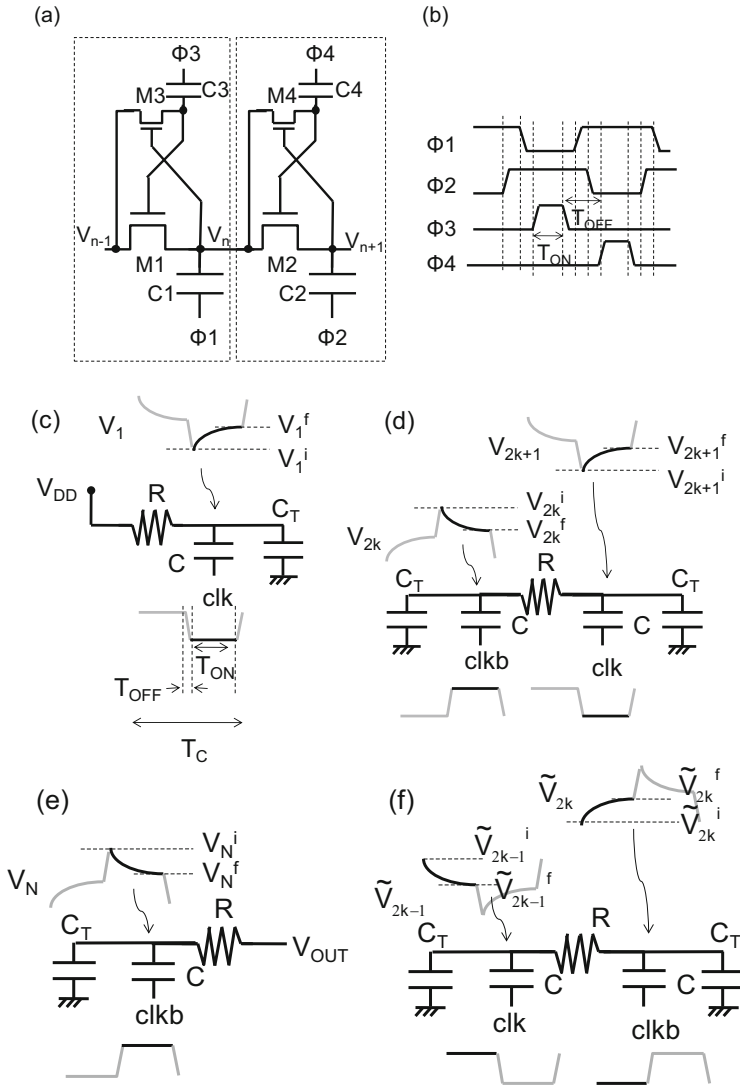
Fig. 3.19 (a)  $V_{GS} - I_{GS}$ , (b)  $V_{GS} - \log(I_{GS})$ , (c)  $V_S - V_{TH}$

**Fig. 3.20** (a)  $I_{OUT}-V_{OUT}$  when C is 1 pF and frequency is 100 MHz, and frequency vs. (b)  $R_{PMP}$  and (c)  $V_{TH}^{EFF}$  when C is 1 pF



### 3.2.3 DC-DC Charge Pump Using Switching MOSFET in Triode Region

In this subsection, a switching resistance aware (SRA) model for a charge pump using switching MOSFETs that operate in triode region is discussed. Figure 3.21a shows a Dickson charge pump circuit driven by four nonoverlapping phases as



**Fig. 3.21** Charge pump driven by four nonoverlapping phases. Four representative combinations between next-neighbor capacitors. (c) First, (d) second to third, and (e) last stages, in the first half of period, followed by (f) first to second stages in the second half of the period

shown in Fig. 3.21b whose operation is discussed in Chap. 5. The capacitors C1,2 driven by  $\Phi_{1,2}$  are main pumping capacitors, and C3,4 driven by  $\Phi_{3,4}$  are auxiliary to eliminate the effect of the threshold voltage of the transfer gates. The transfer transistor M1(2) turns on in a triode region when  $\Phi_{3(4)}$  stays high.

Figure 3.21c–f shows several combinations among the pump in steady state, where  $V_{DD}$  is the supply voltage,  $R$  the channel resistance of switches,  $C$  the main capacitor per stage,  $T_{ON}$  the period when the switch turns on,  $T_{OFF}$  the period when the switch turns off,  $T$  the period of the clock,  $V_k = V_k(t)$  the voltage at node  $k$  ( $1 \leq k \leq N$ ),  $N$  the number of stages,  $V_{ki}$  the initial voltage of  $V_k$  in the first half of the period,  $V_{kf}$  the final voltage of  $V_k$  in the first half of the period,  $\tilde{V}_{ki(f)}$  the initial (final) voltage of  $V_k$  in the second half of the period, and  $V_{OUT}$  the output voltage. Strictly speaking,  $R$  has a voltage dependency per node, but it is assumed that the switch resistance can be treated as a constant averaged in  $T_{ON}$ .  $V_T$  is neglected below assuming  $V_T$  canceling techniques presented in Chap. 5. From Fig. 3.21c–f, Eqs. (3.78a)–(3.78d) hold during  $T_{ON}$ :

$$C(1 + \alpha_T) \frac{dV_1}{dt} = \frac{V_{DD} - V_1}{R} \quad (3.78a)$$

$$C(1 + \alpha_T) \frac{dV_{2k+1}}{dt} = -C(1 + \alpha_T) \frac{dV_{2k}}{dt} = \frac{V_{2k} - V_{2k+1}}{R} \quad (3.78b)$$

$$C(1 + \alpha_T) \frac{dV_{2k}}{dt} = -C(1 + \alpha_T) \frac{dV_{2k-1}}{dt} = \frac{V_{2k-1} - V_{2k}}{R} \quad (3.78c)$$

$$C(1 + \alpha_T) \frac{dV_N}{dt} = -\frac{V_N - V_{OUT}}{R} \quad (3.78d)$$

where  $\alpha_T$  is  $C_T/C$  which represents the gate overdrive loss. The initial and final voltages at each capacitor node in each half period are connected one another as follows:

$$\tilde{V}_{2k-1}^f - \frac{V_{DD}}{1 + \alpha_T} = V_{2k-1}^i \quad (3.79a)$$

$$\tilde{V}_{2k-1}^i = V_{2k-1}^f + \frac{V_{DD}}{1 + \alpha_T} \quad (3.79b)$$

$$\tilde{V}_{2k}^f + \frac{V_{DD}}{1 + \alpha_T} = V_{2k}^i \quad (3.79c)$$

$$\tilde{V}_{2k}^i = V_{2k}^f - \frac{V_{DD}}{1 + \alpha_T} \quad (3.79d)$$

Steady state indicates the following relations at each node, where  $q$  is the charge transferred to the output terminal in a period:

$$V_1^i = V_1^f - \frac{q}{C(1 + \alpha_T)} \quad (3.80a)$$

$$V_{2k}^i = V_{2k}^f + \frac{q}{C(1 + \alpha_T)} \quad (3.80b)$$

$$V_{2k+1}^i = V_{2k+1}^f - \frac{q}{C(1 + \alpha_T)} \quad (3.80c)$$

$$V_N^i = V_N^f + \frac{q}{C(1 + \alpha_T)} \quad (3.80d)$$

where  $N$  is assumed to be an even number here. However, the final results do not depend on whether  $N$  is even or odd.

From the conditions where  $V_1(0) = V_{1i}$  and  $V_1(T_{ON}) = V_{1f}$ , Eq. (3.78a) results in Eq. (3.81a):

$$V_1^f = V_{DD} - \zeta(V_{DD} - V_1^i) \quad (3.81a)$$

$$\zeta = \exp\left(-\frac{T_{ON}}{(1 + \alpha_T)RC}\right) \quad (3.81b)$$

Similarly, Eq. (3.78b) results in the following two relations:

$$V_{2k}^f = \frac{1}{2}(V_{2k}^i + V_{2k+1}^i + \zeta^2(V_{2k}^i - V_{2k+1}^i)) \quad (3.82a)$$

$$V_{2k+1}^f = \frac{1}{2}(V_{2k}^i + V_{2k+1}^i - \zeta^2(V_{2k}^i - V_{2k+1}^i)) \quad (3.82b)$$

Equations (3.78c) and (3.78d) also result in

$$\tilde{V}_{2k-1}^f = \frac{1}{2}(\tilde{V}_{2k-1}^i + \tilde{V}_{2k}^i + \zeta^2(\tilde{V}_{2k-1}^i - \tilde{V}_{2k}^i)) \quad (3.83a)$$

$$\tilde{V}_{2k}^f = \frac{1}{2}(\tilde{V}_{2k-1}^i + \tilde{V}_{2k}^i - \zeta^2(\tilde{V}_{2k-1}^i - \tilde{V}_{2k}^i)) \quad (3.83b)$$

and

$$V_{OUT} = \frac{V_N^f - \zeta V_N^i}{1 - \zeta} \quad (3.84)$$

$V_{1f}$  is calculated to be Eqs. (3.85a) from Eqs. (3.80a) and (3.81a, 3.81b):

$$V_1^f = V_{DD} - \frac{q}{C(1 + \alpha_T)1 - \zeta} \quad (3.85a)$$

From Eqs. (3.79a), (3.79b), (3.79d), and (3.83a),

$$V_{2k-1}^i = \frac{1}{2}(V_{2k-1}^f + V_{2k}^f + \zeta^2(V_{2k-1}^f - V_{2k}^f)) - \frac{V_{DD}}{1 + \alpha_T}(1 - \zeta^2) \quad (3.85b)$$



From Eqs. (3.80c) and (3.85b),

$$V_{2k}^f = V_{2k-1}^f + A \quad (3.85c)$$

$$A = -\frac{2q}{C(1 + \alpha_T)} \frac{1}{1 - \zeta^2} + \frac{2V_{DD}}{1 + \alpha_T} \quad (3.85d)$$

From Eqs. (3.80b), (3.80c), and (3.82b),

$$V_{2k+1}^f = V_{2k}^f + B \quad (3.85e)$$

$$B = -\frac{2q}{C(1 + \alpha_T)} \frac{\zeta^2}{1 - \zeta^2} \quad (3.85f)$$

From Eqs. (3.85c) and (3.85e),

$$\begin{aligned} V_2^f &= V_1^f + A \\ V_3^f &= V_2^f + B \\ V_4^f &= V_3^f + A \\ V_5^f &= V_4^f + B \\ &\vdots \\ +) V_N^f &= V_{N-1}^f + A \\ \hline V_N^f &= V_1^f + \frac{N}{2}A + \left(\frac{N}{2} - 1\right)B \end{aligned} \quad (3.85g)$$

From Eqs. (3.85a), (3.85d), (3.85f), and Eq. (3.85g),

$$V_N^f = V_{DD} \left(1 + \frac{N}{1 + \alpha_T}\right) - \frac{q}{C(1 + \alpha_T)} \left(N \frac{1 + \zeta^2}{1 - \zeta^2} + \frac{\zeta}{1 - \zeta^2}\right) \quad (3.86)$$

From Eqs. (3.80d) and (3.84),

$$V_{OUT} = V_N^f - \frac{q}{C(1 + \alpha_T)} \frac{\zeta}{1 - \zeta^2} \quad (3.87)$$

From Eqs. (3.86) and (3.87),

$$I_{OUT} \equiv \frac{q}{T} = \frac{V_{MAX} - V_{OUT}}{R_{PMP}} \quad (3.88a)$$

where  $R_{PMP}$  and  $V_{MAX}$  are, respectively,

$$R_{PMP} = \frac{T}{C(1 + \alpha_T)} \left(N \frac{1 + \zeta^2}{1 - \zeta^2} + \frac{2\zeta}{1 - \zeta^2}\right) \quad (3.88b)$$

$$V_{\text{MAX}} = V_{\text{DD}} \left( 1 + \frac{N}{1 + \alpha_T} \right) \quad (3.88c)$$

Using the relation of  $T = 2(T_{\text{ON}} + T_{\text{OFF}})$ , one can easily calculate  $I$ - $V$  curves with  $T_{\text{ON}}$  and  $T_{\text{OFF}}$  instead of  $T$  and  $T_{\text{ON}}$ . Equation (3.88c) is same as that of the original model which is given by Eq. (2.77) and  $V_T = 0$ . (3.88b) is reduced to the original model Eqs. (2.78) or (3.89) in case where  $T_{\text{ON}}/RC(1 + \alpha_T)$  is much greater than 1 or  $\zeta$  is much less than 1:

$$R_{\text{PMP}} = \frac{NT}{(1 + \alpha_T)C} \quad (3.89)$$

According as  $T_{\text{ON}}/RC(1 + \alpha_T)$  decreases and becomes comparable to an order of 1,  $R_{\text{PMP}}$  increases, thereby  $I_{\text{OUT}}$  decreases. Since the input current is a sum of the output current multiplied by  $(N + 1)$  and the charging current to another parasitic capacitor per stage at the bottom terminal of the capacitor ( $C_B$ ), Eq. (2.89) holds.

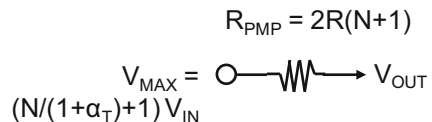
Conversely, when the clock frequency is very high, one can approximate  $T_{\text{ON}}$  and  $\zeta$  to be  $T/2$  and  $1 - T/2RC(1 + \alpha_T)$ , respectively. As a result, Eq. (3.88a) is approximated to be Eq. (3.90):

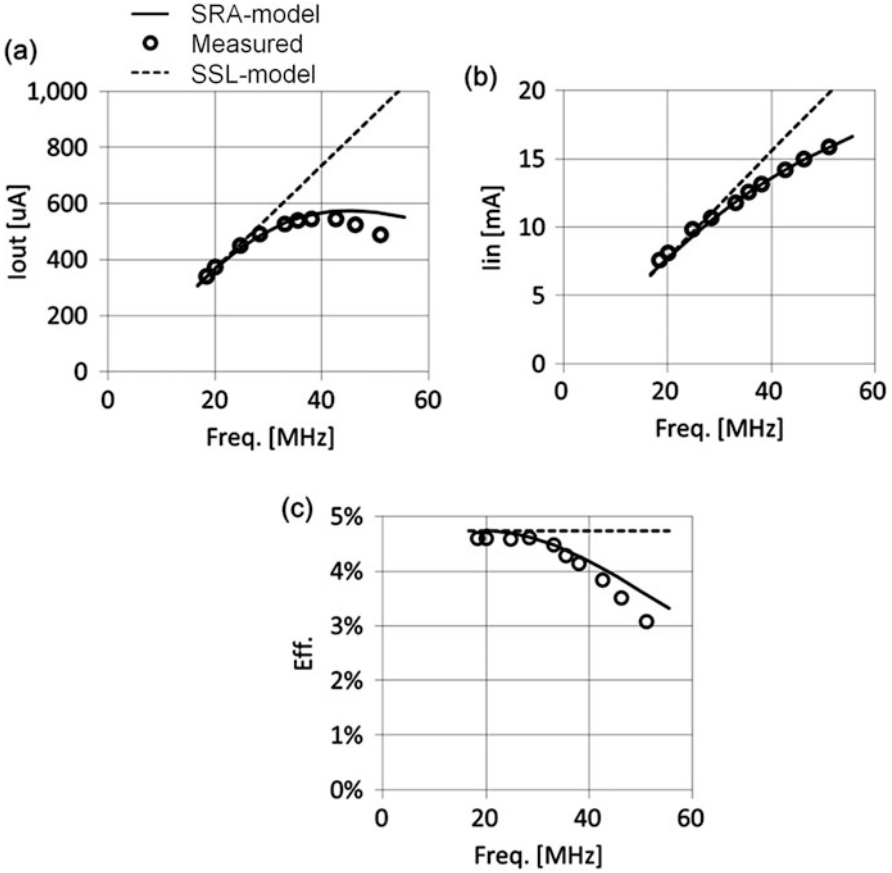
$$I_{\text{OUT}} = \frac{V_{\text{MAX}} - V_{\text{OUT}}}{2R(N + 1)} \quad (3.90)$$

Equation (3.90) indicates an equivalent circuit in FSL as shown in Fig. 3.22. A factor of 2 comes from the fact that the output current flows in half cycle. A factor of  $(N + 1)$  comes from the fact that there are  $(N + 1)$  parasitic resistors connected in series between the input and output terminals. Equation (3.90) has no frequency term. Thus, in FSL, the output current is determined by the characteristics of switching FETs' resistance and the number of stages.

Figure 3.23a, b, respectively, shows measured and modeled output and input current of the pump which has the design parameters shown in Table 3.3, where SRA and SSL model is, respectively, switch resistance aware model expressed by Eqs. (3.88a, 3.88b and 3.88c) and slow switching limit model expressed by Eq. (2.75).  $I_{\text{IN}}$  of model is calculated using Eq. (2.89) as well. The current consumption in an on-chip oscillator was subtracted from the measured input current for comparison with the models. Frequency dependencies of the input and output current of the SRA model with Eqs. (3.88a, 3.88b, and 3.88c) are in good agreement with the measured results up to the peak frequency where the output current has the peak value. It can be considered that the effective channel resistance of switching transistors is no longer treated as a constant because the boosting

**Fig. 3.22** Equivalent circuit of a charge pump with switching FETs in fast switching limit condition



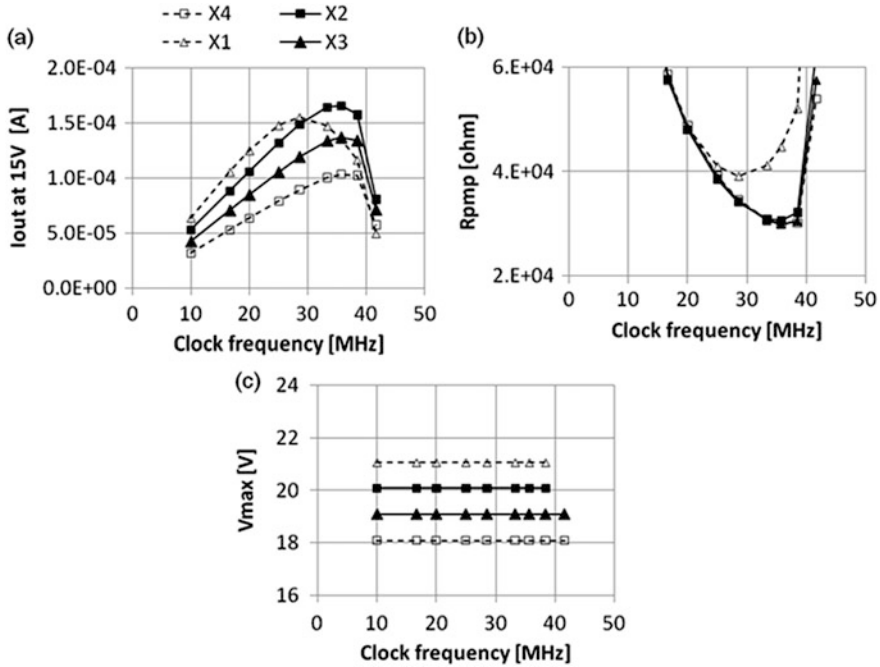


**Fig. 3.23** Measured and calculated output (a) and input (b) current and current efficiency (c) of the pump with the parameters shown in Table 3.3

**Table 3.3** Design parameters used in Fig. 3.23

$N$	8	a.u.
$C$	$4.0E-11$	F
$C_T$	$2.0E-12$	F
$C_B$	$1.2E-11$	F
$R$	180	Ohm
$V_{DD}$	2.15	V
$V_{OUT}$	15	V
$T_{OFF}$	$4.0E-09$	s

capacitors such as  $C_3$  and  $C_4$  in Fig. 3.21a are not fully charged as the clock frequency increases beyond the peak output current. On the other hand, disagreement of the ideal SSL model with Eqs. (2.76), (2.77), and (2.78) with the measured currents starts at much lower clock frequency. Figure 3.23c shows the current

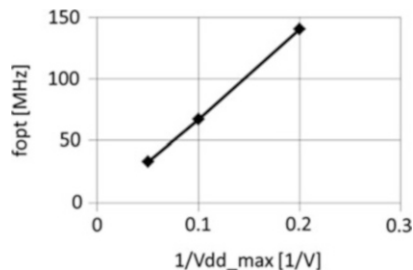


**Fig. 3.24** Performance comparisons between the four design with different transistor sizes using the model with Eqs. (3.88a–c) and parameters in Table 3.4a, b

efficiency. At a sacrifice of efficiency, the output current is maximized at about 40 MHz in this test case.

The optimum operation to maximize the output current depends on the switch resistance. As the channel width of the switching transistors increases, the channel resistance  $R$  decreases. The optimum frequency can be shifted toward a faster side. With the faster clock,  $R_{PMP}$  could be reduced. However, as the transistor size is increased, the parasitic capacitance  $C_T$  is also increased. Thus, the effective voltage amplitude  $V_{MAX}$  is reduced. Therefore, given the ratio of  $R$  and  $C_T$ , one can determine the optimum frequency to maximize the output current. Let's demonstrate design optimization. For simplicity, it is assumed that  $C_T$  is inversely proportional to  $R$  as shown in Table 3.4b. Under the condition that the total pump area is given, increase in the transistor area results in decrease in the main pumping capacitor  $C$ . Table 3.4b includes this consideration as well. Figure 3.24 compares four designs with different transistor sizes using the model with Eqs. (3.88a–c) and parameters in Table 3.4a. As far as the clock frequency is low enough,  $R_{PMP}$  is same between the four cases, however  $V_{MAX}$  is the highest with the smallest transistors and the largest capacitors, resulting in the highest  $I_{OUT}$ . According as the frequency increases,  $I_{OUT}$  also increases until it reaches the peak. One can find the maximum  $I_{OUT}$  at the peak frequency per parameter set of  $R$  and  $C$ . Among the peak values, the optimum frequency,  $R$ , and  $C$  values are identified to have the highest output

**Fig. 3.25** Optimum frequency as a function of maximum drain-to-source voltages of transfer transistors



**Table 3.4** Design parameters used in Fig. 3.24

(a)				
$N$	10	a.u.		
$V_{DD}$	2.0	V		
$T_{OFF}$	1.0E-08	s		
(b)				
	X1	X2	X3	X4
$R$ [ohm]	500	250	167	125
$C_T$ [pF]	0.5	1	1.5	2
$C$ [pF]	10	9.4	8.8	8.2

current. In Fig. 3.24, the parameter set as shown by X2 is the best one among the four cases to have the highest attainable output current and its optimum frequency is determined to be about 33 MHz.

The switch resistance depends on the transistors' dimensions such as channel length and gate oxide thickness which are determined by the maximum voltage applied to the transistors ( $V_{DD\_MAX}$ ). According as the pump output voltage required decreases, transistors used for the switches can be scaled. Figure 3.25 shows optimum frequency  $f_{OPT}$  as a function of  $1/V_{DD\_MAX}$ , where it is assumed that the gate oxide thickness, channel length, and channel width are scaled with a factor of  $1/V_{DD\_MAX}$ . Starting with  $f_{OPT}$  of 33 MHz at  $V_{DD\_MAX}$  of 20 V for the case in Tables 3.4a, b and Fig. 3.24, two cases of  $V_{DD\_MAX}$  of 10 and 5 V are further analyzed with the method used for Fig. 3.24, which results in Fig. 3.25. Thus, the method described above can provide an initial guess for optimum frequency and switch size for detailed SPICE simulations.

In summary, steady-state  $I-V$  characteristic is modeled for the pump operating with a fast frequency clock where the switch resistance cannot be negligible. Using the model, one can optimize the clock frequency where the output current is maximized. For practical design, current efficiency needs to be considered as well to make an optimization.

### 3.3 Power Efficiency

This section describes power efficiency in various views. Combining the equation for  $V_{OUT}-I_{OUT}$  Eq. (2.75) with that for  $V_{OUT}-I_{IN}$  Eq. (2.94) results in  $I_{IN}-I_{OUT}$  Eq. (3.91), which is valid regardless of a value of  $V_{TH}$ :

$$I_{IN} = \left( \frac{N}{1 + \alpha_T} + 1 \right) I_{OUT} + \left( \frac{\alpha_T}{1 + \alpha_T} + \alpha_B \right) fNCV_{IN} \quad (3.91)$$

From Eq. (3.91), one can have the relationship between power efficiency  $\eta$  vs.  $I_{OUT}$  given by Eq. (3.92):

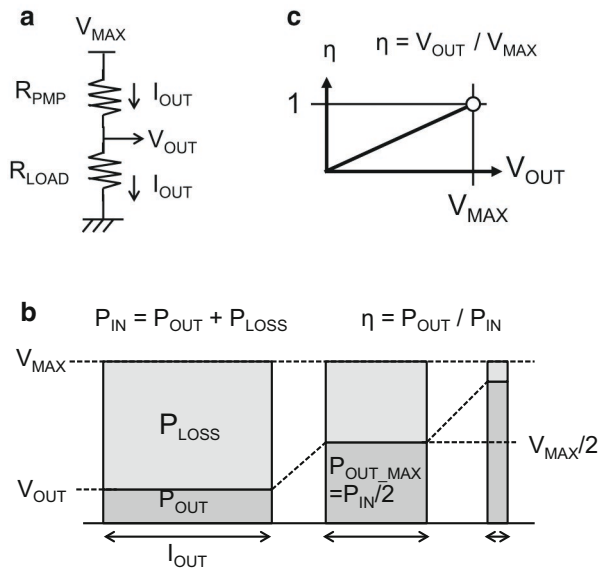
$$\frac{1}{\eta} = \frac{V_{IN}}{V_{OUT}} \left( \left( \frac{N}{1 + \alpha_T} + 1 \right) + \left( \frac{\alpha_T}{1 + \alpha_T} + \alpha_B \right) \frac{fNCV_{IN}}{I_{OUT}} \right) \quad (3.92)$$

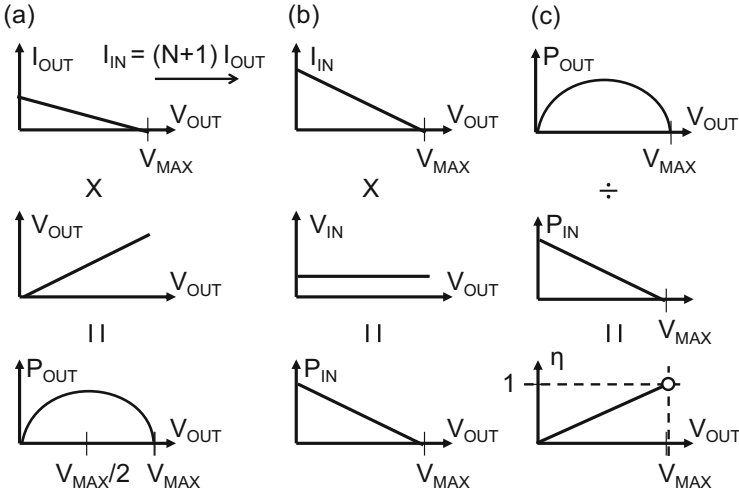
In an ideal case where  $\alpha_T = \alpha_B = 0$ , Eq. (3.92) is reduced to Eq. (3.93):

$$\eta = (V_{OUT}/V_{IN})/(N + 1) = V_{OUT}/V_{MAX} \quad (3.93)$$

Figure 3.26a shows an equivalent circuit, which is given by Eq. (2.75). Input and output power are respectively given by  $V_{MAX}I_{OUT}$  and  $V_{OUT}I_{OUT}$ . As a result,  $\eta$  is given by the voltage ratio as shown in Eq. (3.93) in an ideal case. Figure 3.26b shows power expressed by the area of rectangles, where the base is  $I_{OUT}$  and the height is  $V_{OUT}$  for  $P_{OUT}$  and  $V_{MAX}$  for  $P_{IN}$ . When  $V_{OUT}$  is close to 0 V,  $I_{OUT}$  is close

**Fig. 3.26** (a) Equivalent circuit of a charge pump, (b) power expressed by the area of rectangles, (c)  $\eta$  vs.  $V_{OUT}$



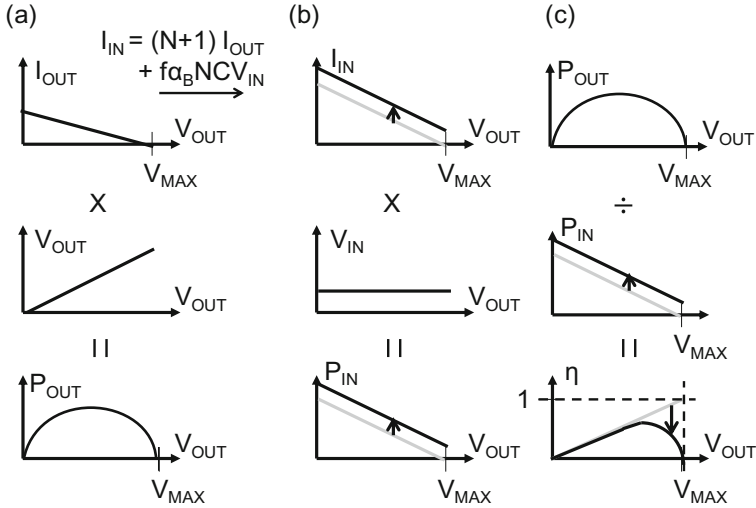


**Fig. 3.27** Characteristics of an ideal charge pump with no nonideal parameters: (a)  $I_{OUT}$ ,  $V_{OUT}$ ,  $P_{OUT}$ , (b)  $I_{IN}$ ,  $V_{IN}$ ,  $P_{IN}$ , (c)  $P_{OUT}$ ,  $P_{IN}$ ,  $\eta$

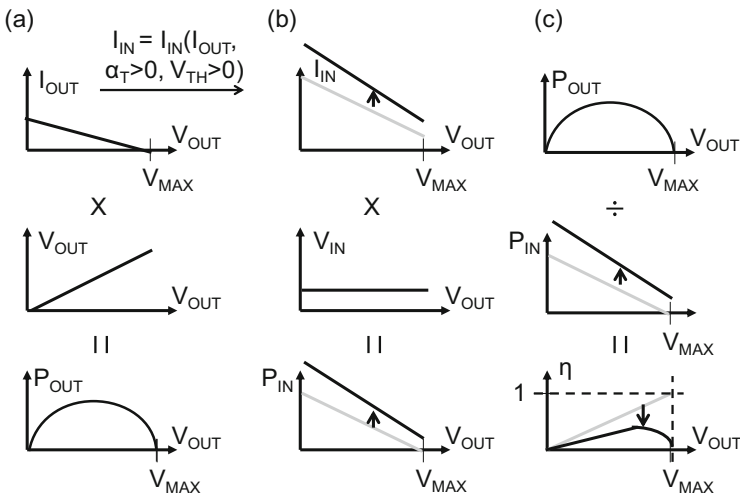
to its maximum, as illustrated in the left most of Fig. 3.26b. Because the power loss is close to 100 %,  $\eta$  has to be close to 0. As  $V_{OUT}$  increases,  $\eta$  improves. When  $V_{OUT}$  is close to  $V_{MAX}$ ,  $I_{OUT}$  is close to 0. Even though  $P_{OUT}$  is close to 0,  $\eta$  is close to 1, as shown in Fig. 3.26c. Thus, power efficiency of an ideal charge pump is very analogue to that of a liner regulator, which is also determined by the ratio of the output voltage to the input voltage, because the equivalent circuit is the same as Fig. 3.26a. As shown in Fig. 3.26b,  $P_{OUT}$  is maximized to be  $P_{IN}/2$  when  $V_{OUT} = V_{MAX}/2$ .

Figure 3.27 provides a different view on how power efficiency is determined.  $P_{OUT}$  is given by the product of  $I_{OUT}$  and  $V_{OUT}$ . As shown by Eq. (2.76),  $I_{OUT}$  is a first order equation of  $V_{OUT}$ . As a result,  $P_{OUT}$  is a second order equation of  $V_{OUT}$ .  $P_{OUT}$  has the maximum at  $V_{MAX}/2$ , as shown in Fig. 3.27a, where  $R_{LOAD} = R_{PMP}$ . This fact is known as impedance matching.  $I_{IN}$  is proportional to  $I_{OUT}$  in an ideal case where  $\alpha_T = \alpha_B = 0$ , as suggested by Eq. (3.91). As a result,  $P_{IN}$  is a first order equation of  $V_{OUT}$ , as shown in Fig. 3.26b. Therefore, as shown in Fig. 3.26c,  $\eta$  is given by a line, as shown by Eq. (3.93):

Figure 3.28 shows how the bottom plate parasitic capacitance affects power efficiency.  $\alpha_B$  does not affect  $V_{OUT}-I_{OUT}$  because it is not included in Eqs. (2.76)–(2.78). On the other hand,  $I_{IN}$  increases by the last term of Eq. (3.91), which does not have a factor of  $V_{OUT}$ . Therefore,  $\alpha_B$  shifts  $I_{IN}$ , i.e.  $P_{IN}$ , in the vertical direction as shown in Fig. 3.28b. As a result,  $\eta$  is affected at high  $V_{OUT}$  most as shown in Fig. 3.28c. Especially,  $\eta$  approaches zero at  $V_{OUT} = V_{MAX}$  rather than one in the ideal case. Thus, a finite  $\alpha_B$  creates the maximum point in  $\eta$ . Similarly, Fig. 3.29 shows how the top plate parasitic capacitance and the threshold voltage of the switching circuit affect power efficiency. The maximum attainable voltage  $V_{MAX}$  is given by Eq. (2.77) when  $V_{TH}$  is not zero.



**Fig. 3.28** Characteristics of a charge pump with a finite  $\alpha_B$ : (a)  $I_{OUT}$ ,  $V_{OUT}$ ,  $P_{OUT}$ , (b)  $I_{IN}$ ,  $V_{IN}$ ,  $P_{IN}$ , (c)  $P_{OUT}$ ,  $P_{IN}$ ,  $\eta$

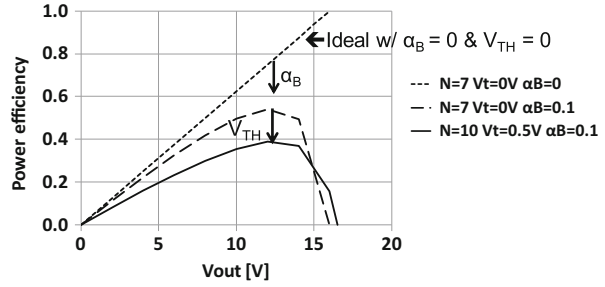


**Fig. 3.29** Characteristics of a charge pump with a finite  $\alpha_T$  and  $V_{TH}$ : (a)  $I_{OUT}$ ,  $V_{OUT}$ ,  $P_{OUT}$ , (b)  $I_{IN}$ ,  $V_{IN}$ ,  $P_{IN}$ , (c)  $P_{OUT}$ ,  $P_{IN}$ ,  $\eta$

In order to keep  $V_{MAX}$  as high as the case where  $V_{TH}$  is zero to have the same  $I_{OUT}$ , one needs to increase  $N$ . Independently, as  $\alpha_T$  increases,  $N$  needs to be increased for compensation. To keep  $R_{PMP}$ , which is given by Eq. (2.78), when  $N$  needs to be increased, one also needs to increase  $C$ . As a result, both first and second terms of the right hand side of Eq. (3.91) increases whereas  $I_{OUT}$  is unchanged. Figure 3.29b shows that  $I_{IN}$  shifts in vertical direction due to the second



**Fig. 3.30** Simulated  $V_{OUT}-\eta$  curves in an ideal case and two realistic cases.  $V_{IN}=2.0\text{ V}$ ,  $V_{OUT}=12\text{ V}$ ,  $I_{OUT}=60\text{ }\mu\text{A}$ ,  $f=10\text{ M-Hz}$ ,  $C=10\text{ pF}$



**Table 3.5** Optimum design of charge pumps

Section	Given parameters	Parameter(s) for optimization
Section 3.4.1	$V_{IN}, V_{OUT}, f, I_{OUT}$	Area to be minimized
Section 3.4.2	$V_{IN}, V_{OUT}, f, Area$	Rise time to be minimized
Section 3.4.3	$V_{IN}, V_{OUT}, f, I_{OUT}$	$I_{IN}$ or $P_{IN}$ to be minimized
Section 3.4.4	$V_{IN}, V_{OUT}, f, I_{OUT}$	Area and $P_{IN}$ to be balanced
Section 3.4.5	$V_{IN}, V_{OUT}, I_{OUT}$	Area, frequency, and $P_{IN}$ to be balanced

term of Eq. (3.91) and its slope has to be steeper due to the first term. Thus,  $\eta$  is degraded even in a low voltage range, as described in Fig. 3.29c.

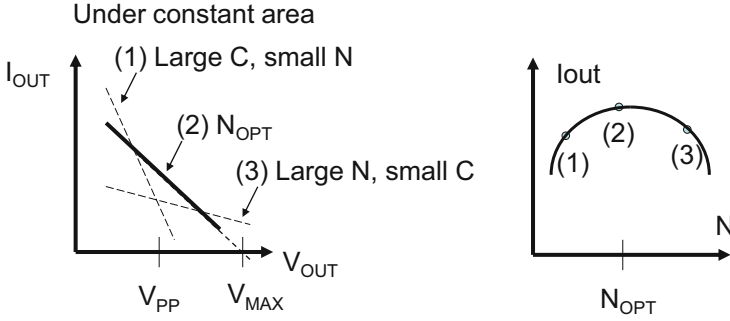
Figure 3.30 shows simulated  $V_{OUT}-\eta$  curves. The charge pump is designed to have  $\eta$  of 80% at  $V_{OUT}=12\text{ V}$  in the ideal case. When  $\alpha_B$  of 10% is included,  $\eta$  is reduced by 25% at 12 V. In addition, when  $V_{TH}$  increases to 0.5 V,  $N$  needs to be increased from 7 to 10. As a result,  $\eta$  is further reduced to 40%. Thus, the parasitic components significantly affect power efficiency.

## 3.4 Optimum Design

This section discusses various optimizations as described in Table 3.5.  $V_{IN}$  and  $V_{OUT}$  are also given parameters for each section.

### 3.4.1 Optimization for Maximizing the Output Current

This subsection discusses optimization for maximizing the output current under a given circuit area. In other words, the optimization is equivalent to a minimum circuit area to output a given current. Because the output voltage is fixed, maximizing the output current is equivalent to maximizing the output power. Figure 3.31 shows three possible options in terms of  $N$  and  $C$  under the product  $NC$  given. Option 1 has a large pump capacitor per stage with a small number of stages,



**Fig. 3.31** Optimization for maximizing  $I_{OUT}$  under a given total pump capacitor area

whereas option 3 has a small pump capacitor per stage with a large number of stages. Option 2 has moderate values for  $C$  and  $N$ . The question here is how to determine the option 2 whose  $I_{OUT}$  is the largest.

One can use Lagrange multiplier introducing functions  $f$  and  $h$ , and a parameter  $\lambda$  as follows:

$$f(C, N) \equiv \frac{C(1 + \alpha_T)}{TN} \left[ \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1)V_T - V_{PP} \right] - I_{PP} = 0 \tag{3.94}$$

$$h(C, N, \lambda) \equiv CN - \lambda f(C, N) \tag{3.95}$$

where Eq. (2.75) is used. One needs to minimize  $CN$  or  $h$  under the constraint where the pump has an output current of  $I_{PP}$  at an output voltage of  $V_{PP}$ .  $\lambda$  and  $I_{PP}$  are eliminated from Eq. (3.94) and  $\partial h / \partial N = \partial h / \partial C = 0$ , resulting in an optimum number of stages to minimize the area as Eq. (3.96):

$$N_{A\_OPT} = 2N_{MIN} \tag{3.96}$$

$$N_{MIN} = \frac{(1 + \alpha_T)(G_V - 1 + v_T)}{1 - v_T - \alpha_T v_T} \tag{3.97a}$$

where  $G_V$  is a voltage gain of  $V_{PP}/V_{DD}$ ,  $v_T$  is a relative threshold voltage of  $V_T/V_{DD}$ , and  $N_{MIN}$  is the number of stage to meet the condition of  $I_{PP} = 0$ . In case where  $V_T$  can be neglected, Eq. (3.97a) is reduced to

$$N_{MIN} = (1 + \alpha_T)(G_V - 1) \tag{3.97b}$$

The condition of Eqs. (3.96) and (3.97b) is equivalent to the condition of the impedance matching which is given by Eq. (2.11). Note that Eq. (3.96) is independent of  $I_{PP}$ .

### 3.4.2 Optimization for Minimizing the Rise Time

This subsection discusses another optimization for minimizing the rise time under a given circuit area. In other word, the optimization is equivalent to specifying the circuit area for a given rise time. As discussed in Sect. 3.1.1, the self-load capacitance  $C_{\text{PMP}}$  is almost constant for a given circuit area which is proportional to the total charge pump capacitance  $NC$ . The output series resistance of the charge pump,  $R_{\text{PMP}} = N/C$ , increases as the square of the number of stages  $N$  in case of a given circuit area because the charge pump capacitance  $C$  is inversely proportional to the number of stages. On the other hand, the maximum output voltage  $V_{\text{MAX}} = (N + 1)V_G$  proportionally increases with the number of stages. As a result, there will be an optimum number of stages to minimize the rise time. If the self-load capacitance of the charge pump,  $C_{\text{PMP}}$ , is set to be just one-third of the total charge pump capacitance,  $NC/3$ , under the condition of a given circuit area,

$$T_{\text{R}} = \frac{\ln\left(1 - \frac{N}{N_{\text{MIN}}}\right)}{\ln\left(1 + \frac{(1+\alpha_{\text{T}})C}{NC_{\text{OUT}}}\right)^{-1}} \quad (3.98)$$

$$C_{\text{OUT}} = C_{\text{L}} + (1 + \alpha_{\text{T}})C_{\text{TOT}}/3 \quad (3.99)$$

where  $C_{\text{L}}$  is the load capacitance of the charge pump,  $C_{\text{TOT}}$  is the total charge pump area  $CN$ , and the second term of Eq. (3.99) is an approximation with an error of 7 % or less as discussed in Sect. 3.1.1. When  $N$  is greater than 4, the denominator of Eq. (3.98) is expanded with an error of 6 % or less so that Eq. (3.98) can be expressed by

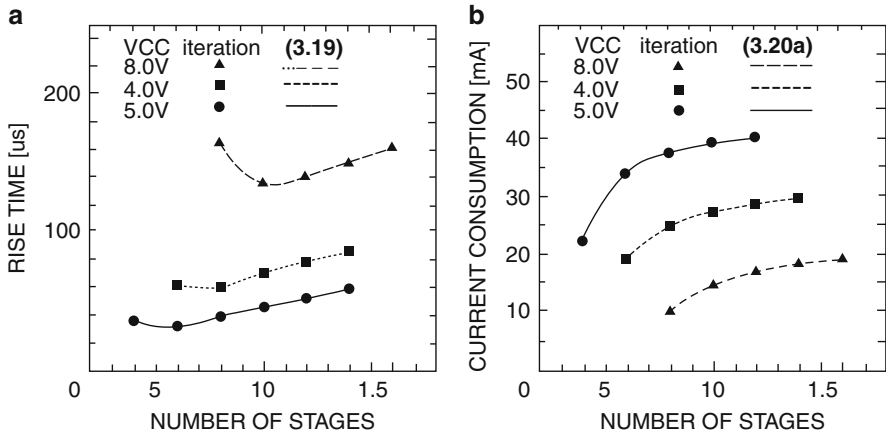
$$T_{\text{R}} = -\frac{N^2 C_{\text{OUT}}}{(1 + \alpha_{\text{T}})C_{\text{TOT}}} \ln\left(1 - \frac{N}{N_{\text{MIN}}}\right) \quad (3.100)$$

Equation (3.100) has the optimum number of stages to minimize the rise time,  $N_{\text{R\_OPT}}$ , is given by

$$N_{\text{R\_OPT}} = 1.40N_{\text{MIN}} \quad (3.101)$$

Note that Eq. (3.101) is independent of  $I_{\text{pp}}$ .

Figure 3.32a, b show the rise time and the current consumption under the condition of a constant circuit area. The rise time proportionally increases with the number of stages in case of a large number of stages. On the other hand, the rise time will be infinite in case of a number of stages as small as  $N_{\text{MIN}}$ . As a result, there is an optimum number of stages in any case. The current consumption increases with the number of stages, so that a charge pump with an excessive number of stages increases not only the rise time but also the current consumption.



**Fig. 3.32** Dependence of the (a) rise time and (b) current consumption on the number of stages under the condition of a constant circuit area  $CN=1$  nF and  $V_{PP}=20.0$  V,  $V_T=0.6$  V,  $\alpha_T=\alpha_B=0$ ,  $T=100$  ns, and  $C_{LOAD}=10$  nF

**Fig. 3.33** Dependence of the optimum number of stages on the boosted voltage under the condition of a constant circuit area and  $V_{DD}=3.0$  V,  $V_T=0.6$  V, and  $\alpha_T=\alpha_B=0$

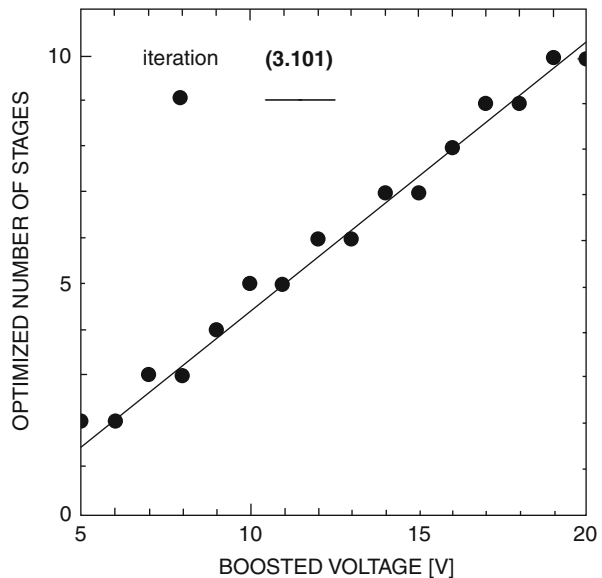
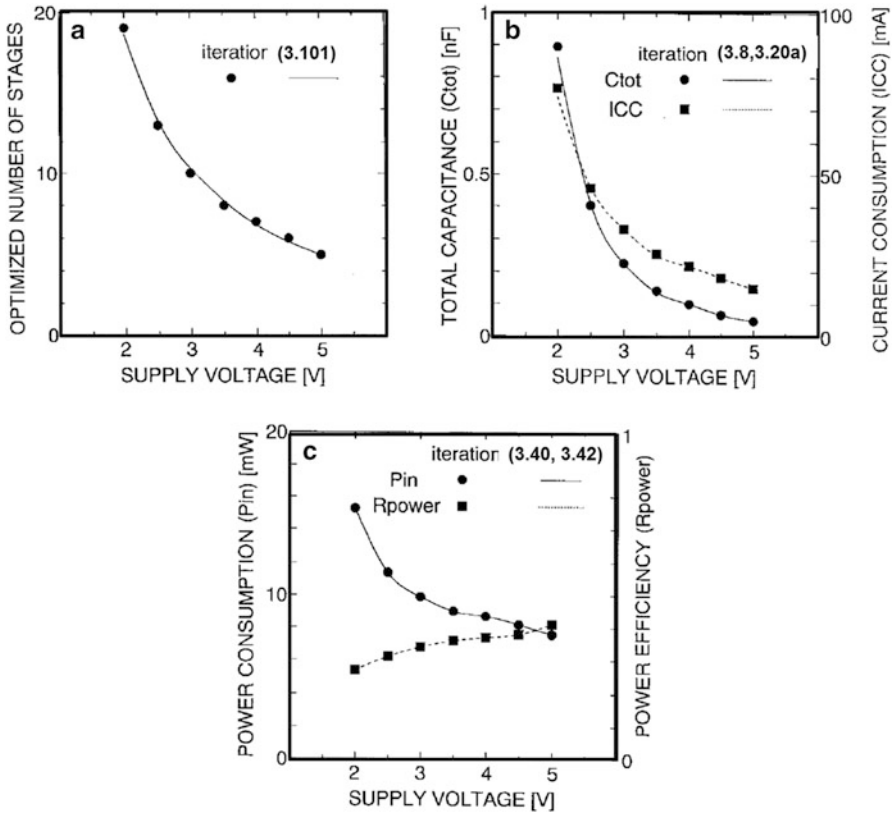


Figure 3.33 shows the dependence of the optimum number of stages on the boosted voltage. The analytical expression represented by the continuous line agrees with the iteration method represented by the discrete dots. The optimum number of stages proportionally increases with the boosted voltage, as represented by Eqs. (3.101) and (3.99). Figure 3.34a shows dependence of the optimum number of stages on the supply voltage. The optimum number of stages increases as the



**Fig. 3.34** Dependence of the optimum number of stages on the supply voltage under the condition of a constant circuit area and  $V_{PP} = 20.0$  V,  $V_T = 0.6$  V,  $\alpha_T = \alpha_B = 0$ ,  $T = 100$  ns (a). Dependence of the total capacitance,  $C_{TOT} = CN$ , and the current consumption on the supply voltage (b), and the power consumption and efficiency (c) under the condition that the rise time at any supply voltage is the constant value of  $63 \mu\text{s}$  and  $C_{OUT} = 10$  nF

supply voltage decreases. As mentioned above, an increase in the number of stages results in an increase in the current consumption. Figure 3.34b shows the total capacitance and the current consumption which are necessary for a constant rise time of  $63 \mu\text{s}$ . The circuit area and the current consumption at a supply voltage of 2 V are 17.9 and 5.1 times larger than those at 5 V, respectively. Figure 3.34c shows dependence of the power consumption and the power efficiency on the supply voltage under the same condition as Figure 3.34b. As a result, not only the circuit area but also the power consumption increase as the supply voltage decreases, unless the boosted voltage is scaled down according to the difference between the supply voltage and the threshold voltage of the transfer diode.

### 3.4.3 Optimization for Minimizing the Input Power

This subsection discusses an optimum design for minimizing the input power under the condition that the pump outputs a given current  $I_{PP}$  at a given output voltage  $V_{PP}$ . Based on the  $V_{OUT} - I_{OUT}$  relation as shown in Eq. (2.75) and the  $I_{DD} - I_{OUT}$  relation as shown in Eq. (2.94), one can use Lagrange multiplier introducing functions  $f$  and  $g$ , and a parameter  $\lambda$  as follows:

$$f(C, N) \equiv \frac{C(1 + \alpha_T)}{TN} \left[ \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1)V_T - V_{PP} \right] - I_{PP} = 0 \quad (3.102a)$$

$$g(C, N, \lambda) \equiv I_{DD}(C, N) - \lambda f(C, N) \quad (3.102b)$$

$\lambda$  and  $I_{PP}$  are eliminated from Eq. (3.102a) and  $\partial g / \partial N = \partial g / \partial C = 0$ , resulting in a quadratic equation in terms of  $N$  in Eq. (3.103):

$$\begin{aligned} & (1 - (1 + \alpha_T)v_T)(1 - v_T + \alpha_B)N^2 \\ & - 2(G_V - 1 + v_T)(1 + \alpha_T)(1 + \alpha_B - v_T)N \\ & + (1 + \alpha_T)(G_V - 1 + v_T)^2 = 0 \end{aligned} \quad (3.103)$$

where  $G_V$  is a voltage gain of  $V_{PP}/V_{DD}$  and  $v_T$  is a relative threshold voltage of  $V_T/V_{DD}$ . Equation (3.103) is accurately solved with no approximation, as shown below:

$$N_{P\_OPT} = N_{MIN} \left( 1 + \sqrt{\frac{\alpha_T + \alpha_B + \alpha_T \alpha_B}{(1 + \alpha_T)(1 + \alpha_B - v_T)}} \right) \quad (3.104a)$$

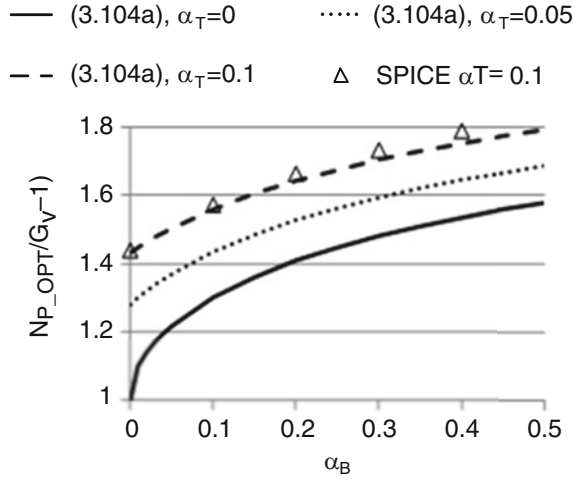
$$N_{MIN} = \frac{(1 + \alpha_T)(G_V - 1 + v_T)}{1 - v_T - \alpha_T v_T} \quad (3.104b)$$

SPICE simulations were performed to validate Eq. (3.104a) for the condition where  $C = 10$  pF,  $V_{DD} = 2.5$  V,  $V_T = 0$  V,  $T = 100$  ns,  $I_{PP} = 100$   $\mu$ A, and  $\alpha_T = 0.1$  as  $N$ ,  $V_{PP}$ , and  $\alpha_B$  are varied. A number of simulations were done with different values of  $N$  for each value of  $\alpha_B$  to determine the optimum value of  $N$  that leads to the lowest input current. Figure 3.35 includes a comparison of SPICE results with Eq. (3.104a).  $N_{P\_OPT}/(G_V - 1)$  was used for the comparison because it is a simple function of  $\alpha_T$  and  $\alpha_B$  according to Eqs. (3.104a) and (3.104b) when  $V_T = 0$  V. The discrepancy was 4 % or less for the  $\alpha_B$  range from 0 to 0.5.

### 3.4.4 Optimization with Area Power Balance

The optimum number of stages with different optimization methods such as minimizing the total pump capacitor area with Eq. (3.96), minimizing the rise

**Fig. 3.35** Dependence of  $N_{P\_OPT}$  on  $\alpha_T$  and  $\alpha_B$ , where  $V_T = 0$  V



time Eq. (3.101), and minimizing the power Eq. (3.104a) are summarized below.  $N_{OPT}$  is commonly expressed by

$$N_{OPT} = \varepsilon N_{MIN} \quad (3.105a)$$

where  $\varepsilon$  is a multiplication factor and is, respectively, given by

$$\varepsilon(A_{MIN}) = 2.0 \quad (3.105b)$$

$$\varepsilon(T_{RISE\_MIN}) = 1.40 \quad (3.105c)$$

$$\varepsilon(I_{DD\_MIN}) = 1 + \sqrt{\frac{\alpha_T + \alpha_B + \alpha_T \alpha_B}{(1 + \alpha_T)(1 + \alpha_B - v_T)}} \quad (3.105d)$$

To see how  $\varepsilon$  affects the total capacitor area, the input power, and the rise time, the normalized parameters  $A_{NORM}$ ,  $I_{DD\_NORM}$ , and  $T_{R\_NORM}$  are calculated as follows. Charge pump area  $C_{TOT}(N)$  as a function of  $N$  is calculated using Eqs. (3.102a), (3.104b), and (3.98):

$$C_{TOT}(N) = CN = \frac{TI_{PP}}{(1 - v_T - \alpha_T v_T)V_{DD}} \frac{N^2}{N - N_{MIN}} \quad (3.106a)$$

Thus, the normalized parameter  $A_{NORM}(\varepsilon)$  is defined by

$$A_{NORM}(\varepsilon) \equiv \frac{C_{TOT}(N)}{C_{TOT}(N_{A\_OPT})} = \frac{1}{4} \frac{\varepsilon^2}{\varepsilon - 1} \quad (3.106b)$$

Similarly, using Eqs. (2.94), (3.102a), (3.104a), and (3.104b),

$$I_{DD\_NORM}(\varepsilon) \equiv \frac{I_{DD}(N)}{I_{DD}(N_{P\_OPT})} = \frac{\varepsilon + \frac{1}{N_{MIN}} + \frac{\varepsilon}{\varepsilon - 1} \left( \frac{\alpha_T}{1 + \alpha_T} + \varepsilon \frac{\alpha_B + \alpha_T v_T}{1 - v_T - \alpha_T v_T} \right)}{\eta + \frac{1}{N_{MIN}} + \frac{\eta}{\eta - 1} \left( \frac{\alpha_T}{1 + \alpha_T} + \eta \frac{\alpha_B + \alpha_T v_T}{1 - v_T - \alpha_T v_T} \right)} \quad (3.106c)$$

$$\eta = 1 + \sqrt{\frac{\alpha_T + \alpha_B + \alpha_T \alpha_B}{(1 + \alpha_T)(1 + \alpha_B - v_T)}} \quad (3.106d)$$

Also, using Eqs. (3.104b), (3.100), and (3.101),

$$T_{R\_NORM}(\varepsilon) \equiv \frac{T_R(N)}{T_R(N_{R\_OPT})} = \frac{\varepsilon^2}{1.4^2} \frac{\ln(1 - \frac{1}{\varepsilon})}{\ln(1 - \frac{1}{1.4})} \quad (3.106e)$$

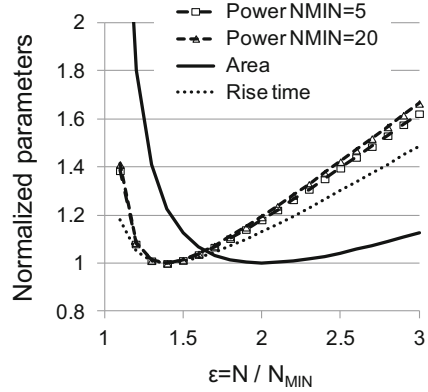
Three normalized parameters  $A_{NORM}$ ,  $I_{DD\_NORM}$ , and  $T_{R\_NORM}$  are overlaid in Fig. 3.36, which shows that the area is minimized at  $\varepsilon = 2$  and that the area is larger by 28 % when  $N_{MIN} = 20$ ,  $N = N_{P\_OPT}$  (i.e.,  $\varepsilon = 1.36$ ) and by 23 % when  $N = N_{R\_OPT}$  (i.e.,  $\varepsilon = 1.40$ ) than when  $N = N_{A\_OPT}$ . On the other hand, Fig. 3.36 also shows that the power is minimized when  $\varepsilon = 1.36$  when  $N_{MIN} = 5$  to 20,  $\alpha_T = 0.05$ ,  $\alpha_B = 0.1$ , and  $v_T = 0.1$  and that  $P_{IN}$  is larger by about 20 % when  $\varepsilon = 2$  (i.e.,  $N = N_{A\_OPT}$ ) than when  $\varepsilon = 1.36$  (i.e.,  $N = N_{P\_OPT}$ ). Note that the impact of  $N_{MIN}$  is not significant. Thus, when one selects the optimum  $N$  in terms of either area or power, one needs more power or area by 20 % or more. Since  $T_{R\_NORM}$  is smaller than  $I_{DD\_NORM}$  at any  $\varepsilon$  when the conditions that  $T_{R\_NORM} \geq 1.1$  and  $I_{DD\_NORM} \geq 1.1$  are considered, the relation between  $A_{NORM}$  and  $I_{DD\_NORM}$  is summarized in Fig. 3.37. One can select the optimum  $\varepsilon$  in range of 1.5 to 1.7 for making a balance between power and area.

Figure 3.38 shows dependence of  $\alpha_T$ ,  $\alpha_B$ , and  $v_T$  on an optimum  $\varepsilon$  to minimize the power under a practical design constraint of  $A_{NORM} \leq 1.1$ .  $N_{MIN}$  was varied from 2 to 50 and had no impact on the optimum  $N/N_{MIN}$  to second decimal places. There are six curves in Fig. 3.38, where one parameter among  $\alpha_T$ ,  $\alpha_B$ , and  $v_T$  is varied while the others are set at 0 or 0.3. As a result, it is validated that  $\varepsilon$  of 1.5 to 1.7 can minimize the power under the condition that  $A_{NORM}$  is smaller than 1.1 and each of  $\alpha_T$ ,  $\alpha_B$ , and  $v_T$  is smaller than 0.2. To study the impact of  $v_T$  on the total capacitor area for three cases using Eq. (3.96), Eqs. (3.104a and 3.104b), and  $N = 1.6N_{MIN}$ , the area and power are calculated per each  $v_T$  and is, respectively, normalized by Eqs. (3.96) and (3.104a, 3.104b), resulting in Fig. 3.39. A design with the number of stages calculated by  $1.6N_{MIN}$  can have a moderate area overhead with an increase of about 7 % or a moderate power overhead with about a 5 % increase.

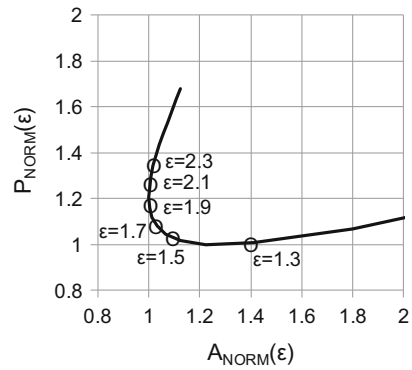
Let's consider the physical aspect on the monotonic increase in  $N_{P\_OPT}$  as a function of  $\alpha_B$ , as shown in Fig. 3.35. Figure 3.36 shows that the total capacitor area decreases as  $\varepsilon$  increases, as long as  $\varepsilon$  is less than 2. Because  $N$  and  $C$  need to be chosen to meet Eq. (2.75), a larger value of  $N$  requires a smaller value of  $C$ . When



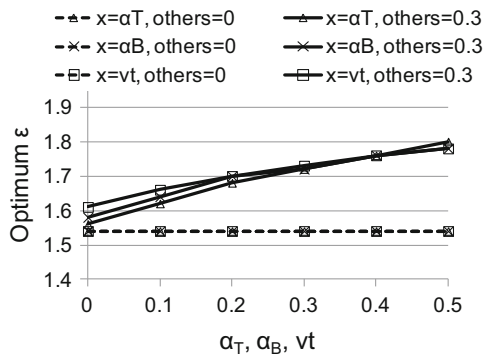
**Fig. 3.36** Dependence of area, power, and rise time on  $\epsilon$  with  $\alpha_T = 0.05$ ,  $\alpha_B = 0.1$ ,  $v_T = 0.1$



**Fig. 3.37** Relationship between input power and area with  $N_{MIN} = 20$ ,  $\alpha_T = 0.05$ ,  $\alpha_B = 0.1$ , and  $v_T = 0.1$

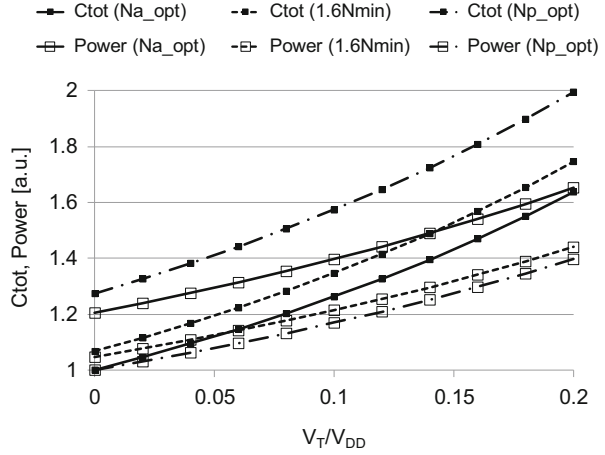


**Fig. 3.38** Dependence of  $\alpha_T$ ,  $\alpha_B$ , and  $v_T$  on optimum  $\epsilon$



$\alpha_B$  increases, an increase in the third term of Eq. (2.94) is mitigated with a larger  $N$  because  $NC$  can decrease. By selecting a larger value of  $N$  for a larger  $\alpha_B$ , a decrease in the third term can be larger than an increase in the sum of the first and second term. As a result, the total input power can be reduced with a larger  $N$ .

**Fig. 3.39** Impact of  $v_T$  on the total capacitor area and power given by Eqs. (3.96) and (3.104a, 3.104b), and  $1.6N_{MIN}$  with  $G_V = 20$ ,  $\alpha_T = 0.05$ ,  $\alpha_B = 0.1$



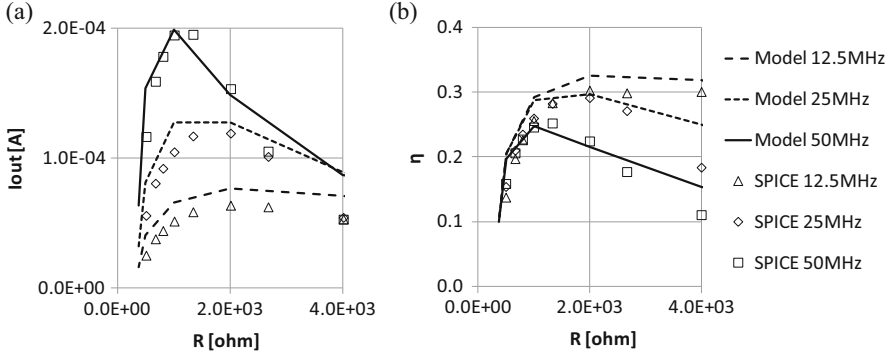
**Table 3.6** Device and circuit parameters used for verification

Parameter	Max	Min
$f$	100 M-Hz	10 M-Hz
$R$	4 k-ohm	250- $\Omega$
$N$	14	8
$V_{IN}$	2.5 V	
$V_{OUT}$	15 V	
$I_{OUT}$	270 $\mu$ A	
$C_{TOT}$	63 pF	
Parameter	Value	
$C_{STG}$	$C_{TOT}/N$	
$tRC_T$	1.4 ns	
$C_T$	$tRC_T/R$	
$C$	$C_{STG}-C_T$	
$\alpha_T$	$C_T/C$	
$\alpha_B$	0.22	
$T_{OFF}$	2.0 ns	

$N_{P\_OPT}$  also increases as  $\alpha_T$  increase, as indicated in Fig. 3.35. This is mainly because  $N_{MIN}$  monotonically increases as  $\alpha_T$ , as given by Eq. (3.104b).

### 3.4.5 Guideline for Comprehensive Optimum Design

An optimum clock frequency and optimum area ratio of pump capacitor to switching circuit can be determined so as to maximize the output current under a given technology as discussed in Sect. 3.2.3. Based on power efficiency expressed by Eq. (3.92) as a general form, one can have a comprehensive optimization



**Fig. 3.40** (a)  $I_{OUT}$  and (b)  $\eta$  as a function of  $R$ .  $N = 10$ ,  $C_{TOT} = 63$  pF

methodology to take both power efficiency and area into consideration for determining the clock frequency, area ratio of pump capacitor to switching circuit, number of stages, and capacitance per stage at the same time. To demonstrate the methodology, design and device parameters shown in Table 3.6 are used for model calculation and SPICE simulation.

Figure 3.40 shows  $I_{OUT}$  and  $\eta$  as a function of  $R$  under the condition where the number of stages and the total capacitance  $C_{TOT}$  are assumed to be 10 and 63 pF, respectively. The errors in the data points with the model from the SPICE simulation were within 20 %. Based on the results as shown in Fig. 3.40, trajectories for  $I_{OUT}-\eta$  are plotted in Fig. 3.41. The top plate parasitic capacitance  $C_T$  is determined by the size of the switching circuit, which is specified by the channel resistance  $R$ , via Eq. (3.107), where a technology-dependent parameter  $tRC_T$  is assumed to be 1.4 ns:

$$C_T = tRC_T/R \quad (3.107)$$

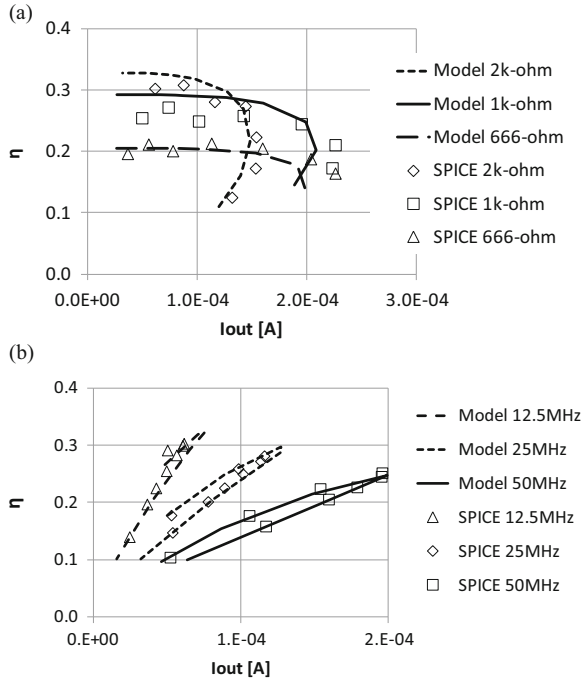
Thus, the top plate parasitic capacitance is inversely proportional to the channel resistance. Then, capacitance per stage is provided by Eq. (3.108) because  $C_{TOT}$  includes both area for pump capacitor and switching circuit:

$$C = C_{TOT}/N - C_T \quad (3.108)$$

Thus, capacitance per stage decreases as the size of switching circuit increases. In Fig. 3.41a, the clock frequency is varied while the area ratio of pump capacitor and switching circuit is unchanged. When the frequency is low,  $I_{OUT}$  is relatively small and  $\eta$  is relatively large. As the frequency increases,  $I_{OUT}$  increases while  $\eta$  does not change significantly. When the frequency is as high as 100 M-Hz, both  $I_{OUT}$  and  $\eta$  become too low. It is difficult to determine the optimum frequency from this behavior.

In Fig. 3.41b, the area ratio of pump capacitor and switching circuit is varied while the clock frequency is unchanged. When the area ratio of pump capacitor to

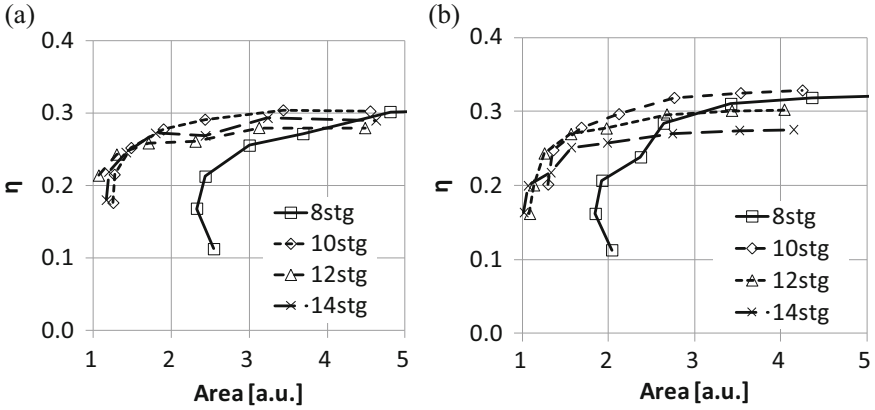
**Fig. 3.41** Comparisons of the model with the SPICE simulation. Variables are (a) the clock frequency and (b) the area ratio of pump capacitor to the switching circuit.  $N = 10$ ,  $C_{TOT} = 63$  pF



switching circuit is large, the pump is in a state where the charge transfer is incomplete under a given frequency. As a result,  $I_{OUT}$  is relatively small. When  $N$  and  $f$  are given, Eq. (3.92) suggests that  $\eta$  is a function of  $I_{OUT}$  mainly because  $C$  is a weak function of the area ratio, and contribution of  $\alpha_B$  and  $\alpha_T$  is minor. Therefore,  $\eta$  is also relatively small when  $I_{OUT}$  is small. As the area ratio decreases, more amounts of charges can be transferred in a half clock period. As a result,  $I_{OUT}$  increases, thereby  $\eta$  also increases through the second term of Eq. (3.92). Thus, it is much easier to determine the optimum area ratio by clock frequency, at which both  $I_{OUT}$  and  $\eta$  are maximized. Then, the required total pump area ( $C_{TOT\_REQ}$ ) is calculated to output a targeted  $I_{OUT}$  ( $I_{OUT\_TAR}$ ) of 270  $\mu$ A in this demonstration using Eq. (3.109):

$$C_{TOT\_REQ} = C_{TOT\_ASS} \times I_{OUT\_CAL} / I_{OUT\_TAR} \quad (3.109)$$

By doing the similar procedure for different number of stages, one can plot the optimum points in a single *area- $\eta$*  plane, as shown in Fig. 3.42. The area is normalized by the minimum among the data points. Figure 3.42a, b respectively show the results using SPICE simulation and model calculation. Both plots show that (1) power efficiency is about 20 % when the charge pump is designed so as to have the minimum area, i.e.,  $N = 14$ ,  $f = 100$  M-Hz, and  $R = 700$  ohm, (2) power efficiency improves to about 30 % when the area is allowed to be twice as large as



**Fig. 3.42** Area vs. power efficiency by SPICE simulation (a) and model calculation (b)

**Table 3.7** Optimization methodology

1. $I_{OUT}$ , $V_{OUT}$ , $V_{IN}$ , given
2. Technology and switching circuit selected
3. Basic parameters of $tRC_T$ , $T_{OFF}$ , $\alpha_T$ , and $\alpha_B$ determined
4. Calculate $I_{OUT}$ and $\eta$ using Eqs. (3.81b), (3.88a–c), (3.91), and (3.92) by varying $f$ under various the area ratios of pump capacitor to switching circuit per $N$ . Area is assumed to be given
5. Determine the optimum area ratio to have the maximum $I_{OUT}$ and $\eta$ per $f$ per $N$
6. Calculate required area for each optimum area ratio per $f$ per $N$ to meet target $I_{OUT}$ using Eq. (3.109)
7. Plot area vs. $\eta$ per $N$
8. Select the best point which determines $f$ , area ratio, $N$ , and $C$ , simultaneously

the minimum, i.e.,  $N = 10$ ,  $f = 25$  M-Hz, and  $R = 2$  k-ohm, and (3) power efficiency is saturated at larger area. Discrepancy of the model from the SPICE result was about 40 % in an extreme case where the number of stages was as small as 8, which barely generated 15 V, but was within 20 % in nominal cases where the number of stages is 10 or larger. As a result, it is validated that the model can provide the initial values for the circuit parameters when the design specification is given. Table 3.7 summarizes the methodology.

### 3.5 Summary of Useful Equations

This section summarizes key design equations for charge pumps with different types of switching devices. General  $V_{OUT}$ – $I_{OUT}$  equations are gathered in Table 3.8.

**Table 3.8** General  $V_{OUT}$ – $I_{OUT}$  equations

Switching device	$V_{OUT}$ – $I_{OUT}$
Diode Eq. (3.59)	$\exp\left(\frac{V_{OUT}}{V_T}\right) = \frac{1}{4} \left( \frac{I_S}{(1 + \alpha_T) f C V_T} \right)^{N+1} \exp\left(\frac{(N + 1 + \alpha_T) C V_{DD} - N I_{OUT} / f}{(1 + \alpha_T) C V_T}\right)$ $\times \left( 1 - \exp\left(-\frac{I_{OUT}}{(1 + \alpha_T) f C V_T}\right) \right)^{-2} \left( 1 - \exp\left(-\frac{2 I_{OUT}}{(1 + \alpha_T) f C V_T}\right) \right)^{-N+1}$
MOSFET in saturated regime Eq. (3.69)	$\left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1) V_{TH}^{EFF} - V_{OUT}$ $= (N - 1) \sqrt{\frac{I_{OUT}^2}{(1 + \alpha_T)^2 f^2 C^2} + \frac{2 I_{OUT}}{k}} + \sqrt{\frac{I_{OUT}^2}{(1 + \alpha_T)^2 f^2 C^2} + \frac{8 I_{OUT}}{k}}$
MOSFET in triode regime Eqs. (3.81b) and (3.88a, 3.88b and (3.88c)	$I_{OUT} = f C (1 + \alpha_T) \left( \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - V_{OUT} \right) / \left( N \frac{1 + \zeta^2}{1 - \zeta^2} + \frac{2 \zeta}{1 - \zeta^2} \right)$ $\zeta \equiv \exp\left(-\frac{T_{ON}}{RC(1 + \alpha_T)}\right)$

**Table 3.9** Slow switching limit approximation

Switching device	$V_{TH}^{EFF}$
Diode Eq. (3.63b)	$V_T \ln \left( 4 \frac{1}{N + 1} \frac{(1 + \alpha_T) f C V_T}{I_S} \right)$
MOSFET in saturated regime Eq. (3.70)	$\sum_{k=1}^{N+1} V_{TH,k} / (N + 1)$
MOSFET in triode regime, cf. Eq. (3.88c)	0

SSL equation is given by Eq. (3.61) et al.:

$$I_{OUT} = (V_{MAX} - V_{OUT}) / R_{PMP} \quad (3.61)$$

where the effective output impedance of the charge pump  $R_{PMP}$  and the maximum attainable voltage  $V_{MAX}$  are given by Eqs. (3.62) and (3.63a) below where the effective threshold voltage is shown in Table 3.9:

$$R_{PMP} = \frac{N}{(1 + \alpha_T) C f} \quad (3.62)$$

$$V_{MAX} = \left( \frac{N}{1 + \alpha_T} + 1 \right) V_{DD} - (N + 1) V_{TH}^{EFF} \quad (3.63a)$$

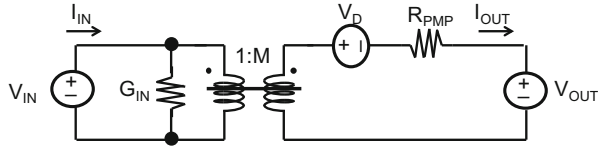
FSL equation has no frequency term, as shown in Table 3.10. The output current is determined by the characteristics of switching devices.

**Table 3.10** Fast switching limit approximation

Switching device	$V_{OUT}$ - $I_{OUT}$
Diode Eq. (3.64)	$I_{OUT} = \frac{I_S}{2} \exp \left[ \frac{(N/(1 + \alpha_T) + 1)V_{DD} - V_{OUT}}{V_T(N + 1)} \right]$
MOSFET in saturated regime Eqs. (3.74)–(3.75)	$I_{OUT} = \frac{k}{2} \left( \frac{(N/(1 + \alpha_T) + 1)V_{DD} - V_{OUT}}{N + 1} - \sum_{k=1}^{N+1} \frac{V_{TH,k}}{(N + 1)} \right)^2$
MOSFET in triode regime Eqs. (3.88c) and (3.90)	$I_{OUT} = \frac{N/(1 + \alpha_T) + 1}{2R(N + 1)} (V_{DD} - V_{OUT})$

**Table 3.11** Optimum number of stages

Parameter to be minimized	$N_{OPT}/N_{MIN}$
Area	2.0
Rise time	1.4
Input power	$1 + \sqrt{\frac{\alpha_T + \alpha_B + \alpha_T \alpha_B}{(1 + \alpha_T)(1 + \alpha_B - v_T)}}$

**Fig. 3.43** Equivalent circuit in SSL

$I_{IN}$ - $I_{OUT}$  is given by Eq. (3.91) which does not depend on  $V_{TH}$  of switching devices:

$$I_{IN} = \left( \frac{N}{1 + \alpha_T} + 1 \right) I_{OUT} + \left( \frac{\alpha_T}{1 + \alpha_T} + \alpha_B \right) f_{NC} V_{IN} \quad (3.91)$$

One can select the number of stages from Table 3.11. The minimum number of stage  $N_{MIN}$  to meet the condition of  $I_{OUT} = 0$  is defined by Eq. (3.104b) where  $G_V$  is a voltage gain of  $V_{OUT}/V_{IN}$ ,  $v_T$  is a relative threshold voltage of  $V_T/V_{IN}$ :

$$N_{MIN} = \frac{(1 + \alpha_T)(G_V - 1 + v_T)}{1 - v_T - \alpha_T v_T} \quad (3.104b)$$

Figure 3.43 shows an equivalent circuit of charge pumps in SSL, which includes Eqs. (3.61)–(3.63a) for  $V_{OUT}$ - $I_{OUT}$  and Eq. (3.91) for  $I_{IN}$ - $I_{OUT}$ . DC conversion coefficient  $M$ , loss at input  $G_{IN}$ , and voltage drop  $V_D$  are, respectively, given by Eqs. (3.110)–(3.112):

$$M = \frac{N}{1 + \alpha_T} + 1 \quad (3.110)$$

$$G_{IN} = \left( \frac{\alpha_T}{1 + \alpha_T} + \alpha_B \right) f_{NCV_{IN}} \quad (3.111)$$

$$V_D = (N + 1)V_{TH}^{EFF} \quad (3.112)$$

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# Chapter 4

## Design of AC–DC Charge Pump

**Abstract** This chapter discusses circuit theory of AC–DC charge pump circuits. The input is a continuous wave with a single frequency or a multi-sine wave with multiple frequencies. An analytical, closed-form AC–DC charge pump voltage multiplier model is described to show the dependency of output current and input power on circuit and device parameters for continuous wave AC–DC charge pump. Then, it is expanded for multi-sine wave AC–DC charge pump. Analysis enables circuit designers to estimate circuit parameters, such as the number of stages and capacitance per stage, and device parameters such as saturation current (in the case of diodes) or transconductance (in the case of MOSFETs). In addition, design optimizations and the impact of AC power source impedance on output power are investigated.

Even though switched-capacitor voltage multipliers were originated with AC–DC, i.e., AC input and DC output, by Greinacher and Cockcroft–Walton, most voltage multipliers for integrated circuits (ICs) have been DC–DC, i.e., DC input and DC output for decades because almost all ICs work with DC input. Recently, wireless sensing nodes and implantable microelectronic devices have been attracting the interest of researchers and engineers. These devices use AC–DC rectifier voltage multipliers to receive power or to harvest energy in AC form. These applications require low power (typically nothing higher than hundreds of  $\mu\text{W}$ ) and have small form factors—features that are well-matched with the features of voltage multipliers with no inductor or any magnetic element required. Section 4.1 discusses continuous wave AC–DC charge pump voltage multiplier which operates at a single frequency. Section 4.1.1 provides a circuit model which only includes DC voltage source, output resistance, and internal capacitance. Each parameter is expressed by circuit and device parameters. Section 4.1.2 investigates design and device parameter sensitivity on the pump performance. Section 4.1.3 discusses optimum design for maximizing output current at a given output voltage and for making a balance between circuit area and input power. Section 4.1.4 studies the impact of AC power source impedance on the pump performance and the dependency of design parameters on the AC power source impedance. Section 4.2 expands into multi-sine wave voltage multipliers where the AC signal has multiple frequencies. Section 4.2.1 provides a circuit model. One can estimate output and input power using the model equations when design and device parameters are given. Section 4.2.2 shows design and device parameter sensitivity on the

pump performance. Section 4.2.3 investigates the effectiveness of multi-sine waves over continuous waves in terms of power efficiency.

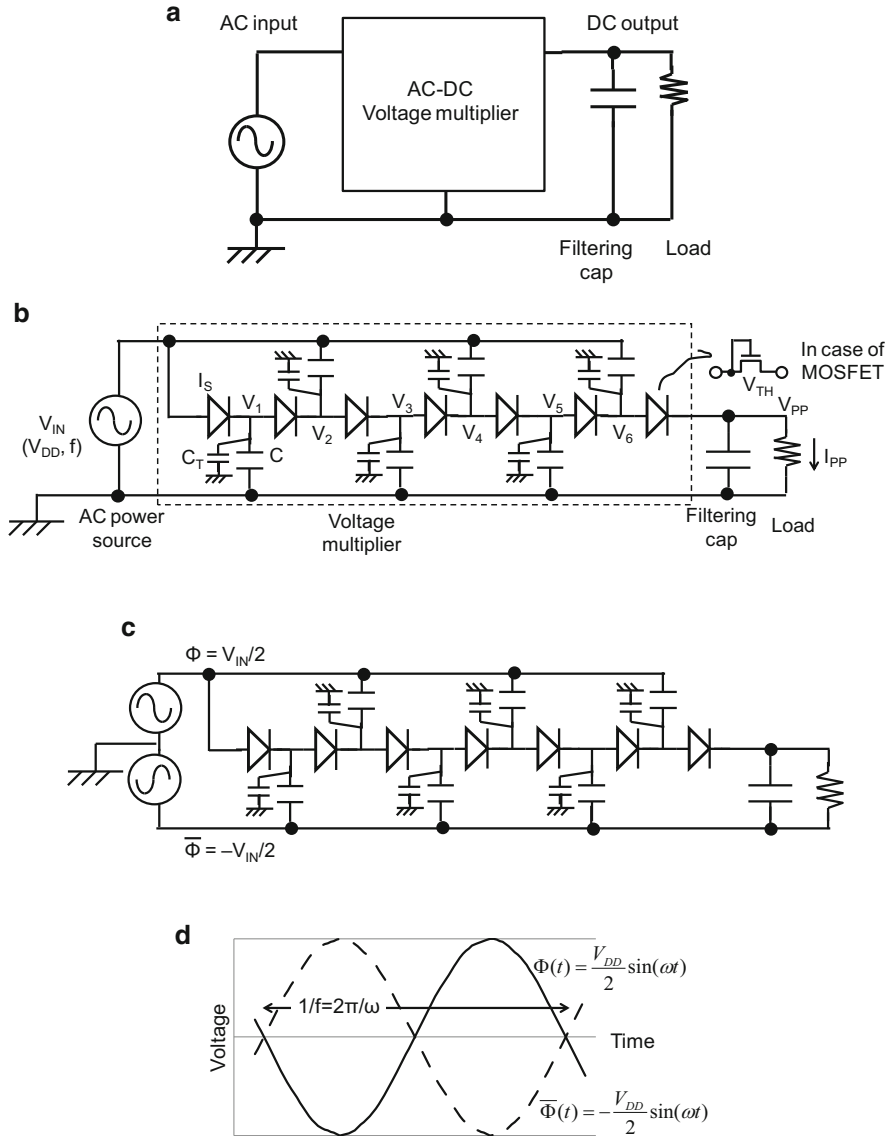
## 4.1 Continuous Wave (CW) AC–DC Charge Pump Voltage Multipliers

### 4.1.1 Circuit Model

Depending on applications, the frequency used for AC–DC power converters varies greatly from 10 Hz to a few kHz for vibrating energy harvesting (Meninger et al. 2001), 868 MHz and 2.45 GHz ISM (Industry—Science—Medical) bands for UHF RFID (Umeda et al. 2005; Papotto et al. 2011), and 10–35 GHz for UPS (Utility Power Satellite) (Gutmann and Borrego 1979; Yoo and Chang 1992) in space. In addition, the voltage gain, i.e., the ratio of output DC voltage to input AC voltage amplitude, needs to be as low as a few in an application or greater than ten in another. Thus, a circuit model for AC–DC voltage multiplier needs to cover wide operating frequency and voltage gain ranges.

Figure 4.1a shows a block diagram of AC–DC charge pump voltage multiplier to generate a DC output voltage from an AC input power source. Figure 4.1b illustrates a six stage AC–DC voltage multiplier. Two terminals for the AC power source are alternately connected with capacitors. The first and last diodes are connected with the AC source and the output terminal, respectively. Capacitors next to each other are connected with a diode. The number of capacitors and diodes are  $N$  and  $N + 1$ , respectively.  $f$  is the frequency of the AC source  $V_{IN}$ .  $I_S$  is the diode saturation current and  $V_T$  is the thermal voltage.  $V_{PP}$  is the output DC voltage and  $I_{PP}$  is the output current.  $N$  can be an odd number, but the last stage does not contribute to a voltage gain. Thus, in this paper, only even number of stages is considered. Figure 4.1c illustrates the voltage multiplier in differential expression using complementary signals  $\Phi$  and  $/\Phi$ , which is identical to Fig. 4.1b in terms of electrical characteristics of the multiplier. Because of its symmetry, Fig. 4.1c is analyzed in this chapter. Figure 4.1d shows the voltage waveform of the AC source in differential expression.  $V_{DD}$  is the amplitude of  $V_{IN}$  in single-ended expression and  $\omega$  is the angular velocity ( $2\pi f$ ).

Assumptions made in this section are: (1) the parasitic resistance of the multiplier capacitors is small enough in comparison with the resistance in the diodes; (2) the parasitic capacitance of the diodes is small enough in comparison with the multiplier capacitors; (3) the output impedance of the AC power source is small enough in comparison with the input impedance of the multiplier; (4) the filtering capacitor is large enough to recognize the output voltage  $V_{PP}$  as a constant under a given output current; (5) the reverse bias leakage current of the diode is small enough in comparison with the forward bias current. When the reverse bias current is not small enough due to very low diode threshold voltage, the net output current,



**Fig. 4.1** Block diagram (a), six stage AC–DC voltage multiplier in (b) single-ended and (c) differential expressions, and (d) waveform of the input voltages in differential expression

and thereby the power efficiency, is reduced from the values that the model predicts.

Figure 4.2a shows the simulated voltage waveform at each capacitor node of a six-stage voltage multiplier as shown in Fig. 4.1c, where  $V_{DD}$  is 0.5 V,  $V_{PP}$  is 1.0 V,

and the clock frequency is 1 GHz. It turns out that the DC offset voltage from  $k$ -th stage to  $(k + 1)$ -th stage is independent of  $k$ , where  $k$  is an integer from one to five, which is equal to that from the input to the first stage and that from sixth stage to the output. To investigate how an internal capacitor voltage is associated with the driving clock, the simulated voltage waveform at the fourth capacitor node  $V_4(t)$ , the driving clock  $\Phi(t)$  added by an offset of  $V_{4\text{DC}}$  which is the DC term of  $V_4(t)$ , and the signal  $\Phi(t + \theta/\omega)$  shifted by  $\theta/\omega$  from  $\Phi(t)$ , are plotted in Fig. 4.2b. Note that  $\theta$  is simply introduced to fit  $V_4(t)$  here, but it will be omitted in the final model equation as shown later. In addition, the voltage waveform  $V_3(t)$  and  $V_4(t)$  at two adjacent capacitor nodes are shown in Fig. 4.2c. The other voltages  $V_{2k-1}(t)$  and  $V_{2k}(t)$  at odd and even capacitors, where  $k$  is an integer from one to  $N/2$  in case of a general  $N$ -stage voltage multiplier, have the same AC terms as  $V_3(t)$  and  $V_4(t)$  and different DC terms  $V_{2k-1\text{DC}}$  and  $V_{2k\text{DC}}$ , respectively. In the period when there is no charge transfer, the capacitor voltage varies at the same rate as the driving clock when the top plate capacitance is negligibly small. This is because the charge stored in the capacitor, which is proportional to the voltage difference between the capacitor and clock voltages, does not change until charge transfer occurs. When the charge transfer occurs, the voltage difference varies. After the charge transfer is complete, the voltage difference becomes a constant again. As a result, it appears that  $\Phi(t + \theta/\omega)$  is more fitted with the actual  $V_4(t)$  than  $\Phi(t)$ , where  $\theta$  is a temporal parameter, as shown in Fig. 4.2b. Thus, a charge transfer affects and modulates the phase of the internal capacitor voltage similar to how the amplitude modulation affects the phase of the AC signal. It is assumed that Eqs. (4.1) and (4.2) hold in general in a no charge transfer period, where  $V_{n\text{DC}}$  is the DC term of  $V_n(t)$ .

$$V_n(t) = \frac{V_{\text{DD}}}{2} \sin(\omega t + \theta) + V_{n\text{DC}} \quad (4.1)$$

$$V_{n+1}(t) = -\frac{V_{\text{DD}}}{2} \sin(\omega t + \theta) + V_{n+1\text{DC}} \quad (4.2)$$

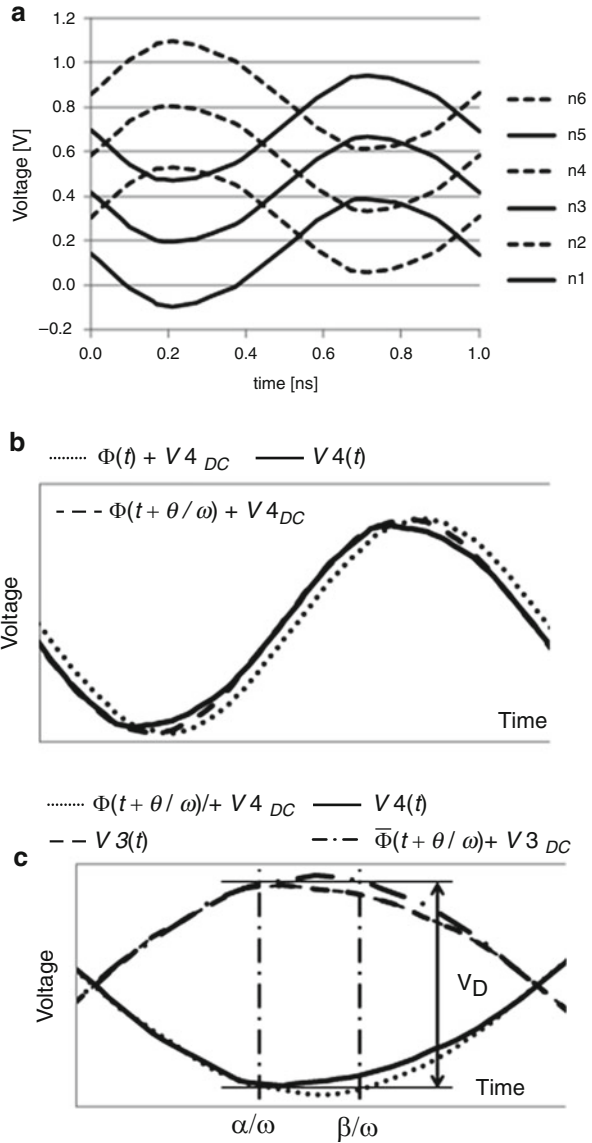
Parameters  $\alpha$  and  $\beta$  in Fig. 4.2c are the angles when the charge transfer starts and ends, respectively. It is assumed that the voltage across the diode between these capacitors is equal to a constant voltage  $V_{\text{D}}$  as a first approximation to have an analytical equation, even though the voltage across the diode varies during the charge transfer. Since  $V_{\text{D}}$  is the voltage difference between  $V_n$  and  $V_{n+1}$  at  $\alpha/\omega$  and  $\beta/\omega$ , subtracting Eq. (4.2) from Eq. (4.1) results in Eq. (4.3) where  $\Delta V$  is given by Eq. (4.4).

$$\begin{aligned} V_{\text{D}} &= V_{\text{DD}} \sin(\alpha + \theta) - \Delta V \\ &= V_{\text{DD}} \sin(\beta + \theta) - \Delta V \end{aligned} \quad (4.3)$$

$$\Delta V = V_{n+1\text{DC}} - V_{n\text{DC}} \quad (4.4)$$

Since  $\Delta V$  is a voltage gain per diode and there are  $(N + 1)$  diodes between the input and output terminals, it is related to the output voltage  $V_{\text{PP}}$  through Eq. (4.5).

**Fig. 4.2** Simulated six-stage rectifier multiplier waveform (Tanzawa 2014)



$$\Delta V = V_{PP}/(N + 1) \tag{4.5}$$

From Eq. (4.3),  $\alpha$ ,  $\beta$  and  $\theta$  are related each other via Eq. (4.6), resulting in Eq. (4.7).

$$\alpha + \theta = \pi - (\beta + \theta) \tag{4.6}$$

$$\theta = (\pi - \alpha - \beta)/2 \tag{4.7}$$

For the diode forward bias voltage  $V_D \gg V_T$ , the diode voltage ( $V_D$ )—current ( $I_D$ ) equation is approximated to be Eq. (4.8), where  $I_S$  is the diode saturation current and  $V_T$  is the thermal voltage.

$$I_D = I_S \exp(V_D/V_T) \quad (4.8)$$

Since  $I_{PP}$  is an average current of  $I_D$  in a period,  $I_{PP}$  is expressed by  $I_D$  as in Eq. (4.9).

$$I_{PP} = \frac{\beta - \alpha}{2\pi} I_D \quad (4.9)$$

The charges transferred from  $n$ -th node to  $(n + 1)$ -th node per cycle,  $I_{PP}/(fC)$ , are given by the difference between the peak voltage difference (i.e.,  $V_{n\_peak} - V_{n+1\_peak} = (V_{DD}/2 + V_{nDC}) - (V_{DD}/2 + V_{n+1DC})$ ) and the voltage difference at  $t = \beta/\omega$  (i.e.,  $V_n(\beta/\omega) - V_{n+1}(\beta/\omega)$ ). Thus,

$$\frac{I_{PP}}{fC} = V_{DD} \left( 1 - \cos \left( \frac{\beta - \alpha}{2} \right) \right)$$

(4.3), (4.5), and (4.7)–(4.9) result in Eq. (4.10), where  $\gamma$  is Eq. (4.11).

$$I_{PP} = \frac{\gamma I_S}{\pi} \exp \left( \frac{V_{DD} \cos(\gamma) - \frac{V_{PP}}{N+1}}{V_T} \right) \quad (4.10)$$

$$\gamma \equiv \frac{\beta - \alpha}{2} = \cos^{-1} \left( 1 - \frac{I_{PP}}{fCV_{DD}} \right) \quad (4.11)$$

$2\gamma$  shows the conduction angle, which is the ratio of the period when charge transfer occurs to one clock cycle. Equation (4.10) is also expressed by Eq. (4.12) to provide  $V_{PP}$  as a function of  $I_{PP}$ .

$$V_{PP} = (N + 1) \left( V_{DD} - \frac{I_{PP}}{fC} - V_T \ln \left( \frac{\pi I_{PP}}{I_S \cos^{-1} \left( 1 - \frac{I_{PP}}{fCV_{DD}} \right)} \right) \right) \quad (4.12)$$

When parasitic capacitance at a top plate of each capacitor  $C_T$  is not negligible in comparison with the main capacitor  $C$ , one needs to add a term for  $\alpha_T$  defined by  $C_T/C$ . In that case, one can use Eq. (4.13) instead of Eq. (4.12) by replacing  $V_{DD}$  in Eqs. (4.1) and (4.2) with  $V_{DD}/(1 + \alpha_T)$  and  $C$  in Eq. (4.11) with  $C(1 + \alpha_T)$ , as discussed for DC–DC charge pump in Chap. 3.

$$V_{PP} = (N + 1) \left( \frac{V_{DD}}{1 + \alpha_T} - \frac{I_{PP}}{fC(1 + \alpha_T)} - V_T \ln \left( \frac{\pi I_{PP}}{I_S \cos^{-1} \left( 1 - \frac{I_{PP}}{fCV_{DD}} \right)} \right) \right) \quad (4.13)$$

When the last term is recognized as an effective threshold voltage of the rectifying diode as given by Eq. (4.14), Eq. (4.13) is rewritten by Eq. (4.15).

$$V_{\text{TH}}^{\text{EFF}} \equiv V_{\text{T}} \ln \left( \frac{\pi I_{\text{PP}}}{I_{\text{S}} \cos^{-1} \left( 1 - \frac{I_{\text{PP}}}{f C V_{\text{DD}}} \right)} \right) \quad (4.14)$$

$$V_{\text{PP}} = (N + 1) \left( \frac{V_{\text{DD}}}{1 + \alpha_{\text{T}}} - V_{\text{TH}}^{\text{EFF}} - \frac{I_{\text{PP}}}{f C (1 + \alpha_{\text{T}})} \right) \quad (4.15)$$

To approximate Eq. (4.13) to have a linear relation between  $V_{\text{PP}}$  and  $I_{\text{PP}}$ , Eq. (4.16) can be used with an error lower than 5 % in a range of 0.1–1.9 for  $\gamma$ , resulting in Eqs. (4.17)–(4.19a, 4.19b).

$$\gamma = 1.1 \frac{I_{\text{PP}}}{f C V_{\text{DD}}} + 0.46 \quad (4.16)$$

$$I_{\text{PP}} = (V_{\text{MAX}} - V_{\text{PP}}) / R_{\text{PMP}} \quad (4.17)$$

$$R_{\text{PMP}} = \frac{(N + 1)}{f C (1 + \alpha_{\text{T}})} \quad (4.18)$$

$$V_{\text{MAX}} = (N + 1) \left( \frac{V_{\text{DD}}}{1 + \alpha_{\text{T}}} - V_{\text{TH}}^{\text{EFF}} \right) \quad (4.19a)$$

$$V_{\text{TH}}^{\text{EFF}} = V_{\text{T}} \ln \left( \frac{f C V_{\text{DD}}}{I_{\text{S}}} \right) \quad (4.19b)$$

One can find other formulations in Vita and Iannaccone 2005; Barnett et al. 2009; Cardoso et al. 2012.

Averaged input power  $P_{\text{IN}}$  is calculated using an averaged  $V_{\text{DD}}$ , described by  $\bar{V}_{\text{DD}}$ , over the time between  $\alpha/\omega$  and  $\beta/\omega$ , as shown in Eq. (4.20), resulting in Eq. (4.21), where  $N$  is an even integer and Eqs. (4.9) and (4.11) are used, and  $\alpha = \pi/2 - \gamma$  and  $\beta = \pi/2 + \gamma$  are assumed as a first approximation.

$$\bar{V}_{\text{DD}} = \int_{\alpha/\omega}^{\beta/\omega} V_{\text{DD}} \sin(\omega t) dt / (\beta/\omega - \alpha/\omega) = V_{\text{DD}} \sin \gamma / \gamma \quad (4.20)$$

$$P_{\text{IN}} = \bar{V}_{\text{DD}} I_{\text{D}} \frac{\beta - \alpha}{2\pi} (N + 1) = (N + 1) \frac{\sin \gamma}{\gamma} V_{\text{DD}} I_{\text{PP}} \quad (4.21)$$

In Eq. (4.21), it is considered that the power is input from  $N$  capacitors driven by the AC source and directly from one input terminal of the first diode. Unlike DC–DC voltage multipliers, the AC power source doesn't consume power to charge parasitic capacitance  $C_{\text{T}}$ ,  $C_{\text{B}}$ . Note that  $P_{\text{IN}}$  is not a direct function of  $V_{\text{PP}}$ , but is indirectly affected by  $V_{\text{PP}}$  via  $N$  since  $N$  is determined by a required voltage gain  $V_{\text{PP}}/V_{\text{DD}}$  and the device and circuit parameters. Power efficiency is then given by Eq. (4.22a).



$$P_{\text{EFF}} \equiv \frac{P_{\text{OUT}}}{P_{\text{IN}}} = \eta_{\text{CA}} \frac{V_{\text{PP}}}{(N+1)V_{\text{DD}}} \quad (4.22a)$$

where the power loss factor associated with a finite conduction angle  $\eta_{\text{CA}}$  is defined by Eq. (4.22b).

$$\eta_{\text{CA}} \equiv \frac{\gamma}{\sin \gamma} \quad (4.22b)$$

As the conduction angle increases,  $\eta_{\text{CA}}$  thereby  $P_{\text{EFF}}$  decreases. With an increased conduction angle, effective voltage amplitude at the capacitor node during charge transfer is reduced, which gives a voltage drop from full swing of  $V_{\text{DD}}$ , resulting in a power loss.

Figure 4.3 shows how  $\gamma$  and  $\eta_{\text{CA}}$  behave as a function of  $I_{\text{PP}}$ . Nominal parameters in Table 4.1 and  $N=20$  are used. In Fig. 4.3a, b,  $V_{\text{PP}} - I_{\text{PP}}$  is also shown. In this example,  $\gamma$  becomes about 0.8 when  $I_{\text{PP}}$  is 150  $\mu\text{A}$ . A conduction angle  $2\gamma$  is then about  $\pi/2$ .  $\eta_{\text{CA}}$  is reduced to 0.9 at this operation condition. As a result, power efficiency decreases by the same factor.

Effective input impedance can be defined and expressed by Eq. (4.23).

$$Z_{\text{IN}}^{\text{EFF}} \equiv \frac{\overline{V_{\text{DD}}}^2}{P_{\text{IN}}} = \frac{1}{\eta_{\text{CA}}} \frac{V_{\text{DD}}}{(N+1)I_{\text{PP}}} \quad (4.23)$$

Equations (4.17)–(4.19a, 4.19b) and (4.23) suggest an AC–DC rectifier multiplier linear model as illustrated in Fig. 4.4. The parasitic capacitance at the top plate of each capacitor  $C_{\text{T}}$  is also included in Fig. 4.4 through the output impedance  $R_{\text{PMP}}$  and the maximum attainable output voltage  $V_{\text{MAX}}$  as given by Eqs. (4.18) and (4.19a, 4.19b), respectively.

When the rectifying device is a metal-oxide-semiconductor (MOS) transistor whose gate and drain are connected as shown in Fig. 4.1b, the drain current is expressed by Eq. (4.24), where  $k$  is transconductance,  $V_{\text{TH}}$  is a threshold voltage, and  $\nu$  is a power factor which changes from 2 to 1 as the carrier velocity saturation gets more severe (Sakurai and Newton 1990).

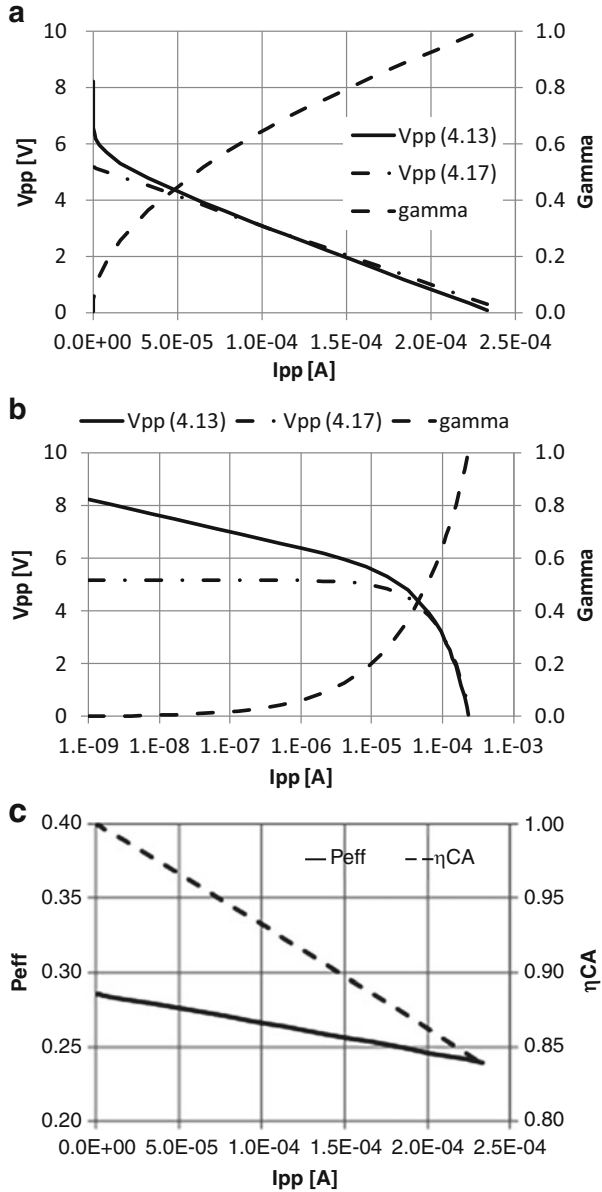
$$I_{\text{D}} = k(V_{\text{D}} - V_{\text{TH}})^{\nu} \quad (4.24)$$

Equations (4.3), (4.5), (4.7), (4.9), and (4.24) result in Eq. (4.25).

$$I_{\text{PP}} = \frac{\gamma k}{\pi} \left( V_{\text{DD}} \cos(\gamma) - \frac{V_{\text{PP}}}{N+1} - V_{\text{TH}} \right)^{\nu} \quad (4.25)$$

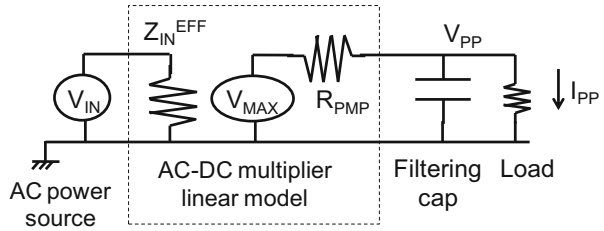
From Eq. (4.25),  $V_{\text{PP}}$  is given as a function of  $I_{\text{PP}}$  by Eq. (4.26). One can find another formulation in Yi et al. 2007.

**Fig. 4.3** (a)  $V_{PP}$  and  $\gamma$  as a function of  $I_{PP}$  in (a) linear and (b) log plots, (c)  $P_{EFF}$  and  $\eta_{CA}$  as a function of  $I_{PP}$



**Table 4.1** Device and circuit parameters used for verification

Parameter	Diode			MOSFET
	Max	Nominal	Min	
$f$	10 G-Hz	1 G-Hz	10 M-Hz	10 M-Hz
$V_{DD}$	0.7 V	0.5 V	0.2 V	0.5 V
$V_{PP}$	3.5 V	2.5 V	0.4 V	2.5 V
$C$	10 pF	1 pF	–	1 pF
$I_S$	100 nA	10 nA	1 nA	$K = 0.9 \text{ mA/V}^2$
$V_T$	30 mV			$V_{TH} = 0.2 \text{ V}$
$\alpha_T$	0.1	0.01	0.001	0.01
$N$	6, 8, 10, ..., 20			

**Fig. 4.4** CW AC–DC rectifier multiplier linear model (Tanzawa 2014)

$$V_{PP} = (N + 1) \left( V_{DD} - \frac{I_{PP}}{fC} - V_{TH} - \left( \frac{\pi I_{PP}}{k \cos^{-1} \left( 1 - \frac{I_{PP}}{fC V_{DD}} \right)} \right)^{1/\nu} \right) \quad (4.26)$$

### 4.1.2 Design and Device Parameter Sensitivity on the Pump Performance

The model is verified by comparing the calculated results based on the model with the SPICE simulation results under the parameters summarized in Table 4.1. SPICE Schottky barrier diode model parameters used for verification are summarized in Table 4.2 (Sze and Ng 2007) so that one can verify the results.  $C_T$  was added as a linear capacitance in the SPICE netlist. Every graph in this section is assumed to use the nominal numbers in Table 4.1, as far as specific numbers are not shown.

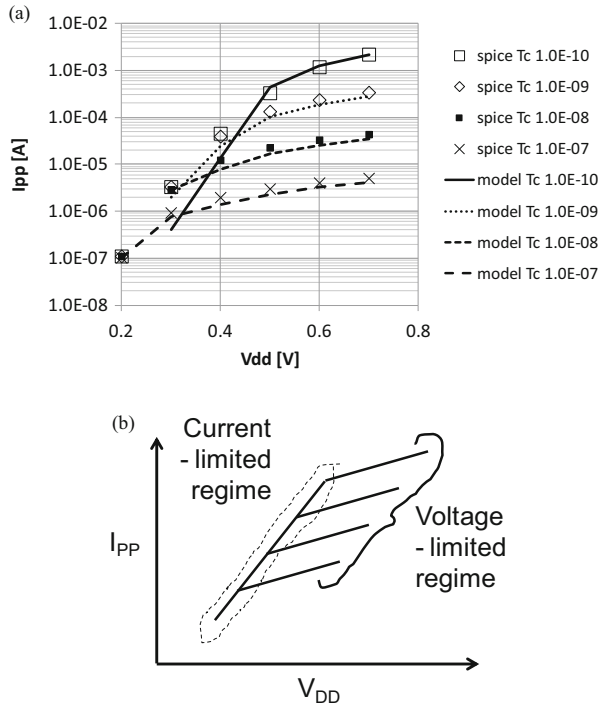
Figure 4.5a compares the model Eq. (4.13) with SPICE results in terms of  $I_{PP}$  as a function of  $V_{DD}$  and  $f$ . Except for very high operating frequency at low  $V_{DD}$ , the model Eq. (4.13) is in good agreement with the SPICE results with an error lower than 30%. Figure 4.5a indicates there are two regimes in  $I_{PP}$ : (1) current-limited and (2) voltage-limited as illustrated in Fig. 4.5b.  $I_{PP}$  has no frequency dependency in the current-limited regime; whereas  $I_{PP}$  is proportional to operating frequency in the voltage-limited regime.

Figure 4.6a, b show an image of the current-limited regime. The frequency in Fig. 4.6b is twice as large as in Fig. 4.6a. These figures assume the transferred

**Table 4.2** SPICE diode model parameters used for verification

Name	Parameter	Value	Unit
Level	Junction diode model	1	No unit
$EG$	Band-gap energy	0.69	eV
$N_P$	Emission coefficient	1	No unit
$IKF$	High-injection knee current	1e-4	A
$RS$	Parasitic resistance	0.1	$\Omega$
$IS$	Saturation current	1e-8	A

**Fig. 4.5** (a)  $V_{DD}$  vs.  $I_{PP}$  under various  $f$  in case of  $N = 20$  and  $V_{PP} = 2.0$  V and (b) its interpretation (Tanzawa 2014)



charge is limited by the diode, even though there is a sufficient voltage difference across the diode because the capacitance of the multiplier capacitor is relatively large and the operating frequency is relatively high. Thus, the conduction angle, or the pulse width in  $I_{PP}$ , is very narrow and a constant current flows during a very limited time. In such a situation,  $I_{PP}$  could not increase by a factor of 2 when the frequency is doubled because the current pulse height is unchanged and the current pulse width is proportional to the clock cycle. Therefore, the averaged  $I_{PP}$  is unchanged with the frequency increase. Conversely, Fig. 4.6c, d show an image of the voltage-limited regime. The frequency in Fig. 4.6d is twice as large as in Fig. 4.6c. These figures assume the conduction angle is relatively wide, and when the current to the next capacitor returns zero at  $\pi/2$ , the voltage across the diode is effectively zero. In this case, the current height and its width are respectively twice

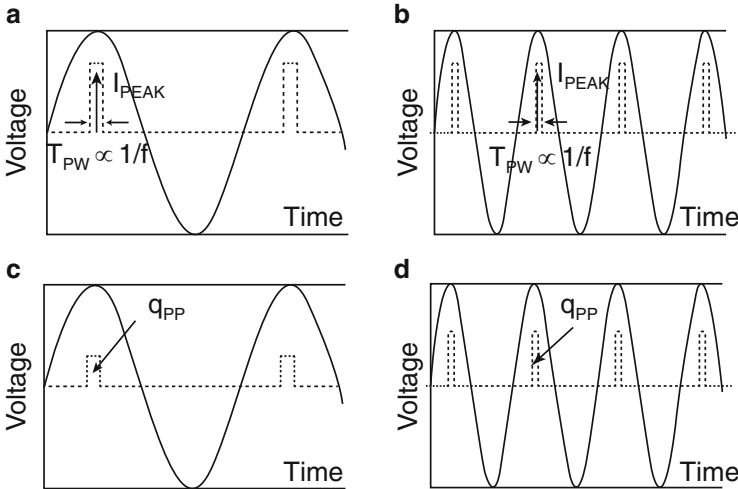


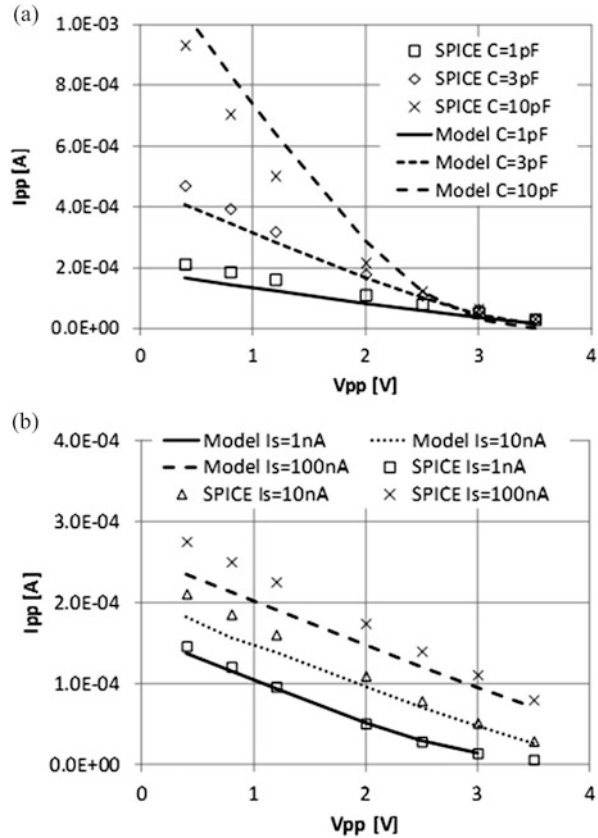
Fig. 4.6 (a, b) Current-limited regime and (c, d) voltage-limited regime

larger and shorter when the clock frequency is doubled. As a result, the transferred charge in a period ( $q_{PP}$ ) can be unchanged with the frequency varied. This means that the averaged  $I_{PP}$  is proportional to the cycle time. The model includes both regimes in the single form. Note that the load current is described as a square in Fig. 4.6a–d based on the model, but the real waveform is not a square.

Figure 4.7a compares the model Eq. (4.13) with the SPICE simulation results in terms of  $C$ . The  $V_{PP} - I_{PP}$  curve is close to a linear line when  $V_{PP}$  is not too high. In this region, one can use a linearized model Eqs. (4.17)–(4.19a, 4.19b). SPICE results show that  $I_{PP}$  does not increase as much as  $C$  increased. For example,  $I_{PP}$  can increase by a factor of 5 when  $C$  increases from 1 to 10 pF with a constant  $I_S$ . The model illustrates this trend. Figure 4.7b compares the model Eq. (4.13) with the SPICE simulation results in terms of  $I_S$ . When  $I_S$  increases by a factor of 10, the  $I-V$  curves shift by about 1 V in X direction. An increase in  $I_S$  by a factor of 10 is equivalent to a decrease in  $V_{TH}^{EFF}$  Eq. (4.19b) by about 60 mV, resulting in a total shift in  $V_{MAX}$  by about 1 V with  $N = 16$ .

Figure 4.8 compares the model Eq. (4.13) with the SPICE simulation results in terms of  $V_{PP}$  and  $N$ . Curves in upper and lower half of Fig. 4.8 refer to  $Y$ -axis at the right- and left-hand side, respectively. Because the model aligns with the SPICE results for the maximum attainable voltage (where  $I_{PP}$  approaches zero) better than it aligns for the maximum attainable current (where  $V_{PP}$  approaches zero), the discrepancy in  $V_{PP}$  in the model from the SPICE results increases as  $I_{PP}$  increases. For example,  $V_{PP}$  is predicted to be 1 V by the model when  $I_{PP}$  is 150  $\mu A$  when  $N = 20$  whereas it is predicted to be about 2 V by the SPICE results. On the contrary,  $V_{PP}$  is predicted to be 3.5 V by the model when  $I_{PP}$  is 50  $\mu A$  when  $N = 20$  whereas it is predicted to be about 3.7 V by the SPICE result. When  $I_{PP}$  is compared at a given  $V_{PP}$ , the discrepancy of the model from the SPICE results is

**Fig. 4.7**  $V_{PP}$  vs.  $I_{PP}$  under various (a)  $C$  and (b)  $I_S$  values in case of  $N = 16$  and  $V_{DD} = 0.5$  V



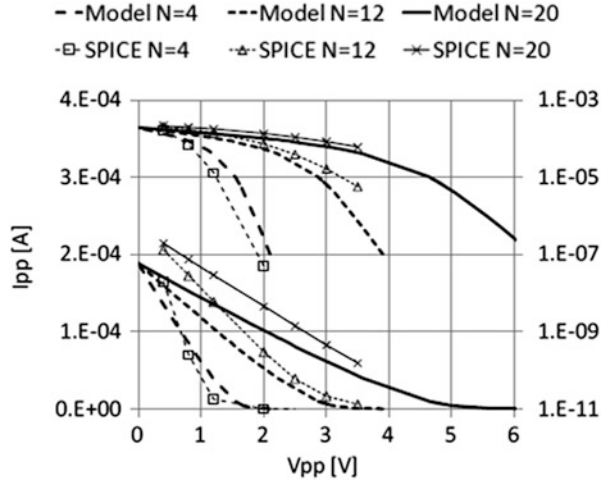
unchanged across  $V_{PP}$  percentage wise. For example, the discrepancy in  $I_{PP}$  of the model from the SPICE results is about 20 % across  $V_{PP}$  when  $N = 20$ .

Figure 4.9 compare the model Eq. (4.25) with the SPICE results in terms of  $V_{PP} - I_{PP}$ , respectively. The discrepancy of the model increases as  $I_{PP}$  increases. The discrepancy is nominally smaller than 50 % across the parameter range described in Table 4.1, except for an extremely low current regime.

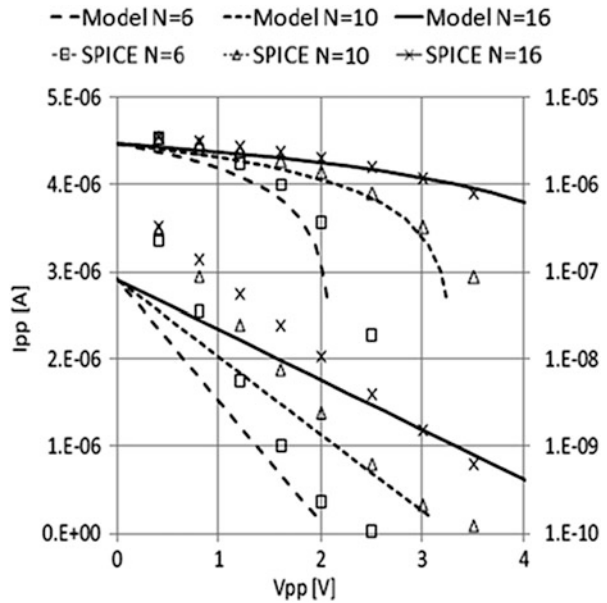
Figure 4.10 compares the model Eq. (4.13) with the SPICE results in terms of (a)  $I_{PP}$  and (b)  $V_{PP}$  on the output power normalized by  $N$ . The curve for  $I_{PP}$  vs. the power normalized by  $N$  is independent of  $N$ . The power density has a peak at about half of the maximum attainable  $I_{PP}$  or  $V_{PP}$ . This trend is the same as DC–DC charge pump as described in Chap. 3.

Figure 4.11a, b compares the model Eq. (4.13), Eq. (4.21) with SPICE results on  $P_{OUT}$  and  $P_{IN}$  over  $V_{PP}$ , respectively. As  $V_{DD}$  increases, the maximum attainable output voltage increases. The peak power point, which is located around a half of the maximum attainable output voltage, also increases accordingly. On the other hand, the input power monotonically decreases as  $V_{PP}$  increases because  $P_{IN}$  is proportional to  $I_{PP}$ , which decreases as  $V_{PP}$  increases as shown by Eq. (4.21). As a

**Fig. 4.8** Comparisons of the model Eq. (4.13) for switching diodes with the SPICE simulation in the case of  $N = 4, 12,$  and  $20$ . Model and SPICE results are respectively shown by *lines* and *symbols*. Lower and upper curves are in linear- and log-plots, respectively (Tanzawa 2014)

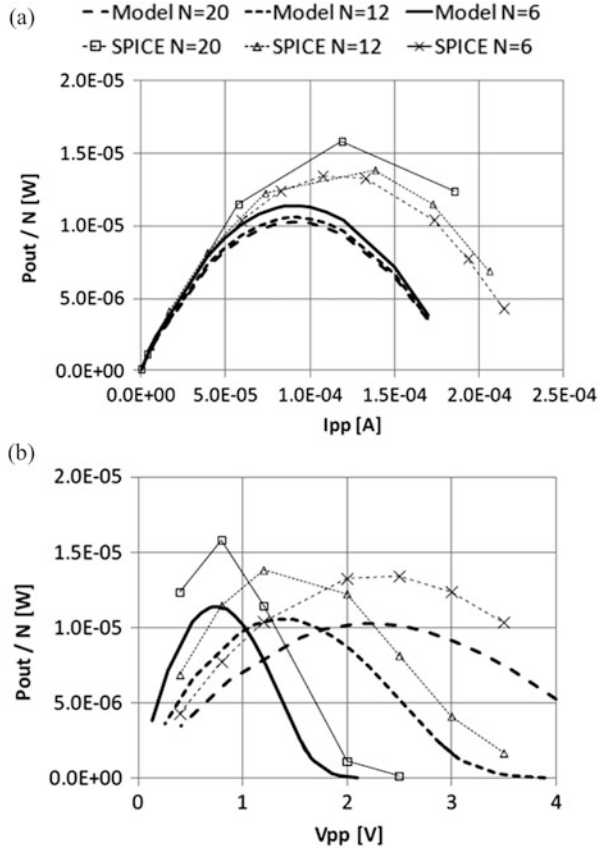


**Fig. 4.9** Comparisons of the model Eq. (4.25) for MOSFETs with the SPICE simulation in the case of  $N = 6, 10,$  and  $16$ . Model and SPICE results are respectively shown by *lines* and *symbols*. Lower and upper curves are in linear- and log-plots, respectively (Tanzawa 2014)



result,  $P_{EFF}$  increases as  $V_{PP}$  increases as shown in Fig. 4.11c. Figure 4.11d shows the same curves as Fig. 4.11c except for the voltage gain  $V_{PP}/V_{DD}$  in X-axis. To the first order,  $P_{EFF}$  is proportional to the voltage gain for a given  $N$ .

**Fig. 4.10** Output power normalized by  $N$  vs. (a)  $I_{PP}$  and (b)  $V_{PP}$  in case of  $N = 6, 12,$  and  $20$



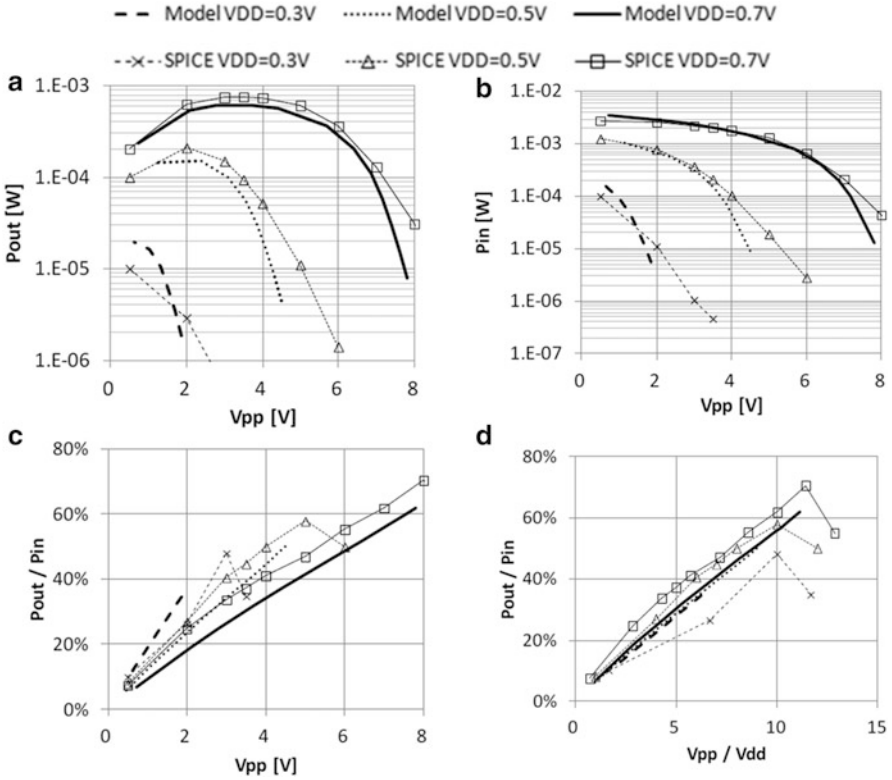
### 4.1.3 Optimum Design

This section discusses two optimizations: (1) maximizing the output power under a given circuit area and a given output voltage and (2) making a balance between the input power and circuit area under a given output voltage and current.

#### 4.1.3.1 Optimization for Maximizing the Output Power

As discussed for DC–DC multipliers in Chap. 3, one can use the Lagrange multiplier to introduce the functions  $f$  and  $g$ , and the parameter  $\lambda$  as follows:  $f(C, N) = CN$ ,  $g(C, N, \lambda) = I_{PP} - \lambda f$ . Using Eqs. (4.17)–(4.19a, 4.19b),  $\partial g/\partial N = \partial g/\partial C = 0$  results in Eq. (4.31).



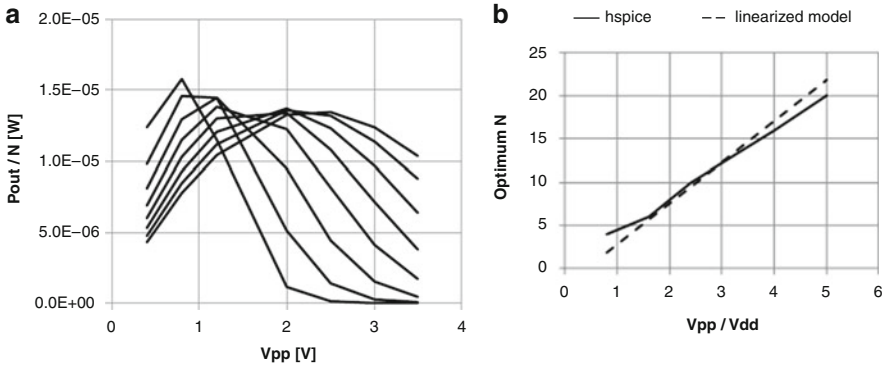


**Fig. 4.11** (a)  $P_{OUT}$  vs.  $V_{PP}$ , (b)  $P_{IN}$  vs.  $V_{PP}$ , (c)  $P_{EFF}$  vs.  $V_{PP}$ , (d)  $P_{EFF}$  vs.  $V_{PP}/V_{DD}$  in case of  $N = 16$ , and  $V_{DD} = 0.3, 0.5$ , and  $0.7$  V (Tanzawa 2014)

$$N_{OPT} = 2V_{PP} / \left( \frac{V_{DD}}{1 + \alpha_T} - V_T \ln \left( \frac{fCV_{DD}}{I_S} \right) \right) \quad (4.31)$$

In this equation, it is assumed that  $N$  is much larger than 1 and  $V_{DD}$  is much larger than  $V_T$ .

Figure 4.12a shows the output power normalized by  $N$  as a function of  $V_{PP}$ . There is the peak power point for each multiplier with a different  $N$ . One can find an optimum  $N$  per  $V_{PP}$  from Fig. 4.12a. The relation of  $N_{OPT}$  to the voltage gain  $V_{PP}/V_{DD}$  is shown in Fig. 4.12b.  $N_{OPT}$  is determined by the number of stages of about four times larger than the voltage gain. For example, when one designs the multiplier to have a voltage gain of 4, one must use  $N$  of about 16.



**Fig. 4.12** (a) Output power normalized by  $N$  vs.  $V_{PP}$  under  $N=6-20$  stages by 2-stage step and (b) optimum  $N$  to maximize the output power normalized by  $N$

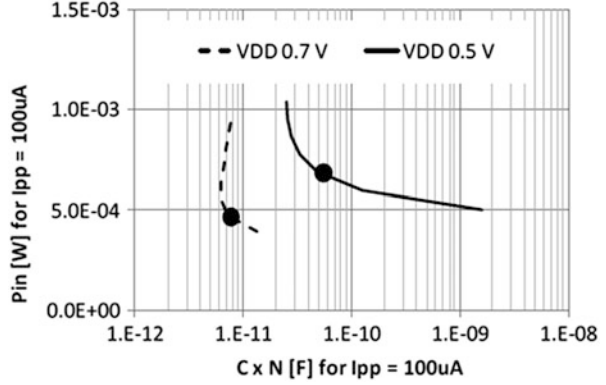
### 4.1.3.2 Optimization for Making a Balance Between the Input Power and the Circuit Area

Another optimization is minimizing the input power. Because there is a significant trade-off between the input power and the circuit area, the actual optimization needs to take the circuit area into account with minimizing the input power.

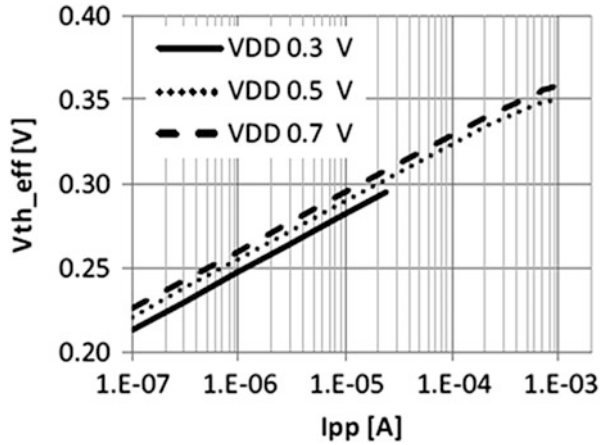
Figure 4.13 demonstrates an area and power trade-off.  $C$  is determined per a given  $N$  when  $I_{PP}$  needs to be  $100 \mu\text{A}$  at  $V_{PP} = 2.5 \text{ V}$ . Then,  $I_S$  is determined to be proportional to  $C$ . For  $N$  below a critical number per  $V_{PP}$ ,  $C$  has to increase rapidly for a multiplier with a small number of stages. This is because the target  $V_{PP}$  comes close to the maximum attainable voltage with a small  $N$ . Conversely, the input power monotonically increases as  $N$  increases as shown by Eq. (4.21). As a result, Fig. 4.13 suggests that one has to select an optimum number of stages to ensure the circuit area does not become too large even for minimizing the input power. As  $V_{DD}$  increases, the capacitors per stage as well as the total number of capacitors decreases significantly.

This is because  $V_{TH}^{EFF}$  is approximately  $0.33 \text{ V}$  at  $I_{PP} = 100 \mu\text{A}$ , as shown in Fig. 4.14. A voltage headroom when  $V_{DD} = 0.7 \text{ V}$  is approximately twice as large as when  $V_{DD} = 0.5 \text{ V}$ . On the other hand, the input power is mainly determined by  $N$  rather than  $V_{DD}$ . As a result, moderate optimum points marked by dots in Fig. 4.13 indicate the input power at  $V_{DD} = 0.7 \text{ V}$  is 30–40 % lower than that at  $V_{DD} = 0.5 \text{ V}$ , whereas the circuit area for  $V_{DD} = 0.7 \text{ V}$  is about 10 times smaller than that for  $V_{DD} = 0.5 \text{ V}$ .

**Fig. 4.13**  $P_{IN}$  vs.  $C \times N$  under a constraint of  $I_{PP} = 100 \mu\text{A}$  at  $V_{PP} = 2.5 \text{ V}$  using the model



**Fig. 4.14**  $I_{PP}$  vs.  $V_{TH}^{EFF}$



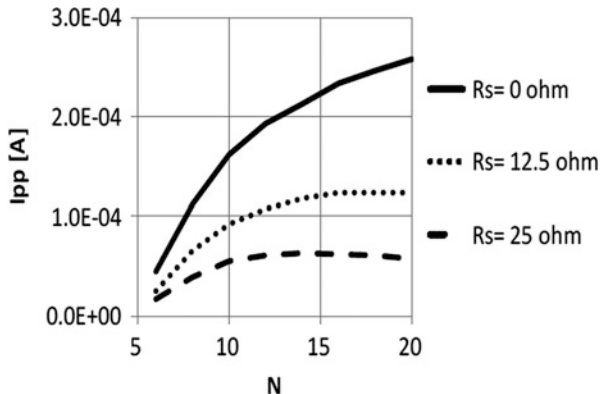
### 4.1.4 Impact of AC Source Impedance

So far, the impedance of the AC power source is ignored. However, antennas in case of RF and UHF energy harvesting and piezoelectric material in case of vibration energy harvesting can have finite impedance. When the impedance of the AC power source is  $R_S$ , the amplitude of the voltage of the source is reduced from  $V_{DD}$  to  $V_{DDS}$ , as given by Eq. (4.32), where it is assumed that the input reactance is canceled out with a proper complex conjugate impedance matching.

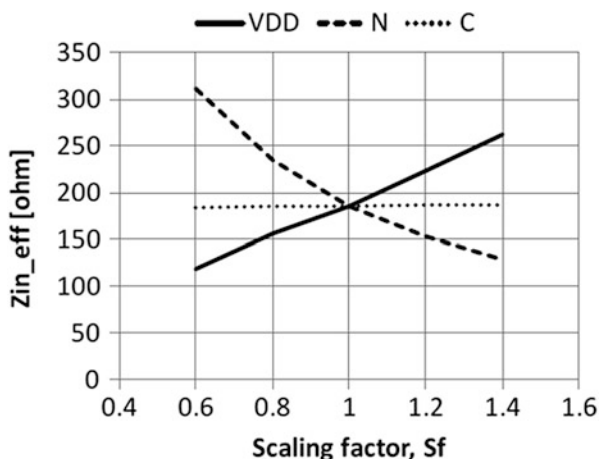
$$V_{DDS} = V_{DD} - R_S I_D (N + 1) \tag{4.32}$$

Using Eqs. (4.9), (4.11), and (4.32), Eq. (4.13) is modified to be Eq. (4.33).

**Fig. 4.15**  $N$  vs.  $I_{PP}$  under various  $Z_S$  with  $V_{DD} = 0.7$  V



**Fig. 4.16** Circuit parameters vs.  $Z_{IN}^{EFF}$



$$V_{PP} = (N + 1) \left( \frac{V_{DD} - \pi R_S I_{PP} (N + 1) / \gamma}{1 + \alpha_T} - \frac{I_{PP}}{fC(1 + \alpha_T)} - V_T \ln \left( \frac{\pi I_{PP}}{I_S \cos^{-1} \left( 1 - \frac{I_{PP}}{fCV_{DD}} \right)} \right) \right) \quad (4.33)$$

Figure 4.15 shows the calculated results using Eq. (4.33). An increase in  $R_S$  can result in more significant reduction in  $I_{PP}$  for the multipliers with a larger number of stages.

Figure 4.16 illustrates the impact of circuit parameters on the effective input impedance  $Z_{IN}^{EFF}$  which is defined by Eq. (4.23). In this plot, starting with a nominal condition of  $f = 1$  G-Hz,  $C = 1$  pF,  $V_{DD} = 0.5$  V,  $V_{PP} = 2.5$  V, and

$N = 20$ , each circuit parameter of  $V_{DD}$ ,  $N$ , and  $C$  is respectively swept by varying a scaling factor  $S_F$  from 0.6 to 1.4, as shown in the horizontal axis. For example, in the case of  $V_{DD}$ ,  $Z_{IN}^{EFF}$  is calculated for  $V_{DD}$  of 0.3–0.7 V whereas the other parameters are unchanged. Likewise, in the case of  $N$ ,  $Z_{IN}^{EFF}$  is calculated for  $N$  of 12–28 whereas the other parameters are unchanged. When  $V_{DD}$  increases,  $Z_{IN}^{EFF}$  increases linearly to  $V_{DD}$ . This is because both  $\bar{V}_{DD}$  and  $P_{IN}$  are approximately proportional to  $V_{DD}$ . Conversely, when  $N$  increases,  $Z_{IN}^{EFF}$  decreases inversely proportional to it. This is because  $\bar{V}_{DD}$  is not a strong function of  $N$  whereas  $P_{IN}$  is approximately proportional to it.  $C$  similarly affects  $(\bar{V}_{DD})^2$  and  $P_{IN}$ , resulting in a very weak impact on  $Z_{IN}^{EFF}$ .

## 4.2 Multi-sine (MS) Wave AC–DC Charge Pump Voltage Multipliers

It was experimentally shown that the communication distance from a radio-frequency identification (RFID) reader to a tag could be extended using multi-sine (MS) waves instead of continuous waves (CW) without increasing the output power of the reader (Trotter et al. 2009; Valenta et al. 2013).

In this section, a closed-form, analytical model of AC–DC rectifier multipliers using MS is described. Using the equations for input and output power, one can estimate the circuit and device parameters as the initial conditions for multiplier design when using MS. This section is organized as follows: Sect. 4.2.1 investigates how modeling of MS multipliers can be connected with that of CW multipliers and describes an analytical model of MS AC–DC multipliers. Section 4.2.2 shows design and device parameter sensitivity on the pump performance. Section 4.2.3 discusses the effectiveness of using MS with respect to power efficiency under a given requirement for output voltage and current.

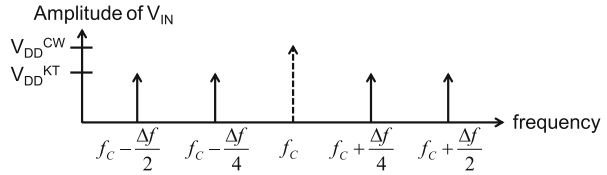
### 4.2.1 Circuit Model

Equally weighted  $N$  subcarriers combined in-phase is considered to have the highest power efficiency because the highest peak voltage is available most frequently with in-phase subcarriers, as described by Eq. (4.34).

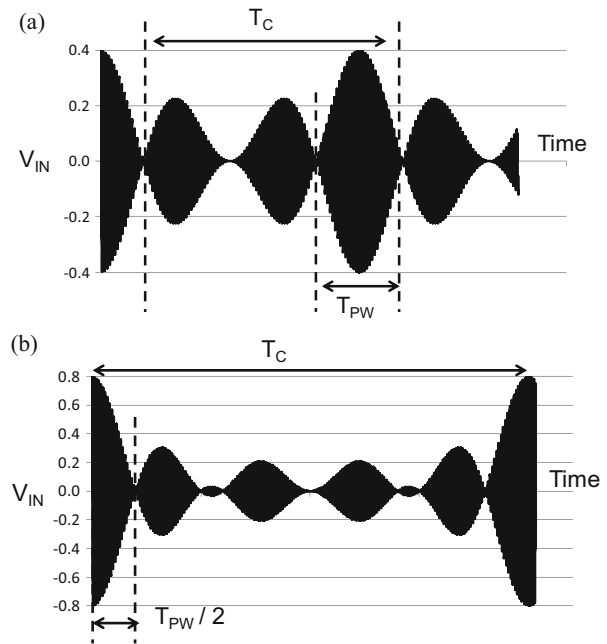
$$\Phi_{2K}(t) = \frac{V_{DD}^{KT}}{2} \sum_{k=1}^K \sin 2\pi \left( f_C \pm k \frac{\Delta f}{2K} \right) t \quad (4.34)$$

$K$  is a positive integer,  $2K$  is the number of subcarriers,  $V_{DD}^{KT}$  is the voltage amplitude of each subcarrier,  $f_C$  is the center frequency (in the case of MS) or the

**Fig. 4.17** Voltage amplitude of subcarriers



**Fig. 4.18** Input voltage waveform in cases of (a) four and (b) eight subcarriers



carrier frequency (in the case of a CW),  $\Delta f$  is the bandwidth of MS. Figure 4.17 shows an example of the multi-sine signal with four subcarriers.

Assumptions made in this section are: (1) the parasitic resistance of the multiplier capacitors is small in comparison with the resistance of the diodes; (2) the parasitic capacitance of the diodes is small in comparison with the multiplier capacitors; (3) the output impedance of the AC power source is small in comparison with the input impedance of the multiplier; (4) the filtering capacitor is large enough to maintain the output voltage  $V_{PP}$  as a constant under a given output current; (5) the reverse bias current of the diode is small in comparison with the forward bias current.

Figure 4.18 shows the voltage waveform Eq. (4.34) for (a) four and (b) eight subcarriers with  $V_{DD}^{KT}$  of 0.1 V. The peak voltage  $V_{DD}^{MAX}$  is 0.4 and 0.8 V, respectively.  $T_C$  and  $T_{PW}$  are the cycle time of  $V_{IN}$  and the pulse width of the highest

pulse of  $V_{IN}$ , respectively. To investigate how many subcarriers are required to be recognized as a representative of multi-sine AC input, SPICE simulation was run under the condition of  $N = 16$ ,  $C = 1$  pF,  $C_T = 10$  fF,  $I_S = 10$  nA,  $V_T = 30$  mV,  $f_C = 1$  GHz,  $\Delta f = 40$  MHz, and  $K = 1, 2, 3, 4$  where  $V_{DD}$  and  $V_{PP}$  are varied. SPICE diode model parameters used in this section are the same as in Sect. 4.1, which is summarized in Table 4.2.

Figure 4.19a shows the power efficiency as a function of the output power with a different number of subcarriers at  $V_{PP} = 2.5$  V. Figure 4.19b shows the number of subcarriers as a function of the duty  $T_{PW}/T_C$ . With the number of subcarriers at four or greater, the power efficiency is unchanged significantly. As a result, a multi-sine AC input with  $K = 2$  is selected in the following studies.

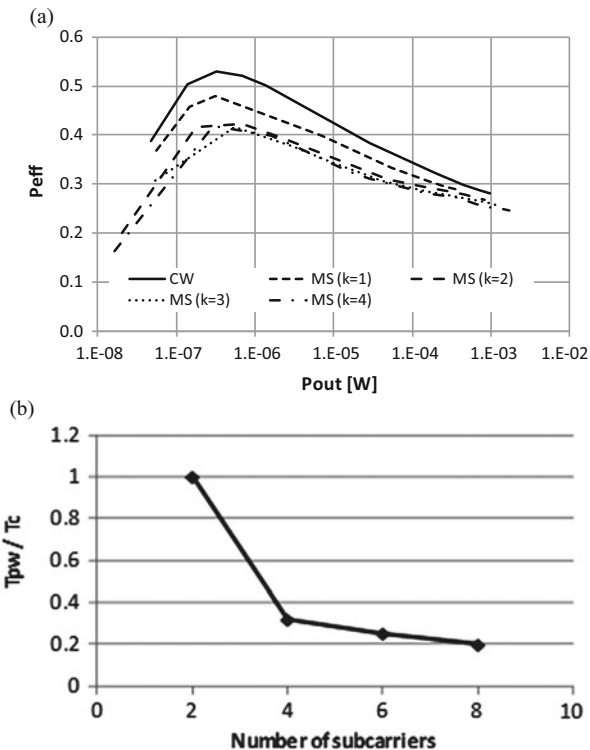
Figure 4.20 compares (a) output power, (b) input power, and (c) power efficiency between CW and MS. In Fig. 4.20a,  $V_{DD}$  was selected for each CW and MS to maximize  $P_{OUT}$  at  $V_{PP} = 3$  V. Then,  $C$  and  $I_S$  were scaled to have  $P_{OUT}$  of  $150 \mu\text{W}$  at  $V_{PP} = 3$  V.  $P_{OUT}$  of MS seems to be less dependent on  $V_{PP}$  than CW. In Fig. 4.20c,  $P_{OUT}$  varied by changing  $V_{DD}$  from 0.3 to 0.65 V for CW and from 0.4 to 0.8 V for MS, and  $P_{EFF}$  was measured as a function of  $P_{OUT}$ . As a result, CW has higher power efficiency by 5–10 % than MS in a range in  $P_{OUT}$  between 1 and  $200 \mu\text{W}$ .

To investigate the characteristics of MS rectifier multipliers, simulated voltage and current is decomposed per pulse. Figure 4.21a illustrates the waveform of  $V_{IN}$ ,  $I_{OUT}$ ,  $I_{IN}$  at  $V_{DD}^{4T} = 0.25$  V and  $V_{PP} = 2.5$  V. The majority of charges is transferred from one capacitor to another with 16 pulses, as shown in Fig. 4.21b.  $j$ -th pulse has its own voltage amplitude of  $V_{DD}(j)$ . The input and output power and the power efficiency are averaged by pulse, as shown in Fig. 4.21c. The power efficiency gradually decreases as the number of pulses increases. However, the variation is relatively small. Figure 4.21d shows  $P_{IN}$  and  $P_{OUT}$  in each pulse as a function of  $V_{DD}(j)$  based on Fig. 4.21c. Figure 4.21d also includes the case of CW where  $V_{DD}$  varies from 0.5 to 1.0 V.  $P_{IN}$  and  $P_{OUT}$  in 9th to 16th pulses are a little lower than those in 1st to 8th pulses. However, both trends are so close that  $P_{IN}$  and  $P_{OUT}$  in both trends can be considered to be as much as those of CW. This suggests that a MS rectifier multiplier can be analyzed using the model of CW, and therefore the model Eq. (4.13) is applicable to MS.  $P_{EFF}$  seems to be constant for pulses with different  $V_{DD}(j)$  in a part of the period of  $V_{IN}$ . This indicates that the power efficiency of MS averaged in  $T_C$  would not be higher than that of CW, which is opposed to some experimental results. Section 4.2.3 discusses more in details.

Figure 4.22 illustrates the highest pulse of  $V_{IN}$ . The envelope is fit by a single cosine whose cycle time is given by  $T_{PW}$ , as shown in Fig. 4.18.

$$V_{DD}(t) = V_{DD}^{\text{MAX}} \cos(2\pi t/T_{PW}) \quad (4.35)$$

**Fig. 4.19** (a) Output power vs. power efficiency with different number of subcarriers under the condition of  $V_{PP} = 2.5$  V and (b) the number of subcarriers vs. the duty  $T_{PW}/T_C$  (Tanzawa 2015)



where  $V_{DD}^{MAX}$  is  $2K V_{DD}^{KT}$ . A linearized AC–DC multiplier model for CW with a constant amplitude of  $V_{DD}$  is given by Eqs. (4.17)–(4.19a, 4.19b). The output power is given by Eq. (4.36).

$$P_{OUT} = I_{PP} V_{PP} \tag{4.36}$$

Because  $V_{DD}$  is a time-varying parameter for MS AC–DC multipliers, Eqs. (4.17) and (4.19a, 4.19b) are rewritten to be Eqs. (4.37) and (4.38), respectively.

$$I_{PP}(t) = (V_{MAX}(t) - V_{PP})/R_{PMP} \tag{4.37}$$

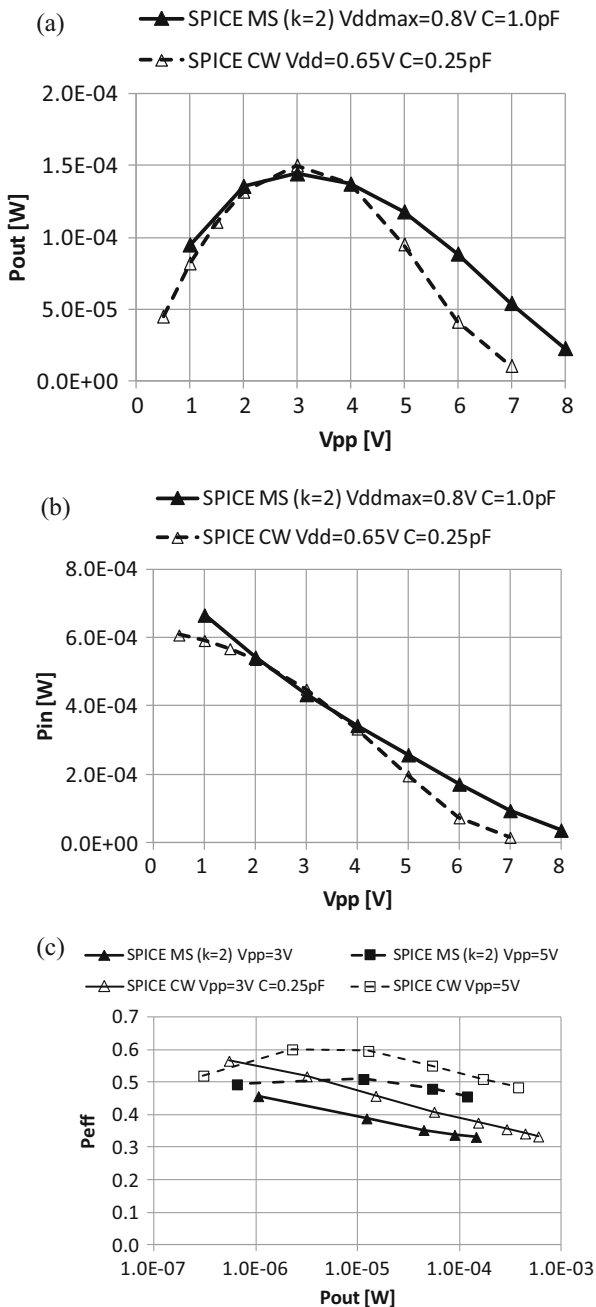
$$V_{MAX}(t) = (N + 1) \left( \frac{V_{DD}(t)}{1 + \alpha_T} - V_T \ln \left( \frac{f C V_{DD}(t)}{I_S} \right) \right) \tag{4.38}$$

As suggested in Fig. 4.22, Eq. (4.35) can be approximated by Eq. (4.39) to the second order because the last few pulses don't contribute to  $I_{PP}$  and  $I_{DD}$  much:

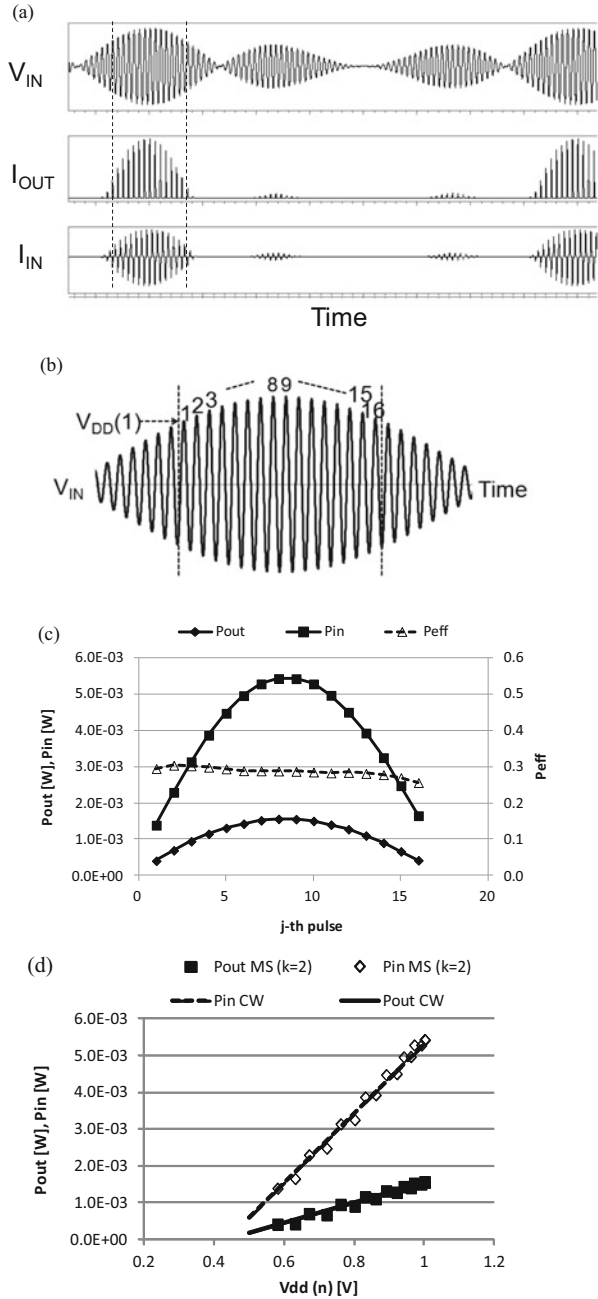
$$V_{DD}(t) = V_{DD}^{MAX} \left( 1 - (2\pi t/T_{PW})^2/2 \right) \tag{4.39}$$



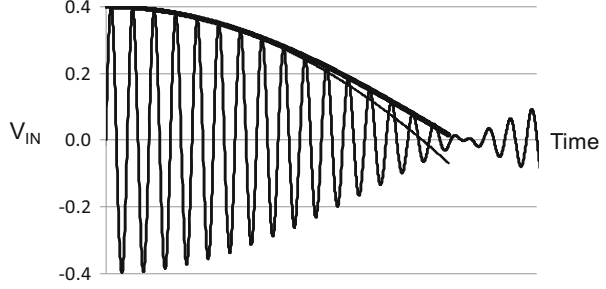
**Fig. 4.20** (a)  $V_{PP}$  vs.  $P_{OUT}$ ,  
 (b)  $V_{PP}$  vs.  $P_{IN}$ , (c)  $P_{OUT}$   
 vs.  $P_{EFF}$  (Tanzawa 2015)



**Fig. 4.21** (a) Definition of  $N$  and  $V_{DD}(N)$ , (b)  $N$ -th pulse vs.  $P_{IN}$ ,  $P_{OUT}$ ,  $P_{EFF}$ , (c)  $V_{DD}(N)$  vs.  $P_{IN}$  and  $P_{OUT}$  at  $V_{DD} = 0.25$  V and  $V_{PP} = 2.5$  V (Tanzawa 2015)



**Fig. 4.22**  $V_{IN}$  with four subcarriers and its envelopes fit by a single cosine Eq. (4.35) in *bold line* and by Eq. (4.39) in *thin line* (Tanzawa 2015)



From Fig. 4.22, Eq. (4.39) is valid as long as the time concerned is shorter than approximately  $2\pi/3 T_{PW}$ . As a result, Eq. (4.38) is approximated by:

$$V_{MAX}(t) = (N + 1) \left( \frac{V_{DD}^{MAX} \left( 1 - (2\pi t/T_{PW})^2/2 \right)}{1 + \alpha_T} - V_T \ln \left( \frac{f C V_{DD}^{MAX}}{I_S} \right) + V_T (2\pi t/T_{PW})^2/2 \right) \quad (4.40)$$

An averaged output current over one cycle  $T_C$ ,  $I_{PP}^{MS}$ , is defined for MS by:

$$I_{PP}^{MS} = \frac{2}{T_C} \int_0^{t_0} I_{PP}(t) dt \quad (4.41)$$

where  $t_0$  is the time when the output current becomes negligibly small in comparison with the peak. For example,  $t_0$  is between 16th and 17th pulses in case of Fig. 4.21a, b. Equation (4.41) is analytically calculated with Eqs. (4.18), (4.37), and (4.40) by ignoring the terms which include  $(t_0/T_{PW})^4$  or higher orders.

$$I_{PP}^{MS} = \frac{4\delta T_{PW}}{3\pi R_{PMP} T_C} (V_A - V_C) \quad (4.42)$$

where  $\delta$ ,  $V_A$ ,  $V_B$ ,  $V_C$  are given by Eqs. (4.33)–(4.36), respectively.

$$\delta = \sqrt{2(V_A - V_C)/(V_A - V_B)} \quad (4.43)$$

$$V_A = (N + 1) V_{DD}^{MAX} / (1 + \alpha_T) \quad (4.44)$$

$$V_B = V_T (N + 1) \quad (4.45)$$

$$V_C = V_T (N + 1) \ln(f C V_{DD}^{MAX} / I_S) + V_{PP} \quad (4.46)$$

Because the output voltage is a constant  $V_{PP}$ , the averaged output power for MS is given together with Eq. (4.42) by:

$$P_{OUT}^{MS} = I_{PP}^{MS} V_{PP} \quad (4.47)$$

Next, the averaged input power is estimated. An AC–DC multiplier model for CW with a constant amplitude of  $V_{DD}$  is given by Eq. (4.21). Using an approximate equation for  $\sin \gamma/\gamma$  as given by Eq. (4.48),

$$\frac{\sin \gamma}{\gamma} = 1 - 0.34 \frac{I_{PP}}{fCV_{DD}} \quad (4.48)$$

Equation (4.18) is rewritten to be Eq. (4.49).

$$P_{IN} = (N + 1) \left( V_{DD} I_{PP} - 0.34 \frac{I_{PP}^2}{fC} \right) \quad (4.49)$$

Because  $V_{DD}$  and  $I_{PP}$  are time-varying parameters in case of MS, Eq. (4.49) is given by Eq. (4.50).

$$P_{IN}(t) = (N + 1) \left( V_{DD}(t) I_{PP}(t) - 0.34 \frac{I_{PP}(t)^2}{fC} \right) \quad (4.50)$$

An averaged input power over one cycle  $T_C$ ,  $P_{OUT}^{MS}$ , is defined for MS by:

$$P_{IN}^{MS} = \frac{2}{T_C} \int_0^{t_0} P_{IN}(t) dt \quad (4.51)$$

Equation (4.51) is analytically calculated with Eqs. (4.18), (4.37), and (4.40) by ignoring the terms which include  $(t_0/T_{PW})^4$  or higher orders.

$$P_{IN}^{MS} = \frac{2(1 + \alpha_T)\delta T_{PW}}{3\pi R_{PMP}} \frac{1}{T_C} \frac{1}{V_A - V_B} \times \quad (4.54)$$

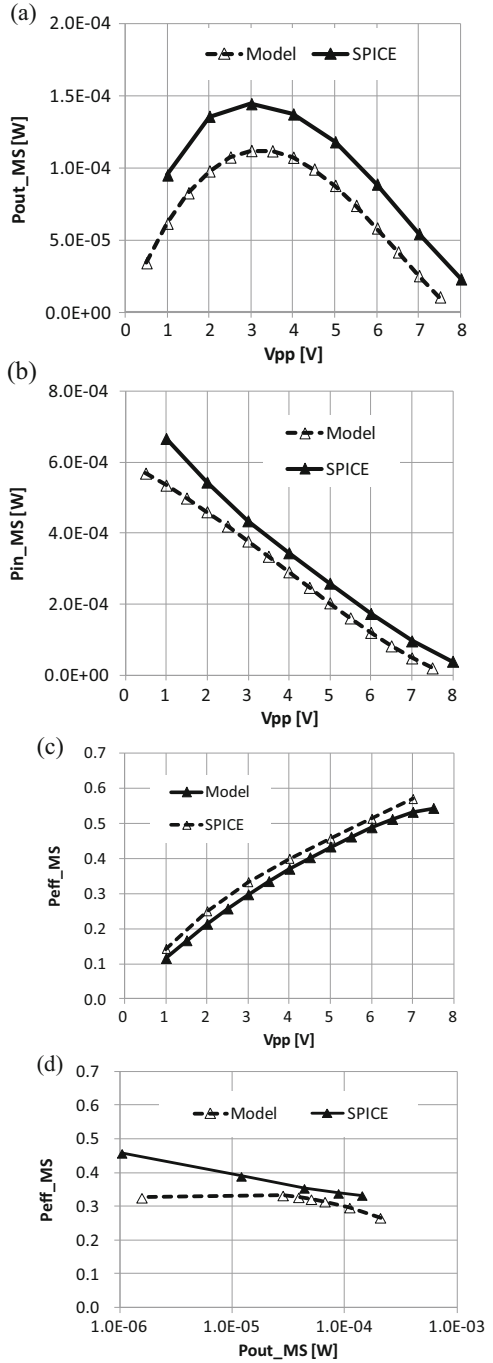
$$\begin{aligned} & (0.66V_A^3 - 1.66V_A^2V_B + 0.68V_A^2V_C \\ & - 1.34V_AV_C^2 + 1.42V_AV_BV_C + 0.34V_BV_C^2) \end{aligned}$$

An averaged power over one cycle  $T_C$ ,  $P_{EFF}^{MS}$ , is given together with Eqs. (4.42), (4.47), and (4.54) by:

$$P_{EFF}^{MS} = \frac{P_{OUT}^{MS}}{P_{IN}^{MS}} \quad (4.55)$$

Figure 4.23 compares the model with the SPICE results for MS. The parameters used are the same as those for Fig. 4.21.  $V_{DD}^{MAX}$  of 0.8 V is used for Fig. 4.23a–c. The differences between the model Eq. (4.47), Eq. (4.54) and SPICE simulation seems to be consistent, differing by a DC offset. This is mainly because a very extreme case where the number of stages is as many as 24 is considered to validate the model. As shown in Fig. 4.24a, the DC offset is reduced as the number of stages is reduced. Even though there is a little difference in the effective threshold voltage of the switching diode between the model and SPICE simulation, it can be

**Fig. 4.23** Comparison of the model and SPICE results in terms of (a)  $V_{PP}$  vs.  $P_{OUT}^{MS}$ , (b)  $V_{PP}^{MS}$  vs.  $P_{IN}^{MS}$ , (c)  $P_{OUT}^{MS}$  vs.  $P_{EFF}^{MS}$ .  $V_{DD}^{MAX}$  of 0.8 V is used for (a–c).  $V_{PP}$  of 3 V is used for (d) (Tanzawa 2015)



accumulated in offsets in  $P_{\text{OUT}}$  and  $P_{\text{IN}}$ . The offsets do not affect the optimum  $V_{\text{PP}}$  where  $P_{\text{OUT}}^{\text{MS}}$  is maximized matches well. Equation (4.47) reproduces a long tail in  $P_{\text{OUT}}^{\text{MS}}$  at a higher  $V_{\text{PP}}$  with MS, as shown in Fig. 4.23a. Equation (4.54) also shows the tendency that  $P_{\text{IN}}^{\text{MS}}$  decreases monotonically as  $V_{\text{PP}}$  increases, as shown in Fig. 4.23b. Because of the discrepancy of the model from the SPICE results in terms of  $P_{\text{OUT}}^{\text{MS}}$  and  $P_{\text{IN}}^{\text{MS}}$ , the deviation in  $P_{\text{EFF}}^{\text{MS}}$  from the SPICE results is about 5 %, as shown in Fig. 4.23c.  $V_{\text{DD}}^{\text{MAX}}$  is swept from 0.4 to 0.8 V at  $V_{\text{PP}} = 3$  V to draw Fig. 4.23d. The model is in agreement with the SPICE results in  $P_{\text{OUT}}^{\text{MS}}$  of 10–100  $\mu\text{W}$ ; however, the model shows much lower  $P_{\text{EFF}}^{\text{MS}}$  at  $P_{\text{OUT}}^{\text{MS}}$  lower than 10  $\mu\text{W}$  compared with the SPICE results. At a lower  $P_{\text{OUT}}^{\text{MS}}$ , the conduction time is so short that the assumption used for the model would not be valid; that is, the average  $P_{\text{OUT}}^{\text{MS}}$  and  $P_{\text{IN}}^{\text{MS}}$  could not be calculated by Eqs. (4.42) and (4.54).

### 4.2.2 Design and Device Parameter Sensitivity on the Pump Performance

Figure 4.24 compares the model results with the SPICE results in terms of  $P_{\text{OUT}}^{\text{MS}}$  and  $P_{\text{EFF}}^{\text{MS}}$  as a function of (a)  $N$ , (b)  $V_{\text{DD}}^{\text{MAX}}$ , where  $N = 24$ , (c)  $C$ , (d)  $f$ , (e)  $I_{\text{S}}$ , and (f)  $\alpha_{\text{T}}$ . The typical condition is  $N = 16$ ,  $C = 1$  pF,  $C_{\text{T}} = 10$  fF,  $I_{\text{S}} = 10$  nA,  $V_{\text{T}} = 30$  mV,  $f_{\text{C}} = 1$  GHz,  $\Delta f = 40$  MHz,  $V_{\text{DD}}^{\text{MAX}} = 0.8$  V, and  $V_{\text{PP}} = 3$  V. The maximum discrepancies in  $P_{\text{OUT}}^{\text{MS}}$  and  $P_{\text{EFF}}^{\text{MS}}$  from the SPICE results are approximately 30 and 5 %, respectively.  $P_{\text{OUT}}^{\text{MS}}$  is proportional to  $N$ ,  $V_{\text{DD}}^{\text{MAX}}$ ,  $C$ , and  $f$ , and  $P_{\text{EFF}}^{\text{MS}}$  is inversely proportional to  $N$  to the first order.

### 4.2.3 On the Effectiveness of Multi-sine Wave Over Continuous Wave

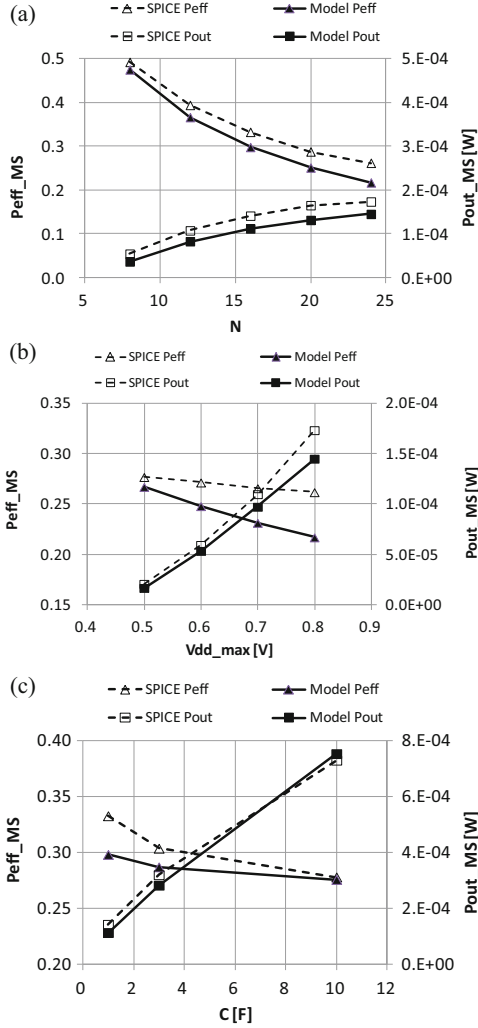
Figure 4.25 compares  $P_{\text{EFF}}$  between CW and MS with the SPICE results and the model under a constraint for a constant  $I_{\text{PP}}$ .

$P_{\text{EFF}}$  of MS is about 5 % lower than that of CW across a wide  $V_{\text{PP}}$  range. Considering one cycle  $T_{\text{C}}$  is divided into two periods, one of which has averaged output (input) power of  $P_{\text{OUT}}^{\text{H}}$  ( $P_{\text{IN}}^{\text{H}}$ ) in the most significant period denoted by  $T_{\text{PW}}$  in Fig. 4.2, and the other of which has averaged output (input) power of  $P_{\text{OUT}}^{\text{L}}$  ( $P_{\text{IN}}^{\text{L}}$ ) in the remaining period, the average  $P_{\text{OUT}}$  and  $P_{\text{IN}}$  in one cycle are written as:

$$P_{\text{OUT}}^{\text{MS}} = DP_{\text{OUT}}^{\text{H}} + (1 - D)P_{\text{OUT}}^{\text{L}} \quad (4.56)$$

$$P_{\text{IN}}^{\text{MS}} = DP_{\text{IN}}^{\text{H}} + (1 - D)P_{\text{IN}}^{\text{L}} \quad (4.57)$$

where  $D$  is the duty cycle  $T_{\text{C}}/T_{\text{PW}}$ . To compare the MS multiplier with CW multiplier,  $P_{\text{OUT}}^{\text{CW}}$  is selected to be as much as  $P_{\text{OUT}}^{\text{MS}}$ .



**Fig. 4.24** Comparison of the model and SPICE results in terms of  $P_{\text{OUT}}^{\text{MS}}$  and  $P_{\text{EFF}}^{\text{MS}}$  as a function of (a)  $N$ , (b)  $V_{\text{DD}}^{\text{MAX}}$ , where  $N = 24$ , (c)  $C$ , (d)  $f$ , (e)  $I_S$ , and (f)  $\alpha_T$ . The typical condition is  $N = 16$ ,  $C = 1$  pF,  $C_T = 10$  fF,  $I_S = 10$  nA,  $V_T = 30$  mV,  $f_C = 1$  GHz,  $\Delta f = 40$  MHz,  $V_{\text{DD}}^{\text{MAX}} = 0.8$  V, and  $V_{\text{PP}} = 3$  V (Tanzawa 2015)

$$P_{\text{OUT}}^{\text{CW}} = P_{\text{OUT}}^{\text{MS}} \quad (4.58)$$

Using other parameters  $\Delta P_{\text{IN}}^{\text{H}}$  and  $\Delta P_{\text{IN}}^{\text{L}}$ ,  $P_{\text{IN}}^{\text{H}}$  and  $P_{\text{IN}}^{\text{L}}$  can be expressed by the following equations:

$$P_{\text{IN}}^{\text{H}} = P_{\text{IN}}^{\text{CW}} + \Delta P_{\text{IN}}^{\text{H}} \quad (4.59)$$

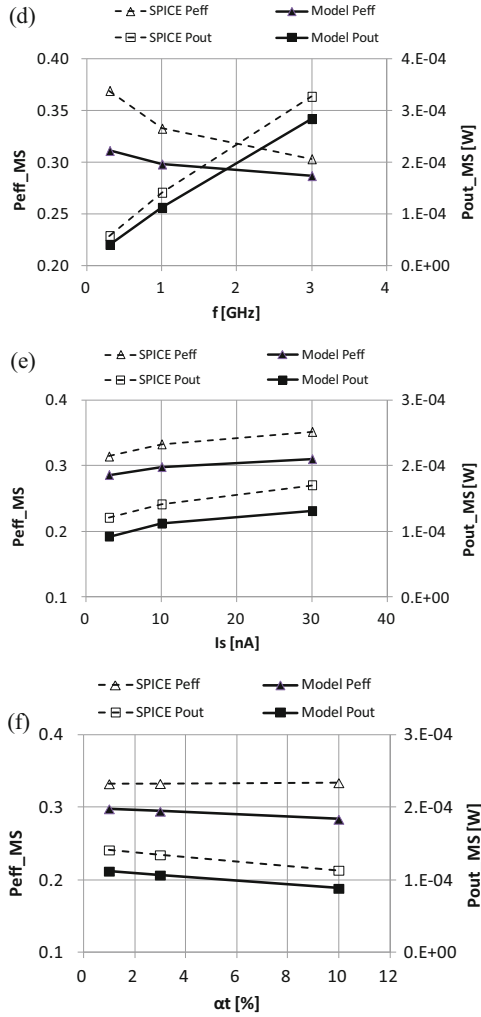
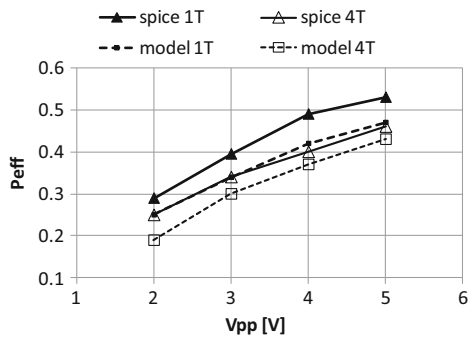


Fig. 4.24 (continued)

Fig. 4.25  $V_{PP}$  vs.  $P_{EFF}$  under a constraint of  $I_{PP} = 100 \mu A$  (Tanzawa 2015)





$$P_{IN}^L = P_{IN}^{CW} - \Delta P_{IN}^L \quad (4.60)$$

Then,  $P_{EFF}^{MS}$  is calculated as follows using Eqs. (4.56)–(4.60):

$$P_{EFF}^{MS} = \left( 1 + \frac{D\Delta P_{IN}^H - (1-D)\Delta P_{IN}^L}{P_{IN}^{CW}} \right)^{-1} P_{EFF}^{CW} \quad (4.61)$$

As shown in Fig. 4.21a, b,  $\Delta P_{IN}^L/\Delta P_{IN}^H \ll 10$  and  $D < 0.3$  for the number of sub-carriers of 4 or more. As a result, Eq. (4.61) is approximated to be Eq. (4.62).

$$P_{EFF}^{MS} \approx \left( 1 + \frac{D\Delta P_{IN}^H}{P_{IN}^{CW}} \right)^{-1} P_{EFF}^{CW} \quad (4.62)$$

Therefore, Eq. (4.62) results in Eq. (4.62).

$$P_{EFF}^{MS} < P_{EFF}^{CW} \quad (4.63)$$

Equation (4.62) is inconsistent with experimental results.

There are several possible reasons on inconsistency with prior measured results. One would have to see if the output was DC. If there is a large ripple at the output terminal, measurements would have been done with different DC in the output. Larger ripple translates into higher DC, which can result in higher power efficiency with MS. The model and SPICE simulations have not considered input impedance, which is different from the condition for measurement. However, if impedance matching between the antenna and rectifying multiplier is effectively perfect, the model and SPICE results should be valid with an input voltage amplitude of half  $V_{DD}$ . The tendency that MS has higher power efficiency than CW wouldn't be affected. Or, if the effect of the input impedance of voltage multiplier on power efficiency is greater with MS than CW because of larger amplitude, MS could be affected by the input impedance more than CW. If the matching has not been perfect, but the input impedance has been capacitive, power factor might have been away from its optimal point. Then, the performance comparison would have to be corrected by matching the input impedance with the source impedance.

As the future work, one will have to study the significance of those on power efficiency for CW and MS. One or some of those may be the reason(s) for the inconsistency with the existing measured results. More work need to be done for determining the root cause of the inconsistency.

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### ***Circuit Model for AC–DC Charge Pump Voltage Multipliers***

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- T. Tanzawa, in *An Analytical Model of Multi-Sine AC-DC Voltage Multiplier*. IEEE International Conference on Circuits and Systems, pp. 1354–1357

# Chapter 5

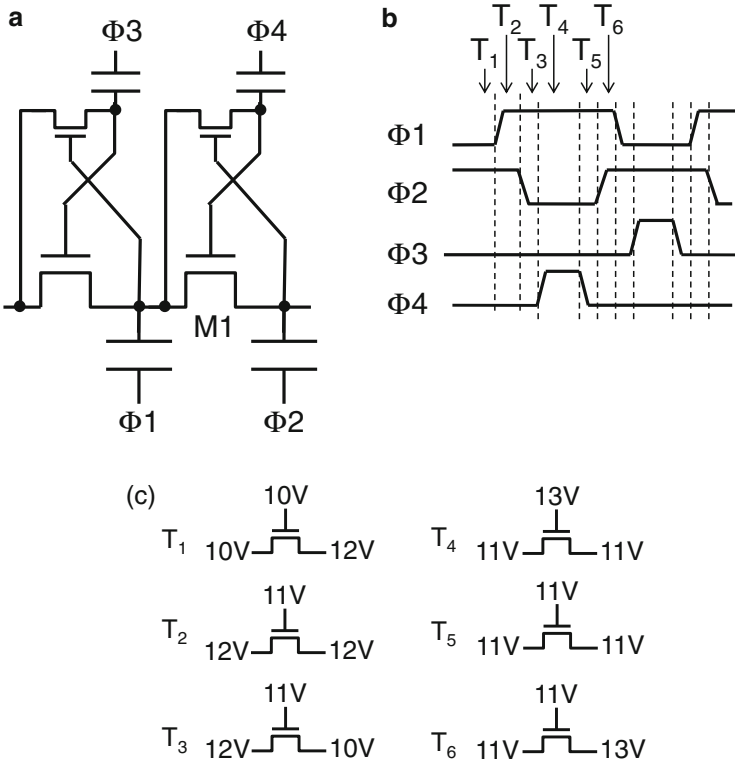
## Charge Pump State of the Art

**Abstract** This chapter discusses design techniques for implementing charge pumps in integrated circuits. Charge pumps are composed of transfer transistors and capacitors. Realistic design needs to take parasitic components such as threshold voltages of the transfer transistors and parasitic capacitance at each of both terminals into account. In order to decrease the pump area and to increase the current efficiency, some techniques such as threshold voltage canceling, stage reconfiguration, and faster clocking are presented. Since the supply current has a frequency component as high as the operating clock, noise reduction technique is another concern for pump design. In addition to design technique for individual pump, system level consideration is also important, since there are usually more than one charge pump in a chip. Area reduction can be also done for multiple charge pump system where all the pumps do not work at the same time. Wide supply voltage range operation and stand-by pump design are also discussed.

This chapter starts with switching diode design in Sect. 5.1, mainly focusing on how the threshold voltage of the transistor and its body effect can be mitigated to increase the output current under a given circuit area. Section 5.2 presents capacitor structures as well as design technique for reducing the top plate parasitic capacitance. Remaining sections discuss control methods for the pumps to operate stably even when the supply voltage can vary widely in Sect. 5.3, to reduce the total pump area in case where two pumps operate in different periods in Sect. 5.4, to decrease noise against the power supply and ripple in the output voltage in Sect. 5.5, and to have stability in the output voltage when stand-by and active pumps are used in Sect. 5.6.

### 5.1 Switching Diode Design

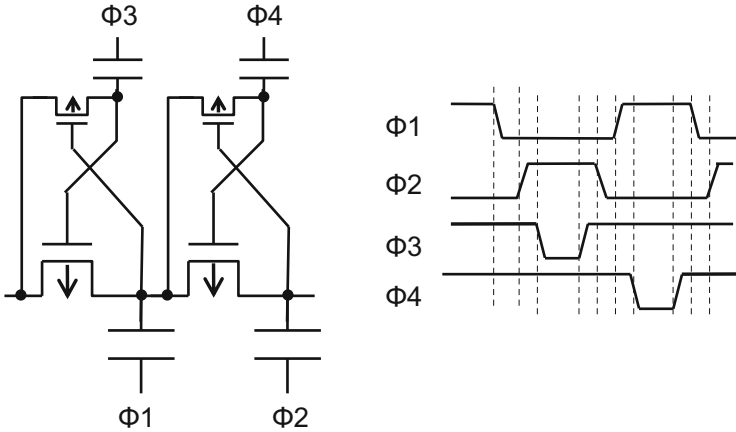
Dickson successfully generated a higher voltage on chip, but the supply voltage was much higher than the threshold voltage of the transfer transistor. According as the supply voltage of an LSI decreases for scaling the transistors and for reducing the power, the impact of the threshold voltage of the transfer transistors on the output



**Fig. 5.1**  $V_T$  cancellation configuration (a) with four nonoverlapping phases (b) and the bias condition (c) (e.g., D'Arrigo et al. 1989; Umezawa et al. 1992)

current becomes significant. This section discusses various types of switching diode implementations to mitigate the impact of the threshold voltage and its body effect.

Figure 5.1a describes a positive voltage multiplier with  $V_T$  cancellation by means of four nonoverlapping phases. In order to increase the gate voltage of the transfer transistors, auxiliary capacitors driven by  $\Phi 3, 4$  and transistors are added to the original devices. Figure 5.1c illustrates the bias condition of the transfer gate M1 at different timings where  $V_{DD}$  is 2 V and  $V_T = 1$  V. The transfer gate M1 fully equalizes the two next neighbor capacitors at  $T_4$  resulting in  $V_T$  cancellation whereas fully turns off at  $T_6$ . Because the timing margins between the phases are needed, the operation frequency is lower than the original Dickson pump with two phases. In case where the threshold voltages of the transfer transistors are the main contributors to limit the output current, the four nonoverlapping phases can improve the pump performance. On the other hand, if a pump running with a faster clock can output a higher current even with a finite threshold voltage, one should choose two phases. Which type is better depends on the threshold voltage, the supply voltage, and the frequency available in a given technology.



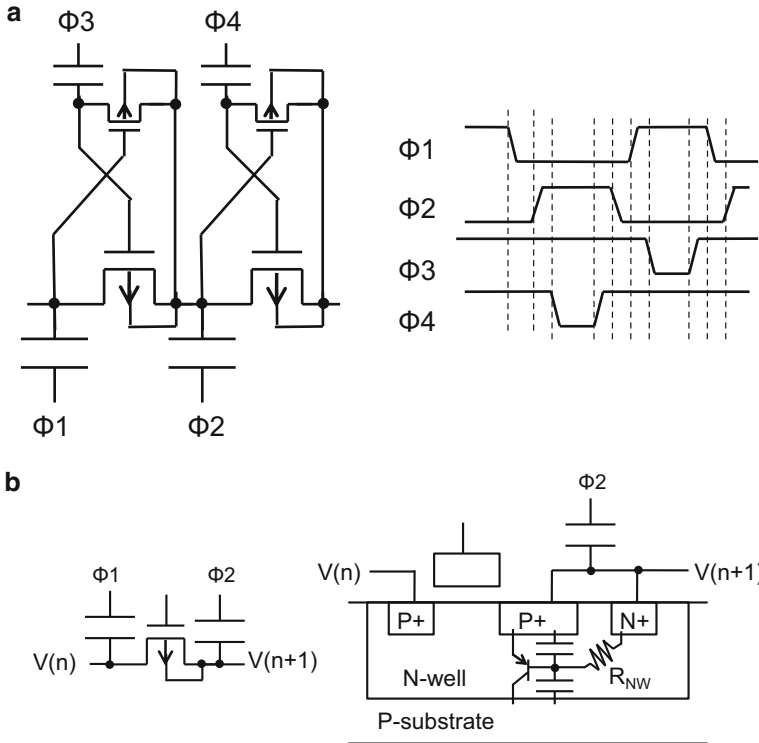
**Fig. 5.2**  $V_T$  cancellation with four nonoverlapping phases for negative voltage generation (Kuriyama et al. 1992)

Figure 5.2 shows a complementary type of Fig. 5.1 outputting a negative voltage. The transfer transistors are pMOSFETs. Their N-well is connected to the supply voltage to prevent source and drain junctions from entering a forward bias regime in entire operation. If that happens, the amount of charges could flow into P-substrate, resulting in reduction in the output current.

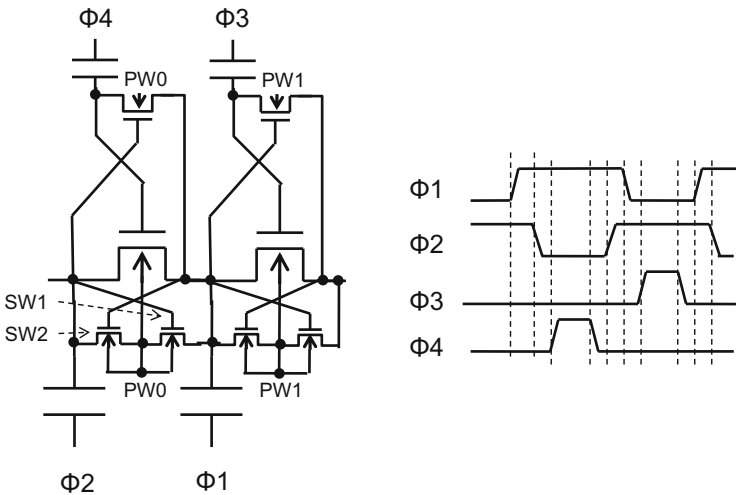
In order to eliminate the body effect of the PMOS transfer transistors, the bodies per stage are connected to a separated N-well, as shown in Fig. 5.3a. Thus,  $V_{GS}$ ,  $V_{DS}$ , and  $V_{BS}$  of every PMOS can be always limited within  $2V_{DD}$ . Because the N-well is connected to the capacitor node, there are several periods when the parasitic bipolar junction transistors could turn on. For instance, when  $\Phi1$  goes high while  $\Phi2$  stays low, the PN junction at the drain could enter in a forward bias regime until the potential of the effective base of the parasitic bipolar junction transistor (bjt) as shown in Fig. 5.3b is recovered after it receives a capacitive coupling of the junction capacitance. A part of the injected current flows into the P-type source and the rest flows into the P-substrate. Device design including the layout of the PMOSFETs and their N-well to reduce the P-substrate current as much as possible is a key to make the pump functional. As the supply voltage decreases, the current via the parasitic bjt decreases.

The body-effect-cancellation with P-well potential control can avoid the potential leakage current due to the parasitic bjt as shown in Fig. 5.4. Isolated P-well is connected to either one of the capacitor nodes with a lower potential. Because the transistors SW1, 2 are added, the parasitic capacitance at the upper terminal of the capacitors increases, resulting in lower boosting ratio. This is the trade-off for this method.

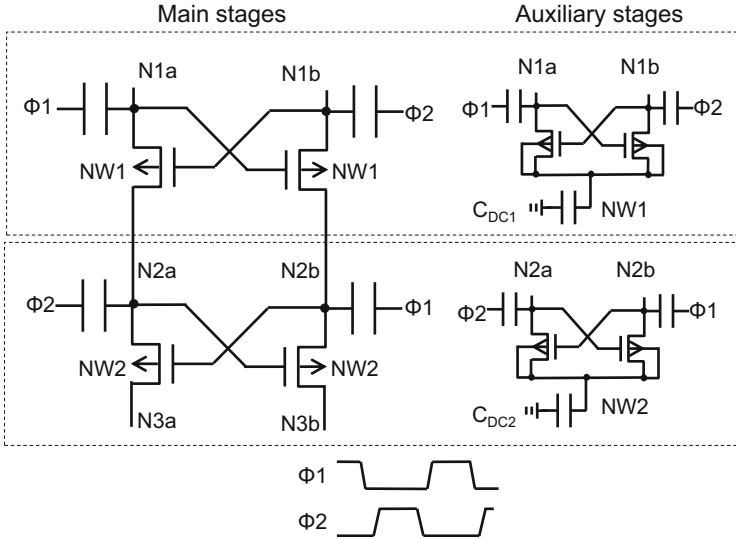
Figure 5.5 illustrates another topology using two phase clock to reduce the number of phases for faster clocking. To stabilize the potential of each isolated N-well, decoupling capacitors are added. The rise time can be affected, but the



**Fig. 5.3** (a) Body-effect-cancellation with divided N-well (Sawada et al. 1995). (b) Parasitic bipolar junction transistor formed in the divided N-well



**Fig. 5.4** Body-effect-cancellation with P-well potential control (Bloch et al. 1998)



**Fig. 5.5** Two phase body-effect-cancellation (Favrat et al. 1998)

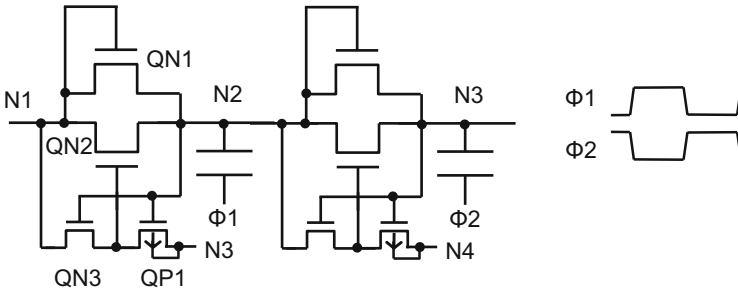
output current wouldn't once the pump enters a steady state because the charging current is not needed for the decoupling capacitors afterward. The transfer gates can have thin gate oxide because any voltage difference between the four terminals is smaller than  $2V_{DD}$ . On the other hand, the main, auxiliary, and decoupling capacitors have to have thick gate oxide because a voltage as high as  $V_{PP}$  is applied to them.

Another  $V_T$  cancelation technique was reported in Fig. 5.6. Besides a diode-connected transfer gate QN1, QN2 is connected in parallel, whose gate voltage is borrowed by the capacitor node N3 of the next stage. Thus, the auxiliary capacitor is not needed. When  $\Phi_1 = L, \Phi_2 = H$ , QP1 turns on to pass the potential at N3 to the gate of QN2. When  $\Phi_1 = H, \Phi_2 = L$ , QN3 connects the potential at N1 to the gate of QN2, resulting in turning QN2 off. During the transition in  $\Phi_1$  and  $\Phi_2$ , there can be the timing when both QN3 and QP1 turn on, resulting in a reverse current flowing from N3 to N1. When this happens, the output current thereby the power efficiency could be reduced. Therefore, the timing margin between  $\Phi_1$  and  $\Phi_2$  is a key design parameter.

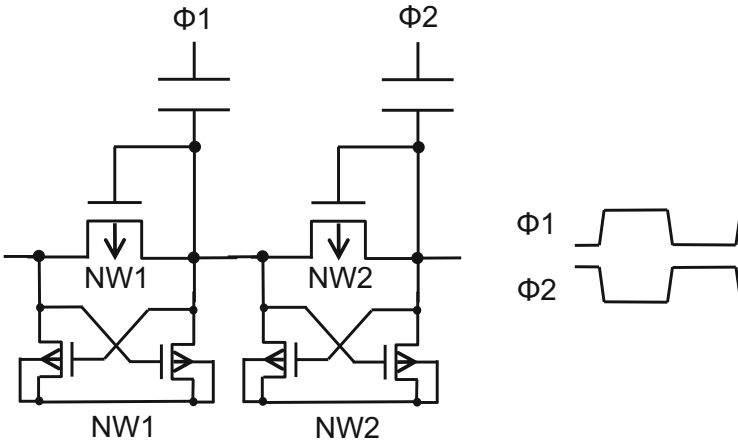
Figure 5.7 illustrates a two phase clock pump with body effect cancelation. Because the transfer gate is connected as a diode, the threshold voltage at  $V_{BS} = 0V$ , the so-called  $V_{T0}$ , does affect the transfer efficiency. However, when transistors with low  $V_{T0}$  is available, this topology may have a lower voltage difference between the gate and source, resulting in a lower stress on the transistor.

To reduce the parasitic capacitance in addition to body effect cancelation, the source terminals of the transfer gates are connected to the P-well by stage as proposed in Fig. 5.8. After  $\Phi_1$  goes high and before  $\Phi_4$  goes high, current flows

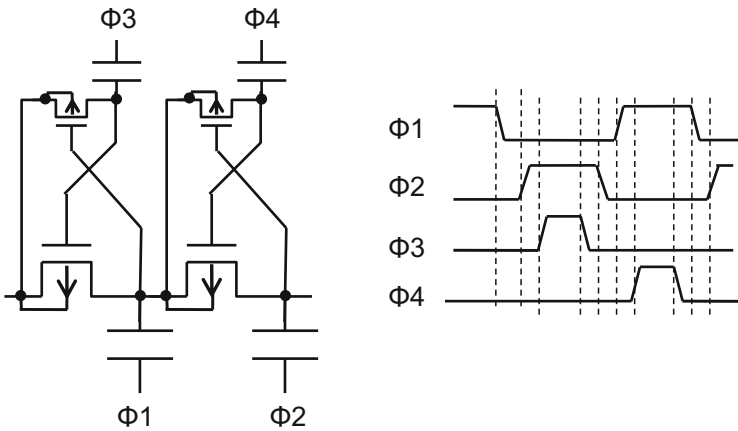




**Fig. 5.6** Two phase CMOS  $V_T$  cancellation (Wu and Chang 1998)

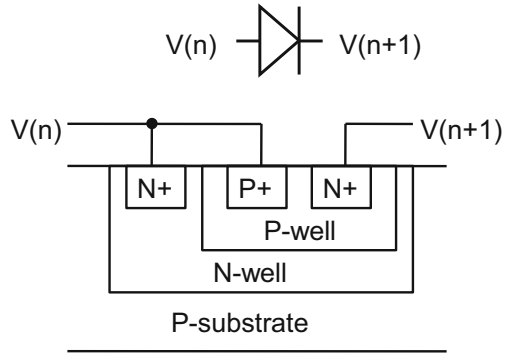


**Fig. 5.7** Body-effect-cancellation with two phase clocks (Shin et al. 2000)

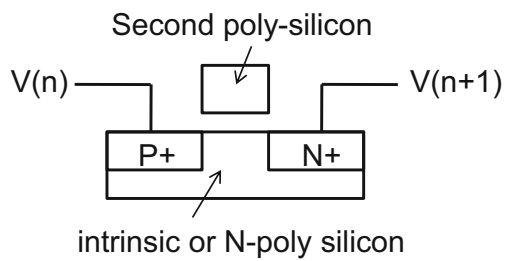


**Fig. 5.8** Body-effect-cancellation with isolated P-well connected to the source of the transfer gates by stage (Javanifard et al. 2008)

**Fig. 5.9** Diode in the substrate (Storti et al. 1988; Kobayashi et al. 1995)



**Fig. 5.10** Poly-Si diode (Mihara et al. 1999)



through a parasitic diode composed of the P-well and the drain N+ junction. As far as the leakage current to P-substrate via the parasitic bjt is sufficiently small in comparison with the current from the pumping capacitor to the next one after  $\Phi_4$  goes high, high-voltage generation is realized.

PN diode is not suffered from the body effect unlike transistor. Figure 5.9 realizes it using triple well structures. To prevent a parasitic bjt from flowing current to the substrate, N-well is connected with P-well. Sheet resistance of P-well is usually lower than that of N-well. In case where the difference is relatively large, the propagation delay from the N-well terminal to the center of P-well is longer than that from the P-well terminal to the center of P-well. If the potential difference reaches its built-in potential, the current can flow from P-well to P-substrate. Therefore, the diode size put in a single P-well needs to be small enough to be able to neglect such a difference in the propagation delay.

Using Flash memory structure, Poly-Si diode is fabricated as shown in Fig. 5.10. Second Poly-Si is used as hard mask to form P+ and N+ at source and drain. Because the source and drain have no connection with P-substrate, there is no parasitic bjt. On the other hand, thin-film poly-silicon devices have much lower mobility than bulk ones do. This is the trade-off for the Poly-Si diode.

## 5.2 Capacitor Design

This section discusses realization of capacitors. N-well capacitor may be able to be fabricated without any significant process cost, as shown in Fig. 5.11a. The N-well terminal can be driven by a clock whose voltage ranges from 0 V to  $V_{DD}$ . The gate oxide is usually thick enough to sustain a high-voltage generated by a pump. There are some parasitic capacitance components associated with the pump capacitor such as a junction capacitance between N-well and P-substrate and a fringe capacitance between the gate edge to the P-substrate. When the interconnection layers pass across the capacitor, it provides another parasitic capacitance to the gate node. When a charge pump is needed in a mixed signal LSI, metal-insulator-metal or polysilicon-insulator-polysilicon capacitor may be available, as shown in Fig. 5.11b. The maximum allowable voltage for the capacitor may restrict using the MIM/PIP capacitor. In advanced silicon technology, many interconnection layers are available. Figure 5.11c shows the cross-sectional view of three interconnection layers. The top and bottom layers are routed in a direction parallel to the sheet and the middle one is routed in the direction perpendicular to the sheet. When the middle portion of the second layer is connected to a terminal of the capacitor and the surrounding portions are connected to another terminal, the capacitance between the two terminals is the sum of the four parasitic capacitors as shown in Fig. 5.11c.

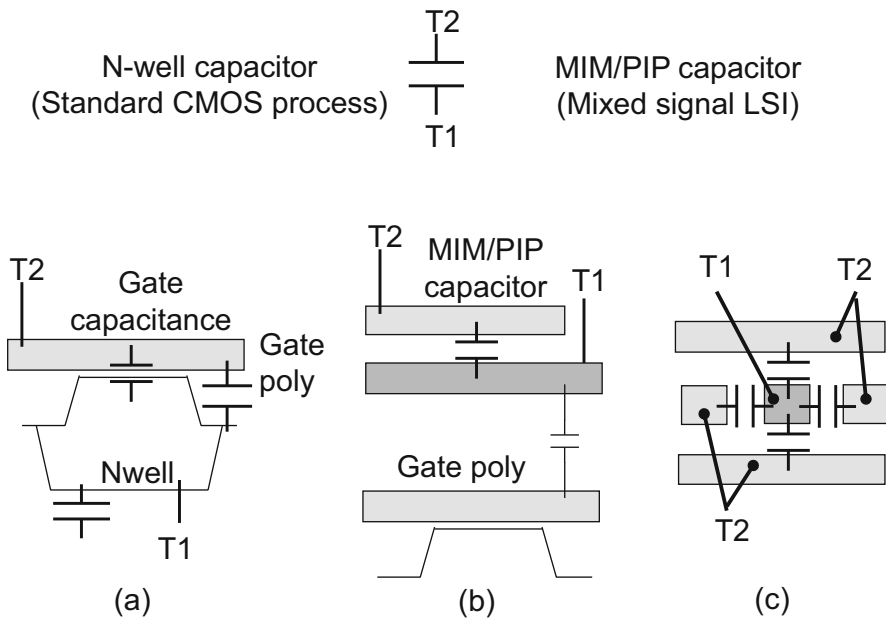


Fig. 5.11 Realization of capacitors

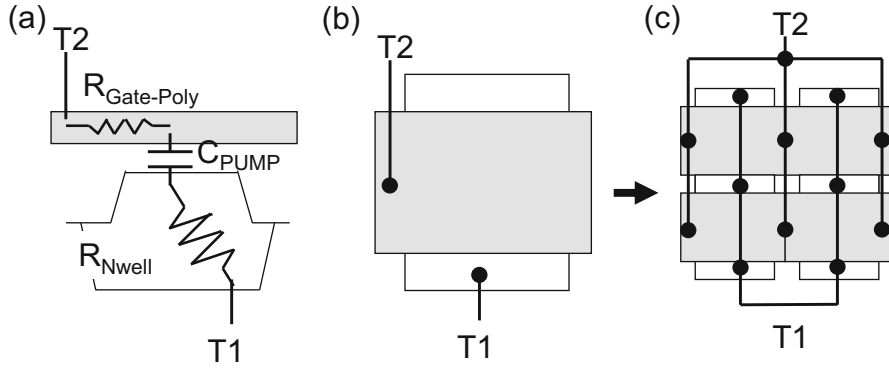


Fig. 5.12 Realization of capacitors with smaller parasitic RC time constant

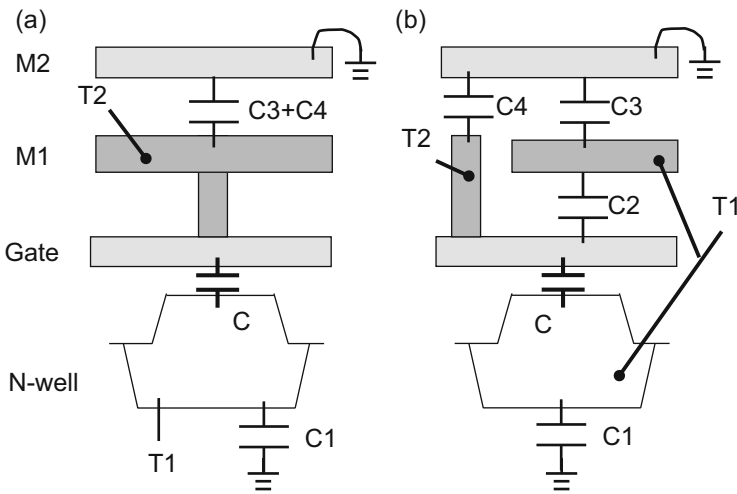


Fig. 5.13 Interconnections to pump capacitors

One needs to make sure that the RC time constant associated with the capacitance of the pump capacitor and parasitic resistance such as gate resistance and well resistance is much smaller than the timing difference between different phases. When the single plate gate is large in terms of the RC time constant as shown in Fig. 5.12a, b, one may have to divide the capacitor into multiple small pieces to make RC time constant of each piece small enough, as shown in Fig. 5.12c.

Figure 5.13 illustrates two different routing to the two terminals of N-well capacitors. The gate is connected with a wide M1 layer in Fig. 5.13a whereas with a narrow M1 layer and another M1 layer over the gate is connected with the terminal  $T_1$  which is connected with N-well in Fig. 5.13b. Table 5.1 compares each capacitance component of the pump capacitor  $C_{\text{CP}}$ , the parasitic capacitance at the

**Table 5.1** Comparison of each capacitance of the pump capacitors of Fig. 5.13

	(a)	(b)
$C_{CP}$	$C$	$C + C_2$
$C_T$	$C_3 + C_4$	$C_4$
$C_B$	$C_1$	$C_1 + C_3$
$\alpha_T$	$(C_3 + C_4)/C$	$C_4/(C + C_2)$

top plate  $C_T$ , and the parasitic capacitance at the bottom plate  $C_B$ . The routing in Fig. 5.13b increases  $C_{CP}$  by  $C_2$  and decreases  $C_T$  by  $C_3$  at a sacrifice of increased  $C_B$  by  $C_3$ , resulting in a smaller ratio of  $C_T$  to  $C_{CP}$ , i.e.,  $\alpha_T$ , than the routing in Fig. 5.13a. This increases the effective clock amplitude and thereby the output current under the same capacitor area. Because of increased  $C_B$ , the power efficiency is equivalent to the first order.

### 5.3 Wide $V_{DD}$ Range Operation Design

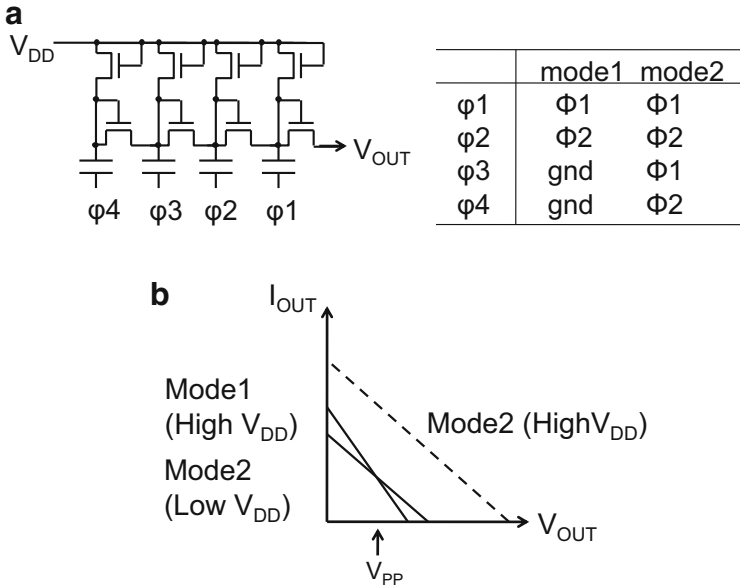
This section investigates the impact of variation in  $V_{DD}$  on  $I_{OUT}$  and  $I_{DD}$  for applications requiring a wide  $V_{DD}$  operation. The design equations for  $I_{OUT}$  and  $I_{DD}$  are given by Eqs. (2.87) and (2.89), respectively. One can extract the derivatives as follows.

$$\frac{dI_{PP}}{dV_{DD}} = \frac{C}{T} \left( 1 + \frac{1 + \alpha_T}{N} \right) \rightarrow \frac{C}{T} \quad (5.1)$$

$$\frac{dI_{DD}}{dV_{DD}} = \frac{C(N+1)}{T} \left( 1 + \frac{1 + \alpha_T}{N} \right) + \frac{NC_B - C_T}{T} \rightarrow \frac{N(C + C_B)}{T} \quad (5.2)$$

The arrows indicate what values are approached to when  $N$  becomes large. The dependence of  $V_{DD}$  on  $I_{PP}$  is not a strong function of  $N$ , but the smaller  $N$  the larger effect on  $I_{PP}$ . The dependence of  $V_{DD}$  on  $I_{DD}$  is a function of  $N^1$ .

Figure 5.14 shows a charge pump with two operational modes in which the number of stages is valuable. Only the last two stages operate in mode 1, whereas all the four stages do in mode 2. Figure 5.14b compares  $I$ - $V$  curves in case of mode 1 at a high  $V_{DD}$ , mode 2 at a high  $V_{DD}$ , and mode 1 at a low  $V_{DD}$ . Two  $I$ - $V$  curves are crossed at  $V_{OUT}$  around  $V_{PP}$ . When the number of stages is reduced as  $V_{DD}$  becomes higher than a critical voltage, the variation in  $I_{OUT}$  across the  $V_{DD}$  operating range can be significantly reduced. Furthermore, in case where the pump is required to output different currents ( $I_{OUT1,2}$ ) at different voltages ( $V_{OUT1,2}$ ) in different period of time, this control method can lower power consumption at  $V_{OUT,1}$  ( $<V_{OUT,2}$ ) because of smaller number of stages.



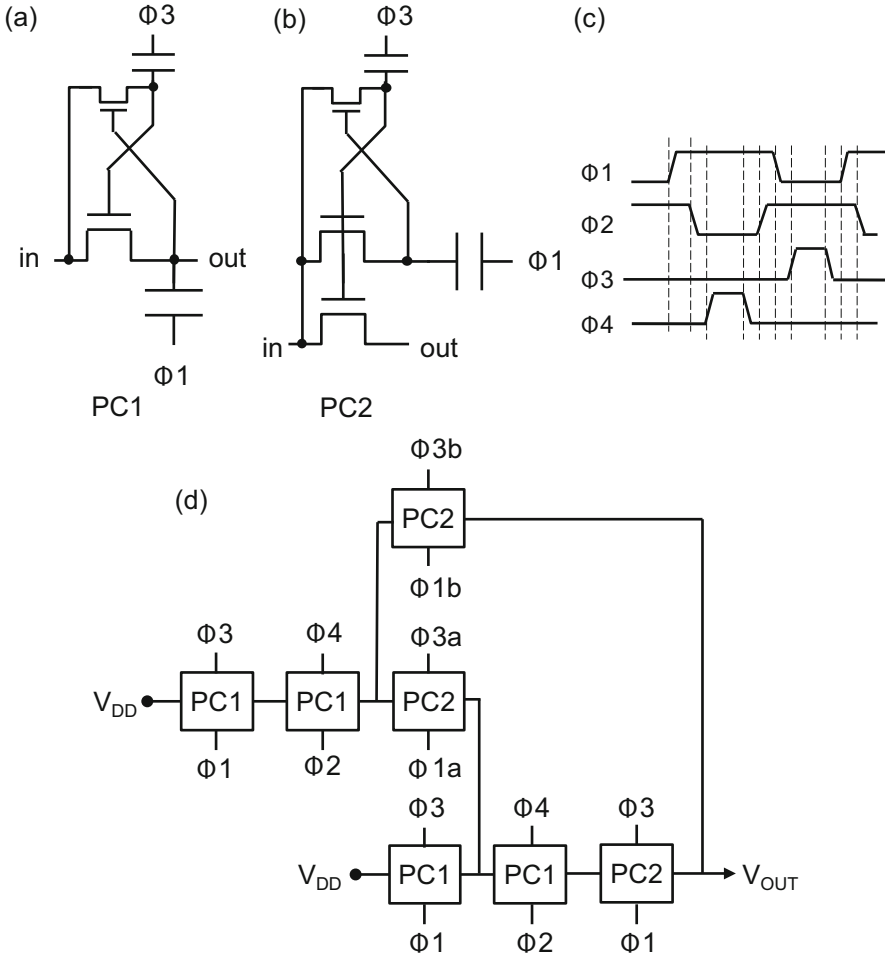
**Fig. 5.14** (a) Low noise pump design for wide  $V_{DD}$  operation with variable number of stages (Gerber et al. 1981). (b) Low noise pump design for wide  $V_{DD}$  operation

### 5.4 Area Efficient Multiple Pump System Design

This section discusses area efficient system design in case where all the multiple charge pumps don't operate simultaneously.

A simple method for generating two different voltages is having two different charge pumps. However, if they are not required to generate at the same time, or in other word, if different high voltages are required in different periods, another method with a single charge pump having additional switches is possible as shown in Fig. 5.15. Figure 5.15a illustrates a unit pump stage cell. The switching circuit shown in Fig. 5.15b is composed of a transfer transistor and a boosting circuit with the same configuration as the unit pump cell. Because the switching circuit doesn't need to transfer large amount of charges, the capacitors used in the switching circuit can be small. Thus, the area for the switching circuit is much smaller than the unit pump cell.

Table 5.2 shows how the additional clocks are given by mode and how the pump is reconfigured. In mode 1, the upper two PC1 stages in Fig. 5.15 are connected with the output terminal in parallel to the lower two PC1 stages. Thus, the pump has a configuration of two arrays of two stages. The output impedance  $R_{PMP}$  and the maximum attainable output  $V_{MAX}$  are respectively given by Eqs. (2.77) and (2.78). The pump in mode 1 has  $R_{PMP}$  of  $T/C$  and  $V_{MAX}$  of  $3V_{DD}$ .  $I_{MAX}$  is defined by  $V_{MAX}/$

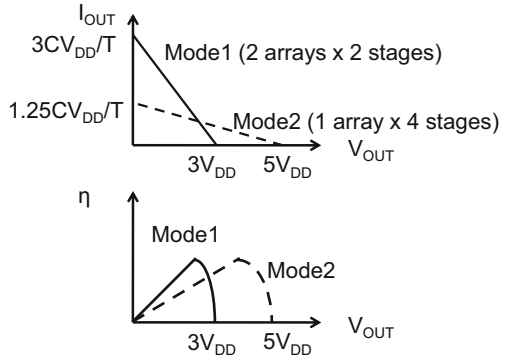


**Fig. 5.15** Pump with variable number of stages and variable effective capacitance per stage (Tanzawa et al. 1997)

**Table 5.2** Clocks for reconfiguring the pump shown in Fig. 5.15

	Mode 1	Mode 2
Φ1a	L	Φ1
Φ3a	L	Φ3
Φ1b	Φ1	L
Φ3b	Φ3	L
# Stages	2	4
# Arrays	2	1
$R_{PMP}$	$T/C$	$4T/C$
$V_{MAX}$	$3V_{DD}$	$5V_{DD}$
$I_{MAX}$	$3CV_{DD}/T$	$1.25CV_{DD}/T$

**Fig. 5.16**  $I$ - $V$ s and  $\eta$ - $V$ s in two modes of the pump of Fig. 5.15



$R_{PMP}$ . On the other hand, in mode 2, the upper two PC1 stages are connected with the lower two PC1 stages. Thus, the pump has a configuration of one array of four stages. The pump in mode 2 has  $R_{PMP}$  of  $4T/C$  and  $V_{MAX}$  of  $5V_{DD}$ .

Figure 5.16 compares  $I$ - $V$  curves between mode 1 and 2. If  $V_{PP}$  is set at  $V_{MAX}/2$ , which maximize the output power as shown in Eq. (2.11), both in mode 1 and 2, the ratio of  $I_{PP}$  of mode 1 to that of mode 2 is equal to  $3/1.25$ . When the output current is not required to be so high, one can simply disable the upper two stages instead of enabling them. The ratio is reduced to  $1.5/1.25$ , but it is still higher than 1. This means that this configuration is also possible when the requirement for the output current in mode 1 is not being smaller than the output current in mode 2. Figure 5.16 also compares  $\eta$ - $V$  curves between mode 1 and 2. Power efficiency can improve by reconfiguring the charge pump in terms of the number of arrays and stages depending on the operating point. The charge pump can also dynamically reconfigure its state as the output voltage increases. This approach can reduce the total rise time and average input power from an initial low voltage to a target high voltage (Tanzawa et al. 1994).

### 5.5 Noise and Ripple Reduction Design

The pump output current  $I_{OUT}$  has a large ripple as shown in Fig. 5.17, resulting in a large ripple in the output voltage  $V_{OUT}$  and in the supply current  $I_{DD}$ . This section discusses design techniques to reduce the ripple. One approach is adding a decoupling capacitor  $C_{DC}$  to the output terminal. When the ripple in  $V_{OUT}$  is required to be  $\Delta V_{PP}$ , the capacitance required for the decoupling capacitor should be

$$C_{DC} > I_{DD}T/\Delta V_{PP} \tag{5.3}$$

The decoupling capacitor can reduce the ripple in output voltage, however, doesn't reduce the ripple in  $I_{DD}$ .

Another method for reducing the ripple in  $V_{PP}$  in case of current load is adding a clamping transistor between the pump output and the load terminals as shown in Fig. 5.18b. Compared with Fig. 5.18a without a clamping transistor, the ripple



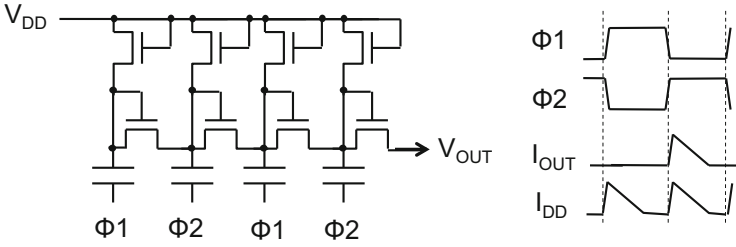


Fig. 5.17 Current profile along with pump operation

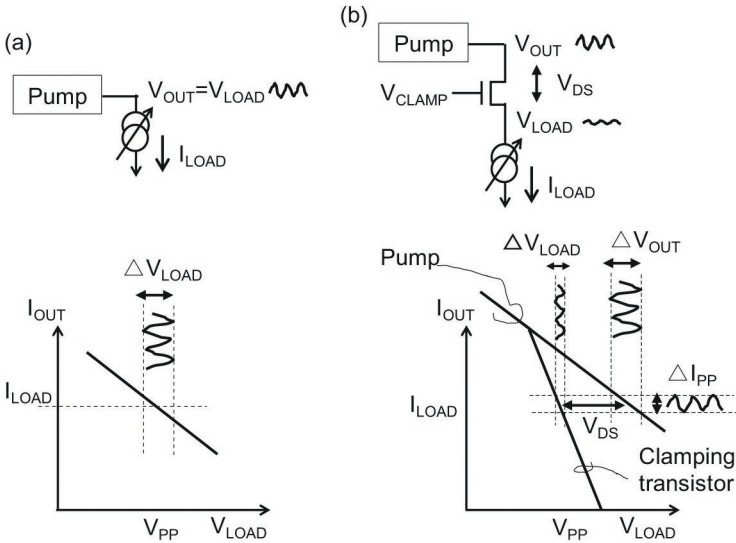


Fig. 5.18 Reduction method in the ripple in  $V_{LOAD}$

voltage at the load terminal can be reduced. As shown in the  $I-V$  graph of Fig. 5.18b, the voltage ripple  $\Delta V_{OUT}$  translates into the current ripple  $\Delta I_{PP}$ , and then it results in the voltage ripple  $\Delta V_{LOAD}$ . Because of the steep slope in  $I-V$  with the clamping transistor,  $\Delta V_{LOAD}$  can be reduced in comparison with  $\Delta V_{OUT}$ . However, to keep the operation point at  $(V_{PP}, I_{PP})$  unchanged, the pump output current needs to be increased to  $(V_{PP} + V_{DS}, I_{PP})$ , where  $V_{DS}$  is the drain to source voltage of the clamping transistor. This requires to increase the output current at  $V_{PP}$ , resulting in a larger pump size. This technique, however, is not effective to reduce the ripple in  $I_{DD}$ .

To reduce the ripple in the supply current, another design technique is needed. Figure 5.19 describes a noise reduction method. A single pump is divided into four arrays. Every array is driven by one of four phases. Thus, both peaks in  $I_{OUT}$  and  $I_{DD}$  can be reduced by a factor of more than 2. The ripple depends on the timing

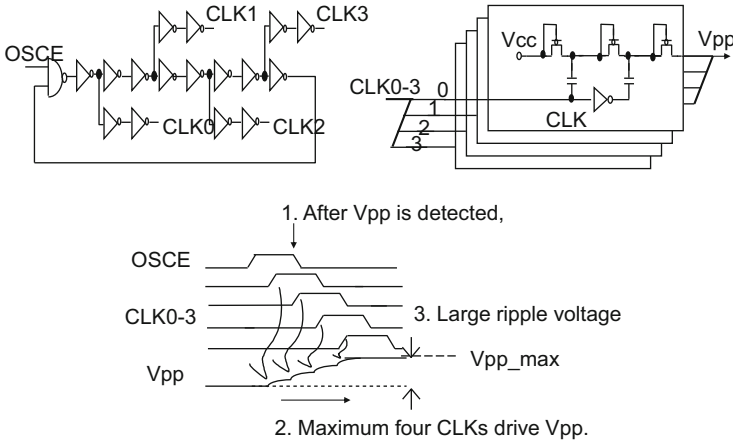


Fig. 5.19 Pump with low noise (Javanifard et al. 1994)

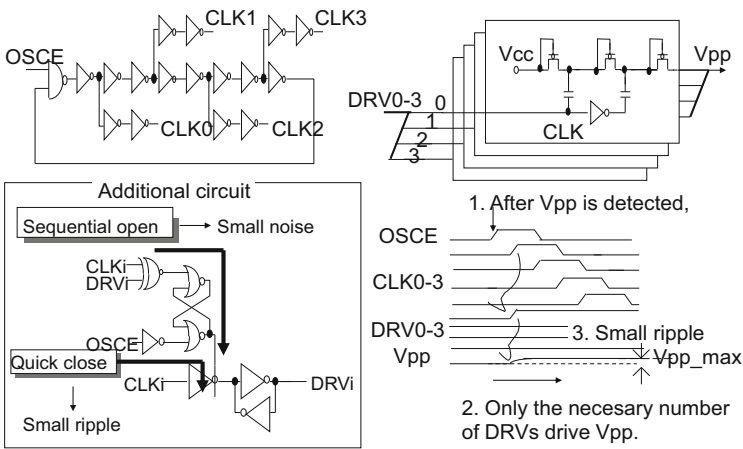


Fig. 5.20 Pump with low ripple (Tanzawa et al. 2002)

when the clock enabling signal  $OSCE$  goes low. Figure 5.19 also shows the worst case in terms of the ripple in  $V_{PP}$ . All the four arrays operate after the oscillator enabling signal  $OSCE$  can go low. In this case, the ripple in  $V_{PP}$  is not reduced in comparison with a single array pump.

Figure 5.20 adds a controlled buffer for the driving signals  $DRV$ . As soon as  $OSCE$  goes low,  $DRVs$  stop changing their logical state and their states are latched. Thus, the ripple in  $V_{PP}$  can be minimized. After  $OSCE$  goes high, transferring  $CLKs$  to  $DRVs$  starts again when the logical state of  $CLKs$  become identical to that of  $DRVs$ . Thus, no simultaneous operation occurs, resulting in averaged current profile in  $I_{OUT}$  and  $I_{DD}$  as well as a low ripple in  $V_{PP}$ .

## 5.6 Stand-by and Active Pump Design

Some applications may need a high-voltage generated on chip with a low stand-by current condition even just after the power supply is input. In addition, the pump output current needs to be sufficient high to supply a load in an active state. When both requirements for a low current consumption in a stand-by state and a high output current in active are made simultaneously, one may have to have two pumps as shown in Fig. 5.21.

When the leakage current at the output node is sufficiently small, the period when the stand-by pump is disabled would be quite long. During this period, all the internal capacitor nodes can be equalized to the output voltage due to the reverse subthreshold current via low- $V_T$  transfer transistors. If the next boosting operation starts with OSCE high under such a situation, only a few clocks may be enough with relatively large pump capacitors to increase the output voltage to a target voltage and the pump operation is disabled again. Assuming the pump needs  $M$  clock cycles to output the current for recovering a reduction in the output voltage of  $\Delta V_{PP}$ ,

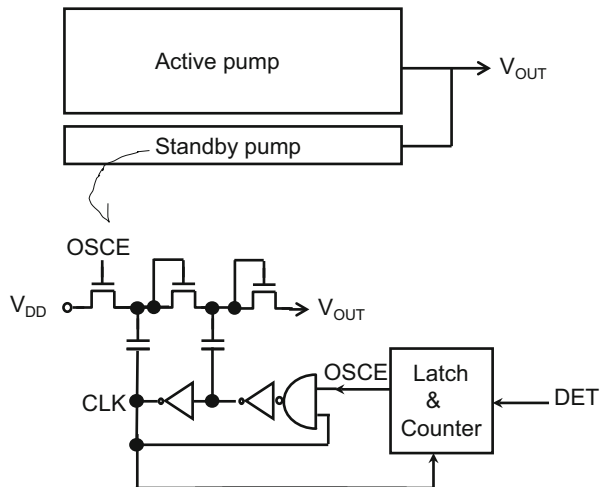
$$\Delta V_{PP} \leq MT_{IP}/C_{LOAD} = \alpha MC/C_{LOAD} \quad (5.4)$$

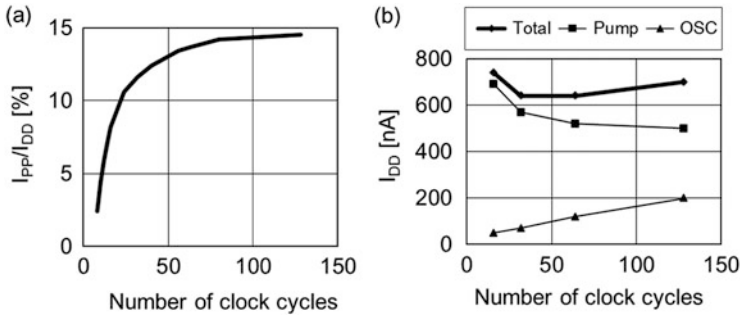
where  $\alpha$  is a proportional coefficient  $[(N+1)V_{DD} - V_{PP}]/N$ . One can simulate the amount of output charges per cycle using similar Eqs. (3.28)–(3.31). One difference is using Eq. (5.5) instead of Eq. (3.27) because all the internal capacitor nodes are equalized to  $V_{PP}$  due to the reverse current.

$$Q(1, j) = Q(1, j+1) = CV_{PP} \quad (5.5)$$

Figure 5.22a shows simulated results under the condition of  $N = 4$ ,  $V_{DD} = 1.5$  V, and  $V_{PP} = 4.5$  V. The graph suggests that one needs a number of clock cycles larger

**Fig. 5.21** Two pump arrays for stand-by and active states (e.g., Sato et al. 1985; Tanzawa et al. 2001)





**Fig. 5.22** Current efficiency  $I_{PP}/I_{DD}$  (a) and stand-by current  $I_{DD}$  (b) vs. number of clock cycles

than 50 to have sufficiently high power efficiency as much as that in a steady state in this example. Figure 5.22b shows the stand-by current as a function of the number of clocks cycles. The input current to the stand-by pump decreases as the number of clocks increases because the power efficiency is improved as shown in Fig. 5.22a whereas the input current to the oscillator increases because the duty ratio of the operation time to the wait time increases thereby the averaged input current increases. Thus, the total input current has a minimum point across the number of clock cycles. In this example,  $M$  of 50–100 should be selected. For given values for  $\Delta V_{PP}$ ,  $C_{LOAD}$ , and  $\alpha$ , Eq. (5.4) constrains the condition for the product  $MC$ . As a result, one can determine optimum values for  $M$  and  $C$ . The counter of Fig. 5.21 is then designed to work with the optimum  $M$ .

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### $V_T$ Canceling Design

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### ***Switching Diode Design***

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# Chapter 6

## Pump Control Circuits

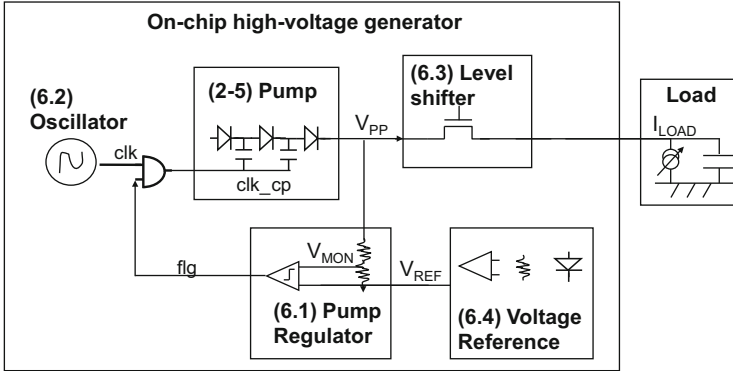
**Abstract** This chapter is devoted to individual circuit block, i.e., pump regulators, oscillators, level shifters, and voltage references, to realize on-chip high-voltage generator together with charge pumps.

Section 6.1 presents pump regulators. Some of the pump output voltages need to be varied to adjust them to the target voltages. This can be done with the voltage gain of the regulator or the reference voltage changed. The voltage divider which is a main component of the regulator has to have small voltage coefficient and fast transient response enough to make the controlled voltage linear to the trim and stable in time. A regulator for a negative voltage has a circuit configuration different from that for a positive voltage. State of the art is reviewed.

Section 6.2 deals with oscillators. Without an oscillator, the charge pump never works. In order to make the pump area small, process, voltage, and temperature variations in oscillator frequency need to be done as small as possible. There is the maximum frequency at which the output current is maximized. If the oscillator is designed to have the maximum frequency under the fastest conditions such as fast process corner, high supply voltage, and low temperature, the pump output current is minimum under the slowest conditions such as slow process, low supply voltage, and high temperature. It is important to design the oscillator with small variations for squeezing the pump area.

Section 6.3 reviews level shifters. The level shifter shifts the voltage for logic high or low of the input signal to a higher or lower voltage of the output signal. Four types of level shifters are discussed (1) high-level NMOS level shifter, (2) high-level CMOS level shifter, (3) high-voltage depletion NMOS+PMOS level shifter, and (4) low-level CMOS level shifter. The trade-offs between the first three high-voltage shifters are mentioned. The negative voltage can be switched with the low-level shifter. As the supply voltage lowers, operation margins of the level shifters decrease. As the supply voltage lowers, the switching speed becomes slower, eventually infinite, i.e., the level shifter does not work. Some design techniques to lower the minimum supply voltage at which the level shifters are functional are shown.

Section 6.4 provides voltage references. Variations in regulated high voltages increase by a factor of the voltage gain of the regulators from those in the reference voltages. Reduction in the variations in voltage references is a key to make the high voltages well controlled. Some innovated designs for low supply voltage operation are presented as well.



**Fig. 6.1** Block diagram of on-chip high-voltage generator

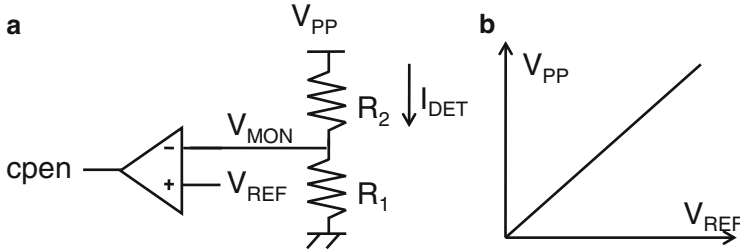
Figure 6.1 shows on-chip high-voltage generator system and each component circuit block discussed in each section of this chapter. The charge pump inputs the supply voltage ( $V_{DD}$ ) and the clock, which is generated by the oscillator, and outputs a voltage ( $V_{PP}$ ) higher than the supply voltage or a negative voltage. The pump regulator enables the charge pump when the absolute value of the output voltage of the charge pump is lower than the target voltage on the basis of the reference voltage  $V_{REF}$ , or disables it otherwise. The output voltage of the pump is determined by the reference voltage and the voltage gain of the regulator. To vary the pump output voltage, either reference voltage or voltage gain of the regulator is varied. The generated high or negative voltage is transferred to a load through high- or low-level shifters. The level shifters are controlled by the input supply voltage. The load is capacitive, resistive, or both.

## 6.1 Regulator

This section presents pump regulators. Some of the pump output voltages need to be varied to adjust them to the target voltages. This can be done with the voltage gain of the regulator or the reference voltage changed. The voltage divider that is a main component of the regulator has to have small voltage coefficient and fast transient response enough to make the controlled voltage linear to the trim and stable in time. A regulator for a negative voltage has a circuit configuration different from that for a positive voltage. State of the art is reviewed.

A pump regulator shown in Fig. 6.2a detecting the output voltage of charge pump contains a voltage divider and a comparator inputting a reference voltage  $V_{REF}$ . The output signal  $cp_{en}$  is a logic signal indicating whether the charge pump needs to operate or not. Design parameters  $R_1$ ,  $R_2$ , and  $V_{REF}$  determine the target  $V_{PP}$ .





**Fig. 6.2** Pump regulator for a positive high voltage  $V_{PP}$  (a) and ideal relation of  $V_{PP}$  to  $V_{REF}$  (b) (Oto et al. 1983)

$$V_{PP} = G_V V_{REF} \quad (6.1)$$

$$G_V = 1 + \frac{R_2}{R_1} \quad (6.2)$$

where  $G_V$  is the voltage gain. Practically, the output voltage can vary due to variations in each design parameter and the offset voltage  $V_{OS}$  of the comparator.

$$\begin{aligned} \delta V_{PP} &= \delta G_V V_{REF} + G_V (\delta V_{REF} + V_{OS}) \\ &= \left( \frac{\delta R_2}{R_1} - \frac{R_2 \delta R_1}{R_1^2} \right) V_{REF} + G_V (\delta V_{REF} + V_{OS}) \end{aligned} \quad (6.3)$$

It is assumed that each variation component is independent of one another. The standard deviation can be given by

$$\sigma V_{PP}^2 = \left( \frac{R_2}{R_1} V_{REF} \right)^2 \left[ \left( \frac{\sigma R_1}{R_1} \right)^2 + \left( \frac{\sigma R_2}{R_2} \right)^2 \right] + G_V^2 (\sigma V_{REF}^2 + \sigma V_{OS}^2) \quad (6.4)$$

To vary  $V_{PP}$  with trimming, there are three methods. The first one is such that  $V_{REF}$  is varied whereas  $G_V$  is constant as shown in Fig. 6.2b. If the input range of the comparator is limited, the operation window would be limited in some portions in  $V_{REF}$ . The second method is such that  $G_V$  is varied whereas  $V_{REF}$  is constant. Figure 6.3a shows a trim-able resistor  $R_1$  as shown in Eq. (6.5) using four signal  $S_i$  ( $1 \leq i \leq 4$ ) to vary  $G_V$  through Eq. (6.2).

$$R_1 = \sum_{i=1}^4 \bar{S}_i r_i \quad (6.5)$$

To reduce the impact of the transistor resistance on  $R_1$ , the transistors need to be large enough or the voltage for logic high of the signal needs to be high enough.

The third one is adding a modulation part to the resistor divider, as shown in Fig. 6.4. Suppose a current source with  $V_{MOD}/R_3$  is connected at the  $V_{MON}$  node,

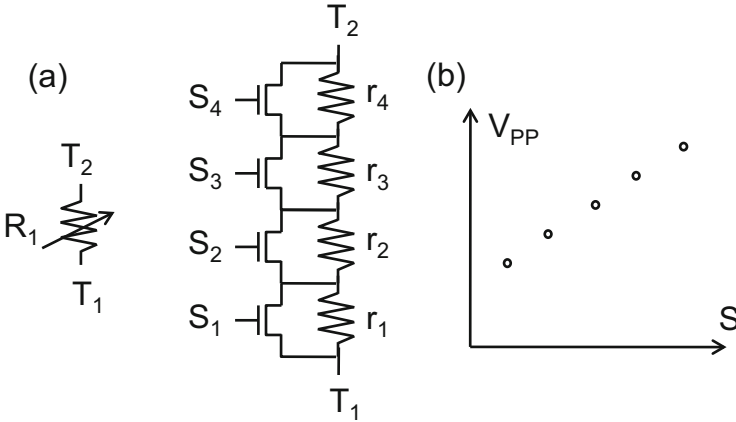


Fig. 6.3 Trim-able resistor (a) and ideal relation of  $V_{PP}$  to  $S$  (b) (Suh et al. 1995)

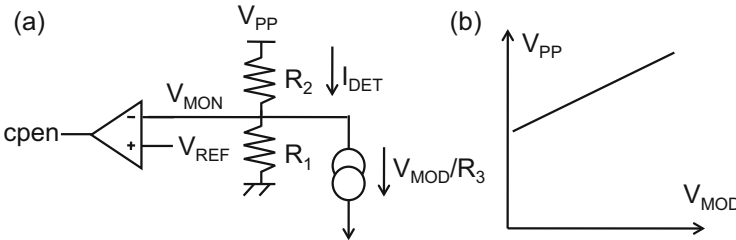


Fig. 6.4 Pump regulator with a voltage modulation path (Tanzawa and Harrington 2010)

$V_{PP}$  is given by two parts of  $V_{REF}$  and  $V_{MOD}$ , as shown by Eq. (6.6). This means that  $V_{PP}$  varies with  $V_{MOD}$  varied while  $V_{REF}$  unchanged. If one can add a modulation component into  $V_{MOD}$ ,  $V_{PP}$  has the characteristic as shown in Fig. 6.4b.

$$V_{PP} = \left(1 + \frac{R_2}{R_1}\right)V_{REF} + \frac{R_2}{R_3}V_{MOD} \quad (6.6)$$

Figure 6.5a shows the current components of a pump and a regulator. As the output voltage of the pump increases, the pump output current  $I_{PUMP}$  decreases whereas the current to the regulator  $I_{REG}$  increases as shown in Fig. 6.5b. The effective current to charge the load,  $I_{LOAD}$ , therefore decreases as  $V_{PP}$ . Thus, the detector current needs to be made low enough not to affect  $I_{LOAD}$  much.

Figure 6.6a shows an n-diffusion resistor fabricated on the p-type substrate. When a terminal of the resistor is applied by a high voltage  $V_{PP}$ , depletion region width increases, resulting in higher resistivity with a thinner conduction layer as shown in Fig. 6.6b. Similar behavior is seen when both terminals are applied by high voltages as shown in Fig. 6.6c. Figure 6.6d indicates that the voltage

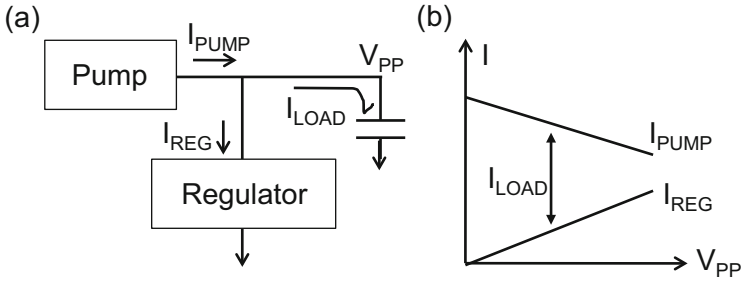


Fig. 6.5 Current components of a pump and a regulator

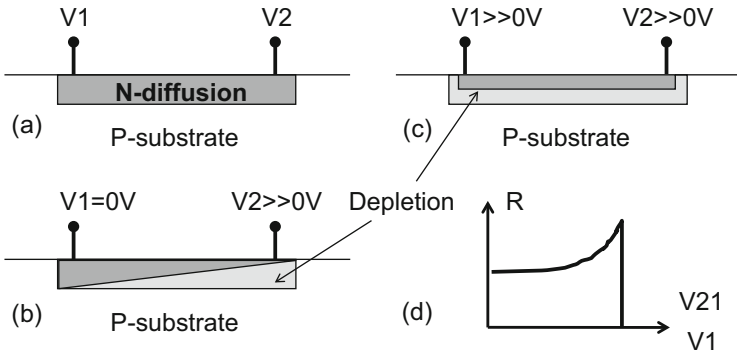
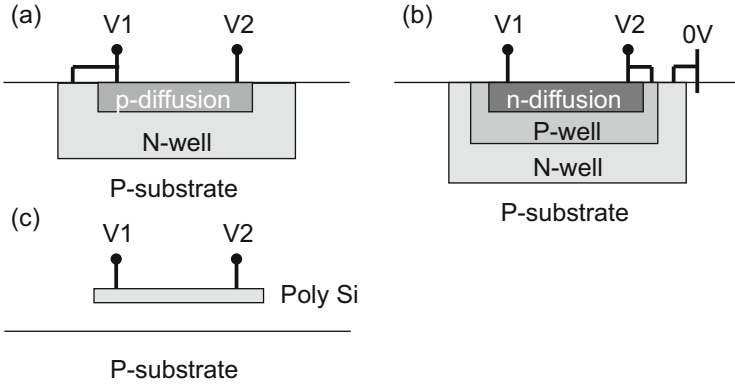


Fig. 6.6 n-diffusion resistor (a), two bias conditions (b, c), resistance vs. bias relation (d)

coefficient of the resistance is small at a low voltage applied and increases as the applied voltages. Over a junction breakdown voltage, it is no longer available as a resistor.

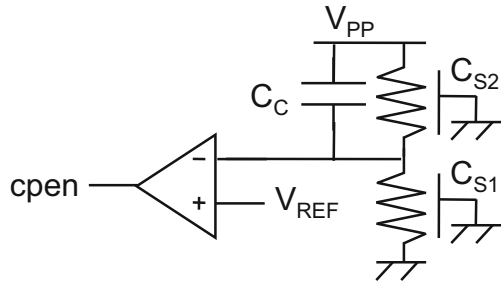
Figure 6.7 illustrates three other types of resistors. When N-well is divided into multiple pieces and every N-well has p-diffusion layers, voltage differences between the p-diffusion layers and N-wells can be reduced to mitigate the nonlinearity of the resistance on the voltages applied in comparison with a single n-diffusion layer on the substrate. To allow a negative voltage to be detected by a pump regulator, n-diffusion layers fabricated on P-well isolated by N-well as shown in Fig. 6.7b. Poly-silicon resistor has benefits of small voltage dependency on the resistivity and of availability of both polarities.

The current used for the regulator is a part of the load current of the pump. In order to reduce the current  $I_{DET}$ , the resistor  $R_1$  is likely high impedance. Assuming  $V_{REF}$  is 1 V and  $I_{DET}$  is 10  $\mu$ A,  $R_1$  is required to be 100 k $\Omega$ . Furthermore, when  $V_{PP}$  is 20 V,  $R_2$  is required to be 1.9 M $\Omega$ . Parasitic capacitance of the resistor depends on the material used. In case of diffusion resistor, its parasitic capacitance per  $\Omega$  is relatively large. Assuming 1 pF/M $\Omega$ , the time constant of  $R_2$  is about 4  $\mu$ s. When the rise time of  $V_{PP}$  is shorter than or compatible to 4  $\mu$ s, the output of the pump can



**Fig. 6.7** Other types of resistors: p-diffusion resistor (a), n-diffusion resistor in a twin well (b), poly silicon resistor (c)

**Fig. 6.8** Pump regulator with a shunt capacitor  $C_C$



have large overshoot. To reduce the propagation delay from  $V_{PP}$  to  $V_{MON}$ , a shunt capacitor  $C_C$  is used as shown in Fig. 6.8. DC operating point is determined by the divider ratio whereas AC signal travels via  $C_C$ .

Figure 6.9 shows a negative voltage detector. The circuit requires a well-controlled regulated voltage  $V_{PP}$  to detect the negative voltage at  $V_{BB}$ , because there is an additional term in Eq. (6.8) compared with Eq. (6.1).

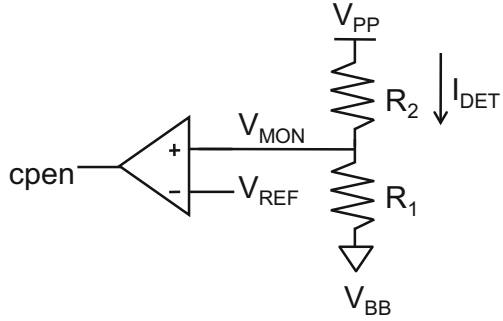
$$\frac{V_{PP} - V_{REF}}{R_2} = \frac{V_{REF} - V_{BB}}{R_1} \tag{6.7}$$

$$V_{BB} = \left(1 + \frac{R_1}{R_2}\right)V_{REF} - \frac{R_1}{R_2}V_{PP} \tag{6.8}$$

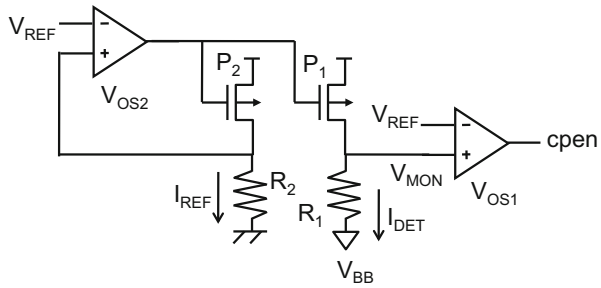
Sensitivity of  $V_{BB}$  on each parameter is calculated by

$$\begin{aligned} \delta V_{BB} = & \left(1 + \frac{R_1}{R_2}\right)(\delta V_{REF} + V_{OS}) - \frac{R_1}{R_2}\delta V_{PP} \\ & + \left(\frac{\delta R_1}{R_2} - \frac{R_1\delta R_2}{R_2^2}\right)(V_{REF} - V_{PP}) \end{aligned} \tag{6.9}$$

**Fig. 6.9** Pump regulator for a negative voltage  $V_{BB}$



**Fig. 6.10** Pump regulator for a negative voltage  $V_{BB}$  with a reduced gain against variations (Mihara et al. 1999)



The standard deviation can be given by

$$\sigma V_{BB}^2 = \left(\frac{R_1}{R_2}\right)^2 (V_{REF} - V_{PP})^2 \left| \left(\frac{\sigma R_1}{R_1}\right)^2 \left(\frac{\sigma R_2}{R_2}\right)^2 \right| + \left(1 + \frac{R_1}{R_2}\right)^2 (\sigma V_{REF}^2 + \sigma V_{OS2}^2) + \left(\frac{R_1}{R_2} \sigma V_{PP}^2\right)^2 \tag{6.10}$$

When  $V_{BB}$  is shifted by  $\Delta V_{BB}$ ,  $V_{MON}$  is shifted by  $\Delta V_{MON}$  as follows.

$$\Delta V_{MON} = \Delta V_{BB} / \left(1 + \frac{R_1}{R_2}\right) \tag{6.11}$$

The amplitude of the input signal to the comparator is scaled from that of  $V_{BB}$  by a factor of  $1 + R_1/R_2$ . The detector shown in Fig. 6.10 increases the input signal amplitude.

In an ideal case with no mismatch in the parameters, the following equations hold.

$$V_{MON} - V_{BB} = I_{DET} R_1 \tag{6.12}$$

$$I_{\text{REF}} = \frac{V_{\text{REF}}}{R_2} = I_{\text{DET}} \quad (6.13)$$

where it is assumed that two PMOSFETs  $P_1$  and  $P_2$  are identical in size. Using the steady state condition of  $V_{\text{MON}} = V_{\text{REF}}$  for Eqs. (6.12) and (6.13),  $V_{\text{BB}}$  can be given by

$$V_{\text{BB}} = G_V V_{\text{REF}} \quad (6.14a)$$

$$G_V = 1 + \frac{R_1}{R_2} \quad (4.14b)$$

Because  $I_{\text{DET}}$  has no  $V_{\text{BB}}$  dependence in Eq. (6.12), the sensitivity of  $V_{\text{MON}}$  on  $V_{\text{BB}}$  is given by

$$\Delta V_{\text{MON}} = \Delta V_{\text{BB}} \quad (6.15)$$

When one uses long channel  $I - V$  equations, the following relations hold.

$$I_{\text{REF}} = K(V_{\text{GS}} - V_{\text{T2}})^2 \quad (6.16a)$$

$$I_{\text{DET}} = K(V_{\text{GS}} - V_{\text{T1}})^2 \quad (6.16b)$$

where  $V_{\text{GS}}$  is the gate-to-source voltage of  $P_1$  and  $P_2$  and  $V_{\text{T1}}$  and  $V_{\text{T2}}$  are the threshold voltages of  $P_1$  and  $P_2$ . When the opamps have input offset voltages of  $V_{\text{OS1}}$  and  $V_{\text{OS2}}$ , as shown in Fig. 6.10, and the device parameters are independently varied,  $I_{\text{REF}}$  varies by Eq. (6.17a).

$$\delta I_{\text{REF}} = \frac{\delta V_{\text{REF}} + V_{\text{OS2}}}{R_2} - \frac{V_{\text{REF}} \delta R_2}{R_2^2} \quad (6.17a)$$

From Eq. (6.16a, 6.16b) and the assumption that  $V_{\text{T2}}$  is mismatched from  $V_{\text{T1}}$  by  $\delta V_{\text{T}}$ ,

$$\delta I_{\text{DET}} = \delta I_{\text{REF}} - 2\sqrt{K I_{\text{REF}}} \delta V_{\text{T}} \quad (6.17b)$$

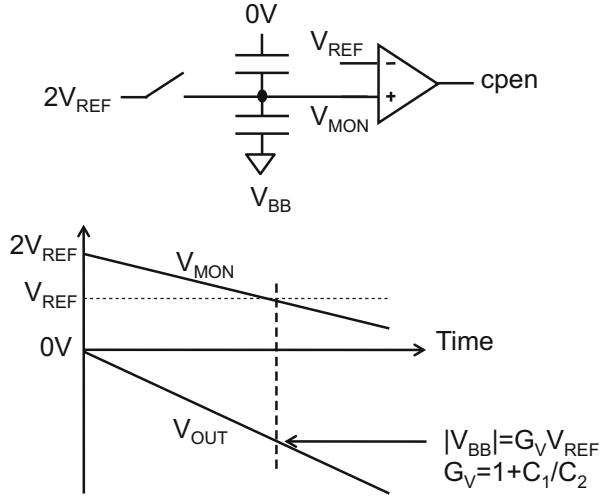
In addition, from Eq. (6.12),

$$V_{\text{OS1}} - \delta V_{\text{BB}} = \delta I_{\text{DET}} R_1 + I_{\text{DET}} \delta R_1 \quad (6.17c)$$

Using Eqs. (6.17a–6.17c) and (6.13), overall variation is given by

$$\begin{aligned} \delta V_{\text{BB}} &= V_{\text{OS1}} - I_{\text{DET}} \delta R_1 - \delta I_{\text{DET}} R_1 \\ &= V_{\text{OS1}} - \frac{V_{\text{REF}}}{R_2} \delta R_1 - \left( \frac{\delta V_{\text{REF}} + V_{\text{OS2}}}{R_2} - \frac{V_{\text{REF}} \delta R_2}{R_2^2} - 2\sqrt{K I_{\text{REF}}} \delta V_{\text{T}} \right) R_1 \end{aligned} \quad (6.18)$$

**Fig. 6.11** Capacitor divider for regulating a negative voltage (Venkatesh et al. 1996)



Assuming each variation component is independent, the standard deviation can be calculated by

$$\begin{aligned} \sigma V_{BB}^2 = & \sigma V_{OS1}^2 + \left(\frac{R_1}{R_2}\right)^2 (\sigma V_{REF}^2 + \sigma V_{OS2}^2) \\ & + \left(\frac{R_1}{R_2} V_{REF}\right)^2 \left[ \left(\frac{\delta R_1}{R_1}\right)^2 + \left(\frac{\delta R_2}{R_2}\right)^2 \right] + 4KI_{REF} R_1^2 \sigma V_T^2 \end{aligned} \quad (6.19)$$

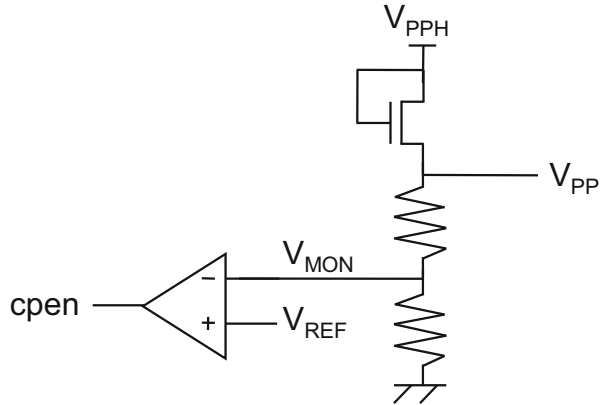
Another interesting design technique is using a capacitor divider as shown in Fig. 6.11. It does not require a resistor, which can be applied by a negative voltage. Initially,  $V_{MON}$  is precharged to  $2V_{REF}$  and then  $V_{MON}$  is made floating. Accordingly as  $V_{OUT}$  goes low,  $V_{MON}$  is also pulled down. Once  $V_{MON}$  reaches  $V_{REF}$ , the detection signal  $cpen$  goes  $L$ . As far as the operation time of the negative voltage generation is short enough so that the leakage current at the floating node is negligibly small, the regulator should function well.

As will be described in Sect. 6.3 for high-voltage switching circuits, a regulator shown in Fig. 6.12 has two output terminals whose voltages are  $V_{PPH}$  and  $V_{PP}$ . A pump is connected with  $V_{PPH}$ . When  $V_{PPH}$  is supplied to the gate of a switching pass NMOS transistor, it can transfer  $V_{PP}$  without any voltage loss as shown in Fig. 6.36.

$$V_{PPH} = V_{PP} + V_T \quad (6.20a)$$

$$V_{PP} = G_V V_{REF} \quad (6.20b)$$

**Fig. 6.12** Regulator with one diode to generate a switching voltage  $V_{PPH}$



## 6.2 Oscillator

This section deals with oscillators. Without an oscillator, the charge pump never works. In order to make the pump area small, process, voltage, and temperature variations in oscillator frequency need to be done as small as possible. There is the maximum frequency at which the output current is maximized. If the oscillator is designed to have the maximum frequency under the fastest conditions such as fast process corner, high supply voltage, and low temperature, the pump output current is minimum under the slowest conditions such as slow process, low supply voltage, and high temperature. It is important to design the oscillator with small variations for squeezing the pump area.

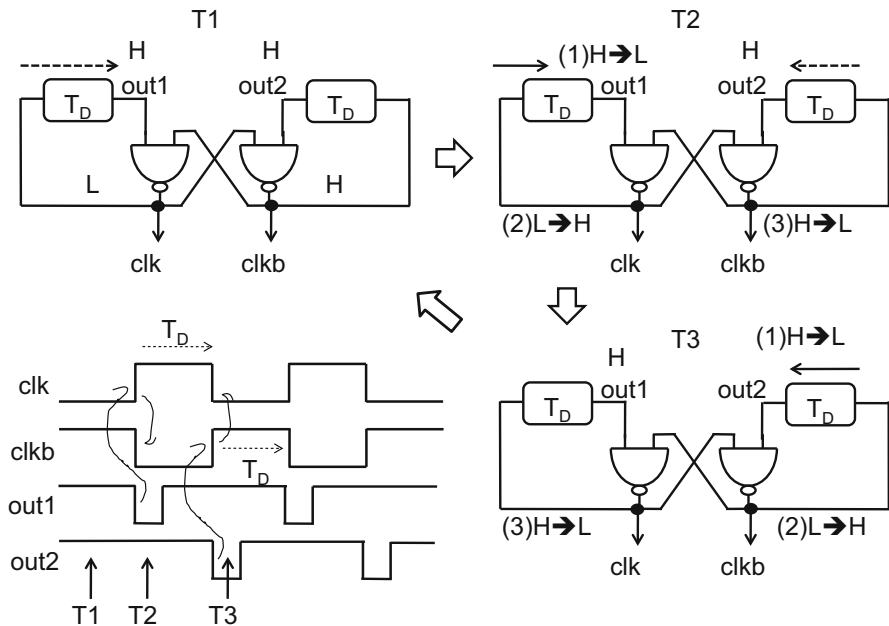
The primal target for oscillators is making  $I_{PP}$  insensitive to process, voltage, and temperature (PVT) variations. What is the parameter which is not varied much?  $C$  is very accurately fabricated within a few percent errors.  $N$  is solid.  $V_T$  can be a weak function of temperature and process variation.  $V_{PP}$  is the solid target.  $V_{DD}$  may be varied a lot without an on-chip voltage regulator or quite solid with it. Thus, there are two cases, use of a linear regulator for  $V_{DD}$  or not. In the former case, all the parameters should be stable to realize the clock frequency or period insensitive to process, voltage, and temperature. To stabilize  $V_{DD}$ , you may need large decoupling capacitors. In case of no  $V_{DD}$  regulator,  $T$  would need to be proportional to the factor  $(N+1)(V_{DD}-V_T)-V_{PP}$  to make  $I_{PP}$  insensitive to PVT variation (Table 6.1).

Figure 6.13 describes a bi-stable oscillator and its operation. It is known that symmetrical bi-stable oscillator generates two phase clock with 50 % high low duties. The half period time is determined by the delay element  $T_D$ . One can start with T1 where out1 and 2 are high and clk and clkb are L and H, respectively. clk L propagates to out1 after  $T_D$ . That flips clk to H, in turn flips clkb to L. clkb L propagates to out2 after  $T_D$ . That flips clkb to H, which flips clk to L. Thus, the oscillator has two states alternately and half period is determined by  $T_D$ . This kind of oscillator is known as bi-stable oscillator. The delay circuit shown in Fig. 6.13



**Table 6.1** Requirement for pump oscillator

Oscillator type	Type 1: use of a linear regulator for $V_{DD}$	Type 2: no use of $V_{DD}$ regulation
Features	$T$ should be insensitive to PVT	$T$ needs to be proportional to $(N + 1)(V_{DD} - V_T) - V_{PP}$
	Decoupling capacitors for $V_{DD}$ regulated is needed	$T$ should be insensitive to PT

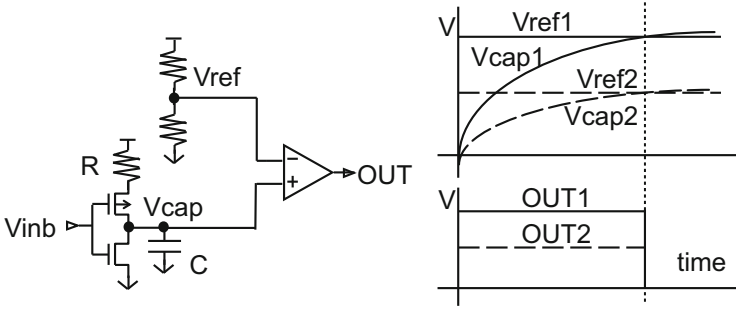


**Fig. 6.13** Bi-stable oscillator generating two phase clock

can be used for the delay element as described by  $T_D$  of Fig. 4.14. The clock period  $T_C$  is simply given by  $2T_D$  as far as the delay of logic gates is negligibly small compared with  $T_D$ . Two delay elements alternately work to have a stable period  $T_C$  given by  $2T_D$ .

Oscillators are composed of multiple delay elements. To have stable oscillators against PVT variations, stable delay elements are essential. When a resistor more stable against PVT variations than channel resistance of a transistor is available, one should use it. Figure 6.14 describes a delay circuit whose delay time is basically determined by the multiple of resistance  $R$  and capacitance  $C$ .

The circuit operates as follows. When the input signal  $V_{inb}$  goes low, the current flows from the supply voltage  $V_{DD}$  to the capacitor node.



**Fig. 6.14** Delay circuit with a delay time proportional to CR (Watanabe et al. 1989)

$$C \frac{dV_{CAP}(t)}{dt} = \frac{V_{DD} - V_{CAP}(t)}{R} \tag{6.21}$$

The reference voltage is proportional to  $V_{DD}$ , where  $\alpha$  is a division ratio.

$$V_{REF} = \alpha V_{DD} \tag{6.22}$$

Under the initial condition where  $V_{CAP}(0)$  is 0 V,  $V_{CAP}(t)$  is solved to be

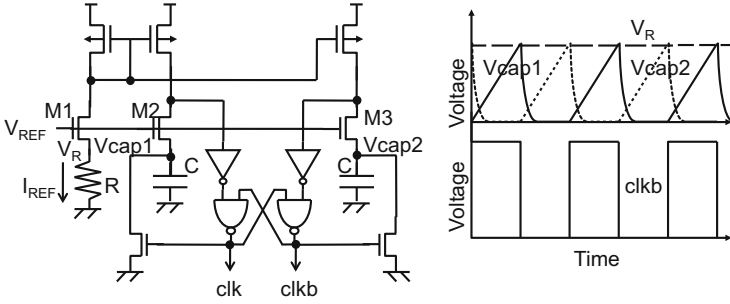
$$V_{CAP}(t) = V_{DD} \left( 1 - e^{-\frac{t}{CR}} \right) \tag{6.23}$$

The output is flipped when  $V_{CAP}$  reaches  $V_{REF}$ . The delay time  $T_D$  is then given by

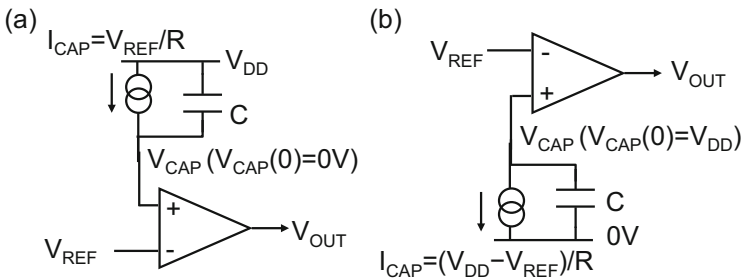
$$T_D = -CR \ln(1 - \alpha) \tag{6.24}$$

Because this equation does not include  $V_{DD}$ ,  $T_D$  is theoretically independent of variation in  $V_{DD}$ . “1” and “2” added to the labels of the waveform of Fig. 6.14 indicate different  $V_{DD}$ . Suppose  $V_{REF1}$  is twice as large as  $V_{REF2}$  due to the variation in  $V_{DD}$ .  $V_{CAP1}$  goes high twice faster than  $V_{CAP2}$  does, resulting in the same delay in OUT1 and OUT2. Nominally the variations in  $R$  against process and temperature are smaller than those in the channel resistance of transistor  $R_{CH}$ . Therefore, overall variation can be small with  $R$  than with  $R_{CH}$ .

Figure 6.15 shows another oscillator with the period that is determined by  $RC$ , where  $V_R$  is the voltage at the upper terminal of the resistor and  $I_{REF}$  is the reference current flowing the resistor and PMOSFETs connected with the clamp NMOSFET M2, 3. The capacitor voltages  $V_{CAP1,2}$  increase linearly to time with  $I_{REF}$ . After the source voltages of M2,3 reach  $V_R$ , the impedance of M2,3 rapidly increases, resulting in rapid increase in the drain voltages of M2,3. The delay time from the time when  $V_{CAP1}$  starts going up to the time when  $V_{CAP1}$  reaches  $V_R$  is given by



**Fig. 6.15** A bi-stable oscillator (Cernea et al. 1989)



**Fig. 6.16** Concept of a delay circuit (Tanzawa and Tanaka 1995)

$CV_R/I_{REF}$ . Even though  $V_R$  and  $I_{REF}$  vary according to the threshold voltage of M1, their ratio is constant as  $R$  as given below.

$$I_{REF} = V_R/R = K(V_{REF} - V_R - V_T)^2 \tag{6.25}$$

$$T_C/2 = CV_R/I_{REF} = RC \tag{6.26}$$

Figure 6.16 illustrates the concept of another delay circuit, which has the delay time with small PVT variations. In Fig. 6.16a, the initial voltage at the capacitor node is set to 0 V. The charging current  $I_{CAP}$  is made to be  $V_{REF}/R$ , where  $V_{REF}$  is the reference voltage for the comparator. The delay time when  $V_{CAP}$  reaches  $V_{REF}$  is given by

$$I_{CAP} = V_{REF}/R \tag{6.27a}$$

$$T_D = CV_{REF}/I_{CAP} = RC \tag{6.28a}$$

In Fig. 6.16b, the initial voltage at the capacitor node is set to  $V_{DD}$ . The discharging current  $I_{CAP}$  is made to be  $(V_{DD} - V_{REF})/R$ . The delay time when  $V_{CAP}$  reaches  $V_{REF}$  is given by

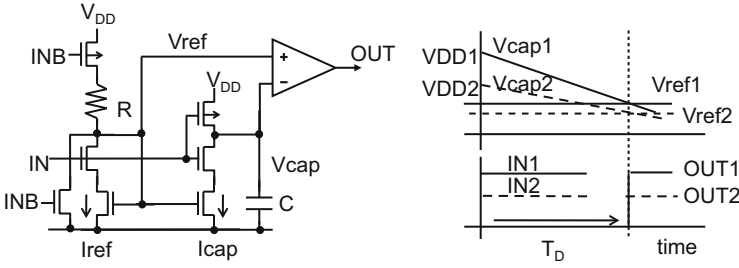


Fig. 6.17 A delay circuit (Tanzawa and Tanaka 1995)

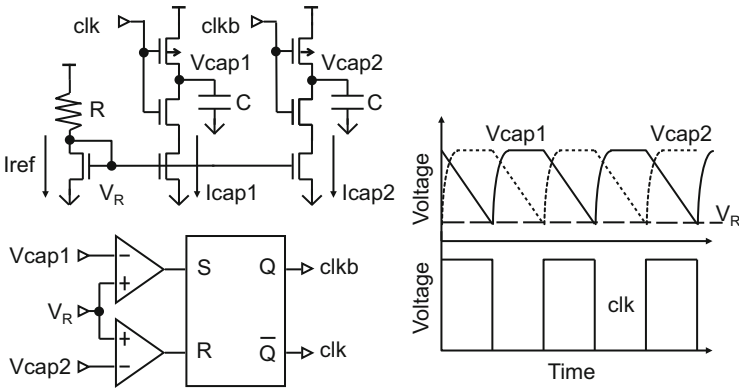


Fig. 6.18 Stable oscillator (type 1) using a delay element described in Fig. 6.17 (Tanzawa and Tanaka 1995)

$$I_{CAP} = (V_{DD} - V_{REF})/R \tag{6.27b}$$

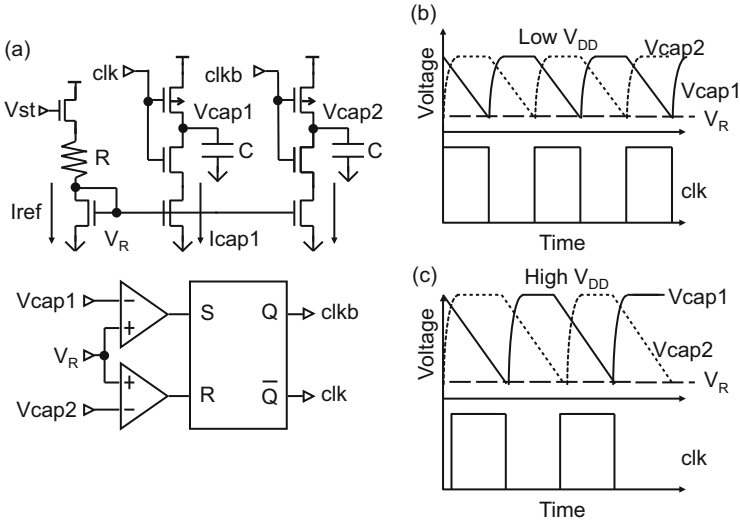
$$T_D = C(V_{DD} - V_{REF})/I_{CAP} = RC \tag{6.28b}$$

Thus, both circuits can have the same delay time with small PVT variations.

Figure 6.17 shows a circuit realizing the concept of Fig. 6.16b. As shown by Eqs. (6.27b) and (6.28b), the delay time is ideally independent of  $V_{DD}$  and  $V_T$  of transistors, resulting in small PVT variations. The key point here is that the voltage swing at the capacitor node  $V_{CAP}$  is proportional to the reference current  $I_{REF}$ .  $V_{DD}$  and  $V_T$  are not included in the ratio of the voltage amplitude of  $V_{CAP}$  and  $I_{REF}$ . Figure 6.18 has two sets of the delay elements of Fig. 6.17. The clock period is given by  $2T_D$ .

To change the type 1 oscillator into type 2,  $I_{REF}$  is made to have less  $V_{DD}$  dependency unlike the type 1, as shown in Fig. 6.19. Because the capacitor voltage has amplitude of  $V_{DD} - V_R$ , the clock cycle is given by

$$T_C/2 = C(V_{DD} - V_R)/I_{REF} \tag{6.29}$$



**Fig. 6.19** Stable oscillator (type 2) using a delay element described in Fig. 6.17 (Tanzawa and Tanaka 1995)

Equation (6.29) indicates the clock period increases as  $V_{DD}$ . To visualize this fact, one can compare a low  $V_{DD}$  case shown in Fig. 6.19b with a high  $V_{DD}$  case shown in Fig. 6.19c. The slopes in  $V_{CAP1,2}$  during the discharging period are same. Thus, as the amplitude increases with  $V_{DD}$ , the delay time also increases.

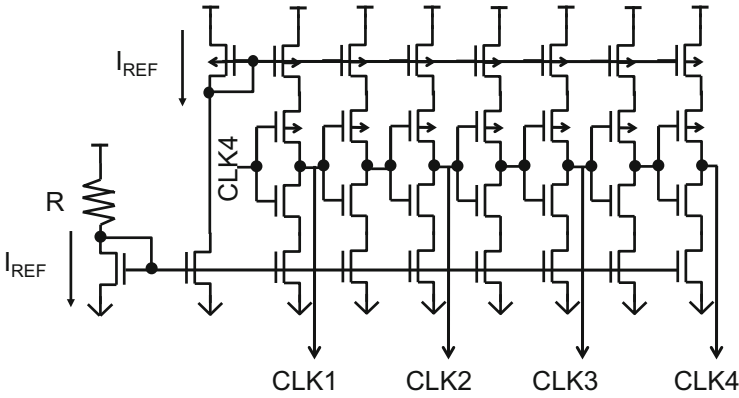
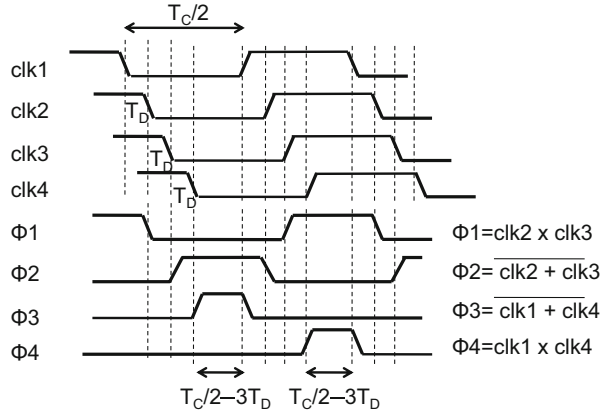
Four nonoverlapping phases  $\Phi 1-4$  are provided by logical addition or multiplication of  $clk1-4$ , each is the clock delayed by a same amount  $T_D$ , as shown in Fig. 6.20. It is noted that  $T_D$  also needs to be stable against PVT variations, because the effective pulse width to transfer the charges from one capacitor to the next one in the charge pump is given by  $T_C/2 - 3T_D$ .

Figure 6.21 illustrates a clock generator to output multiphase clocks. In Sect. 5.5, it was discussed that multiple arrays operating with multiple shifted phases could reduce noise in pump current. The ring oscillator does this. Current sources are connected to both PMOS and NMOS sides to control the operating currents proportional to  $I_{REF}$ . Thus, when  $I_{REF}$  is proportional to  $V_{DD} - V_T$ , the clock cycle time is insensitive to PVT variations. On the other hand, when  $I_{REF}$  is independent of  $V_{DD}$ , the cycle time could be proportional to  $V_{DD}$  but insensitive to PT variation.

### 6.3 Level Shifter

This section reviews level shifters. The level shifter shifts the voltage for logic high or low of the input signal to a higher or lower voltage of the output signal. Four types of level shifters are discussed (1) high-level NMOS level shifter, (2) high-

**Fig. 6.20** Waveform of four nonoverlapping phases  $\Phi 1-4$



**Fig. 6.21** Multiphase clock generator for peak noise reduction

level CMOS level shifter, (3) high-voltage depletion NMOS + PMOS level shifter, and (4) low-level CMOS level shifter. The trade-offs between the first three high-voltage shifters are mentioned. The negative voltage can be switched with the low-level shifter. As the supply voltage lowers, operation margins of the level shifters decrease. As the supply voltage lowers, the switching speed becomes slower, eventually infinite, i.e., the level shifter does not work. Some design techniques to lower the minimum supply voltage at which the level shifters are functional are shown.

### 6.3.1 NMOS Level Shifter

Section 6.3 starts with an NMOS high-level shifter shown in Fig. 6.22. Early days electrically erasable programmable ROM had only NMOS transistor for managing

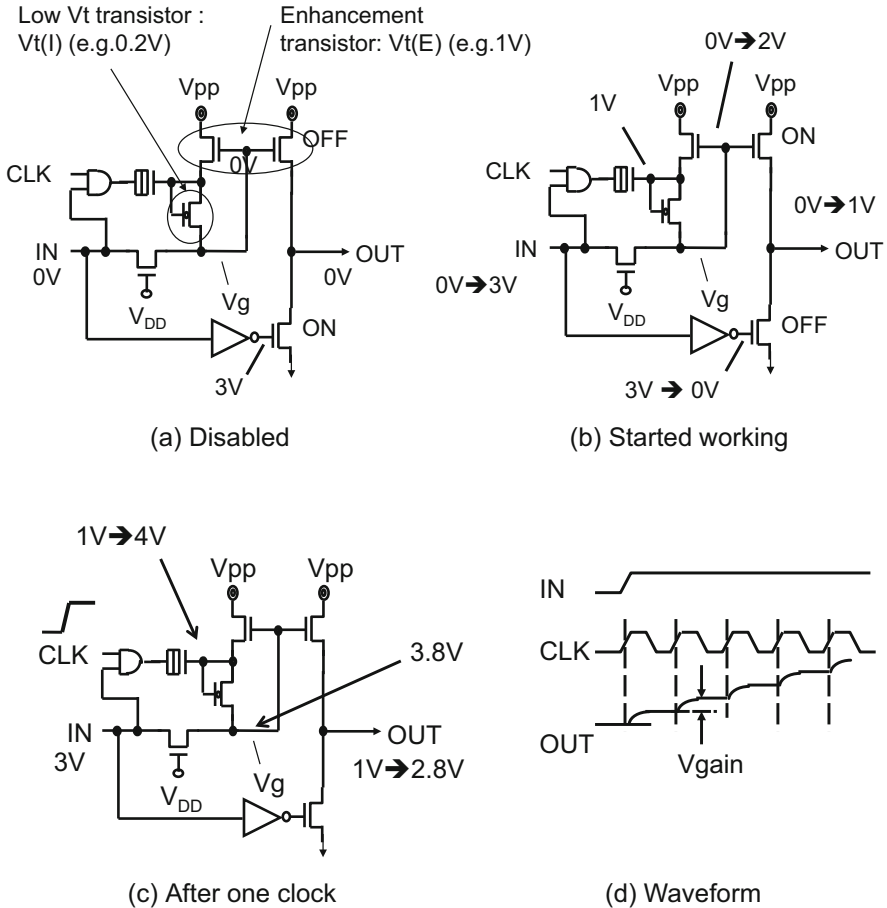


Fig. 6.22 NMOS high-level shifter (Donaldson et al. 1983; Dham et al. 1983)

high voltages. To transfer a high voltage through NMOSFET only without any voltage drop, an overdrive voltage needs to be generated locally. To fully cut off the transfer gate when it is disabled, an enhancement transistor with a high threshold voltage  $V_{TE}$  is used. To operate the local booster at a low supply voltage, a low- $V_t$  transistor is used. Such devices are fabricated without implanting Boron. When the input voltage is 0 V, the grounding NMOS turns on and the high-side NMOSFETs turn off with the gate grounded. When the circuit starts working, an input voltage of 3 V is transferred partially, that is, 2 V to the gates of the high-side NMOSFETs. Thus, the output voltage is 1 V. Then, the clock goes to 3 V, generating a local boosted voltage of 4 V. One diode drop of 3.8 V appears at the gate, resulting in an increase in the output voltage from 1 to 2.8 V.

Unlike CMOS switches with large parasitic capacitance of N-well for PMOSFETs, this NMOS high-level shifter has small gate-, junction-, and wiring-

capacitance, resulting in low power consumption. However, this switch also has a disadvantage in that the minimum operating supply voltage  $V_{DD}$  is mainly limited by the threshold voltage of an enhancement transistor, which prevents the leakage current from flowing in the  $V_{PP}$  switch in an inactive state. A diode-connected intrinsic transistor without channel implantation is used to improve the positive-feedback efficiency of the booster when selected for operation. The minimum operating  $V_{DD}$  is extracted.

Switching operation starts with the input signal  $IN$  high. After that the input clock oscillates to raise the output voltage. The source voltage of the enhancement transistor is lower by the threshold voltage  $V_{IE}$  than the gate voltage  $V_G$  with the clock  $clk$  high. After that, the  $clk$  turns to low and the gate voltage increases by  $V_{DD} - V_{IE} - V_{II}$ . This is the voltage gain per cycle,  $V_{GAIN}$ . Continuing this process alternately, the gate voltage reaches  $V_{PP} + V_{IE}$  and  $V_{PP}$  is output. The necessary condition that the voltage gain be positive at the gate voltage of  $V_{PP} + V_{IE}$  is expressed by

$$V_{GAIN} \equiv V_{DD} - V_{IE} - V_{II} \quad (6.30)$$

at a back bias of  $V_{PP} + V_{IE}$ . Therefore, the minimum operating supply voltage  $V_{DD\_MIN}$  is given by

$$V_{DD\_MIN} \equiv V_{IE} + V_{II} \quad (6.31)$$

When a  $V_{PP}$  of 18 V,  $V_{IE}$  of 1.7 V, and  $V_{II}$  of 0.7 V at a back bias of 18 V are assumed, the minimum operating supply voltage and the maximum voltage for the switching gate are, respectively, 2.4 V and 19.7 V. Thus,  $V_{IE}$  raises the  $V_{DD\_MIN}$  and the maximum  $V_g$  in the NMOS  $V_{PP}$  switch. To decrease  $V_{DD\_MIN}$  for low voltage operation,  $V_{IE}$  needs to be reduced, but the leakage current flowing from  $V_{PP}$  would increase accordingly.

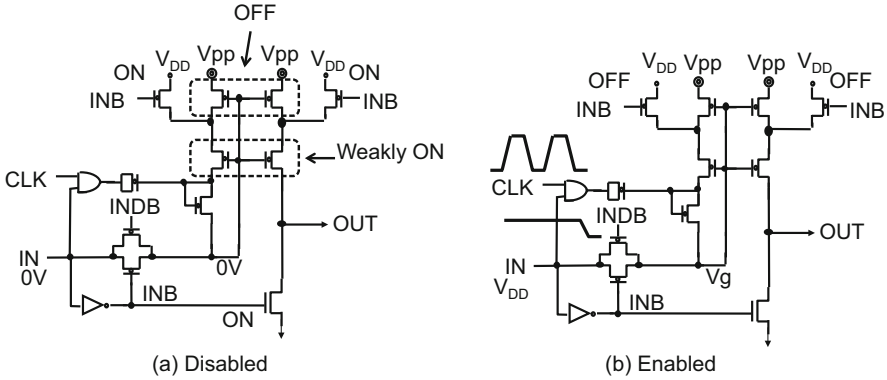
Figure 6.23 overcomes these two contradictory constraints, i.e., reduction of  $V_{DD\_MIN}$  and elimination of the leakage current from  $V_{PP}$  at a sacrifice of a little higher  $I_{DD}$  in active mode. All of the high-voltage transistors except for the pull-down used in the switch are intrinsic ones. Instead of the enhancement transistor in the standard NMOS level shifter, three intrinsic transistors whose  $V_{II}$  at a body bias of 0 V is around 0 V are used. In selected state, the input signal  $IN$  turns to high. In Fig. 6.23, the voltage gain per cycle,  $V_{GAIN}$  and the minimum operating supply voltage  $V_{DD\_MIN}$  are respectively given by

$$V_{GAIN} \equiv V_{DD} - 2V_{II} \quad (6.32)$$

$$V_{DD\_MIN} \equiv 2V_{II} \quad (6.33)$$

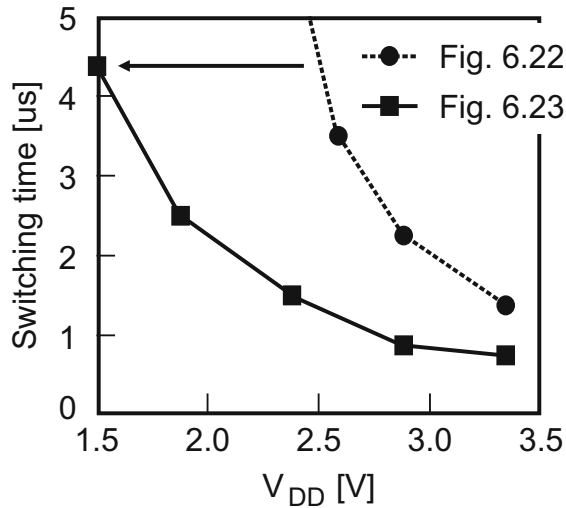
As shown in Fig. 6.23a, in disabled state, the third low-Vt transistor connected between the serially connected low-Vt transistors forces the intermediate node to 1 V or higher. This bias condition creates a negative  $V_{GS}$  of the upper transistor, resulting in no leakage current flowing from  $V_{PP}$ . On the other hand, the lower





**Fig. 6.23** Low voltage NMOS high-level shifter (Tanzawa et al. 1997)

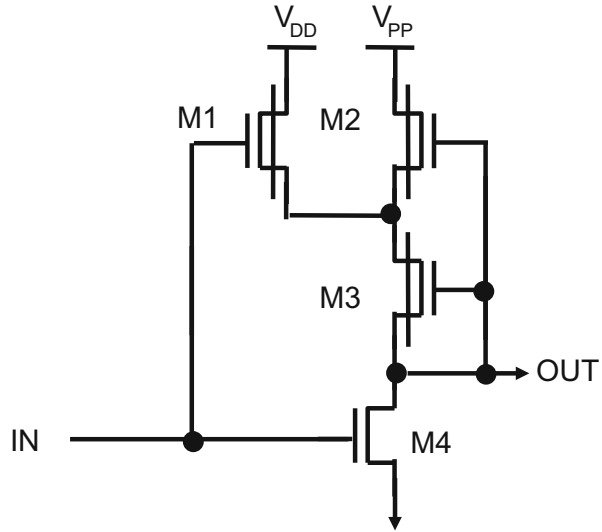
**Fig. 6.24**  $V_{DD}$  vs. switching time (Tanzawa et al. 1997)



transistor can flow a finite leakage current from  $V_{DD}$  even with the gate grounded, resulting in a slight increase in active current. In this example,  $V_{DD\_MIN}$  can be reduced from 2.5 to 1.5 V, as shown in Fig. 6.24. In the case of a  $V_{PP}$  of 18 V,  $V_{IE}$  of 1.7 V, and  $V_{II}$  of 0.7 V at a back bias of 18 V, each of  $V_{GAIN}$  and  $V_{DD\_MIN}$  is reduced by 1 V comparing Eq. (6.32) with Eq. (6.30) and Eq. (6.33) with Eq. (6.31).

Figure 6.25 shows another topology of NMOS level shifter. The circuit uses depletion NMOS M1–3 and enhancement NMOS M4 instead of using low- $V_t$  or enhancement NMOS and driving clock. When the input signal  $IN$  is high, M4 turns on to output low. M1 biases the source terminal of M2, so that M2 is cut off to prevent the leakage current from flowing from  $V_{PP}$  at a sacrifice of an increase in  $I_{DD}$ . The depletion NMOS needs to have the conditions on  $|V_t|$  as given below.

**Fig. 6.25** Depletion and enhancement NMOS high-level shifter (Lucero et al. 1983a)



$$|V_t(V_{BS} = -V_{DD})| < V_{DD} \quad (6.34)$$

$$V_t(V_{BS} = -V_{PP}) < 0 \text{ V} \quad (6.35)$$

Equation (6.34) guarantees that M2 is off when  $IN$  is high. Equation (6.35) shows that the output is as high as  $V_{PP}$  without any voltage drop when  $IN$  is low.

### 6.3.2 CMOS High-Level Shifter

This subsection focuses on CMOS high-level shifter with two cross-coupled PMOS and two complementary pull-down NMOS, as shown in Fig. 6.26.

Figure 6.27 shows level shifter operations. When the input goes from 0 V to  $V_{DD}$  of 2 V, the output is supposed to go from 0 V to  $V_{PP}$  (a). Thus, the high level increases  $V_{OUT}$  from  $V_{DD}$  to  $V_{PP}$ . One can divide the period into three portions (b), (c), and (d). When the input is 0 V as in (b),  $V_{OUT}$  is grounded thereby P1 turns on. Because N1 is off, the drain voltage of N1 is stable at  $V_{PP}$ , which turns off P2. As a result, all the nodes are in a latched state with no DC current flowing. When the input goes to 2 V as in (c), both N1 and P1 flow the current from  $V_{PP}$  to ground. Figure 6.28 shows the behavior in this transition. Suppose the NMOS is much stronger than PMOS as shown in the  $V_{OUT} - I_{DS}$  curves. The initial  $V_{OUT}$  is  $V_{PP}$  as shown by  $V_{INIT}$ . Because the NMOS current  $I_{DN}$  is larger than the PMOS current  $I_{DP}$ , the operating point is moving to  $V_{FIN1}$ . At this point, P2 strongly turns on and so the output node increases up to  $V_{PP}$ , which makes P1 turn off as shown in Fig. 6.27d. Thus, the drain voltage of N1 finally reaches 0 V. Because the circuit has symmetry, the same operation occurs when the input goes down to 0 V. Because

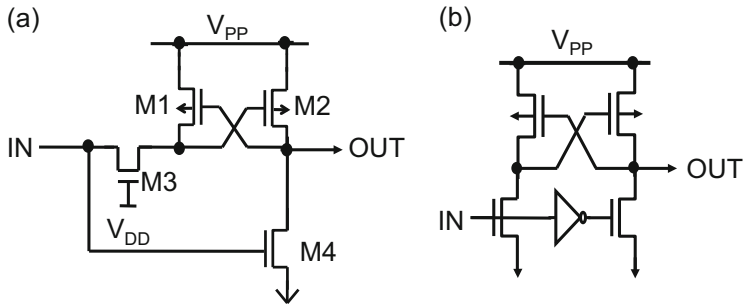


Fig. 6.26 CMOS high-level shifter (e.g., Tanaka et al. 1984a, b; Mehrotra et al. 1984a, b)

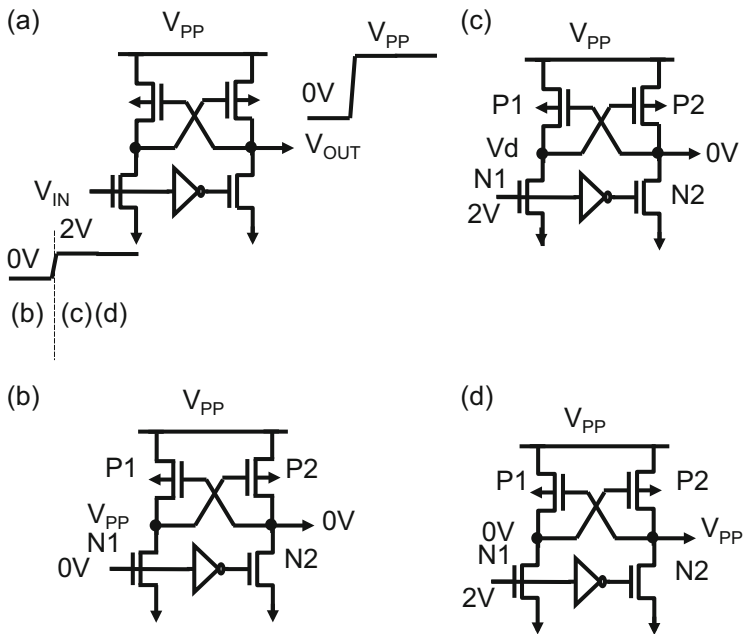


Fig. 6.27 CMOS high-level shifter

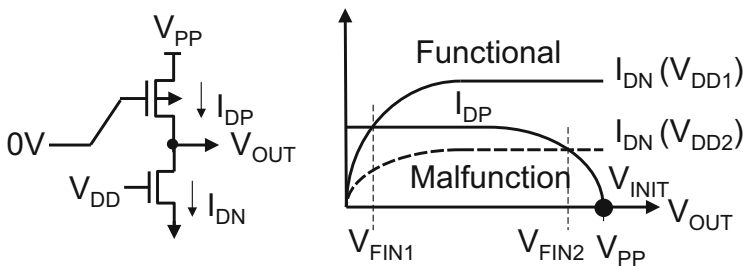


Fig. 6.28 Operating point of the CMOS high-level shifter

$V_{GS}$  of PMOS can be much larger than that of NMOS, the W/L ratio has to be sufficiently imbalanced. To estimate the dimensions, a long channel approximation model is used. The drain current of PMOS and NMOS transistors under the bias condition as shown in Fig. 6.28 is given by

$$I_{DP} = \frac{\mu_h C_{OX} W_P}{2 L_P} (V_{PP} - |V_{tP}|)^2 \quad (6.36)$$

$$I_{DN} = \frac{\mu_e C_{OX} W_N}{2 L_N} (V_{DD} - V_{tN})^2 \quad (6.37)$$

where  $\mu_{h(e)}$  is the mobility of hole (electron),  $C_{ox}$  is the gate capacitance per area,  $W_{P(N)}$  is the channel width of P(N)MOSFET,  $L_{P(N)}$  is the channel length of P(N) MOSFET, and  $V_{tP(N)}$  is the threshold voltage of P(N)MOSFET. To pull down the output node enough to invert the state, the equivalent point where the NMOS current is equivalent to the PMOS current needs to be not as high as  $V_{FIN2}$  but as low as  $V_{FIN1}$  as shown in the waveform of Fig. 6.28.

Thus, the condition where the level shifter works is given by

$$I_{DN\_MIN} > I_{DP\_MAX} \quad (6.38)$$

Assuming

$$\mu_h = \mu_e/2 \quad (6.39)$$

Equations (6.36)–(6.38) are reduced to

$$A_P/A_N \equiv \frac{W_P}{L_P} \bigg/ \frac{W_N}{L_N} < 2(V_{DD\_MIN} - V_{tN})^2 / (V_{PP\_MAX} - |V_{tP}|)^2 \quad (6.40)$$

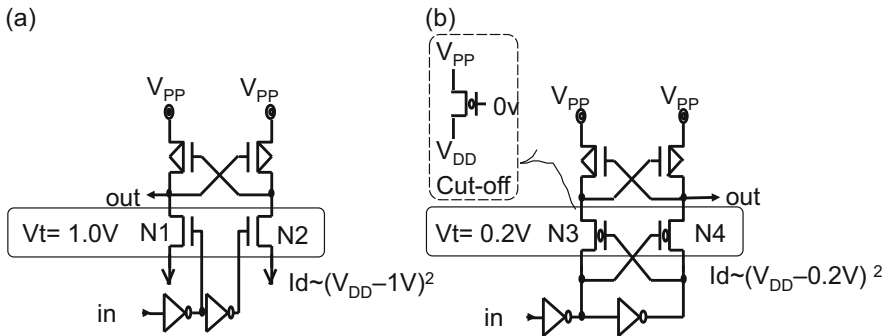
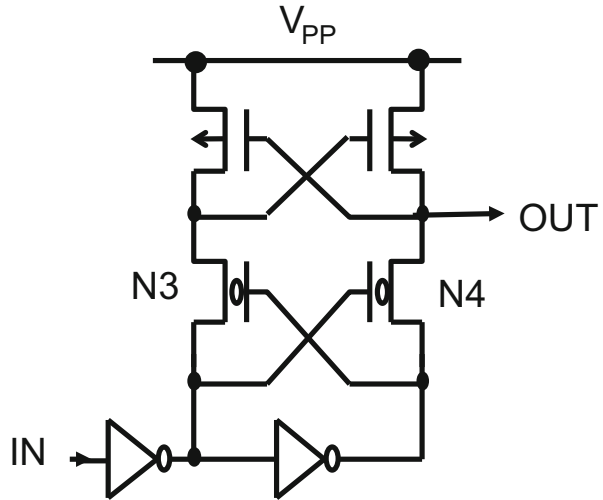
In case of  $V_{DD\_MIN} = 1.5$  V,  $V_{PP\_MAX} = 4$  V, and  $|V_{tP(N)}| = 1$  V, the aspect ratio  $A_P/A_N$  needs to be smaller than 1/18.

To allow lower voltage operation without increasing the switching delay, transistor sizes need to be kept same without  $I_{DN}$  reduced. Figure 6.29 shows a CMOS high-level shifter with low-Vt NMOS N3, N4 with  $V_{tN} \sim 0$  V.

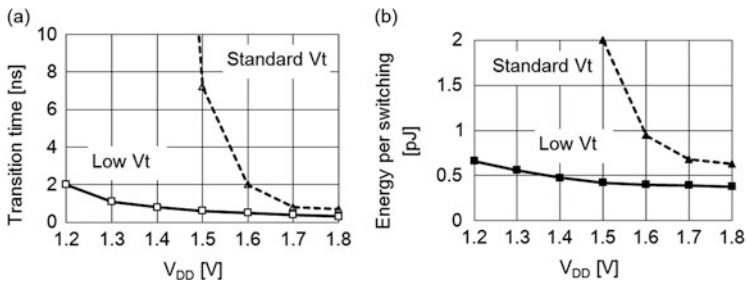
To what extent can the low-Vt NMOS reduce  $V_{DD\_MIN}$ ? In order to not flow a standby current, one only needs to bias the source terminal when the gate is grounded, as shown in 6.40. Thus, both PMOS and NMOS are connected as cross-coupled. When the NMOS needs to strongly turn on, the gate overdrive can be increased with lower Vt, resulting in lower  $V_{DD\_MIN}$ , which has to meet Eq. (6.40) (Fig. 6.30).

Figure 6.31 shows simulation results for the switching time (a) and energy per switching (b) against  $V_{DD}$ . The low- $V_{DD}$  high-voltage level shifter shows significant improvement in reduction in  $V_{DD\_MIN}$  by about 0.5 V.

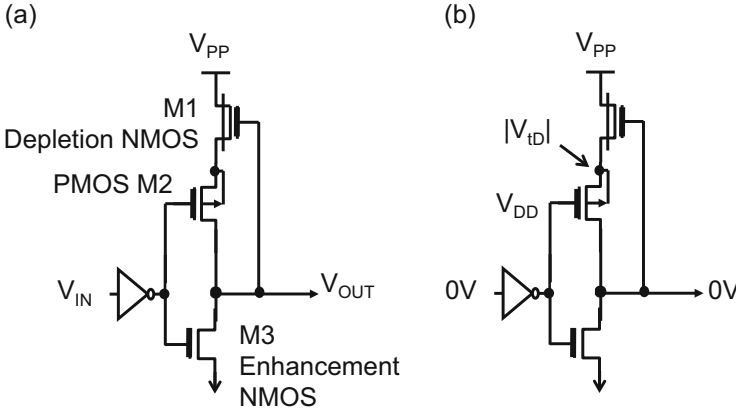
**Fig. 6.29** Low VDD CMOS high-level shifter with low-Vt NMOS N3, N4 (Tanzawa et al. 2001)



**Fig. 6.30** CMOS high-level shifter with standard-Vt (a) and low-Vt (b) NMOS (Tanzawa et al. 2001)



**Fig. 6.31** VDD vs. switching time (a) and energy per switching (b) of the CMOS high-level shifters with standard-Vt and low-Vt NMOS (Tanzawa et al. 2001)



**Fig. 6.32** Depletion NMOS (M1) and enhancement PMOS (M2) high-level shifter (Wada et al. 1989)

### 6.3.3 Depletion NMOS and Enhancement PMOS High-Level Shifter

Figure 6.32a shows another type of high-level shifter. When  $V_{IN}$  stays low, the pull down M3 forces the output node ground. M1 turns off with  $V_S = |V_{tD}|$ , where  $V_{tD}$  is the threshold voltage of M1, as far as M2 turns off with  $V_S = |V_{tD}|$  and  $V_G = V_{DD}$ , as shown in Fig. 6.32b, resulting in Eq. (6.41). When  $V_{IN}$  goes high, M2 turns on as far as  $|V_{tP}|$  is lower than  $|V_{tD}|$ , as shown in Fig. 6.33a, resulting in Eq. (6.42). Theoretically, once the output terminal of the level shifter starts increasing, the loop composed of M1 and M2 becomes positive as shown in Fig. 6.33b. The positive feedback continues until  $V_{tD}(V_{BS} = -V_{OUT})$  becomes 0 V.

Thus, it is necessary that  $V_{tD}(V_{BS} = -V_{PP})$  is negative to make the level shifter functional up to  $V_{PP}$ , resulting in Eq. (6.43). Equations (6.41)–(6.43) define  $V_T$  window to make the level shifter functional under the condition where  $V_{DD}$  is given or  $V_{DDMIN}$  under the condition where  $V_T$ 's are given.

$$|V_{tD}(V_{BS} = V_{tD})| - V_{DD} < |V_{tP}| \quad (6.41)$$

$$|V_{tP}| < |V_{tD}| \quad (6.42)$$

$$|V_{tD}(V_{BS} = -V_{PP})| < 0 \quad (6.43)$$

The requirement for  $V_{DS}$  of PMOS M2 is as low as  $|V_{tD}|$ , which can be much lower than  $V_{PP}$  in case of the CMOS level shifter. Therefore, the process cost may be lower than CMOS level shifter because of no need of specific junction process. Also, the high-voltage device counts can be smaller than CMOS level shifter.

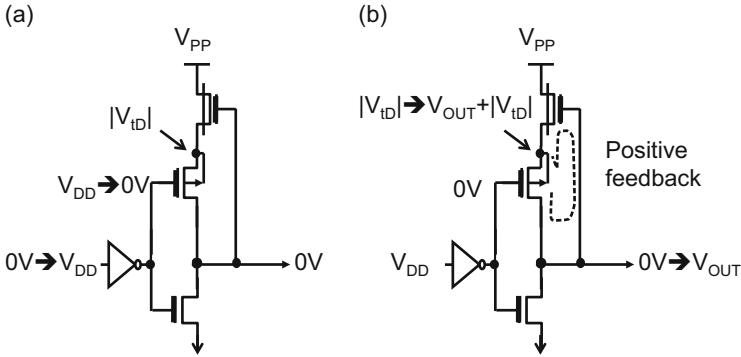


Fig. 6.33 Transient operation (Wada et al. 1989)

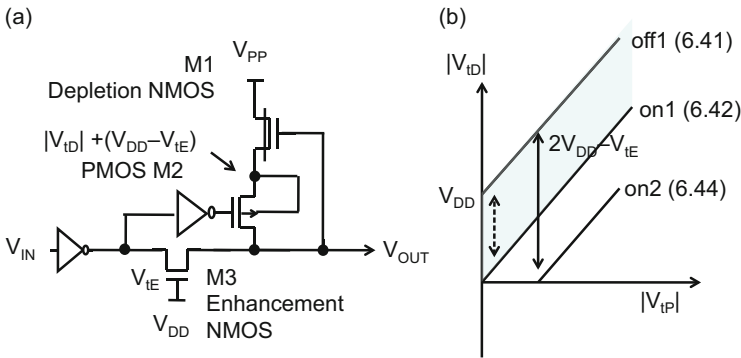
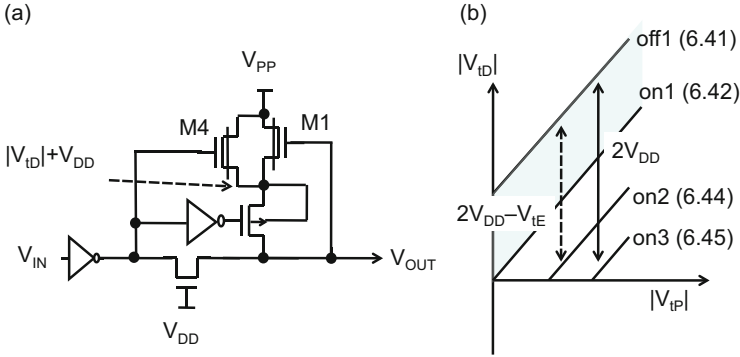


Fig. 6.34 Wider operation window D-NMOS+PMOS high-level shifter (Futatsuyama et al. 2009)

To widen the  $V_T$  window or to reduce  $V_{DDMIN}$ , another circuit shown in Fig. 6.34a adds a precharge path to the depletion NMOS M1. Equation (6.42) is replaced with Eq. (6.44),

$$|V_{IP}| < V_{DD} - V_{tN} + |V_{ID}| \tag{6.44}$$

which is the initial condition where the PMOS becomes conductive. In case of  $V_{tE} = 1$  V and  $V_{DD} = 2$  V, the level shifter as shown in Fig. 6.34 relaxes the constraint for  $|V_{IP}| - |V_{ID}|$  by 1 V. Figure 6.34b shows the  $V_T$  process window to have both the sufficient turn-on and cut-off conditions. The circuit of Fig. 6.32a has the  $V_T$  window between “off1” and “on1” whereas that of Fig. 6.34a has the  $V_T$  window between “off1” and “on2.” Instead of widening the  $V_T$  window, one can reduce  $V_{DD}$ . Assuming that a margin of 2 V is needed between the *off* and *on* conditions,  $V_{DDMIN}$  for the circuits of Figs. 6.32a and 6.34a has to be 2 and 1.5 V, respectively, in case of  $V_{tE} = 1$  V.



**Fig. 6.35** Another wider operation window D-NMOS + PMOS high-level shifter (Tanzawa 2012)

**Fig. 6.36** High-voltage pass gate with  $V_{PPH}$  higher by  $V_{IE}$  than  $V_{PP}$

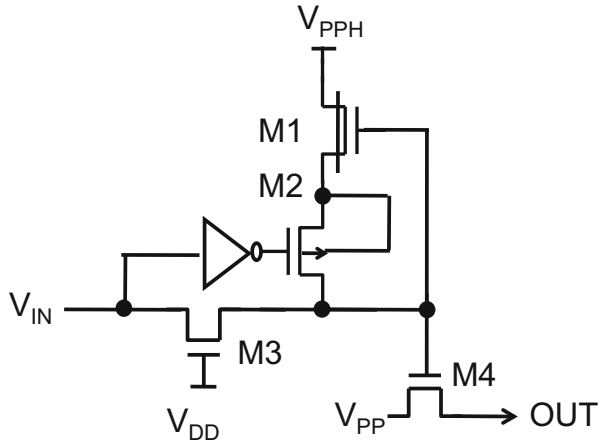


Figure 6.35a shows another high-level shifter with wider operation window. Additional depletion NMOS M4 is connected in parallel with M1, which boosts the source potential of the PMOS at the beginning of the operation. It has

$$|V_{IP}| < V_{DD} + |V_{ID}| \tag{6.45}$$

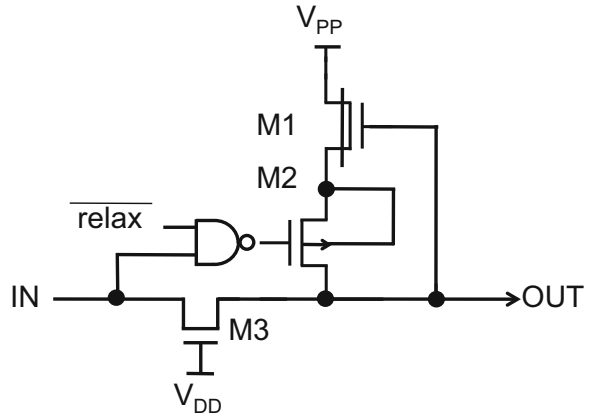
instead of Eq. (6.44). Figure 6.35b shows the window between “off1” and “on3.”  $V_{DD\_MIN}$  can be as low as 1 V under the same assumption as above.

Combining the level shifter of Fig. 6.34a with the regulator of Fig. 6.12, a high-voltage pass gate is obtained as shown in Fig. 6.36.

Because  $V_{PPH}$  is higher by  $V_{IE}$  than  $V_{PP}$ , the pass gate can fully transfer  $V_{PP}$  with a minimal overdrive. The switching speed is determined by the output impedance of M4. When the pass gate M4 is disabled with  $V_{IN}$  low, the drain terminals of M1 and M4 are biased at the high voltages whereas the gate and source terminals are kept low. In this case, there is a gate edge stress from drain to gate. However, because the drain of HV NMOS is usually lightly doped, the voltage stress is low enough. All



**Fig. 6.37** Level shifter with a relaxed gate stress (Tanzawa 2010)

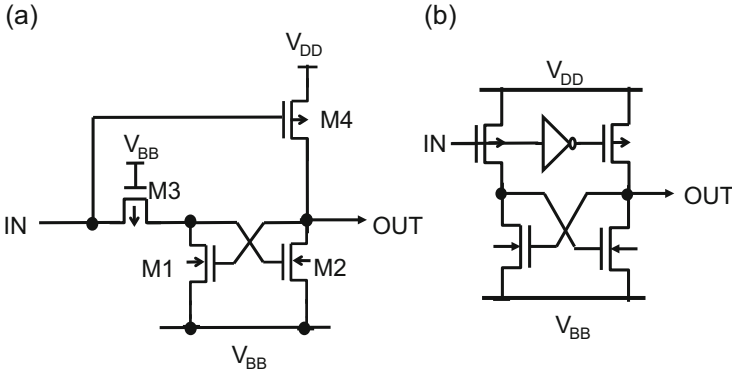


the terminals of the PMOS M2 are biased by low voltages as well. When the pass gate is enabled with  $V_{IN}$  high, all the terminals of M1 and M4 are biased by high voltages, but  $V_{GS}$  of M2 and M4 is much lower than the high voltages. On the other hand, the PMOS M2 is under a gate stress condition with the gate grounded and the source and drain biased with  $V_{PPH}$ . As a result, the HV oxide thickness is determined in a way that  $V_T$  of HV PMOS is not shifted by more than an acceptable amount due to such a Negative Bias Temperature Stability (NBTI) stress.

Figure 6.37 shows a level shifter with a relaxed gate stress. After transferring a part of  $V_{PP}$  to the output terminal, the gate of M2 is biased by  $V_{DD}$ , with the additional control signal/relax low. Even with an input of  $V_{DD}$  to the gate, M2 keeps on-state because the source and drain become high enough. Therefore, the gate oxide thickness can be reduced by roughly  $(V_{PP} - V_{DD})/V_{PP}$  to maintain the NBTI stress. The level shifter of Fig. 6.37 has one logic more than that of Fig. 6.34, but an increase in the area is limited because it only includes low voltage transistors. In addition, there is no timing overhead with the level shifter of Fig. 6.37 over that of Fig. 6.34, because the switching speed is limited by the impedance of the pass transistor such as M4 of Fig. 6.36. Thus, all the HV devices, including the HV capacitors, can be scaled by the ratio  $(V_{PPH} - V_{DD})/V_{PPH}$  with Fig. 6.37 under the condition that the gate electric field is kept the same and the impact of the gate edge stress is still low enough with a thinner gate oxide.

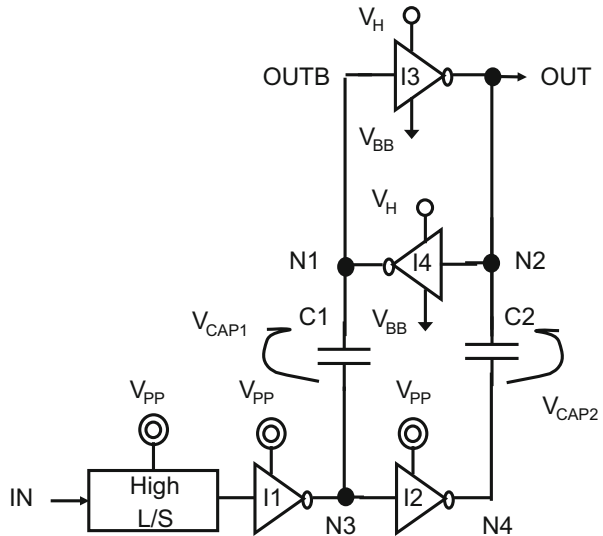
### 6.3.4 CMOS Low-Level Shifter

Low-level shifter converts the low level of the input logic into a negative voltage whereas the high level is unchanged. The circuits of Fig. 6.38 input IN whose voltage amplitude is  $V_{DD}$  or GND and output OUT whose voltage amplitude is  $V_{DD}$  or a negative voltage of  $V_{BB}$ . The topology is fully complementary to the CMOS high-level shifter of Fig. 6.26. Maximum voltage differences between two terminals of each transistor such as  $V_{GS}$ ,  $V_{DS}$ ,  $V_{DB}$ , and  $V_{SB}$  become  $V_{DD} + |V_{BB}|$ .



**Fig. 6.38** CMOS low-level shifter

**Fig. 6.39** CMOS low-level shifter using coupling capacitors (Tanzawa et al. 2002)



In case where  $|V_{BB}|$  is close to  $V_{DD}$ , all the transistors except for the inverter are usually high-voltage ones whose gate oxide is thicker and whose channel length is longer than low-voltage transistors. In case where  $|V_{BB}|$  becomes much larger than  $V_{DD}$ , the gate oxide needs to be much thicker. Under such a condition, reduction in  $V_{DD}$  is limited to make the PMOS strong enough to compulsorily invert the outputs. Thus, scaling the high-voltage transistor is a challenging item for the low-voltage level shifter.

To reduce the voltage for the logic high of the last stage of the low-level shifter, flipping and latching operations are separated using coupling capacitors, as shown in Fig. 6.39. The inverters *I1* and *I2* which, respectively, drive the nodes *N3* and *N4* can have sufficient driving currents to invert the latch via the coupling capacitors

**Table 6.2** Nodal voltages of the low-level shifter of Fig. 6.39 (Tanzawa et al. 2002)

	Initial	Transition	Final
N1	$V_H$	$V_H - (V_{PP} - V_{SS})$	$V_{BB}$
N2	$V_{BB}$	$V_{BB} + (V_{PP} - V_{SS})$	$V_H$
N3	$V_{PP}$	$V_{SS}$	$V_{SS}$
N4	$V_{SS}$	$V_{PP}$	$V_{PP}$
$V_{CAP1}$	$ V_H - V_{PP} $	$ V_H - V_{PP} $	$ V_{SS} - V_{BB} $
$V_{CAP2}$	$ V_{SS} - V_{BB} $	$ V_{SS} - V_{BB} $	$ V_H - V_{PP} $

$C1$  and  $C2$ . The operation voltages of the inverters are  $V_{PP}$  and  $V_{SS}$ , whereas those of the latches are  $V_H$  and  $V_{BB}$ . Table 6.2 shows the nodal voltages of the low-level shifter of Fig. 6.39.

In order to invert the latch, the condition Eq. (6.46) has to hold according to  $V(N1) < V(N2)$  in the transition period.

$$V_H - V_{PP} + V_{SS} < V_{BB} + V_{PP} - V_{SS} \quad (6.46)$$

In addition, because the capacitors have the gate oxide of the high-voltage transistors, the capacitor voltages  $V_{CAP1}$  and  $V_{CAP2}$  are equal to or less than  $V_{MAX}$ , resulting in the following conditions, respectively.

$$V_{PP} - V_H \leq V_{MAX} \quad (6.47)$$

$$V_{SS} - V_{BB} \leq V_{MAX} \quad (6.48)$$

Furthermore, the voltage difference between the logic high and low voltages of the inverters is also equal to or less than the maximum allowable voltage  $V_{MAX}$ .

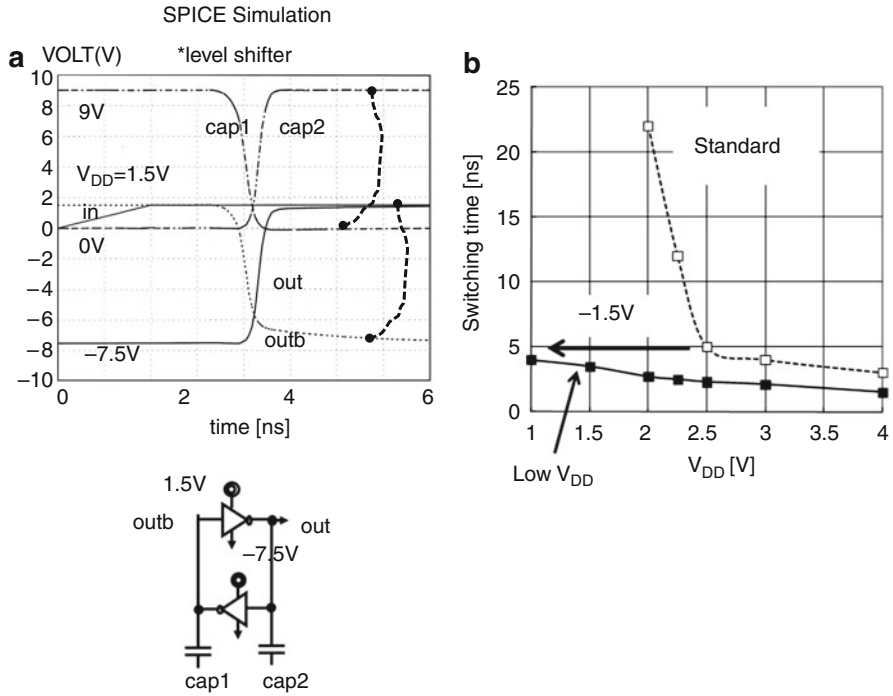
$$V_{PP} - V_{SS} \leq V_{MAX} \quad (6.49)$$

Moreover, the transient voltages at the nodes  $N1$  and  $N2$  have to be between  $V_H$  and  $V_{BB}$ , otherwise the forward bias conditions occur. Thus, the condition should hold as follow.

$$V_{BB} \leq V_H - V_{PP} + V_{SS} \quad (6.50)$$

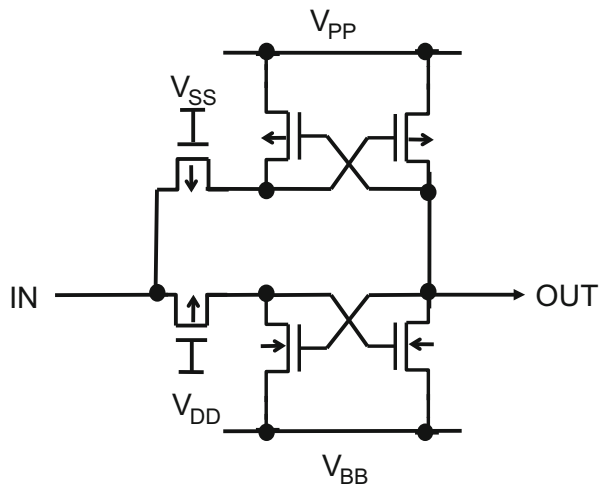
Figure 6.40a shows a simulation waveform of the circuit where  $V_{PP} = 9$  V,  $V_H = 1.5$  V, and  $V_{BB} = -7.5$  V, which meet all the conditions of Eqs. (6.46)–(6.50). The input has 0 and 1.5 V as the two logic levels, which translate into 9 V via a high-voltage shifter. Then, the high amplitude cap1,2 shifts the voltage levels of out and outb as shown without any overstress. Figure 6.40b shows the switching speed vs.  $V_{DD}$ .  $V_{DD\_MIN}$  can be reduced by about 1.5 V.

In case where the high-voltage transistor is determined by another constraint, small circuit area can become the main concern to design the level shifter. Figure 6.41 has just six transistors to convert the voltage level from  $V_{DD}$  ( $V_{SS}$ ) to  $V_{PP}$  ( $V_{BB}$ ). Because the number of gate counts from IN to OUT is much less than the other types of level shifters, the switching delay can be the minimum with this structure.



**Fig. 6.40** Operation waveform (a) and  $V_{DD}$  vs. switching time (b) (Tanzawa et al. 2002)

**Fig. 6.41** CMOS level shifter with both high and low-level shifting (Yamagata et al. 1995)



**Table 6.3** Summary: trade-offs between the level shifters

	High/low	MOS FET	Switching speed	Power	Process cost	Circuit area	$V_{DD\_MIN}$
1.	High	NMOS	Slowest	Highest	Lowest	Large	High
2.	High	CMOS	Fast	High	Highest	Small	Low/ mid
3.	High	D-NMOS + PMOS	Slow	Low	High	Small	Mid
4.	Low	CMOS	Fast	High	Highest	Small to large	Low/ mid
5.	High and Low	CMOS	Fast	High	Highest	Small	Low/ mid

Table 6.3 summarizes the trade-offs in switching speed, switching power, process cost, circuit area, and  $V_{DD\_MIN}$ , among the level shifters discussed. When the applications need to design level shifters where the switching speed is critical, one would have to select a technology supporting high-voltage CMOS with well-controlled  $V_T$ 's, even if that increases the process cost. On the other hand, when the switching speed is not a critical design parameter, one can select either NMOS or D-NMOS + PMOS level shifter depending on the total cost of the process and the die cost. If the level shifter does not affect the die size, the NMOS level shifter should have a lower total cost than the other one. Otherwise, the D-NMOS + PMOS level shifter can be the best choice. Requirement for controllability in  $V_T$ 's of high-voltage depletion NMOS and high-voltage PMOS can be a constraint on  $V_{DD\_MIN}$ .

## 6.4 Voltage Reference

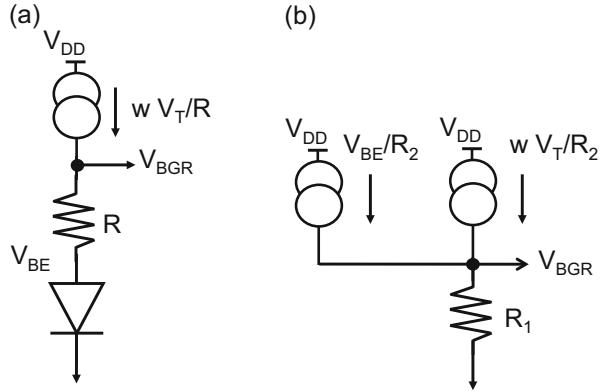
High-voltage generator needs to have a voltage reference to output an accurate high voltage. Bandgap reference outputs an accurate PVT insensitive voltage (Gray et al. 2001; Razavi 2000). Figure 6.42 shows the concept of bandgap reference. In (a),  $V_{BE}$  with a negative temperature coefficient is added with a thermal voltage  $V_T \equiv kT/q$  multiplied by a weight  $w$ .

$$V_{BGR} = V_{BE} + wV_T \quad (6.51)$$

Choosing an appropriate value for  $w$  in Eq. (6.51), one can have a PVT insensitive voltage as known as a bandgap voltage. In (b), two currents are summed with a single resistor  $R_1$ , resulting in another voltage reference.

$$V_{BGR} = (V_{BE} + wV_T)R_1/R_2 \quad (6.52)$$

**Fig. 6.42** Concept of bandgap reference: voltage sum (Widlar 1970) (a) and current sum (Banba et al. 1998) (b)



In addition to  $w$ , one can choose another parameter  $R_1/R_2$  to have a scaled bandgap voltage. Because  $R_1$  and  $R_2$  are made of same material, their ratio should have no temperature and process variations.

This section discusses design equations, sensitivity on device mismatch, and the minimum operation voltage of four types of bandgap references: Kuijk cell, Brokaw cell, Meijer cell, and Banba cell.

### 6.4.1 Kuijk Cell

Figure 6.43 illustrates Kuijk cell composed of two diodes D1–2, three resistors R1–3, PMOS load, and one opamp. D2 has the junction area  $N$  times larger than D1.

$$I_1 = I_S \exp(V_{BE1}/V_T) \tag{6.53}$$

$$I_2 = NI_S \exp(V_{BE2}/V_T) \tag{6.54}$$

From the fact that the two inputs of the opamp are equal,

$$R_1 I_1 = R_2 I_2 \tag{6.55}$$

Because the voltage at the upper terminal of R3 is given by  $V_{BE1}$  with the opamp,

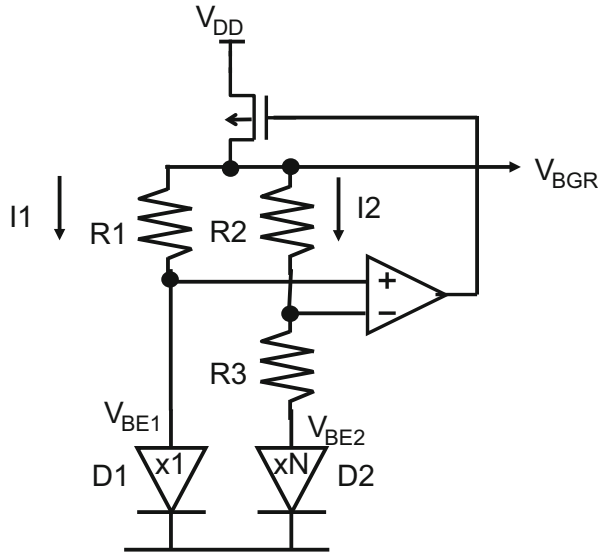
$$V_{BE1} - V_{BE2} = R_3 I_2 \tag{6.56}$$

From Eqs. (6.53) and (6.54),

$$I_2/I_1 = N \exp((V_{BE2} - V_{BE1})/V_T) \tag{6.57}$$

From Eqs. (6.55) and (6.57),

**Fig. 6.43** Kuijk cell bandgap reference (Kuijk 1973)



$$V_{BE1} - V_{BE2} = V_T \ln(NR_2/R_1) \tag{6.58}$$

From Eqs. (6.56) and (6.58),

$$I_2 = V_T \ln(NR_2/R_1)/R_3 \tag{6.59}$$

Therefore,

$$\begin{aligned} V_{BGR} &= V_{BE1} + R_2 I_2 \\ &= V_{BE1} + V_T R_2 / R_3 \ln(NR_2/R_1) \end{aligned} \tag{6.60}$$

Assuming the ratios of  $R$ 's have negligibly small temperature coefficient, the design equation to have zero temperature coefficient in  $V_{BGR}$  at  $T_0$  to the first order is given by

$$R_2/R_3 \ln(NR_2/R_1) = -\frac{q}{k} \frac{dV_{BE1}}{dT} \Big|_{T=T_0} \equiv \alpha \tag{6.61}$$

Without losing generality, one can constrain the following additional equation.

$$R_1 = R_2 \tag{6.62}$$

Equation (6.61) is then reduced to

$$R_2/R_3 = \alpha/\ln(N) \quad (6.63)$$

Next, the impact of mismatches on the reference voltage is considered as follows. In case where there is a finite input offset voltage  $V_{OS}$  of the opamp, it is assumed that the system is stable with  $V_{BE1} + V_{OS}$  at the minus input of the opamp instead of  $V_{BE1}$  and  $I_2 + \Delta I_2$  flowing through  $R_2$  instead of  $I_2$ .

$$\Delta I_2 = V_{OS}/R_3 \quad (6.64)$$

The variation in  $V_{BGR}$  is given by

$$\Delta V_{BGR} = V_{OS} + R_2 \Delta I_2 = (1 + R_2/R_3)V_{OS} \quad (6.65)$$

From Eqs. (6.65) and (6.63),

$$\Delta V_{BGR} = (1 + \alpha/\ln(N))V_{OS} \quad (6.66)$$

To reduce the variation in  $V_{BGR}$ , it is effective to have a large  $N$ . Deviation of  $V_{BGR}$ ,  $\delta V_{BGR}$ , due to each one of the device parameters in Eq. (6.60) is expressed as follows.

$$\begin{aligned} \delta V_{BGR} = & \delta V_{BE1} + V_T [\delta R_2/R_3 \ln(NR_2/R_1) - \delta R_3 R_2/R_3^2 \ln(NR_2/R_1) \\ & + R_2/R_3 (\delta N/N + \delta R_2/R_2 - \delta R_1/R_1)] \end{aligned} \quad (6.67a)$$

Assuming that there is no correlation between any two of the deviations in the device parameters, the standard deviation of  $V_{BGR}$ ,  $\sigma V_{BGR}$ , is calculated together with  $V_{OS}$ .

$$\begin{aligned} (\sigma V_{BGR})^2 = & (\sigma V_{BE1})^2 + V_T^2 [(\sigma R_2)^2/R_3^2 (\ln(NR_2/R_1))^2 \\ & + (\sigma R_3)^2 R_2^2/R_3^4 (\ln(NR_2/R_1))^2 \\ & + R_2^2/R_3^2 ((\sigma N/N)^2 + (\sigma R_2/R_2)^2 + (\sigma R_1/R_1)^2)] \\ & + (1 + R_2/R_3)^2 (\sigma V_{OS})^2 \end{aligned} \quad (6.67b)$$

The minimum operating supply voltage is determined by either one of the load PMOS or the opamp. Assuming the opamp does not limit it,  $V_{DD\_MIN}$  is a sum of the output voltage and  $V_{DS}$  of the load PMOS, i.e.,

$$V_{DD\_MIN} = V_{BGR} + V_{DS} \quad (6.68)$$

which can be as low as about 1.5 V.



### 6.4.2 Brokaw Cell

Figure 6.44 illustrates Brokaw cell composed of two NPN bipolar junction transistors (bjt's), four resistors R1–4, where the left bottom part is counted as one, and one opamp. From Fig. 6.44,

$$I_1 R_3 = I_2 R_4 \tag{6.69}$$

$$I_{B1} = I_1 / \beta_1 = N I_S \exp(V_{BE1} / V_T) \tag{6.70}$$

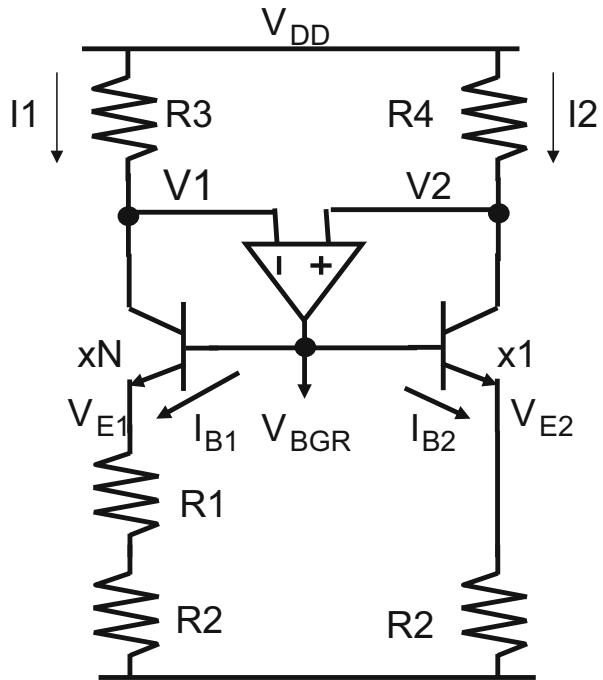
$$I_{B2} = I_2 / \beta_2 = I_S \exp(V_{BE2} / V_T) \tag{6.71}$$

where  $\beta_1$  and  $\beta_2$  are the multiplication factors of the collector currents to the base currents of the left- and right-hand side bjt, respectively, and  $N$  is the area ratio of the two bjt's. In the right-hand side branch,

$$V_{E2} = R_2 I_{b2} (\beta_2 + 1) \tag{6.72}$$

Since the difference between  $V_{BE1}$  and  $V_{BE2}$  appears at the voltage difference between both terminals of  $R_1$ ,

**Fig. 6.44** Brokaw cell bandgap reference (Brokaw 1974)



$$V_{E1} - V_{E2} = V_{BE2} - V_{BE1} = R_1 I_{B1} (\beta_1 + 1) \quad (6.73)$$

From Eqs. (6.69)–(6.71),

$$V_{BE2} - V_{BE1} = V_T \ln \left( N \frac{R_3 \beta_1}{R_4 \beta_2} \right) \quad (6.74)$$

$$I_{B2} = I_{B1} \frac{R_3 \beta_1}{R_4 \beta_2} \quad (6.75)$$

From Eqs. (6.73) and (6.74),

$$I_{B1} = \frac{V_T \ln \left( N \frac{R_3 \beta_1}{R_4 \beta_2} \right)}{R_1 (\beta_1 + 1)} \quad (6.76)$$

Using Eqs. (6.72), (6.75), and (6.76),

$$\begin{aligned} V_{BGR} &= V_{BE2} + V_{E2} = V_{BE2} + R_2 I_2 \\ &= V_{BE2} + V_T \frac{R_2 R_3 \beta_1 (\beta_2 + 1)}{R_1 R_4 (\beta_1 + 1) \beta_2} \ln \left( N \frac{R_3 \beta_1}{R_4 \beta_2} \right) \end{aligned} \quad (6.77a)$$

In case where  $\beta_1 = \beta_2$ ,  $R_3 = R_4$ , Eq. (6.77a) is reduced to Eq. (6.77b).

$$V_{BGR} = V_{BE2} + \frac{R_2}{R_1} V_T \ln(N) \quad (6.77b)$$

Assuming the ratios of  $R$ 's have negligibly small temperature coefficient, the design equation to have zero temperature coefficient in  $V_{BGR}$  at  $T_0$  is given by

$$R_2/R_1 \ln(N) = - \frac{q}{k} \frac{dV_{BE2}}{dT} \Big|_{T=T_0} \equiv \alpha \quad (6.78)$$

Next, the impact of mismatches on the reference voltage is considered as follows. In case where there is a finite input offset voltage  $V_{OS}$  of the opamp, it is assumed that the system is stable with  $V_2 - V_{OS} + \Delta V_1$  at the plus input of the opamp instead of  $V_2$ ,  $V_1 + \Delta V_1$  at the minus input of the opamp instead of  $V_1$ , and  $I_2 + \Delta I_2$  flowing through  $R_3$  instead of  $I_2$ .

$$\Delta I_2 = V_{OS}/R_3 \quad (6.79)$$

Further assuming  $V_{BE2}$  varies by  $\Delta V_{BE2}$  due to  $\Delta I_2$ ,

$$\Delta V_{BGR} = \Delta V_{BE2} = R_2 \Delta I_2 = V_{OS} R_2 / R_3 \quad (6.80)$$

To reduce the variation in  $V_{BGR}$ , it is effective to increase the value for  $R_3$ , which is determined by  $V_{DD\_MIN}$ .

$$V_{DD\_MIN} = V_{E2} + V_{CE2} + R_3 I_2 = V_{BGR} - V_{BE2} + V_{CE2} + R_3 I_2 \quad (6.81)$$

Using Eqs. (6.77a, 6.77b) and (6.81), Eq. (6.80) is written by

$$\Delta V_{BGR} = V_{OS} \frac{V_{BGR} - V_{BE2}}{V_{DD\_MIN} - V_{BGR} + V_{BE2} - V_{CE2}} \frac{\beta_2}{\beta_2 + 1} \quad (6.82)$$

Even for a low supply voltage such as 1.5 V, the variation in  $V_{BGR}$  due to the input offset voltage could be close to the input offset voltage itself. Equation (6.82) is typically much smaller than Eq. (6.66).

### 6.4.3 Meijer Cell

Figure 6.45 shows Meijer cell bandgap reference composed of two resistors, two bjt's, and two PMOS transistors.

$$I_{B1} = N I_S \exp((V_{BE2} - V_{E1})/V_T) \quad (6.83)$$

$$I_{B2} = I_S \exp(V_{BE2}/V_T) \quad (6.84)$$

$$V_{E1} = R_1 (\beta + 1) I_{B1} \quad (6.85)$$

$$I_{C1} = \beta I_{B1} \quad (6.86)$$

$$I_{C2} = \beta I_{B2} \quad (6.87)$$

When two mirror PMOS are identical in size,

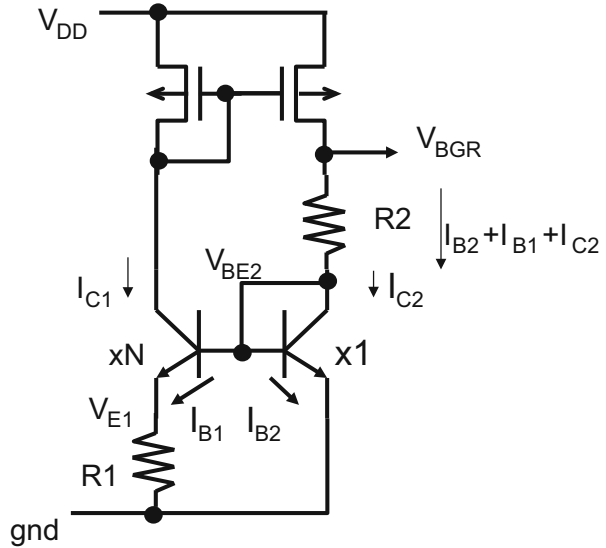
$$I_{C1} = I_{B2} + I_{B1} + I_{C2} \quad (6.88)$$

From Eqs. (6.86)–(6.88),

$$I_{B2} = \frac{\beta - 1}{\beta + 1} I_{B1} \quad (6.89)$$

From Eqs. (6.83), (6.84), and (6.89),

**Fig. 6.45** Meijer cell bandgap reference (Meijer and Verhoeff 1976)



$$N \exp(-V_{E1}/V_T) = \frac{\beta + 1}{\beta - 1} \quad (6.90)$$

From Eqs. (6.85) and (6.90),

$$I_{B1} = \frac{V_T \ln \left[ N \frac{\beta - 1}{\beta + 1} \right]}{R_1 (\beta + 1)} \quad (6.91)$$

$V_{BGR}$  is then

$$\begin{aligned} V_{BGR} &= V_{BE2} + R_2 (I_{B2} + I_{B1} + I_{C2}) \\ &= V_{BE2} + \beta I_{B1} R_2 \\ &= V_{BE2} + \frac{R_2}{R_1} V_T \frac{\beta \ln [N(\beta - 1)/(\beta + 1)]}{\beta + 1} \end{aligned} \quad (6.92)$$

In case where  $\beta$  is much larger than 1, Eq. (6.92) is reduced to

$$V_{BGR} = V_{BE2} + \frac{R_2}{R_1} V_T \ln N \quad (6.93)$$

Next, the impact of the mismatch in the mirror PMOS transistors'  $V_t$  on  $V_{BGR}$  is studied.

$$I_{C1} = K_P(V_{GS} - |V_t|)^2 \tag{6.94}$$

$$\Delta(I_{B1} + I_{B2} + I_{C2}) = 2\sqrt{K_P I_{C1}} \Delta V_t \tag{6.95a}$$

$$\begin{aligned} \Delta V_{BGR} &= R_2 \Delta(I_{B1} + I_{B2} + I_{C2}) \\ &= 2R_2 \sqrt{K_P I_{C1}} \Delta V_t \end{aligned} \tag{6.95b}$$

$V_{DD\_MIN}$  is determined by either lower one of the left Eq. (6.96) or right Eq. (6.97) branch;

$$V_{DD\_MIN} = V_{E1} + V_{CE1} + V_{OD1} + |V_t| \tag{6.96}$$

$$V_{DD\_MIN} = V_{BGR} + V_{OD1} \tag{6.97}$$

where  $V_{CE1}$  is the collector-to-emitter voltage of the left BJT, and  $V_{OD1}$  is the overdrive voltage of the left PMOS.

### 6.4.4 Banba Cell

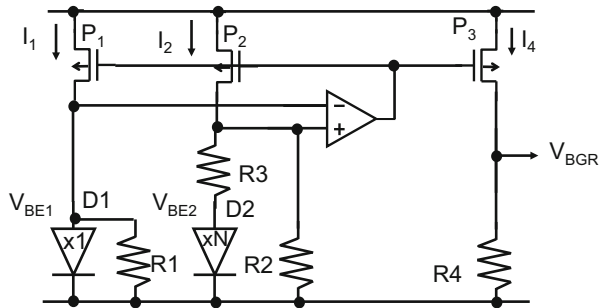
To reduce  $V_{DD\_MIN}$  for low voltage operation in advanced technology, another topology of bandgap reference with folded resistors is proposed as shown in Fig. 6.46, which uses four resistors, two diodes, three load PMOS transistors, and one opamp.

$$I_1 = I_S \exp(V_{BE1}/V_T) + V_{BE1}/R_1 \tag{6.98}$$

$$I_2 = (V_{BE1} - V_{BE2})/R_3 + V_{BE1}/R_2 \tag{6.99}$$

Equation (6.99) is extracted by using the fact that the two input nodes of the opamp are equal with the feedback loop. Because  $I_2$  flows  $R3$  and  $D2$ , the followings hold.

**Fig. 6.46** Banba cell bandgap reference (Banba et al. 1998)



$$(V_{BE1} - V_{BE2})/R_3 = NI_S \exp(V_{BE2}/V_T) \quad (6.100)$$

$$V_{BGR} = I_4 R_4 \quad (6.101)$$

For simplicity,  $R_1$  is equal to  $R_2$  and the  $P1$ ,  $P2$ , and  $P3$  are identical in size. Then, from Eq. (6.98) to Eq. (6.100) and  $I_1 = I_2 = I_4$ ,

$$V_{BE1} - V_{BE2} = V_T \ln N \quad (6.102)$$

From Eqs. (6.99), (6.101), and (6.102),

$$\begin{aligned} V_{BGR} &= I_2 R_4 = R_4 (V_T \ln N / R_3 + V_{BE1} / R_1) \\ &= \frac{R_4}{R_1} \left( V_{BE1} + \frac{R_1}{R_3} V_T \ln N \right) = \frac{R_4}{R_1} V_{BGR-V} \end{aligned} \quad (6.103)$$

where  $V_{BGR-V}$  is a bandgap voltage generated by a type of bandgap references such as Kuijk, Brakow, and Meijer outputting a voltage sum.

Next, the case where the opamp has an input offset voltage of  $V_{OS}$  is considered. Assuming that the voltages at the positive and negative input nodes of the opamp are, respectively, shifted by  $\Delta V_{BE1} + V_{OS}$  and  $\Delta V_{BE1}$  due to  $V_{OS}$ ,  $I_1$  shifts by

$$\Delta I_1 = \Delta V_{BE1} / R_1 + \Delta I_{DIO} = \Delta V_{BE1} (1/R_1 + \ln N / R_3) \quad (6.104)$$

where the following relation is used.

$$\Delta I_{DIO} = \frac{I_S}{V_T} \exp(V_{BE1}/V_T) \Delta V_{BE1} = \ln N / R_3 \Delta V_{BE1} \quad (6.105)$$

Similarly, assuming the voltage at the lower node of  $R3$  shifts by  $\Delta V_{BE2}$ ,  $I_2$  shifts by

$$\Delta I_2 = (\Delta V_{BE1} + V_{OS}) / R_2 + (\Delta V_{BE1} + V_{OS} - \Delta V_{BE2}) / R_3 \quad (6.106)$$

From the fact that the current through  $R3$  is same as that through  $D2$ ,

$$(\Delta V_{BE1} + V_{OS} - \Delta V_{BE2}) / R_3 = \Delta V_{BE2} \ln N / R_3 \quad (6.107)$$

Because the opamp controls the PMOSFETs in such a way that  $\Delta I_1$  is equal to  $\Delta I_2$ , Eqs. (6.104) and (6.106) lead to

$$\Delta V_{BE1} = V_{OS} \frac{\frac{R_3}{R_1} + 1 - \frac{1}{1 + \ln N}}{\ln N - 1 + \frac{1}{1 + \ln N}} \quad (6.108)$$

where Eq. (6.107) and the relation of  $R1 = R2$  are used. The deviation in  $V_{BGR}$  is calculated by Eqs. (6.101), (6.104), and (6.108) as follows:

**Table 6.4** Comparisons in variations and minimum operation voltages

	Kuijk	Brokaw	Meijer	Banba
$V_{DD\_MIN}$	$V_{BGR} + V_{DS} \sim 1.5$ V or $V_{DD\_MIN\_OPAMP}$			$V_{BE} + V_{DS} \sim 1.0$ V or $V_{DD\_MIN\_OPAMP}$
Sensitivity of $V_T$ mismatch (per 1 mV)	10 mV	2 mV	3 mV	10 mV
Sensitivity of $\beta$ variation (10–100)	1 mV	1 mV	100 mV	1 mV
BJT or diode	Diode	BJT	BJT	Diode

$$\Delta V_{BGR} = \Delta I_1 R_4 = (1/R_1 + \ln N/R_2) \frac{\frac{R_2}{R_1} + 1 - \frac{1}{1 + \ln N}}{\ln N - 1 + \frac{1}{1 + \ln N}} V_{OS} \quad (6.109)$$

Table 6.4 summarizes characteristics of four bandgap cells. The values represent typical ones. If bjt is available in a given process, Brokaw cell would be the best among the four types of bandgaps in terms of low  $V_{DD\_MIN}$  and small variation, as far as the supply voltage given is higher than  $V_{DD\_MIN}$  of the bandgap cell. Otherwise, Banba cell would be the best one because it can have lower  $V_{DD\_MIN}$  than the others and similar variation as Kuijk cell does.

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# Chapter 7

## System Design

**Abstract** This chapter provides high voltage generator system design. A gate level hard switching pump model is first presented for designing a single pump block. Multiple pumps are distributed in a die, each of which has wide power ground bus lines. Total area including the charge pump circuits and the power bus lines needs to be paid attention for overall area reduction. This is a nominal case where a charge pump needs to generate a higher voltage from DC energy transducer such as photovoltaic cell and thermo-electronic generator whose output impedance is not small. Design methodology is shown using an example. Another concern on multiple high voltage generator system design is system level simulation time. Even though the switching pump models are used for system verification, simulation run time is still slow especially for Flash memory where the minimum clock period is 20–50 ns whereas the maximum erase operation period is 1–2 ms. In order to drastically reduce the simulation time, another charge pump model together with a regulator model is described which makes all the nodes in the regulation feedback loop analog to eliminate the hard-switching operation.

Figure 7.1 illustrates on-chip high-voltage generator system and summarizes key discussion in each section. Section 7.1 reviews a hard-switching pump model for designing a single pump cell. The pump outputs the current with an enabling signal high and disconnects the output terminal with the signal low. Thus, two logic states in the signal make the pump hardly turn on or off. The pump model can be implemented in a system together with its pump regulator for system simulation. Section 7.2 expands the model to allow the power line resistance to be included as a design parameter rather than a given condition. Thus, one can determine the power line width as well as the pump parameters such as the number of stages and the pump capacitor to minimize the entire area for the pump and the power lines. Section 7.3 shows optimum design when a power line resistance is given, but not a design parameter. For example, a charge pump needs to generate a higher voltage from an extremely low DC voltage from DC energy transducer such as photovoltaic cell and thermo-electronic generator. The design of the charge pump has to take the input resistance into optimization. Section 7.4 then discusses a behavior model supporting to connect the power ground terminal of each pump with its local power ground lines. In case where power ground lines are shared with other pumps and with high power circuit blocks, there can be interference between one pump and the other blocks. Because lower voltage LSIs have larger sensitivity of power ground

noises on performance in terms of speed and variation, the pump behavior model provides high quality on system design. Section 7.5 presents a soft-switching pump model working together with a pump regulator model to avoid a hard-switching for faster system simulation. The soft-switching pump model includes  $I_{DD}$  calculation so that one can get the total  $I_{DD}$  waveform in entire simulation period. Section 7.6 presents system and circuit design and verification procedures using several models to meet the requirement for the system.

## 7.1 Hard-Switching Pump Model

Figure 7.2a shows a high-voltage generator composed of a charge pump circuit and a pump regulator. The regulator detects the output voltage of the pump,  $V_{PP}$ , to output a logical signal  $flg$  to the pump. When  $V_{MON} < V_{REF}$ ,  $flg$  is high, where  $V_{MON}$  is a divided voltage and  $V_{REF}$  is a reference voltage, as shown in Fig. 7.2b. The charge pump outputs the current to the output terminal synchronizing with an input clock  $clk_{cp}$ . When  $V_{MON} > V_{REF}$ ,  $flg$  is low to stop the clock  $clk_{cp}$ . Because the charge pump is operated with a fast continuous clock,  $clk_{cp}$ , which triggers multiple events to a simulator, it takes much time to simulate any system including a pump. A nominal clock frequency is 10 MHz to 1 GHz depending on the voltage conversion ratio or on the technology node.

To reduce the simulation time, especially for a voltage generator system, a modeled pump is used, as shown in Fig. 7.3a, where  $R_{PMP}$  is the effective output resistance of the pump as a function of the clock frequency, the number of stages, and the capacitance of the pump capacitor,  $C_{PMP}$  is the effective internal capacitance to be charged during the ramping period as a function of the number of stages, and the capacitance of the pump capacitor,  $V_{MAX}$  is the maximum attainable output voltage generated by the pump with no current load as a function of the voltage amplitude of the clock and the number of stages, and  $V_{SW}$  is a switching voltage to connect  $V_{MAX}$  to the output terminal via  $R_{PMP}$  and  $C_{PMP}$  with the enable signal  $en$  high. The pump model is disconnected from the output terminal with  $en$  low. The level shifter used in the pump model can be a standard gate-level cell, as shown in Chap. 6. The global clock  $clk$  is forced to high when the model is used for system simulations. This allows to reduce the frequency of the clock  $clk_{cp}$  as low as that of  $flg$ , as shown in Fig. 7.3b. Even though the conventional pump model doesn't require the fast continuous clock, it still needs hard-switching to connect or to disconnect the voltage source to the load synchronized with the feedback signal  $flg$ . Figure 7.3c shows the relation between the output voltage and current. The current  $I_{REG}$  continuously flows in the resistor divider whereas the current  $I_{OUT}$  discontinuously flows into the output terminal from the point p1 to p2 and vice versa. Thus, the simulation time is not fast enough to run the simulations for system-level verification.

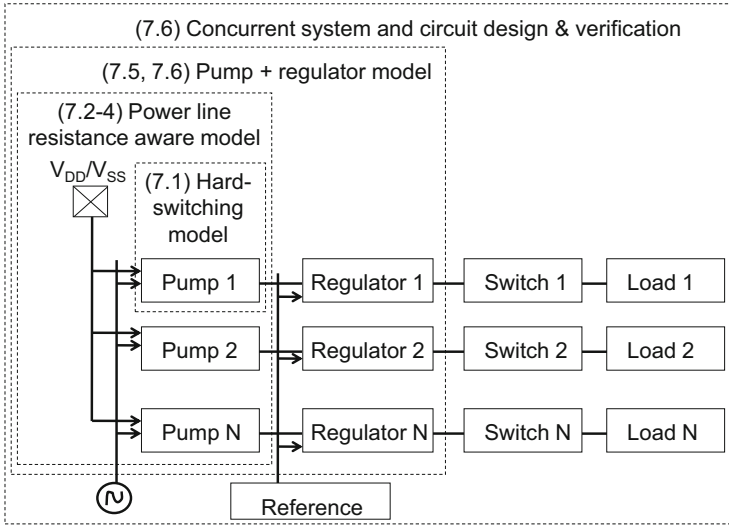
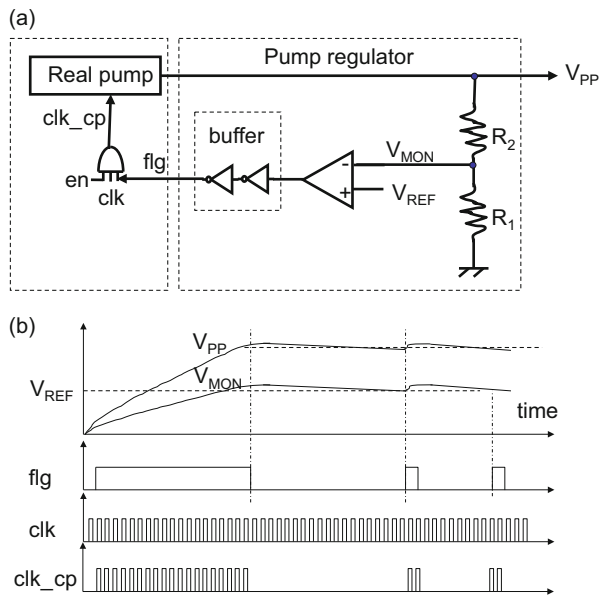


Fig. 7.1 System view and key discussion in each section

Fig. 7.2 Voltage generator composed of a pump and a pump regulator (Tanzawa 2012)



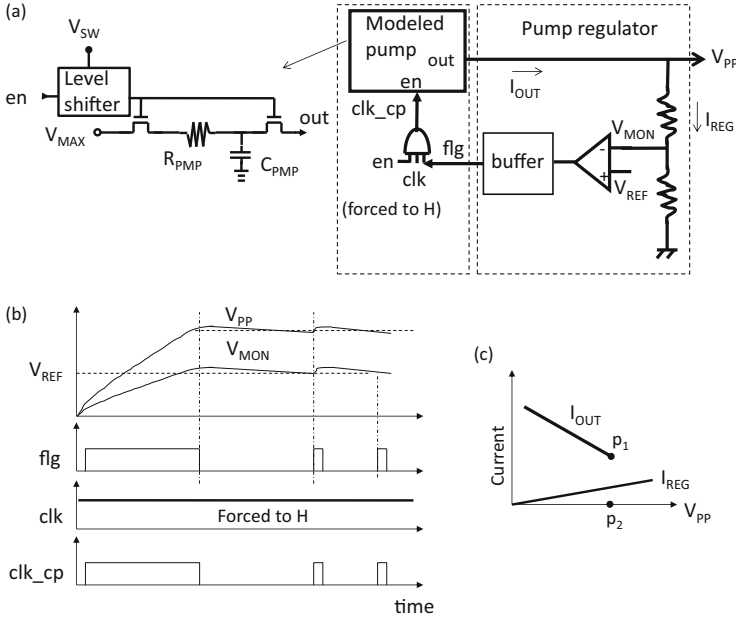


Fig. 7.3 Hard-switching pump model (Tanzawa 2012)

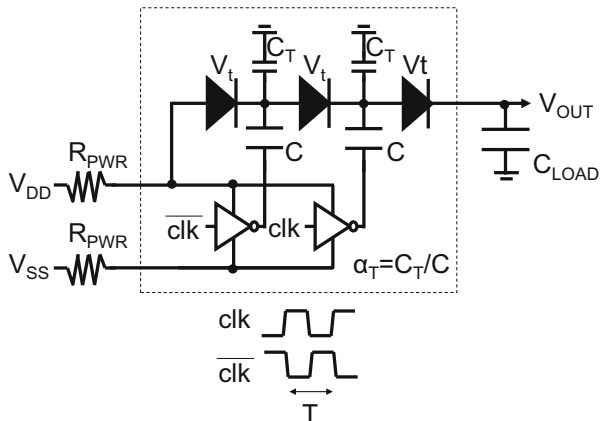
## 7.2 Power Line Resistance Aware Pump Model for a Single Pump Cell

In this section, a finite resistance in power and ground lines is taken into account in the circuit analysis as shown in Fig. 7.4. When the effect of the resistance on the pump performance is low enough to treat it as a perturbation, the amplitude of the clocks,  $V_{DD}$ , can be replaced with  $V_{DD} - 2\Delta V_{DD}$ , where  $\Delta V_{DD}$  is the voltage drop in  $V_{DD}$  line and is assumed to be same as that in ground line, resulting in a factor of 2. This voltage drop is originated from the power supply current  $I_{DD}$  and the wiring resistance  $R_{PWR}$ . Since the former is expressed by  $I_{OUT}/E_{FF}$ , where  $E_{FF}$  is the current efficiency of  $I_{OUT}$  to  $I_{DD}$ ,  $\Delta V_{DD}$  is expressed by  $R_{PWR}I_{OUT}/E_{FF}$ . Approximating a current efficiency in steady state  $E_{FF}$  with  $1/(N + 1)$ , the clock amplitude needs to be replaced with Eq. (7.1).

$$V_{DD} \rightarrow V_{DD} - 2R_{PWR}I_{OUT}(N + 1) \tag{7.1}$$

Following the similar process in Sect. 7.2, the Dickson I-V equation and dynamic behavior of  $V_{OUT}$  are respectively modified by

**Fig. 7.4** Circuit including a pump with power line resistance



$$I_{\text{OUT}} = \left(1 + \frac{2C(N+1)R_{\text{PWR}}}{NT}\right)^{-1} \frac{(1 + \alpha_T)C}{NT} \times \left(N \left(\frac{V_{\text{DD}}}{1 + \alpha_T} - V_T\right) + (V_{\text{DD}} - V_T) - V_{\text{OUT}}\right) \quad (7.2)$$

$$V_{\text{OUT}}(j) = N \left(\frac{V_{\text{DD}}}{1 + \alpha_T} - V_T\right) + (V_{\text{DD}} - V_T) - N \left(\frac{V_{\text{DD}}}{1 + \alpha_T} - V_T\right) \gamma^j \quad (7.3)$$

where

$$\gamma = \left(1 + \frac{1}{1 + \frac{2C(N+1)R_{\text{PWR}}}{T}} \frac{(1 + \alpha_T)C}{NC_T}\right)^{-1} \quad (7.4)$$

The rise time is modified by

$$T_R = T \ln \left(1 - \frac{V_{\text{PP}} - V_{\text{DD}} + V_T}{\frac{V_{\text{DD}}}{1 + \alpha_T} - V_T}\right) / \ln(\gamma) \quad (7.5)$$

Equation (7.3) indicates that the equivalent circuit model parameters are respectively given by

$$V_{\text{MAX}} = N \left(\frac{V_{\text{DD}}}{1 + \alpha_T} - V_T\right) + (V_{\text{DD}} - V_T) \quad (7.6)$$

$$R_{\text{PMP}} = \frac{NT}{(1 + \alpha_T)C_{\text{EFF}}} \quad (7.7)$$

**Table 7.1** Design parameters of three pumps for model verification (Tanzawa 2009)

	$N$	$C$ (pF)	$V_{PP}$ (V)
Pump (a)	2	90	3.5
Pump (b)	5	17	6
Pump (c)	18	8	20

$$C_{\text{EFF}} = C/\eta \quad (7.8)$$

$$\eta = 1 + \frac{2C(N+1)R_{\text{PWR}}}{T} \quad (7.9)$$

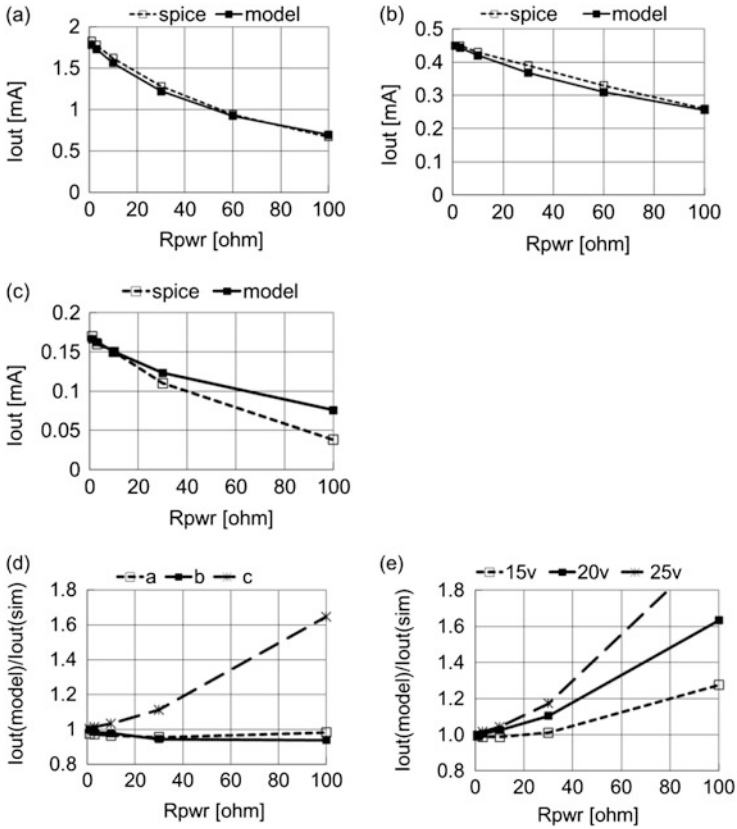
$$C_{\text{PMP}} = \frac{NC}{3}(1 + \alpha_T) \quad (7.10)$$

The difference in the parameters from those with no  $R_{\text{PWR}}$  is that the effective pump capacitor in  $R_{\text{PMP}}$  is reduced by a factor of  $\eta$  given by Eq. (7.9), resulting in an increase in  $R_{\text{PMP}}$ . On the other hand,  $V_{\text{MAX}}$  is unchanged from the case with no  $R_{\text{PWR}}$ . This means that an optimum capacitance per stage  $C_{\text{OPT}}$  needs to be increased as  $R_{\text{PWR}}$  is increased whereas an optimum number of stages  $N_{\text{OPT}}$  doesn't need to be increased no matter what optimization is done.

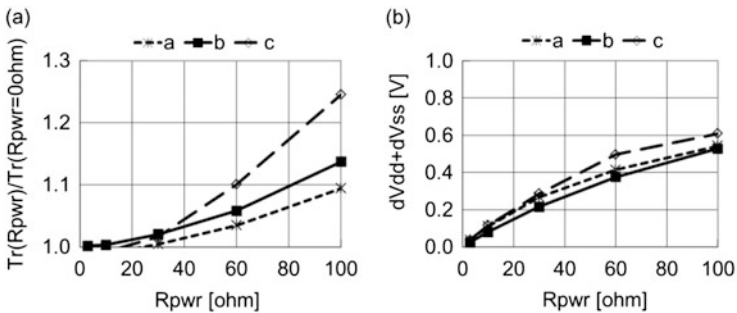
In order to verify the validity of the analysis, SPICE simulations for three different charge pumps shown in Table 7.1 were done under common conditions of  $V_{\text{DD}} = 2.5$  V,  $V_{\text{T}} = 0.35$  V,  $\alpha_{\text{T}} = 0.05$ , and  $T = 50$  ns. Figure 7.5a–c show the comparisons of the output current given by Eq. (7.2) with the simulated results of three pumps listed in Table 7.1. Figure 7.5d summarizes the comparison results. Among the data points, Pump (c) with  $R_{\text{PWR}}$  of 100  $\Omega$  shows a large discrepancy between the simulated and calculated output currents. In order to investigate the discrepancy, the output current at  $V_{\text{PP}}$  of 15 and 25 V are additionally compared in Fig. 7.5e. As  $V_{\text{PP}}$  increases, the body effect of the pass transistors increases. In this case,  $V_{\text{T}}$  of the model needs to be increased accordingly, especially with high  $R_{\text{PWR}}$  or lower effective clock amplitude. This will determine the limitation to be able to apply the model.

Figure 7.6a illustrates the discrepancy between the calculated rise time with Eq. (7.5) and the simulated one. Figure 7.6b shows the average voltage drop of the clock amplitude. A discrepancy of 2–5 % occurs at  $R_{\text{PWR}}$  of 40  $\Omega$  in Fig. 7.6a, where the voltage drop is 0.25 V or more in Fig. 7.6b. This indicates that the analysis made in this paper is in good agreement with the simulation, with less than 10 % discrepancy as long as  $R_{\text{PWR}}$  drops the clock amplitude by 10 % of  $V_{\text{DD}}$ .

Figure 7.7 shows the impact of  $R_{\text{PWR}}$  on  $C_{\text{OPT}}$ . As a factor  $C(N+1)$  increases, the increase rate  $\eta$  given by Eq. (7.9) also increases. In other words, one needs to design the charge pump circuits and/or the power line resistance so as to meet the following equation, in order to ensure that the effect of  $R_{\text{PWR}}$  on the pump performance is negligibly small.



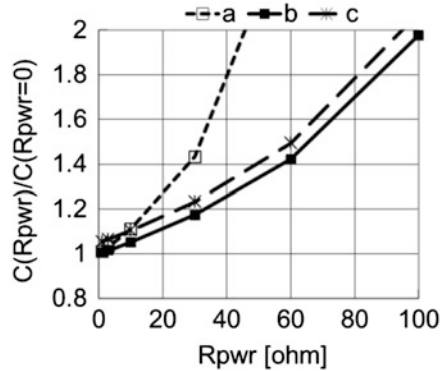
**Fig. 7.5** Comparisons of the output current given by (Fig. 7.2) with the simulated results of three pumps (a–c) listed in Table 5.1. (d) Summary of the comparison results. (e) Errors in the output current of Pump (c) (Tanzawa 2009)



**Fig. 7.6** Dependency of the rise time on  $R_{PWR}$  (a) and the voltage drop in the clock amplitude (b) under the same conditions as Fig. 5.5 (Tanzawa 2009)



**Fig. 7.7** Impact of  $R_{PWR}$  on  $C_{OPT}$  (Tanzawa 2009)



$$R_{PWR} \ll \frac{T}{2C(N + 1)} \tag{7.11}$$

When the design violates Eq. (7.11), one needs to increase the pump capacitor by a factor of  $\eta$  given by Eq. (7.9) with the number of stages unchanged to meet the requirement for the design. This becomes more important especially in lower supply voltage LSIs because  $C$  and  $N$  tend to increase at lower  $V_{DD}$  conditions.

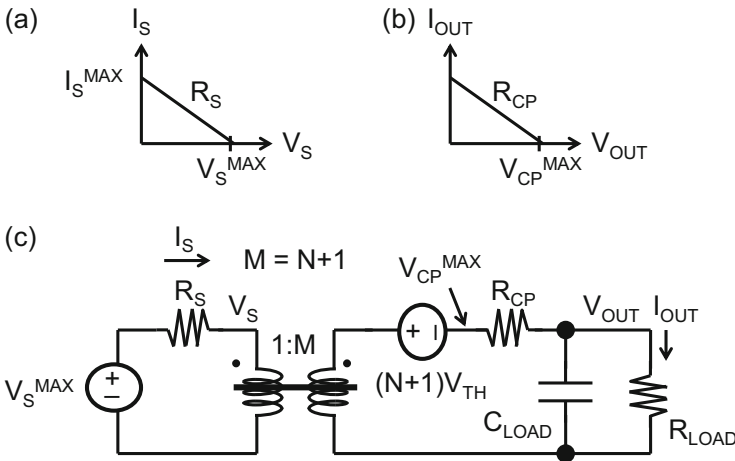
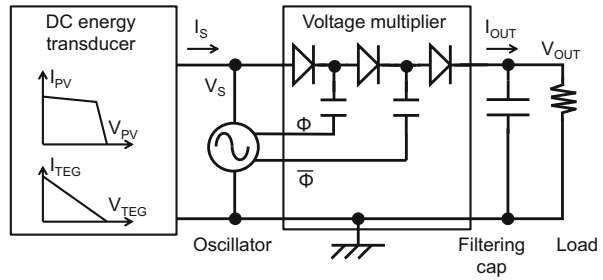
### 7.3 Optimum Design for a Given Power Line Resistance

Wireless sensing nodes, which have attracted the interest of several researchers and engineers, also use rectifier voltage multipliers for energy harvesting. These applications require low power (typically no higher than 1 mW) and have small form factors—features that are well-matched with the features of voltage multipliers. Some applications use DC energy transducers like photovoltaic (PV) and thermoelectric generator (TEG).

This section describes design optimization to maximize the power efficiency of a system that contains a DC energy transducer with high output impedance and a DC–DC voltage multiplier as shown in Fig. 7.8. An equivalent circuit model is presented to help identify the relationship between the circuit parameters and the output voltage/current characteristics of the system. A maximum power point is then identified. When the parasitic capacitance in the charge pump is significant, optimization can be expanded.

Figure 7.9 shows the output voltage and current relationships of (a) a DC energy transducer—where  $I_S^{MAX}$  is the short circuit current or the maximum output current when the output voltage is grounded,  $V_S^{MAX}$  is the open circuit voltage or the maximum output voltage when the output current is zero, and  $R_S$  is the output impedance—and (b) a charge pump—where  $V_{CP}^{MAX}$  is the maximum output voltage when the output current is zero and  $R_{CP}$  is the output impedance. The

**Fig. 7.8** DC–DC voltage multiplier supplied by a DC energy transducer, such as photovoltaic (PV) and thermoelectric generator (TEG). (Tanzawa 2014)



**Fig. 7.9** Output voltage/current relationships of (a) DC energy transducer, (b) charge pump, and (c) DC equivalent circuit for the system (Tanzawa 2014)

output impedance of PV is not linear throughout the entire output voltage range; however, it can be treated as linear around an operating point. Conversely, the output impedance of TEG can be treated as linear throughout the entire output voltage range. To simplify the discussion, first, every parasitic capacitance is neglected.

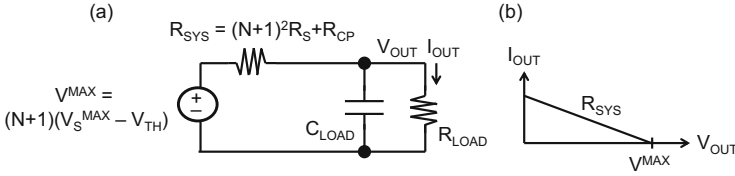
From Fig. 7.9a, b,

$$I_S = (V_S^{\text{MAX}} - V_S) / R_S \tag{7.12}$$

$$I_{\text{OUT}} = (V_{\text{CP}}^{\text{MAX}} - V_{\text{OUT}}) / R_{\text{CP}} \tag{7.13}$$

where

$$R_{\text{CP}} = TN / C \tag{7.14}$$



**Fig. 7.10** (a) Simplified DC equivalent circuit and (b) output voltage/current characteristics of the system (Tanzawa 2014)

$$V_{\text{CP}}^{\text{MAX}} = (N + 1)(V_S - V_{\text{TH}}) \quad (7.15)$$

$N$  is the number of stages or capacitors,  $C$  is the capacitance of each multiplier capacitor,  $T$  is the clock period, and  $V_{\text{TH}}$  is the threshold voltage of the switching diode. In a steady state, the output current of the transducer  $I_S$  is equally divided between the input of the first diode and  $N$  capacitors. Therefore,

$$I_S = (N + 1)I_{\text{OUT}} \quad (7.16)$$

The DC equivalent circuit of the system, which is consistent with Eqs. (7.12)–(7.16), is shown in Fig. 7.9c using a transformer symbol. When reducing the transformer using a simplification method, the simplified DC equivalent circuit of the system results, as shown in Fig. 7.10a, where  $V_{\text{MAX}}$  is the maximum attainable output voltage when the output current is zero, and  $R_{\text{SYS}}$  is the output impedance of the system.

$$I_{\text{OUT}} = (V_{\text{MAX}} - V_{\text{OUT}})/R_{\text{SYS}} \quad (7.17)$$

where

$$R_{\text{SYS}} = (N + 1)^2 R_S + R_{\text{CP}} \quad (7.18)$$

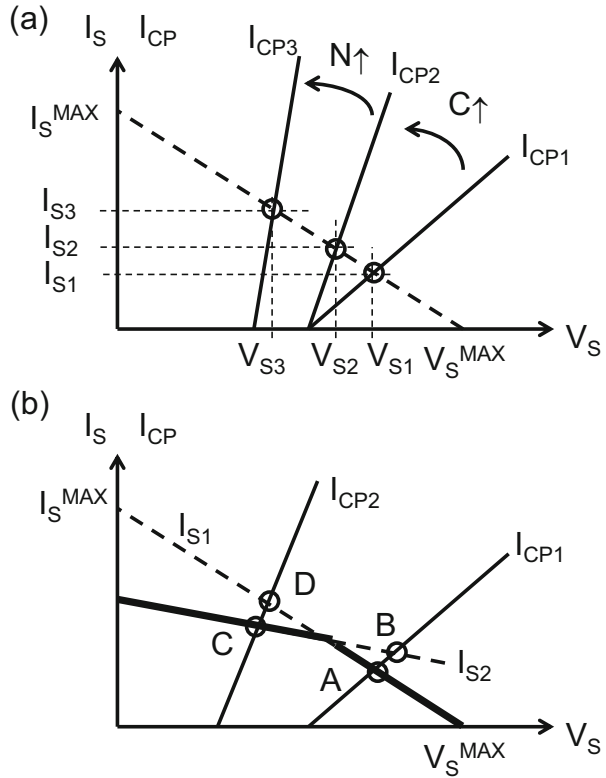
$$V_{\text{MAX}} = (N + 1)(V_S^{\text{MAX}} - V_{\text{TH}}) \quad (7.19)$$

Figure 7.10b shows the output voltage and current characteristics of the system. When one does not have any strict constraint on the charge pump circuit area, one can design the charge pump to meet the condition of  $(N + 1)^2 R_S \gg R_{\text{CP}}$  with sufficiently large pump capacitors and fast clock. In this case, Eq. (7.17) is reduced to Eq. (7.20).

$$I_{\text{OUT}} = \frac{(N + 1)(V_S^{\text{MAX}} - V_{\text{TH}}) - V_{\text{OUT}}}{(N + 1)^2 R_S} \quad (7.20)$$

From Eq. (7.20), the optimum number of stages is given by Eq. (7.21).

**Fig. 7.11** (a) Output voltage/current relationships of (a) TEG and (b) PV (Tanzawa 2014)



$$N_{OPT} = \frac{2V_{OUT}}{V_S^{MAX} - V_{TH}} - 1 \tag{7.21}$$

Equation (7.21) gives the same result as the case where the output impedance of the power source is sufficiently small and  $V_S^{MAX}$  is replaced with the power supply voltage. When the charge pump is designed to have Eq. (7.21), the output voltage and current of the transducer are respectively given by Eqs. (7.22) and (7.23).

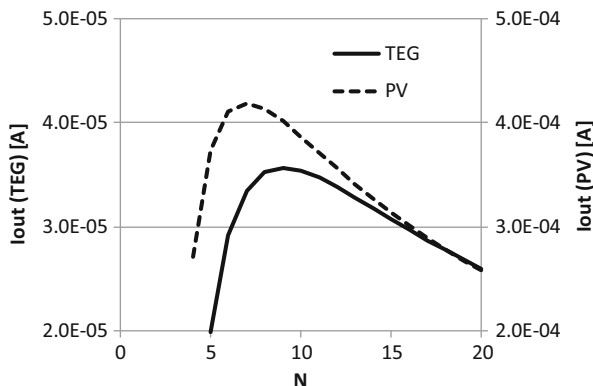
$$V_{S\_OPT} = (V_S^{MAX} + V_{TH})/2 \tag{7.22}$$

$$I_{S\_OPT} = (V_S^{MAX} - V_{TH})/(2R_S) \tag{7.23}$$

Figure 7.11 shows operating points in terms of  $I_S$  and  $V_S$  when the transducer is TEG (a) and PV (b). As  $C$  or  $N$  increases, the input current to the charge pump also increases from  $I_{S1}$  to  $I_{S2}$  or from  $I_{S2}$  to  $I_{S3}$ , respectively. When PV is used as the DC transducer, one needs to calculate Eqs. (7.22) and (7.23) using two different  $V_S^{MAX}$  and  $R_S$  values for  $I_{S1}$  and  $I_{S2}$ , as shown in Fig. 7.11b, where  $I_{S1}$  and  $I_{S2}$  are linearized lines in low output current and low output voltage ranges, respectively. When the input current of the pump is  $I_{CP1}$ , the system has two possible operating points,

**Table 7.2** Linearized model parameters (Tanzawa 2014)

	TEG	PV	
		$I_{S1}$	$I_{S2}$
$I_S^{\text{MAX}}$	1.0 mA	1.0 mA	8.8 mA
$R_S$	700 $\Omega$	10 k $\Omega$	94 $\Omega$

**Fig. 7.12** Output current as a function of  $N$  given by Eq. (7.20) (Tanzawa 2014)

$A$  and  $B$ . One needs to consider  $A$  rather than  $B$  because  $A$  is the actual operating point. One can select the actual operating point by selecting a lower current state in  $A$  and  $B$ .

When linearized model parameters are given by the values shown in Table 7.2, and  $V_{\text{TH}} = 0.2$  V, one can calculate Eq. (7.20) to draw the output current as a function of  $N$ , as shown in Fig. 7.12. In this case, Eq. (7.21) estimates the optimum number of 9 and 7 for TEG and PV, respectively, to provide the maximum output current. These values are well matched with the optimum numbers extracted from Fig. 7.12.

Next, this section discusses the case where the parasitic capacitance in the charge pump is not small. It is assumed that the current to the oscillator is sufficiently small. If needed, one can replace  $I_S$  with  $I_S - I_{\text{OSC}}$ , where  $I_{\text{OSC}}$  is the DC current of the oscillator.  $I_{\text{OUT}}$  and  $I_S$  of a charge pump are given by Eqs. (2.75) and (2.94), respectively, which are replicated as Eqs. (7.24) and (7.25).

$$I_{\text{OUT}} = \frac{C(1 + \alpha_T)}{TN} \left[ \left( \frac{N}{1 + \alpha_T} + 1 \right) V_S - (N + 1)V_{\text{TH}} - V_{\text{OUT}} \right] \quad (7.24)$$

$$I_S = (N + 1)I_{\text{OUT}} + \alpha_T C (V_{\text{OUT}} - V_S + (N + 1)V_{\text{TH}}) / T + N\alpha_B C V_S / T \quad (7.25)$$

where  $\alpha_T$  and  $\alpha_B$  are the ratios of the parasitic capacitance at the top ( $C_T$ ) and bottom ( $C_B$ ) plates of the pump capacitance, i.e.,  $\alpha_T = C_T/C$  and  $\alpha_B = C_B/C$ . From Eqs. (7.12), (7.24), and (7.25), the output current is expressed by Eq. (7.26), only including the circuit design parameters,

$$I_{\text{OUT}} = \frac{(N+1)(V_S^{\text{MAX}} - V_{\text{TH}}) - V_{\text{OUT}} - \delta_1}{(N+1)^2 R_S + TN/C - \delta_2} \quad (7.26)$$

where  $\delta_1$  and  $\delta_2$  are the second order factors including  $\alpha_T$  and  $\alpha_B$ .

$$\delta_1 = \alpha_T N V_S^{\text{MAX}} + N(\alpha_T + \alpha_B)(V_{\text{OUT}} + (N+1)V_{\text{TH}})CR_S/T \quad (7.27)$$

$$\delta_2 = \alpha_T(2N+1)R_S + \alpha_T T/C - \alpha_B N^2 R_S \quad (7.28)$$

For a given  $N$ , when  $C$  is designed too small,  $I_{\text{OUT}}$  will also be small because the charge transferred from one stage to the next is limited by  $C$ ; however, when  $C$  is designed too large,  $I_{\text{OUT}}$  will also be small because the power efficiency is reduced to charge large parasitic capacitance. Thus, an optimum  $C$  must be identified. One can rewrite Eq. (7.26) as Eq. (7.29) using Eqs. (7.30)–(7.33).

$$I_{\text{OUT}} = (A - BC)/(D/C + E) \quad (7.29)$$

$$A = (N+1)I_S^{\text{MAX}} - \frac{(N+1)V_{\text{TH}} + V_{\text{OUT}}}{R_S} - N\alpha_T I_S^{\text{MAX}} \quad (7.30)$$

$$B = N(V_{\text{OUT}} + (N+1)V_{\text{TH}})(\alpha_T + \alpha_B)/T \quad (7.31)$$

$$D = TN(1 - \alpha_T)/R_S \quad (7.32)$$

$$E = (N+1)^2 - N(N+2)\alpha_T + N^2\alpha_B \quad (7.33)$$

From  $\partial I_{\text{OUT}} / \partial C = 0$ , the optimum  $C$  needs to meet the second order equation (7.34).

$$BEC^2 + 2BDC - AD = 0 \quad (7.34)$$

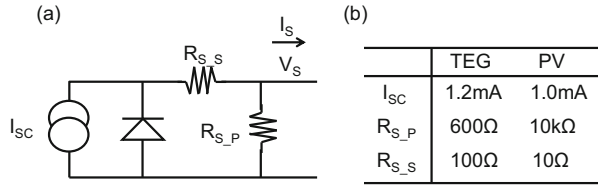
One can identify the optimum  $C$  for a given  $N$  by numerically solving Eq. (7.34).

In the following part, the model is verified by comparing the calculated results based on the model with SPICE simulation results. Figure 7.13a illustrates a model for a DC transducer; its circuit parameters used in simulations are shown in Fig. 7.13b. Figure 7.14 shows the output voltage and current characteristics of the PV and TEG with SPICE and linearized models. The parameters used for model calculations and SPICE simulations are respectively summarized in Tables 7.2 and 7.3. The effective threshold voltage of the switching diode is calculated by Eq. (3.63b) which is reproduced as Eq. (7.35).

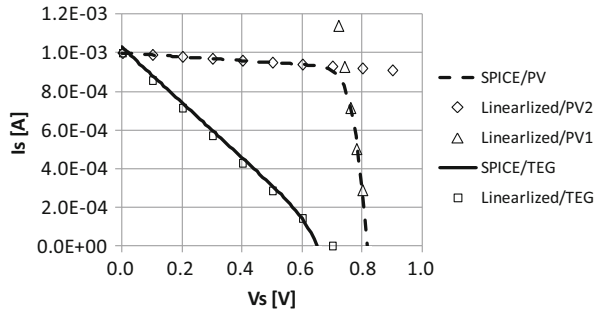
$$V_{\text{TH}} = V_T \ln \left( 4^{\frac{1}{N+1}} \frac{(1 + \alpha_T)CV_T}{T I_{S\_DIO}} \right) \quad (7.35)$$

where  $I_{S\_DIO}$  is the saturation current of switching diodes and  $V_T$  is the parameter determining the slope in  $I_{DIO} = I_{S\_DIO} \exp(V_{DIO}/V_T)$ . The diode model used in the SPICE simulation is based on a Schottky barrier diode shown in Table 7.4. In this

**Fig. 7.13** Equivalent circuit for DC energy transducer (a) and SPICE parameters used for verification (Tanzawa 2014)



**Fig. 7.14** SPICE and linearized models of PV and TEG (Tanzawa 2014)



**Table 7.3** Circuit parameters used for verification (Tanzawa 2014)

Name	Value	Unit
$T$	100	ns
$V_{OUT}$	2.5	V
$\alpha_T$	0.1	No unit
$\alpha_B$	0.05	No unit
$C$	10–100	pF
$N$	4–14	No unit

**Table 7.4** SPICE diode model parameters used for verification (Tanzawa 2014)

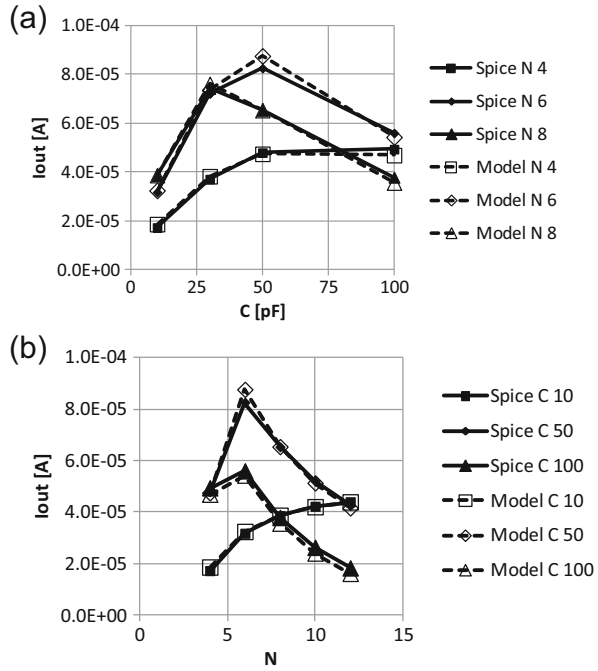
Name	Parameter	Value	Unit
Level	Junction diode model	1	No unit
$EG$	Bandgap energy	0.69	eV
$N_p$	Emission coefficient	1	No unit
$RS$	Parasitic resistance	0.1	Ω
$IS$	Saturation current	1e-8	A

case, one can fit the  $V_{DIO} - I_{DIO}$  curve of the SPICE model with  $I_{S\_DIO} = 20$  nA and  $V_T = 26$  mV. Strictly speaking, Eq. (7.34) and  $\partial I_{OUT} / \partial C = 0$  result in a higher-order equation because  $V_{TH}$  is a function of  $C$ ; however,  $V_{TH}$  can be treated as a constant as long as  $C$  does not vary too much because  $V_{TH}$  is a weak function of  $C$ .

Figure 7.15 shows  $C$  versus  $I_{OUT}$  (a) and  $N$  versus  $I_{OUT}$  (a). The model fits well with the SPICE results. The output current with a given  $N$  or  $C$  has a maximum at a different value of  $C$  or  $N$ , respectively.

Using the matrix data in terms of  $C$  and  $N$ , one can draw the contour plots as shown in Fig. 7.16a, b for the SPICE and model. Because the SPICE simulations run only with the combination of  $N = 4, 6, 8, 10, 12$ , and  $C = 10, 30, 50, 100$  pF, the

**Fig. 7.15** Output current as a function of (a)  $C$  and (b)  $N$  in case of the PV (Tanzawa 2014)



contour of the SPICE seems a little different from the model results, which run with  $N=4-14$  and  $C=10-100$  pF by 10 pF. However, the optimum parameters set for  $C$  and  $N$  is estimated to be 50 pF and 6 in both cases.

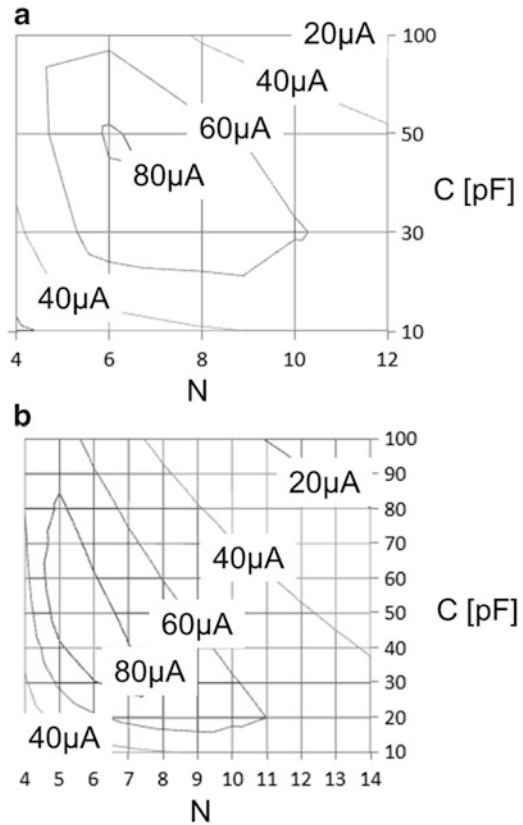
Similarly, Figs. 7.17 and 7.18 show the results for TEG. Unlike PV, the window for an optimum  $N$  is wide and an optimum  $C$  is not a strong function of  $N$ . Thus, the model can be considered well-characterized to reproduce the SPICE results without running several combinations of  $C$  and  $N$ . One can make an initial guess in terms of circuit parameters for detailed design with SPICE. To see the sensitivity of these parameters on the output power, one can draw a similar contour plot under different conditions for diode parameters, such as the saturation current and the parasitic junction capacitance as a part of the parasitic capacitance of the top plate of the pump capacitor, as well as for circuit parameters like the short circuit current and output impedance of the transducer, clock period, and output voltage.

## 7.4 Pump Behavior Model for Multiple Pump System

This section discusses top-down charge pump circuit design with charge pump behavior models, including not only circuit parameters such as the number of stages, the capacitance per stage, and the supply voltage, but also the parasitic power wiring resistance as shown in Fig. 7.19. System designers can determine



**Fig. 7.16** Contour plots for the PV with SPICE (a) and model (b) (Tanzawa 2014)



floor plan for replacement of individual charge pump, required power and ground width and length, and individual pump design parameters once the total area and power meet their design targets. Then, the charge pump circuit designers can start designing each pump with each design parameter determined.

In order to generalize the model for multiple charge pump circuits distributed in LSIs, one can start with the following equations;

$$V_{MAX} = (N/(1 + \alpha_T) + 1)(V_{DD\_LOCAL} - V_{SS\_LOCAL}) + V_{OS} \tag{7.36}$$

$$V_{OS} = -(N + 1)V_T \tag{7.37}$$

$$I_{DD} = I_{SS} = I_{OUT}/E_{FF} = (N + 1)I_{OUT} \tag{7.38}$$

These equations are translated into a behavior model with several elements such as a voltage controlled voltage source (*exvmax*), current controlled current sources (*fxivcc*, *fxivss*), and voltage sources (*vos*, *vxiout*), as shown in Fig. 7.20 and in

**Fig. 7.17** Output current as a function of (a) C and (b) N in case of the TEG (Tanzawa 2014)

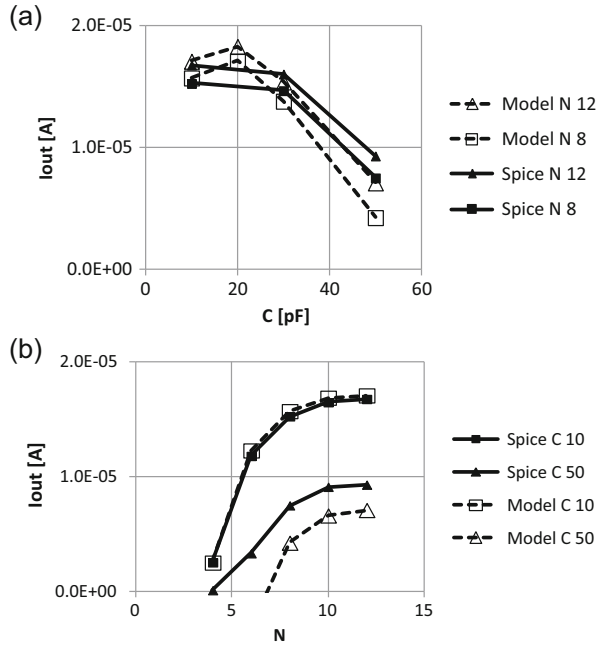
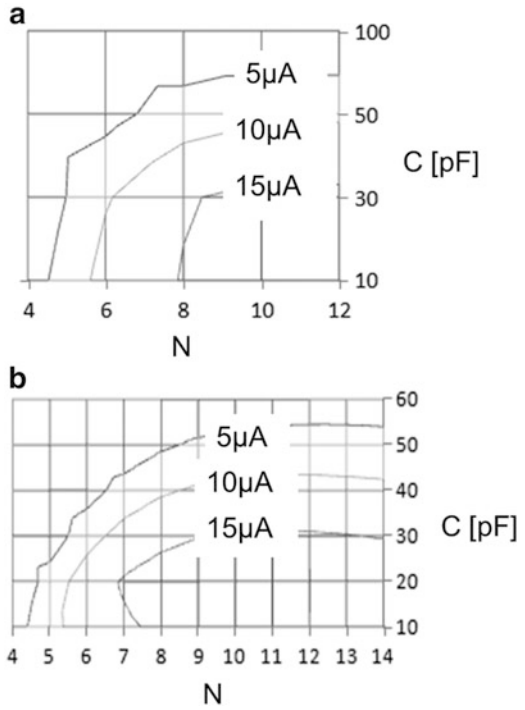


Table 7.5, where  $N$ ,  $\alpha_T$ ,  $\text{eff}$ , and  $V_T$  are design parameters; respectively the number of stages, the ratio of the parasitic capacitance at the top plate to that of the pump capacitor, the current efficiency defined by the ratio of the output current to the input current, and the voltage drop via the switching diode. The voltage controlled voltage source,  $exvmax$ , represents the first term of Eq. (7.36) and the voltage sources,  $vos$ , represents the second term of Eq. (7.36). Also, the current controlled current sources,  $fxivcc$  and  $fxivss$ , and the voltage sources,  $vxiout$ , are related each other through Eq. (7.38). Thus, the input current,  $I_{DD}$  and  $I_{SS}$ , are calculated by monitoring the output current with  $vxiout$ . This behavior model is schematically expressed by each box described in Fig. 7.21.

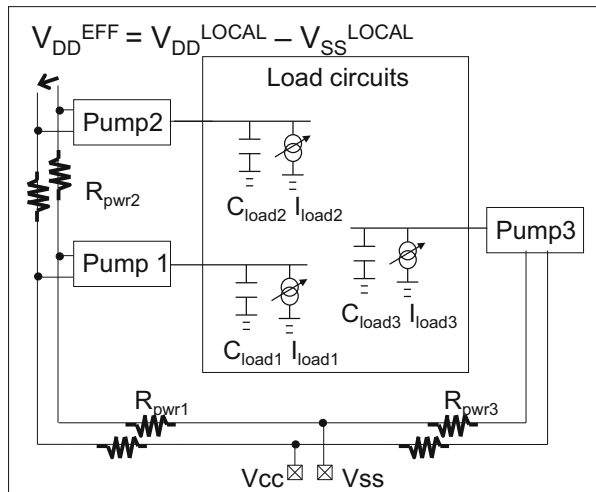
The terminal sw of Fig. 7.20 is synchronized with an output of a regulator. The charge pump and the regulator are configured to be a feedback system to stabilize the output voltage of the pump. The switching elements M1, 2 should be so ideal that their channel resistance is much lower than an output resistance of  $R_{PMP}$ . Since every charge pump can be defined by its own behavior model, it is available in a system level simulation as shown in Fig. 7.21. When each of the terminals  $V_{DD\_LOCAL}$  and  $V_{SS\_LOCAL}$  is simply connected to the parasitic resistor network for power and ground lines, it is reduced to the original model. Thus, this behavior model includes the original one.

In order to verify the behavior model and to see the impact of the common impedance of  $R_{PWR}$  on the circuit performance in the pump system described in Fig. 7.21. SPICE simulations were done together with the real pumps with gate level net list, as shown in Fig. 7.22.

**Fig. 7.18** Contour plots for the TEG with SPICE (a) and model (b) (Tanzawa 2014)



**Fig. 7.19** LSI with distributed multiple charge pump circuits



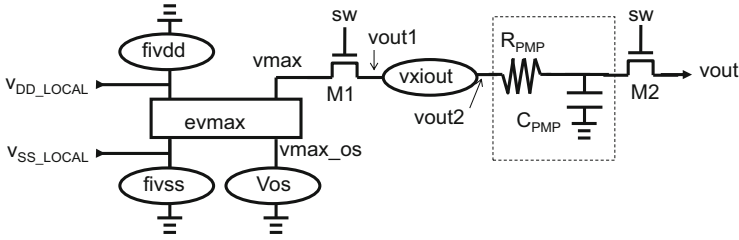


Fig. 7.20 Behavior model used in top-level design (Tanzawa 2010)

Table 7.5 Behavior model of the charge pump circuit (Tanzawa 2010)

<i>Exvmax</i>	<i>vmax</i>	<i>vmax_os</i>	<i>vdd_local</i>	<i>vss_local</i>	$(N)/(1 + \alpha_T) + 1$
<i>Fxivss</i>	<i>gnd</i>	<i>vss_local</i>	<i>vxiout</i>	$1/eff$	
<i>fxivcc</i>	<i>vdd_local</i>	<i>gnd</i>	<i>vxiout</i>	$1/eff$	
<i>vos</i>	<i>vmax_os</i>	<i>gnd</i>	<i>dc</i>	$-(N + 1)V_T$	
<i>vxiout</i>	<i>vout0</i>	<i>vout</i>	<i>dc</i>	$0 V$	

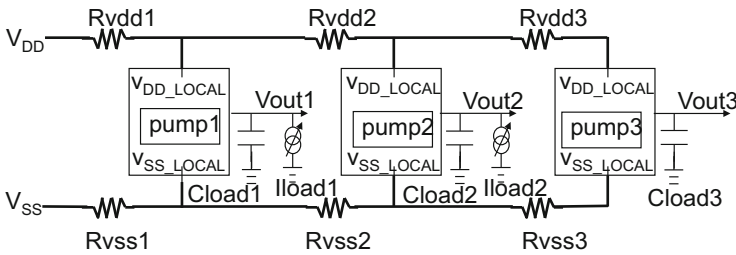
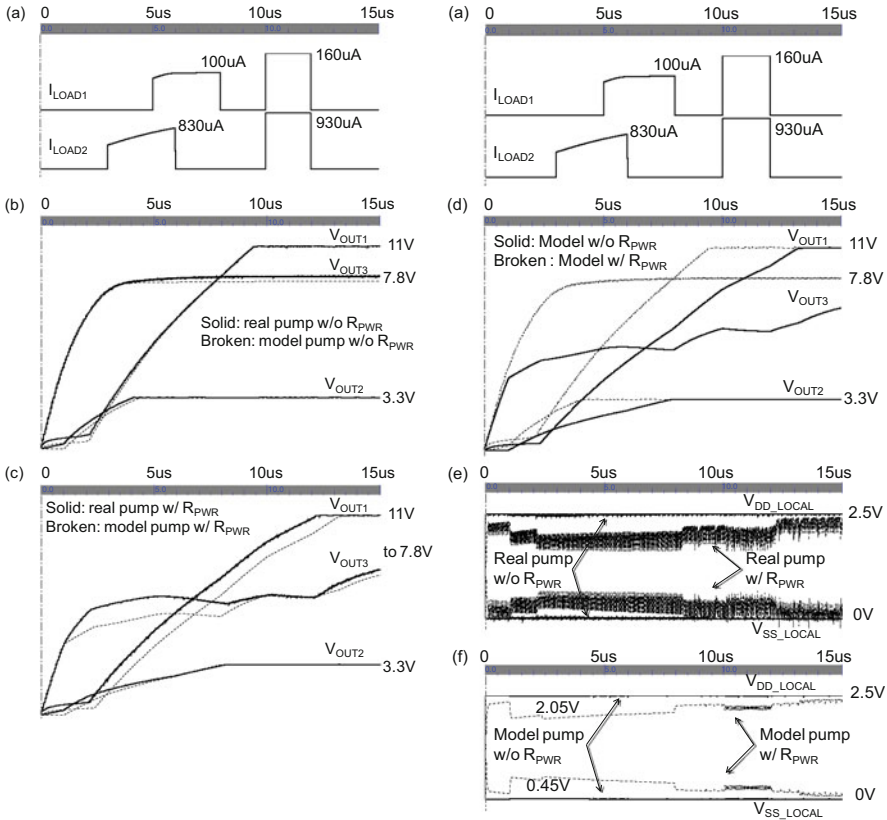


Fig. 7.21 Test bench for pump system (Tanzawa 2010)

The circuit parameters used in the simulations are shown in Table 7.6, where  $V_{MAX}$  is the maximum output voltage in case of no power and ground line resistance.  $V_{DD}$  and  $T$  are 2.5 V and 60 ns, respectively. Each pump is regulated so that  $V_{OUT1-3}$  are stabilized at 11.0, 3.3, and 7.8 V, respectively. Figure 7.22a shows the input load current waveforms. Figure 7.22b compares the modeled pumps with the real ones in the case where no power and ground wiring resistance is considered. The waveforms are in good agreement with an error of less than 3%. Figure 7.22c shows the comparison between the modeled and real pumps in the case where a finite power and ground wiring resistance is considered with the values shown in Table 7.6. Even though the simulated condition was so large that the local power and ground bounces were as high as about 0.45 V at the peak points, respectively, as shown in Fig. 7.22f, the rise time is in agreement within less than 10%. Figure 7.22d compares the waveforms between the cases with and without  $R_{PWR}$ .  $V_{OUT3}$  suffered most from the other pumps such as pump 1 and 2, which share all the power and ground lines. Figure 7.22e shows the local  $V_{DD}$  and  $V_{SS}$  at pump 3 using the real pump net list, which include high frequency components as fast as the clock frequency. The local power ground waveforms in case that the modeled pumps



**Fig. 7.22** Simulated waveform of real and modeled charge pumps with and without power ground line resistance (Tanzawa 2012)

**Table 7.6** Design parameters used for Fig. 7.21 (Tanzawa 2012)

	Pump 1	Pump 2	Pump 3
$N$	12	2	5
$C$ (pF)	50	150	25
$\alpha_T$	0.05	0.05	0.05
$V_T$ (V)	0.4	0	0.3
$C_{LOAD}$	500 pF	5 nF	10 pF
$R_{VDD}$ ( $R_{VSS}$ ) ( $\Omega$ )	12	18	12
$V_{MAX}$ (V)	36.6	7.3	13.8
$R_{PMP}$	13.7 k $\Omega$	750 $\Omega$	9.1 k $\Omega$

are used as shown in Fig. 7.22f behave filtering and averaging ones. Thus, the behavior model is shown to be accurate enough to reproduce the real pump behavior. In addition, the impact of the common impedance in power and ground lines on the pump performance was shown. The simulation time in case with the

modeled pump system was reduced to less than 1/20 of that in case with the real one, in this example.

Assumed values of the worst-case  $V_{DD\_LOCAL}$  and  $V_{SS\_LOCAL}$  are conventionally given as the input parameters such as the clock frequency for designing individual charge pump circuits. However, since the system simulation, including all the charge pump circuits and power and ground wiring resistance, is impractical with respect to the simulation time, dynamic behavior of power and ground noises is hardly reflected to the pump performance. This kind of unknown sometimes results in over design or in larger circuit than necessary.

On the other hand, by using the behavior model, one can use the power ground resistance as parameters to minimize the total area for the power ground wirings and charge pump circuits. With the power and ground line resistance extracted from an initial floor plan for the voltage generator system and load conditions given, the initial solutions for the pump design parameters, such as the number of stages and capacitance per stage, are obtained. If the resultant total area and power don't meet the requirements, the power and ground line resistance has to be updated. Under the updated condition, the circuit parameters are reduced again and checked to be fulfilled with the target values for the total area and power. Thus, the feedback between floor plan and pump design is available to minimize the total area and power. After such a top-down procedure, an individual charge pump design can be started which takes the power and ground voltage drops due to the operation currents of itself and the rest of the circuits into consideration.

## 7.5 Concurrent Pump and Regulator Models for Fast System Simulation

This section discusses modeling of the pump and the pump regulator to make the system simulation much faster than the pump models shown in Sects. 7.1 and 7.4.

Figure 7.23a illustrates models for pump and regulator. The voltage source  $V_{MAX}$  is connected to the resistor  $R_{PMP}$  via the current mirror. It is designed such that the output impedance is sufficiently small compared with  $R_{PMP}$  to keep the total impedance of the pump the same as the original model, i.e., to have  $I_{OUT1}$  in Fig. 7.23c as high as  $I_{OUT}$  in Fig. 7.3c.

In case that the open loop gain of the system is too high to make the pump plus regulator system unstable, a diode optionally needs to be added to reduce the gain, as shown in Fig. 7.23a, especially for a high-voltage generator system where the pole of the resistor divider is not quite far from that of the pump with its load included. Thus, the AC performance of the opamp usually doesn't affect the stability of the entire system. One needs to make sure that the diode added doesn't affect  $I_{OUT1}$  especially at low output voltages, which could slightly increase the output impedance. Because the opamp is one of the circuit components which compose the entire regulator block and which cannot be changed in the model, it is

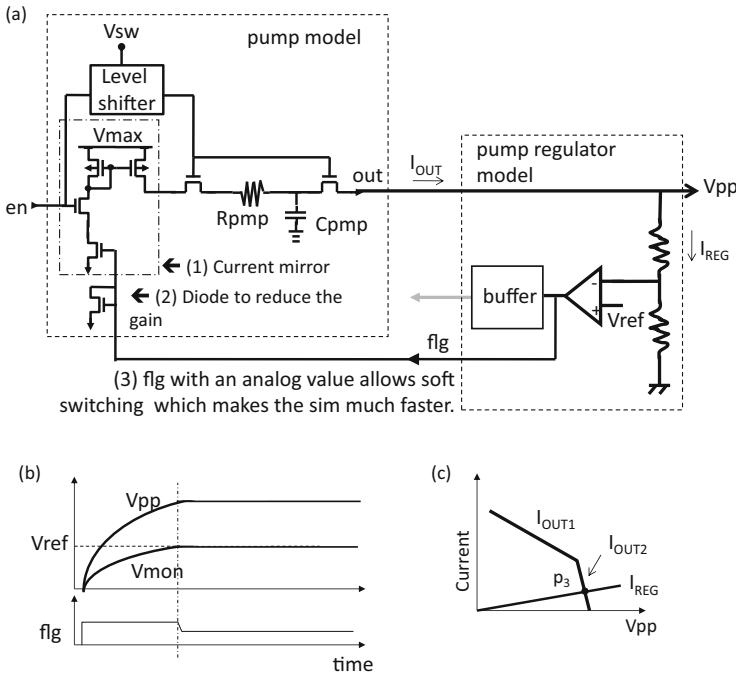


Fig. 7.23 Soft-switching pump and regulator models (Tanzawa 2012)

only the current mirror and the diode optionally that can adjust the I-V characteristics as shown by  $I_{OUT1}$  and  $I_{OUT2}$  in Fig. 7.23c. In addition, the pump regulator is modeled to convert the logic signal  $flg$  to an analog one. This is done with a simple change of the terminal from the output of the buffer to that of the opamp. The buffer is required to transfer the signal with a small slew rate for a real circuit. But, it is not required for a simulation purpose in case that the wiring parasitic resistance and capacitance are not considered. Figure 7.23b shows how the feedback signal  $flg$  behaves. At the beginning of the operation,  $flg$  is higher than the level in a stable state to output the current from the pump. When  $V_{MON}$  gets close to  $V_{REF}$ ,  $flg$  starts decreasing.  $flg$  becomes stable once the feedback system becomes stable unlike the conventional model. Figure 7.23c shows the regulation point  $P_3$  at which the pump output current  $I_{OUT}$  is balanced with the regulator current  $I_{REG}$ . Thus, every node in the loop becomes analog so that hard-switching can be fully eliminated, resulting in much faster simulation time.

Figure 7.24 explains how the cell views for the regulator are implemented into the design. The pump regulator cell has two different cell views: schematic for physical design and *pump\_model* for system simulation or verification. The output terminal  $flg$  is differently connected to an output terminal of the pump regulator core block *cpregcore*, which has two output terminals,  $flg_l$  and  $flg_a$ . Thus,  $flg$  is connected with  $flg_l$  for physical design to drive its heavy load and with  $flg_a$  for

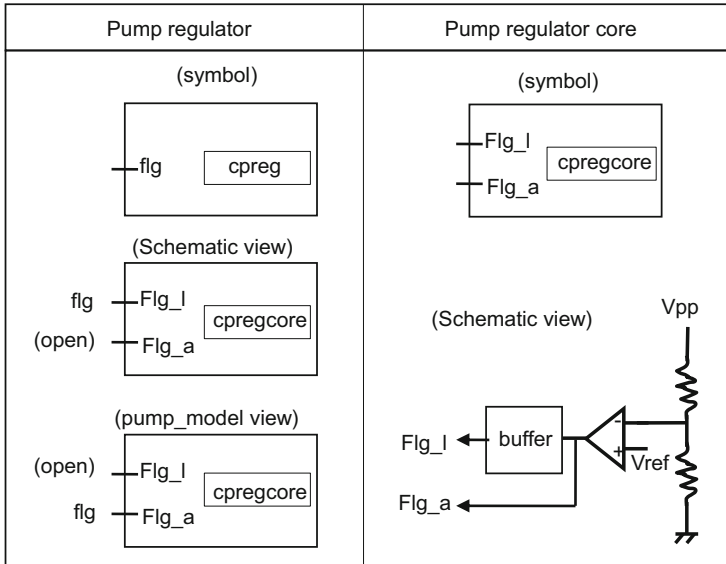


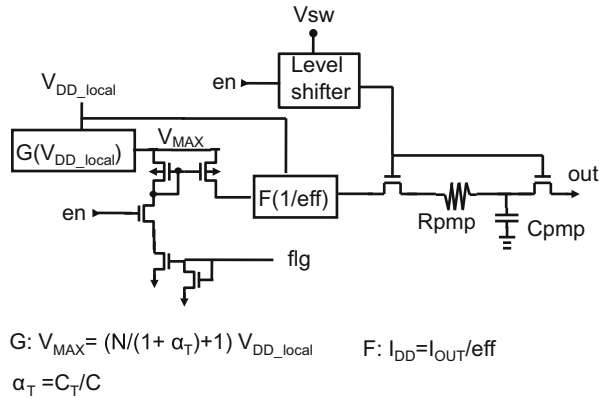
Fig. 7.24 Regulator model with *schematic* and *pump\_model* views (Tanzawa 2012)

verification to make the feedback node analog. One can generate a gate level net list using the schematic view and a net list including the model using the *pump\_model*. This approach enables us to use a single physical block for both physical design and verification. Even though there is a design update in the pump regulator core, one doesn't need to update either the schematic or *pump\_model* views of the pump regulator. Thus, this method has no risk of a potential mismatch between the real and model regulator. One drawback of the soft-switching pump model over the hard-switching pump model is that the soft-switching model doesn't reproduce any ripple in the output voltage unlike the real and hard-switching pump model pumps do. If concerns on the system-level simulations include the ripple, one needs to run some simulations with the hard-switching pump model additionally.

Another pump model is shown in Fig. 7.25 to take the impact of the power line resistance on the pump performance into account, which is represented by G and F. A voltage controlled voltage source G and a current controlled current source F are available in HSPICE and other simulators. These are combined into the pump model of Fig. 7.23, resulting in Fig. 7.25.  $V_{MAX}$  is actually a function of  $V_{DD\_LOCAL}$ , the power supply for the pump, such as  $V_{MAX} = (N/(1 + \alpha_T) + 1) V_{DD\_LOCAL}$ , where  $N$  is the number of stages,  $\alpha_T$  is the ratio of the parasitic capacitance at the top node of the pumping capacitor ( $C_T$ ) to that of the pumping capacitor ( $C$ ), and  $V_{DD\_LOCAL}$  is the local power. Using the output current  $I_{OUT}$  and the current efficiency  $eff$ , the input current  $I_{DD}$  can be given by  $I_{OUT}/eff$ , as shown by F in Fig. 7.25. When the power line resistance is added to the local  $V_{DD}$  terminal for a system simulation, the IR drop in  $V_{DD\_LOCAL}$  is reproduced self-consistently.



**Fig. 7.25** Soft-switching pump model including  $I_{DD}$  calculation (Tanzawa 2012)



**Fig. 7.26** Comparison of I-V curves between a modeled 18 V generator and a real one under various  $V_{DD}$  (Tanzawa 2012)

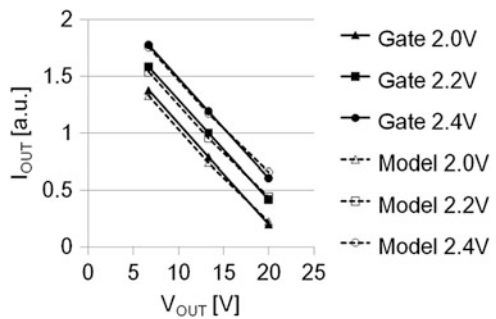


Figure 7.26 compares the  $V_{OUT} - I_{OUT}$  characteristics of a real pump and a model. Because the model shown in Fig. 7.25 includes the clock amplitude ( $V_{DD}$ ) as an input parameter, the model can have the I-V curves close to the real ones under the wide  $V_{DD}$  operation conditions with an error of 5%.

Figure 7.27 shows Bode plots of the 18 V generator composed of the pump, regulator without (a) and with (b) a diode, and current load. Adding the diode, the generator system gets stable with a phase margin of  $5^\circ - 100^\circ$ . Thus, the dimension of the diode can be adjusted according to the gain and output voltage range of the opamp given.

Figure 7.28 compares the waveforms for the output voltage  $V_{OUT}$  and the monitor nodal voltage  $V_{MON}$ .  $V_{OUT}$  is in good agreement each other, but  $V_{MON}$  is different. Because  $V_{MON}$  of the system with the diode added is much smoother than that without the diode, the former is considered to have less simulation time than the latter.

Figure 7.29 compares the waveform with a real pump and regulator with that with the modeled pump and regulator. Due to hard- and soft-switching operation with the real and modeled generator, the current waveform is sawtooth with the real

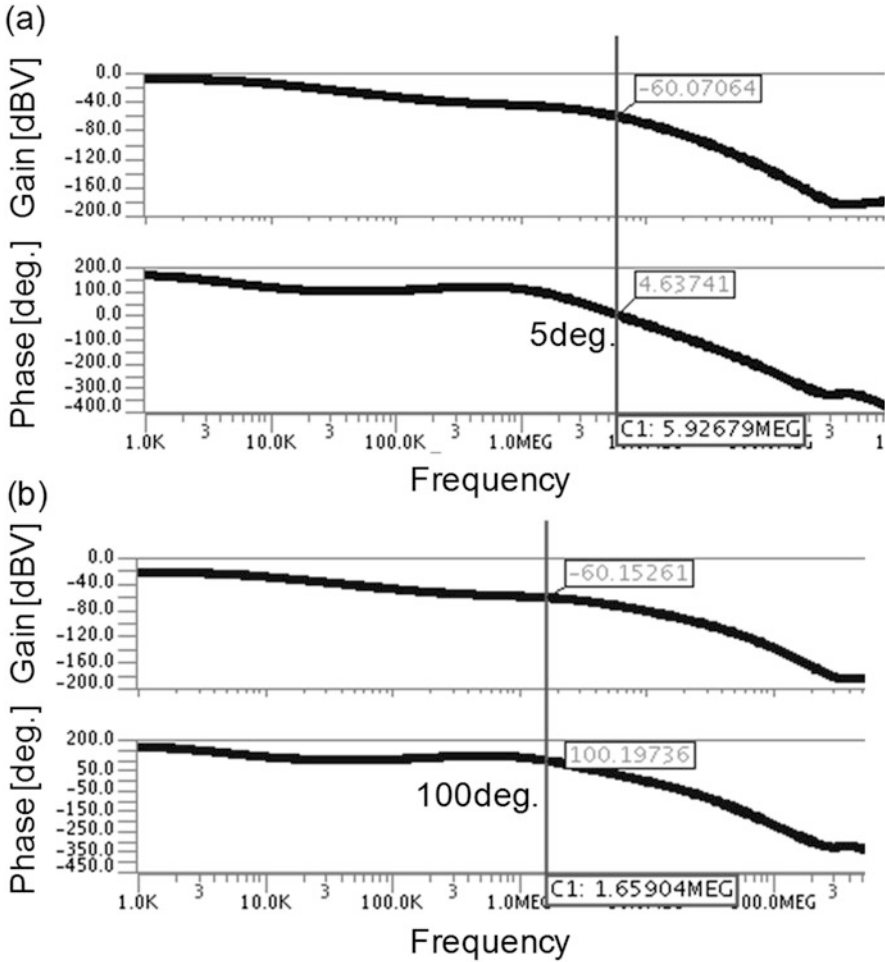
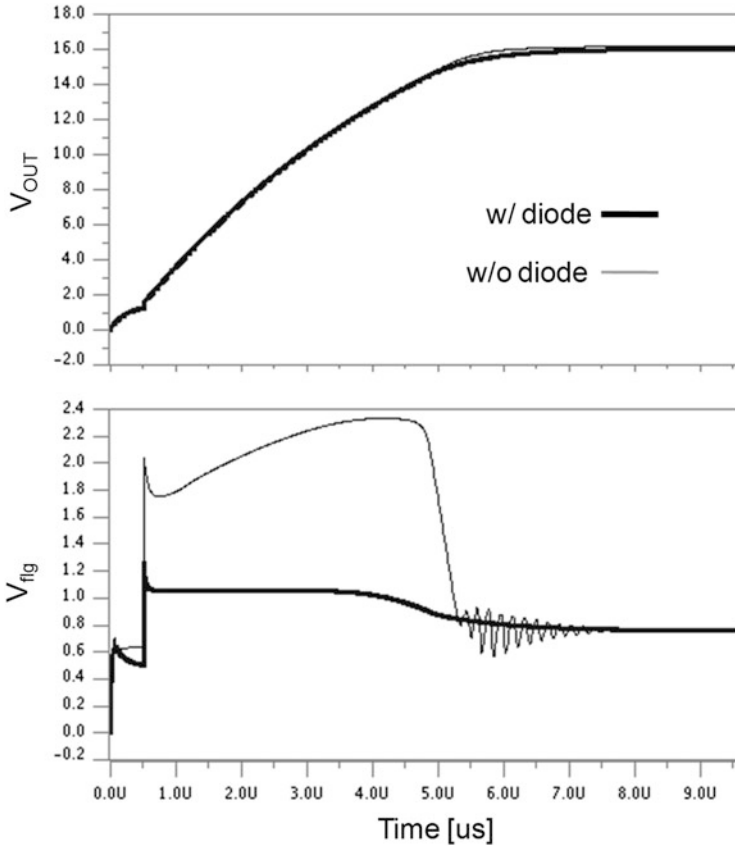


Fig. 7.27 Phase margin without (a) and with (b) a diode connected with *flg* node for the 18 V generator (Tanzawa 2012)

one, whereas smooth with the modeled one. The generator is designed to output 18 V. The voltage waveform with the models is in good agreement with the real one in spite of the different current waveforms. The HSPICE run time with the modeled generator was 75 times shorter than that with the real one for a 10  $\mu$ s transient simulation. At a sacrifice of the accuracy in high frequency components in  $I_{OUT}$  and  $I_{DD}$ , faster simulation was achieved with an error of 5% in  $V_{OUT}$ . As far as the voltage waveform is concerned, the accuracy seems to be enough. Note that the reduction rate depends on the simulator used as well as the simulation net list and simulation period of time.

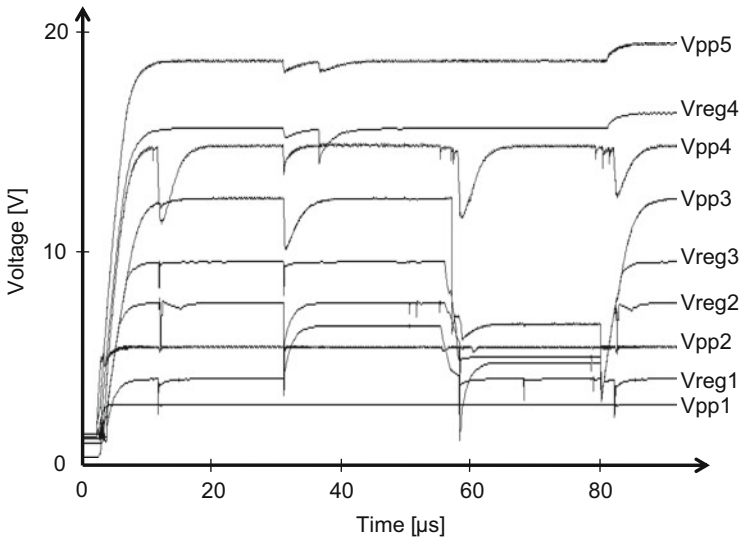
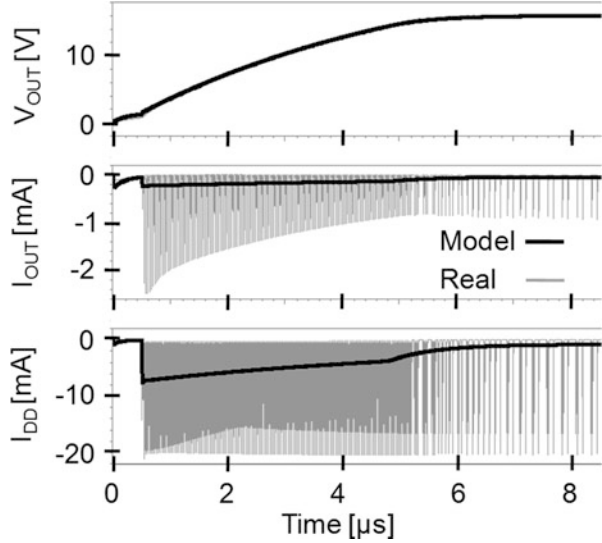


**Fig. 7.28** Transient waveforms with and without the diode added to the modeled 18 V generator (Tanzawa 2012)

Figure 7.30 shows nine simulated waveforms of five output voltages generated by five pumps,  $V_{PP1-5}$ , and four regulated voltages regulated from the pump outputs,  $V_{REG1-4}$ , for a programming operation in 200  $\mu\text{s}$ . The simulated net list includes not only voltage generators, but also switches and loads in NAND Flash memory. The number of devices in the net list is about 50 k. To validate the effectiveness of the models on the system-level simulation time, mixed-signal simulations were done.

Figure 7.31 compares the simulation time for the generator system with the number of devices at 30 k, which only includes the voltage generator, and the full-chip with the number of devices at 50 k between the cases with real generators, the hard-switching modeled ones, and the soft-switching ones. Regarding the voltage generator system, the hard-switching and soft-switching models reduced the simulation time by about 5 and 75, respectively, in comparison with the gate level net list. Regarding the full-chip, the soft-switching model reduced the simulation time

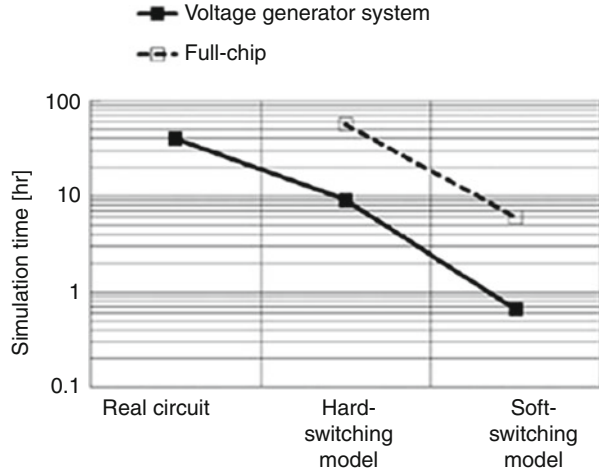
**Fig. 7.29** Comparison of a modeled 18 V generator with a real one (Tanzawa 2012)



**Fig. 7.30** Full-chip simulation waveforms (Tanzawa 2012)

by about 10 in comparison with the hard-switching model. The reduction rate depends on the simulation net list and simulation period of time, but in this example about  $\times 10$  reduction in simulation time was realized with the soft-switching model compared with the hard-switching one.

**Fig. 7.31** Comparison in mixed-signal simulation time between real pumps; hard-switching models; and soft-switching models (Tanzawa 2012)



## 7.6 System Design Methodology

Figure 7.32 shows a design flow for an on-chip high-voltage generator system. One has design requirements for the system such as the output current  $I_{OUT}$  or the rise time  $T_R$  for each voltage source and the total area and the peak and average operation current  $I_{DD}$  for entire voltage sources under the power line resistance  $R_{PWR}$  assumed (Step 1). Pump design parameters such as the clock period  $T$ , each capacitor  $C$ , and the number of stages  $N$  are determined using the design formulas (Step 2). Figure 7.33 shows a flow in Step 2 of entire flow shown in Fig. 7.32 in case where the power line resistance needs to be taken into consideration for low voltage ICs. In addition to the load conditions given by current, resistive, and capacitive load, the power line resistance is taken as a design parameter. Once the design parameters of the number of stages and the capacitance per stage is determined per pump which entirely meet the area and power budget, one can proceed Step 3 of Fig. 7.32 to start physical design. At Step 3, gate level design is done with the schematic and layout for each pump and regulator. In parallel, the pumps are modeled as the hard-switching models (model 1) and the soft-switching ones (model 2). The gate level design is verified with respect to the pump I-V characteristic for each pump block and to the transient simulation for combination of each pump, regulator, and load. The system level design is verified with respect to the voltage ripple with model 1 and to the output current, the rise time, and the operation current with model 2 (Step 4).

When all the simulation results meet the original target, all the on-chip high-voltage system design and their component design are completed. Otherwise, one may need to update some of the original targets because there could be

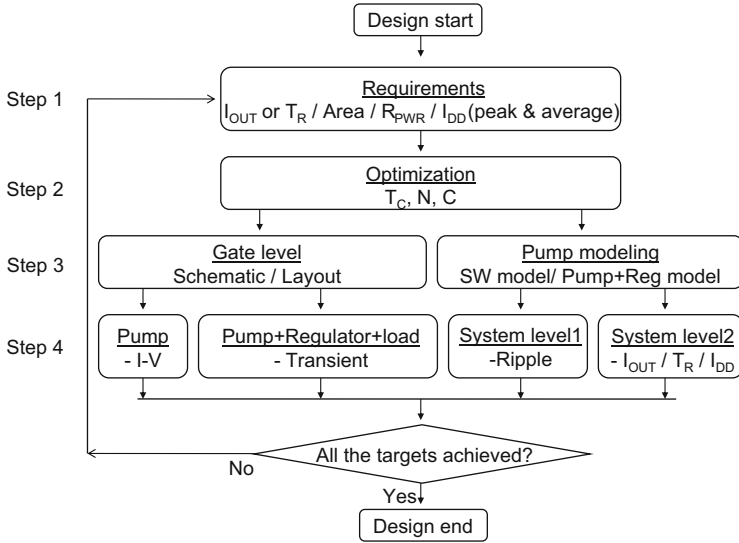


Fig. 7.32 Design and verification flow

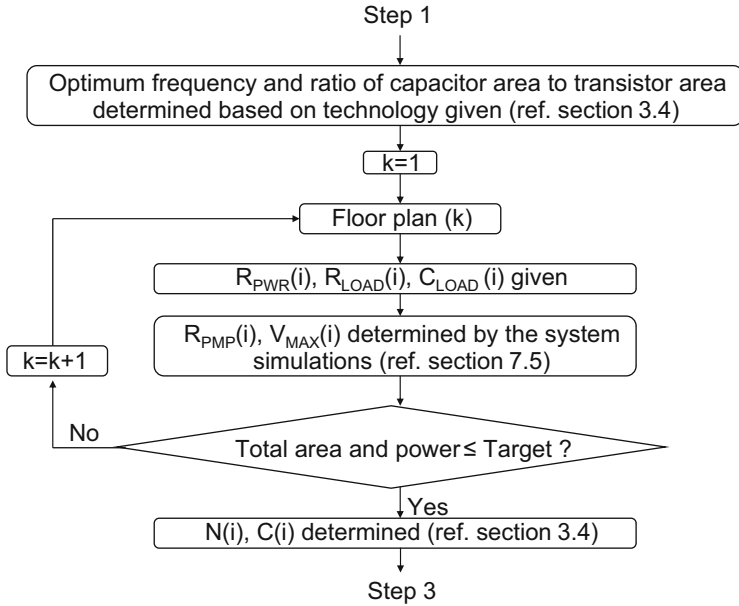


Fig. 7.33 Flow in Step 2 of Fig. 7.32 for an individual cell considering the power line resistance

inconsistency between the design parameters. The verification categorized into “system level 2” conventionally takes more time than the rest in Step 4. Therefore, the soft-switching model can reduce the time for Step 4, resulting in faster entire design and verification periods.

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