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Khac Duc Do • Jie Pan

# Control of Ships and Underwater Vehicles 

Design for Underactuated and Nonlinear Marine Systems

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ISSN 1430-9491
ISBN 978-1-84882-729-5 e-ISBN 978-1-84882-730-1
DOI 10.1007/978-1-84882-730-1
Springer Dordrecht Heidelberg London New York
British Library Cataloguing in Publication Data
A catalogue record for this book is available from the British Library
Library of Congress Control Number: 2009931136
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Cover design: eStudioCalamar, Figueres/Berlin
Printed on acid-free paper
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The first author dedicates this book to his parents, Do Thi Duyen and Do Khac Yen, and his little daughter, Do Thu Trang.

The second author dedicates this book to his parents, Liyou Pan and Hungxiu Diao.

## Series Editors’ Foreword

The series Advances in Industrial Control aims to report and encourage technology transfer in control engineering. The rapid development of control technology has an impact on all areas of the control discipline. New technology, new controllers, actuators, sensors, new industrial processes, computer methods, new applications, new philosophies..., new challenges. Much of this development work resides in industrial reports, feasibility study papers, and the reports of advanced collaborative projects. The series offers an opportunity for researchers to present an extended exposition of such new work in all aspects of industrial control for wider and rapid dissemination.

When advances are made in industrial control technology, for example, sensors, actuators, controllers, communications, and computing power, there are at least three possible consequences for system control. The control engineering may decide to use the new hardware to make an existing control system perform better. Another possibility is that the new hardware advances make a proposed but previously impractical control method feasible. Alternatively, it may be necessary to devise a completely new control technique in order to exploit the new hardware. However, in many cases, it is often economically impossible to advance an aspect of the system's control hardware until there is a major upgrade of the system and the industrial control engineer has to grapple with the limitations of the system as it exists. (This is where control engineering science becomes an art!) Underactuated marine vessels are a case in point. In most configurations, a vessel's main actuators are propellers and rudders, yet a marine vessel has six degrees of freedom in its motion, and the marine control engineer simply has to work with the control surfaces and sensors available. One area that may advance control performance is the use of better control designs, and recently control engineers have become more interested in what nonlinear control might have to offer.

Researchers Khac Duc Do and Jie Pan have published a sequence of journal and conference papers on new control algorithms for underactuated systems. Most of this work has used models of marine systems (surface ships and underwater vehicles), but some of the work has been with other mechanical system models (wheeled mobile robots, vertical take-off and landing (VTOL) aircraft). Now they have taken the opportunity to capture this research and development work in a monograph en-
titled Control of Ships and Underwater Vehicles: Design for Underactuated and Nonlinear Marine Systems for the Advances in Industrial Control series. This will enable industrial control engineers and control researchers to read and study a systematic presentation of their ideas and work. The book opens with chapters that introduce the appropriate nonlinear control theory and marine vessel models used in the research, proceeds to controller derivation and development, and finally details simulation results for a series of nonlinear control schemes devised for marine vessel control problems like point-to-point navigation and path-following. Also presented are some field results of a laboratory-scale vessel on a local river. In one chapter other applications fields are considered and the nonlinear control results for a simple wheeled robots and a simple VTOL aircraft model are given. Each of the applications chapters contains illustrative simulation studies and control results.

The monograph will interest control researchers, graduate students, and industrial control engineers alike, particularly those involved in marine and wheeled robotic motion control problems. The Series Editors, being based in Glasgow, Scotland, have always had an interest in marine control problems and have endeavoured to ensure that the Advances in Industrial Control series has useful volumes from this field of control application. Past volumes have included: Ship Motion Control by Tristan Perez (ISBN 978-1-85233-959-3, 2005), Compressor Surge and Rotating Stall by Jan Tommy Gravdahl and Olav Egeland (ISBN 978-1-85233-067-5, 1998), and Robust Control of Diesel Ship Propulsion by Nikolaos Xiros (ISBN 978-1-85233-543-4, 2002), and we are pleased to welcome this new volume, Control of Ships and Underwater Vehicles into the series.

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2009

## Preface

Control of ocean vessels including ships and underwater vehicles is an active field due to its theoretical challenges and important applications such as passenger and goods transportation, environmental surveying, undersea cable inspection, offshore oil installations, and many others. Most ocean vessels are underactuated meaning that they have more degrees of freedom to be controlled than the number of independent control inputs. Ships do not usually have an independent sway actuator while for underwater vehicles there are often no independent sway and heave actuators. As a result, motion control of underactuated ocean vessels opened a new territory in applied nonlinear control, and attracted special attention from both marine technology and control engineering communities.

If classical motion control systems designed for fully or overactuated vessels are directly used on underactuated vessels, the resulting performance of controlled systems is very poor or control objectives cannot be achieved. For example, the traditional approach, in which a combination of a conventional autopilot and a line-ofsight algorithm is used to steer an underactuated ship from one point to another on a straight line, does not impose on minimizing the lateral distance. Consequently, the shortest traveling distance is not achieved. Another example is that underactuated ocean vessels cannot be stabilized by any time-invariant continuous state feedback controllers although they are open loop controllable. This fact resulted from a direct application of the Brockett necessary condition to feedback stabilization of underactuated ocean vessels.

Inspired by progress in the field, we present this monograph on control of underactuated ocean vessels including ships and underwater vehicles to senior and postgraduate students, researchers and practitioners of marine technology, control engineering, mechanical engineering, electrical engineering, and mechatronics. This is the first book in the literature that offers various solutions to advanced feedback control topics of practical importance including stabilization, trajectory-tracking, path-tracking, and path-following for underactuated ocean vessels. In the control development and stability analysis of the controlled systems, practical motivations as well as nontrivial techniques are carefully detailed. The techniques presented in
the book can be readily applied to other underactuated mechanical systems such as aircraft, spacecraft, mobile robots, and robot arms.

## Acknowledgements

The first author would like to thank the Rectoral Board and his colleagues at Thai Nguyen University of Technology and the University of Western Australia for providing him with a friendly and efficient working environment during his time of writing this book. He is grateful to H.L. Nguyen for her support and encouragement in different ways. He also thanks Z.P. Jiang for inviting him to Polytechnic University for a period of three months from October 2001 to January 2002. Both authors thank the anonymous reviewers and editorial staff of various journals and conferences for their helpful comments on the authors' research papers, which are the main source of research and results for this book. The authors' special thanks go to M.A. Johnson, Joint Editor of the Advances in Industrial Control series, O. Jackson, Editor (Engineering), and A. Bunning, Editorial Assistant (Engineering) of Springer, UK, for their support and constructive comments in the process of publishing the book.

The writing of this book was supported in part by the Australian Research Council under grants DP0453294, LP0219249, DP0774645, and DP0988424.

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May, 2009

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## Chapter 1 <br> Introduction

The subject of this book is the application of nonlinear control theory to control underactuated ocean vessels and to analyze stability of their controlled systems.

A vessel is said to be underactuated if it has more degrees of freedom to be controlled than the number of independent control inputs. Many ocean ships are equipped with propellers and rudders for surge and yaw motions only, but without any actuators for direct control of their sway motion. The control of underactuated ocean vessels has become an active area of research in recent years not only because it poses many challenging questions in applied nonlinear control theory, but also because of its practical importance.

This chapter starts with a brief review of the development in nonlinear control theory and its applications. Next, difficulties in control of underactuated ocean vessels are discussed. The chapter ends with a description of the organization of the book.

### 1.1 Overview of Nonlinear Control Developments

Over the last two decades nonlinear control theory has received considerable research effort and has undergone a period of significant progress. This is motivated from the fact that real world problems are often inherently nonlinear. Conventionally, control engineers studied a nonlinear system by using a linearized model around selected operating points, and suitable techniques for linear control systems founded on the basis of the superposition principle. A fundamental limitation of this linearization approach is that stability and performance can only be guaranteed in the neighborhood of the selected operating points. Consequently, studying nonlinear systems based on linear models is severely limited since real world systems are usually required to operate over a large number of operating points. In addition, many phenomena such as the existence of multiple equilibria or operating points, periodic variation of state variables or limit cycles, finite escape time, and bifurcations cannot be described or predicted by using linear models. The above-mentioned
limitations led control researchers to directly study nonlinear systems. Moving from linear to nonlinear systems, we are faced with a much more difficult situation. The superposition principle no longer holds and analysis tools involve more advanced mathematics. In addition, the separation principle does not hold for nonlinear systems in general.

Currently, the techniques that are available for nonlinear control system design include linearization and gain-scheduling techniques, feedback linearization, circle Popov criteria, small-gain theorems, passivity, averaging, singular perturbation, sliding mode control, Lyapunov design and redesign, backstepping and forwarding, adaptive control, input-to-state stability, integral input-to-state stability, and differential geometric approaches. For details of the above-mentioned approaches, the reader is referred to [1-10]. Nonlinear control methods are increasingly implemented in practice with great success such as electrical motors, diesel engines, ocean vessels, jet engine compressors and electromechanical systems. For some of these applications, the reader is referred to [3,11-15].

### 1.2 Difficulties in Control of Underactuated Ocean Vessels

The main difficulty in the control of an underactuated system is that the system has more degrees of freedom (outputs) to be controlled than the number of independent actuators (inputs). A direct application of the Brockett necessary condition indicates that an underactuated system cannot be stabilized by any time-invariant continuous state feedback controllers although it is open loop controllable. This topic is detailed in Section 2.7.4. Moreover, if classical motion control systems designed for fully or overactuated systems are directly used on the underactuated ones, the resulting performance of the controlled systems is very poor or the control objectives cannot be achieved. The above observations are also true for underactuated ocean vessels. Below we illustrate the aforementioned difficulties through a simple example.

In Chapter 3, the equations of motion of an underactuated ocean vessel without environmetal disturbances are given by (see (3.31)):

$$
\begin{align*}
& \dot{\eta}=J(\eta) v, \\
& M \dot{v}=-C(v) v-D(v) v-g(\eta)+\tau \tag{1.1}
\end{align*}
$$

where

$$
\begin{align*}
\eta & =[x, y, z, \phi, \theta, \psi]^{T}, \\
v & =[u, v, w, p, q, r]^{T} \tag{1.2}
\end{align*}
$$

are the vectors of position/Euler angles and velocities, respectively. The matrices $\boldsymbol{M}, \boldsymbol{C}(\boldsymbol{v})$, and $\boldsymbol{D}(\boldsymbol{v})$ denote inertia, Coriolis, and damping, respectively, $\boldsymbol{g}(\boldsymbol{\eta})$ is the vector of gravitational/buoyancy forces, $\boldsymbol{\tau} \in \mathbb{R}^{6}$ is the vector of the forces and moments provided by the actuators. Depending on the vessel's configuration, several
elements of $\boldsymbol{\tau}$ are zero. For example, for the case where an underwater vehicle does not have independent actuators in the heave and sway axes, the second and third elements of $\boldsymbol{\tau}$ are zero. This gives an underactuated situation. We will show that an underactuation causes serious problems for classical control design techniques.

We now assume that we want to design the control input vector $\boldsymbol{\tau}$ to force the vector of position/Euler angles $\eta$ to track a reference vector of position/Euler angles $\eta_{\boldsymbol{d}}$ with bounded $\dot{\eta}_{\boldsymbol{d}}$ and $\ddot{\boldsymbol{\eta}}_{\boldsymbol{d}}$. As a standard application of the control design techniques in $[3,6]$ for such a strict feedback system like (1.1), we define the tracking error vector $\eta_{\boldsymbol{e}}$ as

$$
\begin{equation*}
\eta_{e}=\eta-\eta_{d} \tag{1.3}
\end{equation*}
$$

and differentiate $\boldsymbol{\eta}_{\boldsymbol{e}}$ twice along the solutions of the system (1.1) to obtain

$$
\begin{equation*}
\ddot{\eta}_{e}=\dot{J}(\eta) v+J(\eta) M^{-1}(-C(v) v-D(v) v-g(\eta)+\tau)-\ddot{\eta}_{d} \tag{1.4}
\end{equation*}
$$

Letting

$$
\begin{equation*}
\dot{J}(\eta) v+J(\eta) M^{-1}(-C(v) v-D(v) v-g(\eta)+\tau)-\ddot{\eta}_{d}=-K_{1} \eta_{e}-K_{2} \dot{\eta}_{e} \tag{1.5}
\end{equation*}
$$

where $K_{\mathbf{1}} \in \mathbb{R}^{6 \times 6}$ and $K_{\mathbf{2}} \in \mathbb{R}^{6 \times 6}$ are positive definite symmetric control gain matrices, results in

$$
\begin{equation*}
\ddot{\eta}_{e}=-K_{1} \eta_{e}-K_{2} \dot{\eta}_{e}, \tag{1.6}
\end{equation*}
$$

where $\dot{\boldsymbol{\eta}}_{\boldsymbol{e}}=\boldsymbol{J}(\boldsymbol{\eta}) \boldsymbol{v}-\dot{\boldsymbol{\eta}}_{\boldsymbol{d}}$, which is globally exponentially stable (GES) at the origin. We now need to solve (1.5) for $\boldsymbol{\tau}$ for implementation. From (1.5), we have

$$
\begin{align*}
\tau & =M J^{-1}(\eta)\left(-K_{1} \eta_{e}-K_{2} \dot{\eta}_{e}-\dot{J}(\eta) v+\ddot{\eta}_{d}\right)+C(v) v+D(v) v+g(\eta) \\
& :=\Psi\left(\eta_{d}, \dot{\eta}_{d}, \ddot{\eta}_{d}, \eta, v\right) \tag{1.7}
\end{align*}
$$

Let us assume that the actuators available are independent and assume that they number $m_{c}$, where $m_{c}<6$. Consequently we can propose a decomposition of $\boldsymbol{\tau}$ as:

H1. $\boldsymbol{\tau}_{\boldsymbol{a}} \in \mathbb{R}^{6}$, where $\boldsymbol{\tau}_{\boldsymbol{a}}$ has zeros in the element positions identifying that there are no actuator inputs and where $\boldsymbol{\tau}_{\boldsymbol{a}}$ contains the elements of $\boldsymbol{\tau}$ associated with actuator locations.
$\mathrm{H} 2 . \boldsymbol{\tau}_{\boldsymbol{n} \boldsymbol{a}} \in \mathbb{R}^{6}$. This vector complements $\boldsymbol{\tau}_{\boldsymbol{a}}$ by having zeros in the locations associated with actuator inputs and contains the elements of $\boldsymbol{\tau}$ in the locations associated with no actuator inputs.
H3. Similarly, the vector $\boldsymbol{\Psi}(\cdot)$ can also be decomposed into vectors $\boldsymbol{\Psi}_{\boldsymbol{a}}(\cdot)$ and $\boldsymbol{\Psi}_{\boldsymbol{n} \boldsymbol{a}}(\cdot)$ according to the rules of schemes (H1) and (H2).
Thus, the computed control vector $\boldsymbol{\tau}$ can be decomposed as

$$
\begin{equation*}
\tau=\tau_{a}+\tau_{n a}=\Psi_{a}(\cdot)+\Psi_{n a}(\cdot) \in \mathbb{R}^{6} \tag{1.8}
\end{equation*}
$$

and

$$
\begin{align*}
& \boldsymbol{\tau}_{\boldsymbol{a}}=\boldsymbol{\Psi}_{\boldsymbol{a}}(\cdot) \in \mathbb{R}^{6}, \\
& \boldsymbol{\tau}_{\boldsymbol{n} \boldsymbol{a}}=\boldsymbol{\Psi}_{\boldsymbol{n} \boldsymbol{a}}(\cdot) \in \mathbb{R}^{6} . \tag{1.9}
\end{align*}
$$

In practice the actuator configuration only allows $\boldsymbol{\tau}_{\boldsymbol{a}}$ to be implemented, and the response of the vessel (system) will result from this input alone.

If the vessel response is given by $\boldsymbol{\eta}$, where

$$
\begin{equation*}
\eta=\boldsymbol{H}(\tau) \tag{1.10}
\end{equation*}
$$

the full response would have been $\boldsymbol{\eta}_{\text {full }}$, where

$$
\begin{equation*}
\eta_{\text {full }}=\boldsymbol{H}(\boldsymbol{\tau})=\boldsymbol{H}\left(\tau_{\boldsymbol{a}}+\tau_{\boldsymbol{n} \boldsymbol{a}}\right) \tag{1.11}
\end{equation*}
$$

but the actual response is $\eta_{\text {partial }}$, where

$$
\begin{equation*}
\eta_{\text {partial }}=\boldsymbol{H}\left(\boldsymbol{\tau}_{\boldsymbol{a}}\right)=\boldsymbol{H}_{\boldsymbol{s}}\left(\boldsymbol{\tau}_{\boldsymbol{r} \boldsymbol{a}}\right) \tag{1.12}
\end{equation*}
$$

and $\boldsymbol{\tau}_{\boldsymbol{a}} \in \mathbb{R}^{6}, \boldsymbol{\eta}_{\text {partial }} \in \mathbb{R}^{6}, \boldsymbol{H}(\cdot): \mathbb{R}^{6} \rightarrow \mathbb{R}^{6}$ but $\boldsymbol{\tau}_{\boldsymbol{r} \boldsymbol{a}}$ is the $m_{\boldsymbol{c}}$ vector of nonzero actuator inputs, $\boldsymbol{\tau}_{\boldsymbol{r} \boldsymbol{a}} \in \mathbb{R}^{m_{c}}$, and $\boldsymbol{H}_{\boldsymbol{s}}(\cdot)$ is the reduced functional representation with $\boldsymbol{H}_{\boldsymbol{s}}: \mathbb{R}^{m_{c}} \rightarrow \mathbb{R}^{6}$.

Therefore the fully actuated vessel is represented by $\boldsymbol{H}(\cdot)$ and the underactuated vessel is represented by $\boldsymbol{H}_{\boldsymbol{s}}(\cdot)$. These two representations may have very different properties. The implications of not being able to implement the control action $\boldsymbol{\tau}_{\boldsymbol{n a}}$, and having a vessel whose representation is $\boldsymbol{H}_{\boldsymbol{s}}(\cdot)$ are potentially:

1. Degraded performance achieved by $\boldsymbol{\eta}_{\text {partial }}$ when compared to $\boldsymbol{\eta}_{\text {full }}$.
2. Total loss of performance and inability to meet the control objectives in any useful way.
3. Loss of guaranteed properties like closed loop stability with possible catastrophic system failure.

Ways to overcome these problems include:

1. Reduction of performance specification as a control objective that can be achieved by $\eta_{\text {partial }}$.
2. Modification of the control design procedure to accommodate the reduced degrees of freedom available for control.
3. Procedures to use the available actuators to minimize the deleterious effects caused by not being able to implement the $\boldsymbol{\tau}_{\boldsymbol{n} \boldsymbol{a}}$ component.
4. The use of advanced control procedures like those described in this book.
5. The installation of additional actuators able to create access to effect and control the missing degrees of freedom.

### 1.3 Organization of the Book

The rest of the book is organized in 14 chapters, see Table 1.1. A brief summary of each chapter is given below.

Chapter 2 presents mathematical tools like Lyapunov stability theory, Barbalatlike lemmas, and the backstepping technique to be used in control design and stability analysis in the book.

Chapter 3 sets out the basic material that will be used in the subsequent chapters. Basic motion control tasks for ocean vessels and modeling of vessels are given.

Chapter 4 presents control properties of ocean vessels. Then, the existing literature on the control of underactuated ocean vessels is reviewed. Through the review of the previous work in the area of stabilization, trajectory-tracking, path-following, and output feedback control of underactuated ocean vessels, challenging questions are raised. These questions motivate contributions of the book to new solutions for the motion control of underactuated ships and underwater vehicles.

Chapter 5 addresses the problem of trajectory-tracking control of underactuated ships. These ships do not have independent actuators in the sway axis. The reference trajectory is generated by a suitable virtual ship. The control development is based on an elegant coordinate transformation, Lyapunov's direct method, and the backstepping technique, and utilizes passivity properties of the ship dynamics and their interconnected structure.

Chapter 6 examines the problem of designing a single controller that achieves stabilization and trajectory-tracking simultaneously for underactuated ships. In comparison with the preceding chapter, a path approaching the origin and a set-point can also be included in the reference trajectory, i.e., stabilization/regulation is also considered. The control development is based on several special coordinate transformations plus the techniques in the preceding chapter.

Chapter 7 presents global partial-state feedback and output feedback control schemes for trajectory-tracking control of underactuated ships. For the case of partial-state feedback, measurements of the ship sway and surge velocities are not needed, while for the case of output feedback, no ship velocities are required for feedback. Global nonlinear coordinate changes are introduced to transform the ship dynamics to a system affine in ship velocities to design observers to globally exponentially estimate unmeasured velocities. These observers plus the techniques in Chapters 5 and 6 facilitate the controllers' development.

Chapter 8 deals with the problem of path-tracking control of underactuated ships. In comparison with Chapters 5, 6, and 7, the requirement of the reference trajectory generated by a suitable virtual ship is relaxed. Both full state feedback and output feedback cases are considered. The control development is based on a series of nontrivial coordinate transformations plus the techniques in the previous chapters.

Chapter 9 addresses the problem of way-point tracking control of underactuated ships. Both full state feedback and output feedback controllers are designed. The controllers in this chapter can be considered as an advanced version of the conventional course-keeping controllers in the sense that in addition to maintaining the
desired heading, both the sway displacement (lateral distance) and sway velocity are controlled.

Chapter 10 develops full state feedback and output feedback controllers that force underactuated ships to follow a predefined path. The control development is motivated by an observation that it is practical to steer a vessel such that the vessel is on the reference path and its total velocity is tangent to the reference path, and that the vessel's forward speed is controlled separately by the main thruster control system. The techniques in the previous chapters plus the use of the Serret-Frenet frame facilitate the results.

Chapter 11 presents a different approach from Chapter 10 to solve a pathfollowing problem for underactuated ships. Unlike in Chapter 10, here the control development is based on the method of generating reference paths by the helmsman. The path-following errors are first interpreted in polar coordinates, the techniques used in the previous chapters are then used to design path-following controllers. Interestingly, some situations of practical importance such as parking and point-to-point navigation are covered in this chapter as a by-product of the developed path-following system.

Chapter 12 addresses the problem of trajectory-tracking control of underactuated underwater vehicles. These vehicles do not have independent actuators in the sway and heave axes. The control development is built on the techniques developed for underactuated ships in Chapters 5, 6, and 7. Due to complex dynamics of underwater vehicles in comparison with that of the surface ships, the control design and stability analysis require more complicated coordinate transformations and techniques than those developed for underactuated ships in Chapters 5, 6, and 7.

Chapter 13 extends the results of Chapter 11 to the design of a path-following control system for underactuated underwater vehicles. A series of path-following strategies for the vehicles is first discussed. A practical approach is then chosen to facilitate the control development. We also address the parking and point-to-point navigation problems of the vehicles in this chapter.

Chapter 14 illustrates several applications of the observer and control design techniques developed in the previous chapters to control of other underactuated mechanical systems, which are common in practice. These systems include mobile robots and VTOL aircraft. For mobile robots, a global exponential observer is first designed based on the observer design for underactuated ships in Chapter 7. Output feedback simultaneous stabilization and trajectory-tracking and path-following controllers are then developed using the control design techniques proposed for underactuated ships in Chapters 5 and 10. For VTOL aircraft, the observer and control design strategies used for underactuated ships in Chapters 5 and 6 are utilized to design a global output feedback trajectory-tracking controller.

Finally, Chapter 15 concludes the book by briefly summarizing the main results presented in the previous chapters and presenting related open problems for further investigation.

Table 1.1 Organization of the book

| 1. Introduction |  |
| :--- | :--- |
| Part I: Mathematical Tools | 2. Mathematical Preliminaries |
| Part II: Modeling and Control <br> Properties of Ocean Vessels | 3. Modeling of Ocean Vessels <br> 4. Control Properties and Previous Work on Control of <br> Ocean Vessels |
| Part III: Control of | 5. Trajectory-tracking Control of Underactuated Ships <br> 6. Simultaneous Stabilization and Trajectory-tracking Con- <br> Underactuated Ships <br> trol of Underactuated Ships <br> 7. Partial-state and Output Feedback Trajectory-tracking <br> Control of Underactuated Ships <br> 8. Path-tracking Control of Underactuated Ships <br> 9. Way-point Tracking Control of Underactuated Ships <br> 10. Path-following of Underactuated Ships Using Serret- <br> Frenet Coordinates <br> 11. Path-following of Underactuated Ships Using Polar Co- <br> ordinates |
| Part IV: Control of |  |
| Underactuated Underwater <br> Vehicles | 12. Trajectory-tracking Control of Underactuated Under- <br> water Vehicles <br> 13. Path-following of Underactuated Underwater Vehicles |
| Part V: Control of Other <br> Underactuated Mechanical <br> Systems | 14. Control of Other Underactuated Mechanical Systems |

## Part I Mathematical Tools

## Chapter 2 Mathematical Preliminaries

This chapter presents mathematical tools, which will be used in control design and stability analysis in the subsequent chapters. Some standard theorems, lemmas and corollaries, which are available in references, are sometimes given without a proof.

### 2.1 Lyapunov Stability

Stability theory is important in system theory and engineering. For a controlled system to be usable in practice, the least requirement is that the system is stable under unknown disturbances or noise. There are various types of stability problems that arise in the study of dynamical systems such as input-output stability, Lyapunov stability, absolute stability and stability of periodic solutions. These stability concepts have been studied extensively over the last 100 years. For control design and stability analysis in this book, Lyapunov stability plays an important role for the following reasons. First, Lyapunov's direct method uses an energy-like function called the Lyapunov function to study the behaviors of dynamical systems analytically. This function reflects physical properties of the systems under study. Second, Lyapunov's second method is applicable to both linear and nonlinear systems. Third, many results in input-output stability can be obtained directly using Lyapunov stability theory. Lyapunov stability theory generally includes Lyapunov's first and second methods. The first method and the center manifold theory developed later are techniques based on the lowest-order approximation around a given point or a nominal trajectory. The stability results obtained using these two approximation methods are inherently local, and the stability regions are hard to estimate and often quite small. Since the objective of this book is to study nonlinear dynamics of ocean vessels, we omit the aforementioned two approximation techniques. Our primary interest is in stability based on Lyapunov's second method for systems described by ordinary differential equations.

This section is concerned with stability of equilibrium points in the sense of Lyapunov. An equilibrium point is stable if all solutions starting at nearby points
stay nearby; otherwise it is unstable. It is asymptotically stable if in addition, all solutions tend to the equilibrium point as time approaches infinity. These notions will be made mathematically rigorous in this section.

Consider the following nonautonomous system:

$$
\begin{equation*}
\dot{x}=f(t, x) \tag{2.1}
\end{equation*}
$$

where $f:[0, \infty) \times D \rightarrow \mathbb{R}^{n}$ is piecewise continuous in $t$ and locally Lipschitz in $x$ on $[0, \infty) \times D$ and $D \subseteq \mathbb{R}^{n}$ is a domain that contains the origin $x=0$.

### 2.1.1 Definitions

Definition 2.1. The origin $x=0$ is the equilibrium point of (2.1) if

$$
\begin{equation*}
f(t, 0)=0, \forall t \geq 0 \tag{2.2}
\end{equation*}
$$

Definition 2.2. A continuous function $\alpha:[0, a) \rightarrow \mathbb{R}^{+}$is said to belong to class $K$ if it is strictly increasing and $\alpha(0)=0$. It is said to belong to class $K_{\infty}$ if $a=\infty$ and $\alpha(r) \rightarrow \infty$ as $r \rightarrow \infty$.

Definition 2.3. A continuous function $\beta:[0, a) \times \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$is said to belong to class $K L$ if, for each fixed $s$, the mapping $\beta(r, s)$ belongs to class $K$ with respect to $r$ and, for each fixed $r$, the mapping $\beta(r, s)$ is decreasing with respect to $s$ and $\beta(r, s) \rightarrow 0$ as $s \rightarrow \infty$. It is said to belong to class $K L_{\infty}$ if, in addition, for each fixed $s$ the mapping $\beta(r, s)$ belongs to class $K_{\infty}$ with respect to $r$.

Definition 2.4. The equilibrium point $x=0$ of (2.1) is

1. stable if, for each $\varepsilon>0$, there is $\delta=\delta\left(\varepsilon, t_{0}\right)>0$ such that

$$
\begin{equation*}
\left\|x\left(t_{0}\right)\right\|<\delta \Rightarrow\|x(t)\|<\varepsilon, \forall t \geq t_{0} \geq 0 \tag{2.3}
\end{equation*}
$$

2. uniformly stable if, for each $\varepsilon>0$, there is $\delta=\delta(\varepsilon)>0$ independent of $t_{0}$ such that (2.3) is satisfied,
3. unstable if it is not stable,
4. asymptotically stable if it is stable and there is a positive constant $c=c\left(t_{0}\right)$ such that $x(t) \rightarrow 0$ as $t \rightarrow \infty$, for all $\left\|x\left(t_{0}\right)\right\|<c$,
5. uniformly asymptotically stable if it is uniformly stable and there is a positive constant $c$, independent of $t_{0}$, such that for all $\left\|x\left(t_{0}\right)\right\|<c, x(t) \rightarrow 0$ as $t \rightarrow \infty$, uniformly in $t_{0}$; that is, for each $\eta>0$, there is $T=T(\eta)>0$ such that

$$
\begin{equation*}
\|x(t)\|<\eta, \forall t \geq t_{0}+T(\eta), \forall\left\|x\left(t_{0}\right)\right\|<c \tag{2.4}
\end{equation*}
$$

6. globally uniformly asymptotically stable (GUAS) if it is uniformly stable, $\delta(\varepsilon)$ can be chosen to satisfy $\lim _{\varepsilon \rightarrow \infty} \delta(\varepsilon)=\infty$, and, for each pair of positive numbers $\eta$ and $c$, there is $T=T(\eta, c)>0$ such that

$$
\begin{equation*}
\|x(t)\|<\eta, \forall t \geq t_{0}+T(\eta, c), \forall\left\|x\left(t_{0}\right)\right\|<c \tag{2.5}
\end{equation*}
$$

Definition 2.5. The equilibrium point $x=0$ of (2.1) is exponentially stable if there exist positive constants $c, k$, and $\lambda$ such that

$$
\begin{equation*}
\|x(t)\| \leq k\left\|x\left(t_{0}\right)\right\| e^{-\lambda\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0, \forall\left\|x\left(t_{0}\right)\right\|<c \tag{2.6}
\end{equation*}
$$

and GES if (2.6) is satisfied for any initial state $x\left(t_{0}\right)$.
Definition 2.6. The equilibrium point $x=0$ of (2.1) is $K$-exponentially stable if there exist positive constants $c$ and $\lambda$ and a class $K$ function $\alpha$ such that

$$
\begin{equation*}
\|x(t)\| \leq \alpha\left(\left\|x\left(t_{0}\right)\right\|\right) e^{-\lambda\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0, \forall\left\|x\left(t_{0}\right)\right\|<c \tag{2.7}
\end{equation*}
$$

and globally $K$-exponentially stable if (2.7) is satisfied for any initial state $x\left(t_{0}\right)$.
Definition 2.7. The solutions of (2.1) are as follows:

1. Uniformly bounded if there exists a positive constant $c$, independent of $t_{0} \geq 0$, and for every $a \in(0, c)$, there is $\beta=\beta(a)>0$, independent of $t_{0}$, such that

$$
\begin{equation*}
\left\|x\left(t_{0}\right)\right\| \leq a \Rightarrow\|x(t)\| \leq \beta, \forall t \geq t_{0} \tag{2.8}
\end{equation*}
$$

2. Globally uniformly bounded if (2.8) holds for an arbitrarily large $a$.
3. Uniformly ultimately bounded with ultimate bound $b$ if there exist positive constants $b$ and $c$, independent of $t_{0} \geq 0$, and for every $a \in(0, c)$ there is $T=T(a, b) \geq 0$, independent of $t_{0}$, such that

$$
\begin{equation*}
\left\|x\left(t_{0}\right)\right\| \leq a \Rightarrow\|x(t)\| \leq b, \forall t \geq t_{0}+T \tag{2.9}
\end{equation*}
$$

4. Globally uniformly ultimately bounded if (2.9) holds for an arbitrarily large $a$.

### 2.1.2 Lemmas and Theorems

The following lemma provides equivalent definitions of uniform stability and uniform asymptotic stability by using class $K$ and class $K L$ functions.

Lemma 2.1. The equilibrium point $x=0$ of (2.1) is

1. uniformly stable if and only if there exist a class $K$ function $\alpha$ and a positive constant $c$, independent of $t_{0}$, such that

$$
\begin{equation*}
\|x(t)\| \leq \alpha\left(\left\|x\left(t_{0}\right)\right\|\right), \forall t \geq t_{0} \geq 0, \forall x\left(t_{0}\right) \mid\left\|x\left(t_{0}\right)\right\|<c \tag{2.10}
\end{equation*}
$$

2. uniformly asymptotically stable if and only if there exist a class $K L$ function $\beta$ and a positive constant $c$, independent of $t_{0}$, such that

$$
\begin{equation*}
\|x(t)\| \leq \beta\left(\left\|x\left(t_{0}\right)\right\|, t-t_{0}\right), \forall t \geq t_{0} \geq 0, \forall x\left(t_{0}\right) \mid\left\|x\left(t_{0}\right)\right\|<c \tag{2.11}
\end{equation*}
$$

3. GUAS if and only if inequality (2.11) is satisfied with $\beta \in K L_{\infty}$ for any initial state $x\left(t_{0}\right)$.

Proof. See [6].
Lemma 2.2. Assume that $d: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ satisfies

$$
\begin{equation*}
P\left[\frac{\partial d}{\partial x}\right]+\left[\frac{\partial d}{\partial x}\right]^{T} P \geq 0, \forall x \in \mathbb{R}^{n}, \tag{2.12}
\end{equation*}
$$

where $P=P^{T}>0$. Then

$$
\begin{equation*}
(x-y)^{T} P(d(x)-d(y)) \geq 0, \quad \forall x, y \in \mathbb{R}^{n} . \tag{2.13}
\end{equation*}
$$

Proof. See [16].
Lemma 2.3. The following nonlinear interconnected system:

$$
\begin{align*}
& \dot{x}_{1}=f_{1}\left(t, x_{1}, x_{2}\right)+g_{1}\left(t, x_{1}, x_{2}\right) u,  \tag{2.14}\\
& \dot{x}_{2}=f_{2}\left(t, x_{1}, x_{2}\right)+g_{2}\left(t, x_{1}, x_{2}\right) u,
\end{align*}
$$

where $x_{i} \in \mathbb{R}, i=1,2, f_{i}\left(t, x_{1}, x_{2}\right)$ are locally Lipschitz in $x_{i}$ and piecewise continuous in $t ; u \in \mathbb{R}$ is the control input, and $g_{2}\left(t, x_{1}, x_{2}\right) \neq 0, \forall t \geq 0, x_{i} \in \mathbb{R}$, can be transformed to the following system

$$
\begin{align*}
& \dot{z}_{1}=\gamma_{1}\left(t, z_{1}, x_{2}\right),  \tag{2.15}\\
& \dot{x}_{2}=\gamma_{2}\left(t, z_{1}, x_{2}\right)+\varphi_{2}\left(t, z_{1}, x_{2}\right) u .
\end{align*}
$$

Proof. Define

$$
\begin{equation*}
z_{1}=x_{1}+\pi\left(t, x_{1}, x_{2}\right), \tag{2.16}
\end{equation*}
$$

where $\pi\left(t, x_{1}, x_{2}\right)$ is to be determined. Differentiating both sides of (2.16) along the solutions of (2.14) yields

$$
\begin{align*}
\dot{z}_{1}= & \left(\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{1}}+1\right) f_{1}\left(t, x_{1}, x_{2}\right)+ \\
& \frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{2}} f_{2}\left(t, x_{1}, x_{2}\right)+\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial t}+  \tag{2.17}\\
& {\left[\left(\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{1}}+1\right) g_{1}\left(t, x_{1}, x_{2}\right)+\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{2}} g_{2}\left(t, x_{1}, x_{2}\right)\right] u . }
\end{align*}
$$

Now choosing the function $\pi\left(t, x_{1}, x_{2}\right)$ such that

$$
\begin{equation*}
\left(\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{1}}+1\right) g_{1}\left(t, x_{1}, x_{2}\right)+\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{2}} g_{2}\left(t, x_{1}, x_{2}\right)=0 \tag{2.18}
\end{equation*}
$$

results in (2.15), where the functions $\gamma_{1}\left(t, z_{1}, x_{2}\right), \gamma_{2}\left(t, z_{1}, x_{2}\right)$, and $\varphi_{2}\left(t, z_{1}, x_{2}\right)$ are defined as

$$
\begin{align*}
\gamma_{1}\left(t, z_{1}, x_{2}\right):= & \left(\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{1}}+1\right) f_{1}\left(t, x_{1}, x_{2}\right)+ \\
& \frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial x_{2}} f_{2}\left(t, x_{1}, x_{2}\right)+\frac{\partial \pi\left(t, x_{1}, x_{2}\right)}{\partial t}  \tag{2.19}\\
\gamma_{2}\left(t, z_{1}, x_{2}\right):= & f_{2}\left(t, x_{1}, x_{2}\right), \varphi_{2}\left(t, z_{1}, x_{2}\right):=g_{2}\left(t, x_{1}, x_{2}\right)
\end{align*}
$$

with $x_{1}$ being solved from (2.16) and substituted in.

## Remark 2.1.

1. The success of the above lemma depends on whether a solution to the partial differential equation (2.18) can be found. Solving this partial differential equation is difficult in general but might be possible in some specific cases such as the vehicle systems in this book.
2. In some cases, designing a control input $u$ for the transformed system (2.15) is simpler than that for the original system (2.14).

The main Lyapunov stability theorem, which has a number of applications in studying stability of (2.1), is given below.

Theorem 2.1. Let $D=\left\{x \in \mathbb{R}^{n} \mid\|x\|<r\right\}$ and $x=0$ be an equilibrium point of (2.1). Let $V: D \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{+}$be a continuously differentiable function such that $\forall t \geq 0, \forall x \in D$,

$$
\begin{align*}
& \gamma_{1}(\|x\|) \leq V(x, t) \leq \gamma_{2}(\|x\|) \\
& \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x) \leq-\gamma_{3}(\|x\|) . \tag{2.20}
\end{align*}
$$

Then the equilibrium point $x=0$ is

1. uniformly stable, if $\gamma_{1}$ and $\gamma_{2}$ are class $K$ functions on $[0, r)$ and $\gamma_{3} \geq 0$ on $[0, r)$,
2. uniformly asymptotically stable, if $\gamma_{1}, \gamma_{2}$ and $\gamma_{3}$ are class $K$ functions on $[0, r)$,
3. exponentially stable if $\gamma_{i}(\rho)=k_{i} \rho^{\alpha}$ on $[0, r), k_{i}>0, \alpha>0, i=1,2,3$,
4. globally uniformly stable if $D=\mathbb{R}^{n}, \gamma_{1}$ and $\gamma_{2}$ are class $K_{\infty}$ functions, and $\gamma_{3} \geq 0$ on $\mathbb{R}^{+}$,
5. GUAS if $D=\mathbb{R}^{n}, \gamma_{1}$ and $\gamma_{2}$ are class $K_{\infty}$ functions, and $\gamma_{3}$ is a class $K$ function on $\mathbb{R}^{+}$,
6. GES, if $D=\mathbb{R}^{n}, \gamma_{i}(\rho)=k_{i} \rho^{\alpha}$ on $\mathbb{R}^{+}, k_{i}>0, \alpha>0, i=1,2,3$.

Proof. See [3].
Theorem 2.2. Let $x=0$ be an equilibrium point of (2.1) and suppose that $f$ is locally Lipschitz in $x$ and uniformly continuous in $t$. Let $V: \mathbb{R}^{n} \times \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$be a continuously differentiable function such that

$$
\begin{align*}
& \gamma_{1}(\|x\|) \leq V(x, t) \leq \gamma_{2}(\|x\|) \\
& \dot{V}=\frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x) \leq-W(x) \leq 0, \tag{2.21}
\end{align*}
$$

for all $t \geq 0$ and $x \in \mathbb{R}^{n}$, where $\gamma_{1}$ and $\gamma_{2}$ are class $K_{\infty}$ functions, and $W$ is a continuous function. Then all solutions of (2.1) are globally uniformly bounded and satisfy

$$
\begin{equation*}
\lim _{t \rightarrow \infty} W(x(t))=0 \tag{2.22}
\end{equation*}
$$

In addition, if $W(x)$ is positive definite, then the equilibrium point $x=0$ is GUAS. Proof. See [3].

The following Lyapunov-like theorem is useful for showing uniform boundedness and ultimate boundedness.

Theorem 2.3. Let $V:[0, \infty) \times D \rightarrow \mathbb{R}$ be a continuously differentiable function and $D \subset \mathbb{R}^{n}$ be a domain that contains the origin such that

$$
\begin{align*}
& \alpha_{1}(\|x\|) \leq V(x, t) \leq \alpha_{2}(\|x\|) \\
& \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x) \leq-W(x), \forall\|x\| \geq \mu>0 \tag{2.23}
\end{align*}
$$

for all $t \geq 0$ and $x \in D$ where $\alpha_{1}$ and $\alpha_{2}$ are class $K$ functions, and $W$ is a continuous positive definite function. Take $r>0$ such that $B_{r} \subset D$ and suppose that

$$
\begin{equation*}
\mu<\alpha_{2}^{-1}\left(\alpha_{1}(r)\right) \tag{2.24}
\end{equation*}
$$

Then, there exists a class KL function $\beta$ and for every initial state $x\left(t_{0}\right)$, satisfying $\left\|x\left(t_{0}\right)\right\|<\alpha_{2}^{-1}\left(\alpha_{1}(r)\right)$, there is $T \geq 0$ (dependent on $x\left(t_{0}\right)$ and $\mu$ ) such that the solutions of (2.1) satisfies

$$
\begin{align*}
& \|x(t)\| \leq \beta\left(\left\|x\left(t_{0}\right)\right\|, t-t_{0}\right), \quad \forall t_{0} \leq t \leq t_{0}+T, \\
& \|x(t)\|<\alpha_{1}^{-1}\left(\alpha_{2}(\mu)\right), \forall t \geq t_{0}+T . \tag{2.25}
\end{align*}
$$

Moreover, if $D=\mathbb{R}^{n}$ and $\alpha_{1}$ belongs to class $K_{\infty}$, then (2.25) holds for any initial state $x\left(t_{0}\right)$ with no restriction on how large $\mu$ is.

Proof. See [6].

### 2.1.3 Stability of Cascade Systems

Consider the following cascade system:

$$
\begin{align*}
& \dot{z}_{1}=f_{1}\left(t, z_{1}\right)+g\left(t, z_{1}, z_{2}\right) z_{2}  \tag{2.26}\\
& \dot{z}_{2}=f_{2}\left(t, z_{2}\right)
\end{align*}
$$

where $z_{1} \in \mathbb{R}^{n}, z_{2} \in \mathbb{R}^{m}, f_{1}\left(t, z_{1}\right)$ is continuously differentiable in $\left(t, z_{1}\right)$, and $f_{2}\left(t, z_{2}\right)$ and $g\left(t, z_{1}, z_{2}\right)$ are continuous and locally Lipschitz in $z_{2}$ and $\left(z_{1}, z_{2}\right)$, respectively.

If we set $z_{2}=0$, then the first equation of (2.26) becomes $\dot{z}_{1}=f_{1}\left(t, z_{1}\right)$. Therefore we can view the first equation of (2.26) as the system

$$
\begin{equation*}
\Omega_{1}: \dot{z}_{1}=f_{1}\left(t, z_{1}\right) \tag{2.27}
\end{equation*}
$$

which is perturbed by the output of the system

$$
\begin{equation*}
\Omega_{2}: \dot{z}_{2}=f_{2}\left(t, z_{2}\right) \tag{2.28}
\end{equation*}
$$

Now assume that the systems $\Omega_{1}$ and $\Omega_{2}$ are asymptotically stable at the origin, i.e., (2.27) and (2.28) yield $\lim _{t \rightarrow \infty} z_{1}(t)=0$ and $\lim _{t \rightarrow \infty} z_{2}(t)=0$, respectively. Based on these assumptions, it is plausible to conclude that the system (2.26) is asymptotically stable at the origin in general. In many cases, the solution $z_{1}(t)$ of the system (2.26) goes to infinity in finite time as can be seen from the simple cascade system:

$$
\begin{align*}
& \dot{z}_{1}=-k_{1} z_{1}+z_{1}^{2} z_{2} \\
& \dot{z}_{2}=-k_{2} z_{2} \tag{2.29}
\end{align*}
$$

where $k_{1}$ and $k_{2}$ are strictly positive constants. It is obvious that the subsystems $\dot{z}_{1}=-k_{1} z_{1}$ and $\dot{z}_{2}=-k_{2} z_{2}$ are GES at the origin. It is straightforward to show that the solution of (2.29) is

$$
\begin{align*}
& z_{1}(t)=\frac{z_{1}\left(t_{0}\right)\left(k_{1}+k_{2}\right)}{z_{1}\left(t_{0}\right) z_{2}\left(t_{0}\right) e^{-k_{2}\left(t-t_{0}\right)}+\left(k_{1}+k_{2}-z_{1}\left(t_{0}\right) z_{2}\left(t_{0}\right)\right) e^{k_{1}\left(t-t_{0}\right)}}, \\
& z_{2}(t)=z_{2}\left(t_{0}\right) e^{-k_{2}\left(t-t_{0}\right)} \tag{2.30}
\end{align*}
$$

It can be seen from (2.30) that if $z_{1}\left(t_{0}\right) z_{2}\left(t_{0}\right)<k_{1}+k_{2}$, then both $z_{1}(t)$ and $z_{2}(t)$ are bounded and converge exponentially to zero, respectively. If $z_{1}\left(t_{0}\right) z_{2}\left(t_{0}\right)=k_{1}+$ $k_{2}$, then $z_{2}(t)$ is still bounded and converges exponentially to zero but $z_{1}(t)$ tends to infinity exponentially fast. If $z_{1}\left(t_{0}\right) z_{2}\left(t_{0}\right)>k_{1}+k_{2}$, the situation is catastrophic, that is $z_{1}(t)$ tends to infinity when $t \rightarrow t_{0}+t_{f}$ with

$$
\begin{equation*}
t_{f}=\frac{1}{k_{1}+k_{2}} \ln \left(\frac{z_{1}\left(t_{0}\right) z_{2}\left(t_{0}\right)}{z_{1}\left(t_{0}\right) z_{2}\left(t_{0}\right)-\left(k_{1}+k_{2}\right)}\right) . \tag{2.31}
\end{equation*}
$$

The following theorem gives sufficient conditions of stability of the cascade system (2.26) based on stability of (2.27) and (2.28), and the connected term $g\left(t, z_{1}, z_{2}\right)$.
Theorem 2.4. Consider the following assumptions:

1. The systems (2.27) and (2.28) are both GUAS and we know explicitly a $C^{1}$ Lyapunov function $V\left(t, z_{1}\right)$, two class- $K_{\infty}$ functions $\alpha_{1}$ and $\alpha_{2}$, a class- $K$ function $\alpha_{4}$, and a positive semidefinite function $W\left(z_{1}\right)$ such that

$$
\begin{align*}
& \alpha_{1}\left(\left\|z_{1}\right\|\right) \leq V\left(t, z_{1}\right) \leq \alpha_{2}\left(\left\|z_{1}\right\|\right) \\
& \frac{\partial V}{\partial t}+\frac{\partial V}{\partial z_{1}} f_{1}\left(t, z_{1}\right) \leq-W\left(z_{1}\right),\left\|\frac{\partial V}{\partial z_{1}}\right\| \leq \alpha_{4}\left(\left\|z_{1}\right\|\right) . \tag{2.32}
\end{align*}
$$

2. For each fixed $z_{2}$, there exists a continuous function $\lambda: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$such that

$$
\begin{align*}
& \lim _{s \rightarrow \infty} \lambda(s)=0 \\
& \left\|\frac{\partial V}{\partial z_{1}} g\left(t, z_{1}, z_{2}\right)\right\| \leq \lambda\left(\left\|z_{1}\right\|\right) W\left(z_{1}\right) . \tag{2.33}
\end{align*}
$$

3. There exist continuous functions $\theta: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$and $\alpha_{5}: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$such that

$$
\begin{equation*}
\left\|g\left(t, z_{1}, z_{2}\right)\right\| \leq \theta\left(\left\|z_{2}\right\|\right) \alpha_{5}\left(\left\|z_{1}\right\|\right) \tag{2.34}
\end{equation*}
$$

and a continuous nondecreasing function $\alpha_{6}: \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$, and a nonnegative constant a such that

$$
\begin{align*}
& \alpha_{6}(s) \geq \alpha_{4}\left(\alpha_{1}^{-1}(s)\right) \alpha_{5}\left(\alpha_{1}^{-1}(s)\right), \\
& \int_{a}^{\infty} \frac{\mathrm{d} s}{\alpha_{6}(s)}=\infty \tag{2.35}
\end{align*}
$$

4. For each $r>0$, there exist constants $\chi>0$ and $\eta>0$ such that for all $t \geq 0$ and all $\left\|z_{2}\right\|<r$

$$
\begin{equation*}
\left\|\frac{\partial V}{\partial z_{1}} g\left(t, z_{1}, z_{2}\right)\right\| \leq \chi W\left(z_{1}\right), \forall\left\|z_{1}\right\| \geq \eta . \tag{2.36}
\end{equation*}
$$

5. There exists a class $K$ function $\phi$ such that the solution $z_{2}(t)$ of (2.28) satisfies

$$
\begin{equation*}
\int_{t_{0}}^{\infty}\left\|z_{2}(t)\right\| \mathrm{d} t \leq \phi\left(\left\|z_{2}\left(t_{0}\right)\right\|\right) \tag{2.37}
\end{equation*}
$$

Then we can conclude that if

- Assumptions 1 and 2, or
- Assumptions 1, 3 and 4, or
- Assumptions 1, 3 and 5
hold then the cascade system (2.26) is GUAS.
Proof. See [17].


### 2.2 Input-to-state Stability

Definition 2.8. The system

$$
\begin{equation*}
\dot{x}=f(t, x, u), \tag{2.38}
\end{equation*}
$$

where $f$ is piecewise continuous in $t$ and locally Lipschitz in $x$ and $u$, is said to be input-to-state stable (ISS) if there exist a class $K L$ function $\beta$ and a class $K$ function
$\gamma$, such that, for any $x\left(t_{0}\right)$ and for any input $u(\cdot)$ continuous and bounded on $[0, \infty)$, the solution exists for all $t \geq t_{0} \geq 0$ and satisfies

$$
\begin{equation*}
\|x(t)\| \leq \beta\left(\left\|x\left(t_{0}\right)\right\|, t-t_{0}\right)+\gamma\left(\sup _{t_{0} \leq \tau \leq t}\|u(\tau)\|\right) \tag{2.39}
\end{equation*}
$$

The following theorem establishes the equivalence between the existence of a Lyapunov-like function and the input-to-state stability.

Theorem 2.5. Suppose that for the system (2.38) there exists a $C^{1}$ function $V: \mathbb{R}^{+} \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{+}$such that for all $x \in \mathbb{R}^{n}$ and $u \in \mathbb{R}^{m}$,

$$
\begin{align*}
& \gamma_{1}(\|x\|) \leq V(t, x) \leq \gamma_{2}(\|x\|) \\
& \|x\| \geq \rho(\|u\|) \Rightarrow \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x, u) \leq-\gamma_{3}(\|x\|) \tag{2.40}
\end{align*}
$$

where $\gamma_{1}, \gamma_{2}$, and $\rho$ are class $K_{\infty}$ functions and $\gamma_{3}$ is a class- $K$ function. Then the system (2.38) is ISS with $\gamma=\gamma_{1}^{-1} \circ \gamma_{2} \circ \rho$.

Proof. If $x\left(t_{0}\right)$ is in the set

$$
\begin{equation*}
R_{t_{0}}=\left\{x \in \mathbb{R}^{n} \mid\|x\| \leq \rho\left(\sup \|u(\tau)\|_{\tau \geq t_{0}}\right)\right\} \tag{2.41}
\end{equation*}
$$

then $x(t)$ remains within the set

$$
\begin{equation*}
S_{t_{0}}=\left\{x \in \mathbb{R}^{n} \mid\|x\| \leq \gamma_{1}^{-1} \circ \gamma_{2} \circ \rho\left(\sup \|u(\tau)\|_{\tau \geq t_{0}}\right)\right\}, \tag{2.42}
\end{equation*}
$$

for all $t \geq t_{0}$. Define $B=\left[t_{0}, T\right)$ as the time interval before $x(t)$ enters $R_{t_{0}}$ for the first time. In view of the definition of $R_{t_{0}}$ we have

$$
\begin{equation*}
\dot{V} \leq-\gamma_{3} \circ \gamma_{2}^{-1}(V), \forall t \in B \tag{2.43}
\end{equation*}
$$

Then, there exists a class- $K L$ function $\beta_{v}$ such that $V(t) \leq \beta_{v}\left(V\left(t_{0}\right), t-t_{0}\right), \forall t \in$ $B$, which implies

$$
\begin{equation*}
\|x(t)\| \leq \gamma_{1}^{-1}\left(\beta_{v}\left(\gamma_{2}\left(\left\|x\left(t_{0}\right)\right\|, t-t_{0}\right)\right):=\beta\left(\left\|x\left(t_{0}\right)\right\|, t-t_{0}\right), \quad \forall t \in B\right. \tag{2.44}
\end{equation*}
$$

On the other hand, by (2.42), we conclude that

$$
\begin{equation*}
\|x(t)\| \leq \gamma_{1}^{-1} \circ \gamma_{2} \circ \rho\left(\sup \|u(\tau)\|_{\tau \geq t_{0}}\right):=\gamma\left(\sup \|u(\tau)\|_{\tau \geq t_{0}}\right) \tag{2.45}
\end{equation*}
$$

for all $t \in\left[t_{0}, \infty\right] \backslash B$. Then by (2.44) and (2.45),

$$
\begin{equation*}
\|x(t)\| \leq \beta\left(\left\|x\left(t_{0}\right)\right\|, t-t_{0}\right)+\gamma\left(\sup \|u(\tau)\|_{\tau \geq t_{0}}\right), \quad \forall t \geq t_{0} \geq 0 \tag{2.46}
\end{equation*}
$$

By causality, we have

$$
\begin{equation*}
\|x(t)\| \leq \beta\left(\left\|x\left(t_{0}\right)\right\|, t-t_{0}\right)+\gamma\left(\sup \|u(\tau)\|_{t_{0} \leq \tau \leq t}\right), \quad \forall t \geq t_{0} \geq 0 \tag{2.47}
\end{equation*}
$$

A function $V$ satisfying conditions (2.40) is called an ISS Lyapunov function.

### 2.3 Control Lyapunov Functions

This section presents an extension of the Lyapunov function concept, which is a useful tool to design an adaptive controller for nonlinear systems. Assuming that the problem is to design a feedback control law $\alpha(x)$ for the time-invariant system:

$$
\begin{equation*}
\dot{x}=f(x, u), x \in \mathbb{R}^{n}, u \in \mathbb{R}, f(0,0)=0, \tag{2.48}
\end{equation*}
$$

such that the equilibrium $x=0$ of the closed loop system:

$$
\begin{equation*}
\dot{x}=f(x, \alpha(x)) \tag{2.49}
\end{equation*}
$$

is globally asymptotically stable (GAS). We can take a function $V(x)$ as a Lyapunov candidate function, and require that its derivative along the solutions of (2.49) satisfy $\dot{V}(x) \leq-W(x)$, where $W(x)$ is a positive definite function. We therefore need to find $\alpha(x)$ to guarantee that for all $x \in \mathbb{R}^{n}$ such that

$$
\begin{equation*}
\frac{\partial V(x)}{\partial x} f(x, \alpha(x)) \leq-W(x) . \tag{2.50}
\end{equation*}
$$

This is a difficult problem. A stabilizing control law for (2.48) may exist but we may fail to satisfy (2.50) because of a poor choice of $V(x)$ and $W(x)$. A system for which a good choice of $V(x)$ and $W(x)$ exists is said to possess a control Lyapunov function (CLF). For systems affine in the control:

$$
\begin{equation*}
\dot{x}=f(x)+g(x) u, \quad f(0)=0, \tag{2.51}
\end{equation*}
$$

the CLF inequality (2.50) becomes

$$
\begin{equation*}
\frac{\partial V}{\partial x} f(x)+\frac{\partial V}{\partial x} g(x) \alpha(x) \leq-W(x) . \tag{2.52}
\end{equation*}
$$

If $V(x)$ is a CLF for (2.51), then a particular stabilizing control law $\alpha(x)$, smooth for all $x \neq 0$, is given by

$$
u=\alpha(x)= \begin{cases}-\frac{\frac{\partial V}{\partial x} f(x)+\sqrt{\left(\frac{\partial V}{\partial x} f(x)\right)^{2}+\left(\frac{\partial V}{\partial x} g(x)\right)^{4}}}{\frac{\partial V}{\partial x} g(x)}, & \frac{\partial V}{\partial x} g(x) \neq 0  \tag{2.53}\\ 0, & \frac{\partial V}{\partial x} g(x)=0\end{cases}
$$

It should be noted that (2.52) can be satisfied only if

$$
\begin{equation*}
\frac{\partial V}{\partial x} g(x)=0 \Rightarrow \frac{\partial V}{\partial x} f(x)<0, \quad \forall x \neq 0 \tag{2.54}
\end{equation*}
$$

and that in this case (2.53) gives

$$
\begin{equation*}
W(x)=\sqrt{\left(\frac{\partial V}{\partial x} f\right)^{2}+\left(\frac{\partial V}{\partial x} g\right)^{4}}>0, \quad \forall x \neq 0 \tag{2.55}
\end{equation*}
$$

The main drawback of the CLF concept as a design tool is that for most nonlinear systems a CLF is not known. The task of finding an appropriate CLF may be as complex as that of designing a stabilizing feedback law. The backstepping technique [3] presented in the next section gives a constructive way to construct an appropriate CLF for various systems.

### 2.4 Backstepping

Assumption 2.1. Consider the system

$$
\begin{equation*}
\dot{x}=f(x)+g(x) u, \quad f(0)=0 \tag{2.56}
\end{equation*}
$$

where $x \in \mathbb{R}^{n}$ is the state and $u \in \mathbb{R}$ is the control input. There exist a continuously differentiable feedback control law

$$
\begin{equation*}
u=\alpha(x), \alpha(0)=0 \tag{2.57}
\end{equation*}
$$

and a smooth, positive definite, radially unbounded function $V: \mathbb{R}^{n} \rightarrow \mathbb{R}$ such that

$$
\begin{equation*}
\frac{\partial V}{\partial x}[f(x)+g(x) \alpha(x)] \leq-W(x) \leq 0, \forall x \in \mathbb{R}^{n} \tag{2.58}
\end{equation*}
$$

where $W: \mathbb{R}^{n} \rightarrow \mathbb{R}$ is positive semidefinite.
It should be noted that under this assumption, the control law (2.57) applied to the system (2.56) guarantees global boundedness of $x(t)$, and the regulation of $W(x)$ : $\lim _{t \rightarrow \infty} W(x)=0$. If $W(x)$ is positive definite, the control law (2.57) renders the global asymptotic stable equilibrium of (2.56).

Theorem 2.6. Let the system (2.56) be augmented by an integrator:

$$
\begin{align*}
& \dot{x}=f(x)+g(x) \xi \\
& \dot{\xi}=u \tag{2.59}
\end{align*}
$$

and suppose that the first equation of (2.59) satisfies Assumption 2.1 with $\xi$ as its control.

1. If $W(x)$ is positive definite then

$$
\begin{equation*}
V_{a}=V(x)+\frac{1}{2}[\xi-\alpha(x)]^{2} \tag{2.60}
\end{equation*}
$$

is a CLF for the system (2.59), that is, there exists a feedback control law $u=$ $\alpha_{a}(x, \xi)$, which renders $x=0, \xi=0$ the GAS equilibrium of (2.59). One such control law choice is

$$
\begin{equation*}
u=-c(\xi-\alpha(x))+\frac{\partial \alpha}{\partial x}[f(x)+g(x) \xi]-\frac{\partial V}{\partial x} g(x), c>0 . \tag{2.61}
\end{equation*}
$$

2. If $W(x)$ is only positive semidefinite, then there exists a feedback control law that renders $\dot{V}_{a} \leq-W_{a}(x, \xi) \leq 0$, such that $W_{a}(x, \xi)>0$ whenever $W(x)>0$ or $\xi \neq \alpha(x)$. This guarantees global boundedness and convergence of $[x(t) \xi(t)]^{T}$ to the largest invariant set $M_{a}$ contained in the set

$$
E_{a}=\left\{\left.\left[\begin{array}{l}
x \\
\xi
\end{array}\right] \in \mathbb{R}^{n+1} \right\rvert\, W(x)=0, \xi=\alpha(x)\right\}
$$

Proof. We only prove the first part of the theorem. Proof of the second part is trivial. Introducing an error variable

$$
\begin{equation*}
z=\xi-\alpha(x) \tag{2.62}
\end{equation*}
$$

and differentiating with respect to time, (2.59) can be written as

$$
\begin{align*}
\dot{x} & =f(x)+g(x)[\alpha(x)+z], \\
\dot{z} & =u-\frac{\partial \alpha}{\partial x}[f(x)+g(x)(\alpha(x)+z)] . \tag{2.63}
\end{align*}
$$

Using (2.58), the first time derivative of (2.60) along the solutions of (2.63) satisfies

$$
\begin{equation*}
\dot{V}_{a} \leq-W(x)+z\left[u-\frac{\partial \alpha}{\partial x}(f(x)+g(x)(\alpha(x)+z))+\frac{\partial V}{\partial x} g(x)\right] . \tag{2.64}
\end{equation*}
$$

Any control law, such as (2.61), which renders $\dot{V}_{a} \leq-W_{a}(x, \xi) \leq-W(x)$, with $W_{a}$ positive definite in $z$, guarantees global boundedness of $x$ and $z$, and regulation of $W(x)$ and $z(t)$.

Corollary 2.1. Let the system (2.56) satisfy Assumption 2.1 with $\alpha(x)=\alpha_{0}(x)$ being augmented by a chain of $k$ integrators so that $u$ is replaced by $\xi_{1}$, the state of the last integrator in the chain is

$$
\begin{align*}
& \dot{x}=f(x)+g(x) \xi_{1}, \\
& \dot{\xi}_{1}=\xi_{2} \\
& \quad \vdots  \tag{2.65}\\
& \dot{\xi}_{k-1}=\xi_{k} \\
& \dot{\xi}_{k}=u .
\end{align*}
$$

For this system, repeated application of Theorem 2.6 with $\xi_{1}, \ldots, \xi_{k}$ as virtual controls, results in the Lyapunov function

$$
\begin{equation*}
V_{a}\left(x, \xi_{1}, \ldots, \xi_{k}\right)=V(x)+\frac{1}{2} \sum_{i=1}^{k}\left[\xi_{i}-\alpha_{i-1}\left(x, \xi_{1}, \ldots, \xi_{i-1}\right)\right]^{2} \tag{2.66}
\end{equation*}
$$

Any choice of feedback control which renders $\dot{V}_{a} \leq-W_{a}\left(x, \xi_{1}, \ldots, \xi_{k}\right) \leq 0$, with $W_{a}\left(x, \xi_{1}, \ldots, \xi_{k}\right)=0$ only if $W(x)=0$ and $\xi_{i} \neq \alpha_{i-1}\left(x, \xi_{1}, \ldots, \xi_{i-1}\right), i=1, \ldots, k$, guarantees that $\left[x^{T}(t), \xi_{1}(t), \ldots, \xi_{k}(t)\right]^{T}$ is globally bounded and converges to the largest invariant set $M_{a}$ contained in the set

$$
E_{a}=\left\{\left[x^{T}, \xi_{1}, \ldots, \xi_{k}\right]^{T} \in \mathbb{R}^{n+k} \mid W(x)=0, \xi_{i}=\alpha_{i-1}\left(x, \xi_{1}, \ldots, \xi_{i-1}\right)\right\}
$$

for all $i=1, \ldots, k$. Furthermore, if $W(x)$ is positive definite, that is, if $x=0$ can be rendered GAS through $\xi_{1}$, then (2.66) is a CLF for (2.65) and the equilibrium $x=0, \xi_{1}=0, \ldots, \xi_{k}=0$ can be rendered GAS through $u$.

Proof. See [3].

### 2.5 Stabilization Under Uncertainties

The power of adaptive control is exhibited in the presence of uncertain nonlinearities and unknown parameters. Such uncertainties in linear systems make the control design procedure difficult and become more serious in control of nonlinear systems. For nonlinear systems, the states can easily escape to infinity in a finite time. The following theorem introduces the use of a term in the control law called nonlinear damping to stabilize the system (2.56) in the presence of disturbance.

Theorem 2.7. Consider the system (2.56) satisfying Assumption 2.1 which is perturbed as

$$
\begin{equation*}
\dot{x}=f(x)+g(x)\left[u+\varphi(x)^{T} \Delta(x, u, t)\right], \tag{2.67}
\end{equation*}
$$

where $\varphi(x)$ is a $p \times 1$ ) vector of known smooth nonlinear functions, and $\Delta(x, u, t)$ is a $(p \times 1)$ vector of uncertain nonlinearities, which are uniformly bounded for all values of $x, u$, and $t$. If Assumption 2.1 satisfies with $W(x)$ being positive definite and radially unbounded, then the control

$$
\begin{equation*}
u=\alpha(x)-k \frac{\partial V}{\partial x}(x) g(x)\|\varphi(x)\|^{2}, \quad k>0 \tag{2.68}
\end{equation*}
$$

when applied to the system (2.67) renders the closed loop system ISS with respect to the disturbance input $\Delta(x, u, t)$ and hence guarantees global uniform boundedness of $x(t)$ and convergence to the residual set

$$
\begin{equation*}
R=\left\{x:\|x\| \leq \gamma_{1}^{-1} \circ \gamma_{2} \circ \gamma_{3}^{-1}\left(\frac{\|\Delta\|_{\infty}^{2}}{4 k}\right)\right\}, \tag{2.69}
\end{equation*}
$$

where $\gamma_{1}, \gamma_{2}, \gamma_{3}$ are class $K_{\infty}$ functions such that

$$
\begin{align*}
& \gamma_{1}(\|x\|) \leq V(x) \leq \gamma_{2}(\|x\|), \\
& \gamma_{3}(\|x\|) \leq W(x) . \tag{2.70}
\end{align*}
$$

Proof. By using (2.68) and (2.70), the first time derivative of $V(x)$ is

$$
\begin{align*}
\dot{V} & =\frac{\partial V}{\partial x}[f+g u]+\frac{\partial V}{\partial x} g \varphi^{T} \Delta \\
& =\frac{\partial V}{\partial x}[f+g \alpha]-k\left(\frac{\partial V}{\partial x} g\right)^{2}\|\varphi\|^{2}+\frac{\partial V}{\partial x} g \varphi^{T} \Delta \\
& \leq-W(x)-k\left(\frac{\partial V}{\partial x} g\right)^{2}\|\varphi\|^{2}+\frac{\partial V}{\partial x} g \varphi^{T} \Delta  \tag{2.71}\\
& \leq-W(x)-k\left(\frac{\partial V}{\partial x} g\right)^{2}\|\varphi\|^{2}+\left\|\frac{\partial V}{\partial x} g\right\|\|\varphi\|\|\Delta\|_{\infty} \leq-W(x)+\frac{\|\Delta\|_{\infty}^{2}}{4 k} .
\end{align*}
$$

From (2.71), it follows that $\dot{V}$ is negative whenever $W(x) \geq \frac{\|\Delta\|_{\infty}^{2}}{4 k}$. Combining this with the second equation of (2.70), we conclude that

$$
\begin{equation*}
\|x(t)\|>\gamma_{3}^{-1}\left(\frac{\|\Delta\|_{\infty}^{2}}{4 k}\right) \Rightarrow \dot{V}<0 \tag{2.72}
\end{equation*}
$$

This means that if $\|x(0)\| \leq \gamma_{3}^{-1}\left(\frac{\|\Delta\|_{\infty}^{2}}{4 k}\right)$, then

$$
\begin{equation*}
V(x(t)) \leq \gamma_{2} \circ \gamma_{3}^{-1}\left(\frac{\|\Delta\|_{\infty}^{2}}{4 k}\right), \tag{2.73}
\end{equation*}
$$

which implies that

$$
\begin{equation*}
\|x(t)\| \leq \gamma_{1}^{-1} \circ \gamma_{2} \circ \gamma_{3}^{-1}\left(\frac{\|\Delta\|_{\infty}^{2}}{4 k}\right) . \tag{2.74}
\end{equation*}
$$

If, on the other hand, $\|x(0)\| \geq \gamma_{3}^{-1}\left(\frac{\|\Delta\|_{\infty}^{2}}{4 k}\right)$, then $V(x) \leq V(x(0))$, which implies

$$
\begin{equation*}
\|x(t)\| \leq \gamma_{1}^{-1} \circ \gamma_{2}(\|x(0)\|) \tag{2.75}
\end{equation*}
$$

Combining (2.74) and (2.75) leads to global uniform boundedness of $x(t)$ :

$$
\begin{equation*}
\|x\|_{\infty} \leq \max \left\{\gamma_{1}^{-1} \circ \gamma_{2} \circ \gamma_{3}^{-1}\left(\frac{\|\Delta\|_{\infty}^{2}}{4 k}\right), \gamma_{1}^{-1} \circ \gamma_{2}(\|x(0)\|)\right\} \tag{2.76}
\end{equation*}
$$

while (2.72) and the first equation of (2.70) prove the convergence of $x(t)$ to the residual set defined in (2.69).

Combining the above theorem with Theorem 2.6 leads to the following corollary.
Corollary 2.2. Consider the following system

$$
\begin{equation*}
\dot{x}=f(x)+g(x) u+F(x) \Delta_{1}(x, u, t) \tag{2.77}
\end{equation*}
$$

where $x \in \mathbb{R}^{n}, u \in \mathbb{R}, F(x)$ is an $(n \times q)$ matrix of known smooth nonlinear functions, and $\Delta_{1}(x, u, t)$ is a $(q \times 1)$ vector of uncertain nonlinearities, which is uniformly bounded for all values of $x, u$ and $t$. Suppose that there exists a feedback control law $u=\alpha(x)$ that renders $x(t)$ globally uniformly bounded, and that this is established via positive definite and radially unbounded functions $V(x), W(x)$, and $a$ constant $b$ such that

$$
\begin{equation*}
\frac{\partial V}{\partial x}\left[f(x)+g(x) \alpha(x)+F(x) \Delta_{1}(x, u, t)\right] \leq-W(x)+b . \tag{2.78}
\end{equation*}
$$

Now consider the augmented system

$$
\begin{align*}
& \dot{x}=f(x)+g(x) \xi+F(x) \Delta_{1}(x, u, t) \\
& \dot{\xi}=u+\varphi^{T}(x, \xi) \Delta_{2}(x, \xi, u, t) \tag{2.79}
\end{align*}
$$

where $\varphi(x, \xi)$ is a $(p \times 1)$ vector of known smooth nonlinear functions, $\Delta_{2}(x, u, \xi, t)$ is a $(p \times 1)$ vector of uncertain nonlinearities, which are uniformly bounded for all values of $x, u$, $\xi$ and $t$. For this system, the feedback control law

$$
\begin{align*}
u= & -c(\xi-\alpha(x))+\frac{\partial \alpha}{\partial x}[f(x)+g(x) \xi]- \\
& \frac{\partial V}{\partial x} g(x)-k(\xi-\alpha(x))\left\{\|\varphi(x, \xi)\|^{2}+\left\|\frac{\partial \alpha}{\partial x} F(x)\right\|^{2}\right\} \tag{2.80}
\end{align*}
$$

guarantees global uniform boundedness of $x(t)$ and $\xi(t)$ with any $c>0$ and $k>0$.
Proof. See [3].

### 2.6 Barbalat-like Lemmas

This section presents lemmas that are useful in investigating the convergence of time-varying systems.

Lemma 2.4. (Barbalat's lemma) Consider the function $\phi: \mathbb{R}^{+} \rightarrow \mathbb{R}$. If $\phi$ is uniformly continuous and $\lim _{t \rightarrow \infty} \int_{0}^{t} \phi(\tau) \mathrm{d} \tau$ exists and is finite, then

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \phi(t)=0 . \tag{2.81}
\end{equation*}
$$

Proof. See [6].
Lemma 2.5. Assume that a nonnegative scalar differentiable function $f(t)$ enjoys the following conditions

$$
\begin{align*}
& \text { 1. }\left|\frac{\mathrm{d}}{\mathrm{~d} t} f(t)\right| \leq k_{1} f(t), \\
& \text { 2. } \int_{0}^{\infty} f(t) \mathrm{d} t \leq k_{2} \tag{2.8}
\end{align*}
$$

for all $t \geq 0$, where $k_{1}$ and $k_{2}$ are positive constants, then $\lim _{t \rightarrow \infty} f(t)=0$.
Proof. Integrating both sides of 1) in (2.82) gives

$$
\begin{align*}
& f(t) \leq f(0)+k_{1} \int_{0}^{t} f(s) \mathrm{d} s \leq f(0)+k_{1} k_{2}, \\
& f(t) \geq f(0)-k_{1} \int_{0}^{t} f(s) \mathrm{d} s \geq f(0)-k_{1} k_{2} \tag{2.83}
\end{align*}
$$

These inequalities imply that $f(t)$ is a uniform bounded function. From (2.83) and the second condition in (2.82), we have that $f(t)$ is also bounded on the half axis $[0, \infty]$, i.e., $f(t) \leq k_{3}$ with $k_{3}$ a positive constant. Hence $\left|\frac{\mathrm{d}}{\mathrm{d} t} f(t)\right| \leq k_{1} k_{3}$. Now assume that $\lim _{t \rightarrow \infty} f(t) \neq 0$. Then there exists a sequence of points $t_{i}$ and a positive constant $\epsilon$ such that $f\left(t_{i}\right) \geq \epsilon, t_{i} \rightarrow \infty, i \rightarrow \infty,\left|t_{i}-t_{i-1}\right|>2 \epsilon /\left(k_{1} k_{3}\right)$ and moreover $f(s) \geq \epsilon / 2, s \in L_{i}=\left[t_{i}-\epsilon /\left(2 k_{1} k_{3}\right), t_{i}+\epsilon /\left(2 k_{1} k_{3}\right)\right]$. Since the segments $L_{i}$ and $L_{j}$ do not intersect for any $i$ and $j$ with $i \neq j$, we have

$$
\begin{equation*}
\int_{0}^{\infty} f(t) \mathrm{d} t \geq \int_{0}^{T} f(t) \mathrm{d} t \geq \sum_{t_{i} \leq T} \int_{L_{i}} f(t) \mathrm{d} t \geq \frac{\epsilon}{2} \frac{\epsilon}{k_{1} k_{3}} M(T) \tag{2.84}
\end{equation*}
$$

where $M(T)$ is the number of points $t_{i}$ not exceeding $T$. Since $\lim _{T \rightarrow \infty} M(T)=\infty$, the integral $\int_{0}^{\infty} f(t) \mathrm{d} t$ is divergent. This contradicts Condition 2 in (2.82). This contradiction proves the lemma.

Remark 2.2. Lemma 2.5 is different from Barbalat's lemma 2.4. While Barbalat's lemma assumes that $f(t)$ is uniformly continuous, Lemma 2.5 assumes that $\left|\frac{\mathrm{d}}{\mathrm{d} t} f(t)\right|$ is bounded by $k_{1} f(t)$.

Lemma 2.6. Consider a scalar system

$$
\begin{equation*}
\dot{x}=-c x+p(t) \tag{2.85}
\end{equation*}
$$

where $c>0$ and $p(t)$ is a bounded and uniformly continuous function. If, for any initial time $t_{0} \geq 0$ and any initial condition $x\left(t_{0}\right)$, the solution $x(t)$ is bounded and converges to 0 as $t \rightarrow \infty$, then

$$
\begin{equation*}
\lim _{t \rightarrow \infty} p(t)=0 \tag{2.86}
\end{equation*}
$$

Proof. See [18].
Lemma 2.7. Consider a first-order differential equation of the form

$$
\begin{equation*}
\dot{x}=-\left(a(t)+f_{1}(\xi(t))\right) x+f_{2}(\xi(t)) \tag{2.87}
\end{equation*}
$$

where $f_{1}$ and $f_{2}$ are continuous functions, and $\xi:[0, \infty) \rightarrow \mathbb{R}^{m}$ is a time-varying vector-valued signal that exponentially converges to zero and, for all $t \geq t_{0} \geq 0$, satisfies

$$
\begin{equation*}
\left|f_{i}(\xi(t))\right| \leq \gamma_{i}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{i}\left(t-t_{0}\right)} \tag{2.88}
\end{equation*}
$$

where $\sigma_{i}>0, i=1,2$ and $\gamma_{i}$ are class- $K$ functions. If a $(t)$ enjoys the property that there is a constant $\sigma_{3}$ such that

$$
\begin{equation*}
\int_{t_{1}}^{t_{2}} a(\tau) \mathrm{d} \tau \geq \sigma_{3}\left(t_{2}-t_{1}\right), \forall t_{2} \geq t_{1} \geq 0 \tag{2.89}
\end{equation*}
$$

then there exist a class-K function $\gamma$ and a constant $\sigma>0$ such that

$$
\begin{equation*}
|x(t)| \leq \gamma\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right) e^{-\sigma\left(t-t_{0}\right)} \tag{2.90}
\end{equation*}
$$

Proof. See [19].

### 2.7 Controllability and Observability

### 2.7.1 Controllability and Observability of Linear Time-invariant Systems

This section deals with the controllability and observability properties of systems described by linear time-invariant state-space representations. In particular, consider a linear and time-invariant system defined by the state-space representation:

$$
\begin{align*}
\dot{x}(t) & =\boldsymbol{A} x(t)+\boldsymbol{B} u(t), \\
y(t) & =\boldsymbol{C} x(t)+\boldsymbol{D} u(t), \tag{2.91}
\end{align*}
$$

where $\boldsymbol{A}$ has $n$ distinct eigenvalues $\lambda_{1}, \lambda_{2}, \ldots, \lambda_{n}$, which define the poles and the corresponding modes $e^{\lambda_{i} t}$.

### 2.7.1.1 Controllability

Conditions for controllability of the system (2.91) are given in the following lemmas.

Lemma 2.8. The system (2.91) is controllable if and only if any one of the following (equivalent) conditions holds:

1. The controllability matrix

$$
\begin{equation*}
\boldsymbol{C}_{\boldsymbol{c}}:=\left[\boldsymbol{B}, \boldsymbol{A} \boldsymbol{B}, \ldots, \boldsymbol{A}^{n-1} \boldsymbol{B}\right] \tag{2.92}
\end{equation*}
$$

has full rank $n$.
2. The $n$ rows of $e^{\boldsymbol{A} t} \boldsymbol{B}$, where $e^{\boldsymbol{A} t}$ represents the unique state transition matrix of the system, are linearly independent over the real field $\mathbb{R}$ for all $t$.
3. The controllability Grammian

$$
\begin{equation*}
\boldsymbol{G}_{\boldsymbol{c}}\left(t_{0}, t_{f}\right):=\int_{t_{0}}^{t_{f}} e^{-\boldsymbol{A} \tau} \boldsymbol{B} \boldsymbol{B}^{\boldsymbol{T}} e^{-\boldsymbol{A}^{T} \tau} \mathrm{~d} \tau \tag{2.93}
\end{equation*}
$$

is nonsingular for all $t_{f}>t_{0}$.
4. The $n \times(n+m)$ matrix $[\lambda \boldsymbol{I}-\boldsymbol{A}, \boldsymbol{B}]$ has rank $n$ at all eigenvalues $\lambda_{i}$ of $\boldsymbol{A}$ or equivalently, $\lambda \boldsymbol{I}-\boldsymbol{A}$ and $\boldsymbol{B}$ are left comprime polynomial matrices.

Proof. See [20], Chapter 8.
Since the solution to the system (2.91) is given by

$$
\begin{equation*}
x(t)=e^{\boldsymbol{A}\left(t-t_{0}\right)} x\left(t_{0}\right)+\int_{t_{0}}^{t} e^{\boldsymbol{A}(t-\tau)} \boldsymbol{B} u(\tau) \mathrm{d} \tau \tag{2.94}
\end{equation*}
$$

it follows that the controllability Grammian-based control input

$$
\begin{equation*}
u(t)=\boldsymbol{B}^{\boldsymbol{T}} e^{-\boldsymbol{A}^{\boldsymbol{T}} t} \boldsymbol{G}_{\boldsymbol{c}}^{-1}\left(t_{0}, t_{f}\right)\left[e^{-\boldsymbol{A} t_{f}} x\left(t_{f}\right)-e^{-\boldsymbol{A} t_{0}} x\left(t_{0}\right)\right] \tag{2.95}
\end{equation*}
$$

transfers any initial state $x\left(t_{0}\right)$ to any arbitrarily chosen final state $x\left(t_{f}\right)$ at any arbitrary $t_{f}>0$. This observation is consistent with the more traditional definition of controllability.

### 2.7.1.2 Observability

Conditions for observability of the system (2.91) are given in the following lemmas.
Lemma 2.9. The system (2.91) is observable if and only if any one of the following (equivalent) conditions holds:

## 1. The observability matrix

$$
\boldsymbol{C}_{\boldsymbol{o}}:=\left[\begin{array}{c}
C  \tag{2.96}\\
C A \\
\vdots \\
C A^{n-1}
\end{array}\right]
$$

has full rank n.
2. The $n$ columns of $\boldsymbol{C} e^{\boldsymbol{A} t}$, where $e^{\boldsymbol{A} t}$ represents the unique state transition matrix of the system, are linearly independent over the real field $\mathbb{R}$ for all $t$.
3. The observability Grammian

$$
\begin{equation*}
\boldsymbol{G}_{\boldsymbol{o}}\left(t_{0}, t_{f}\right):=\int_{t_{0}}^{t_{f}} e^{\boldsymbol{A}^{T} \tau} \boldsymbol{C}^{T} \boldsymbol{C} e^{\boldsymbol{A} \tau} \mathrm{d} \tau \tag{2.97}
\end{equation*}
$$

is nonsingular for all $t_{f}>t_{0}$.
4. The $(n+p) \times n$ matrix $\left[\begin{array}{c}\lambda \boldsymbol{I}-\boldsymbol{A} \\ \boldsymbol{C}\end{array}\right]$ has rank $n$ at all eigenvalues $\lambda_{i}$ of $\boldsymbol{A}$ or equivalently $\lambda \boldsymbol{I}-\boldsymbol{A}$ and $\boldsymbol{C}$ are right comprime polynomial matrices.

Proof. See [20], Chapter 8.
If the system (2.91) is observable, it then follows that the system initial state can be determined by

$$
\begin{equation*}
x\left(t_{0}\right)=e^{\boldsymbol{A} t_{0}} \boldsymbol{G}_{\boldsymbol{o}}^{-1}\left(t_{0}, t_{f}\right) e^{\boldsymbol{A}^{T} t_{0}} \int_{t_{0}}^{t_{f}} e^{\boldsymbol{A}^{T}\left(t-t_{0}\right)} \boldsymbol{C}^{T} f(t) \mathrm{d} t \tag{2.98}
\end{equation*}
$$

where

$$
\begin{equation*}
f(t):=y(t)-\boldsymbol{C} \int_{t_{0}}^{t} e^{\boldsymbol{A}(t-\tau)} \boldsymbol{B} u(\tau) \mathrm{d} \tau-\boldsymbol{D} u(t) \tag{2.99}
\end{equation*}
$$

This observation is consistent with the more traditional definition of observability.

### 2.7.2 Controllability and Observability of Linear Time-varying Systems

This section deals with controllability and observability properties of the following $n$-dimensional linear time-varying system

$$
\begin{align*}
\dot{x}(t) & =\boldsymbol{A}(t) x(t)+\boldsymbol{B}(t) u(t), \\
y(t) & =\boldsymbol{C}(t) x(t)+\boldsymbol{D}(t) u(t) . \tag{2.100}
\end{align*}
$$

We assume that the matrices $\boldsymbol{A}(t), \boldsymbol{B}(t), \boldsymbol{C}(t)$, and $\boldsymbol{D}(t)$ are at least continuous functions of $t$.

### 2.7.2.1 Controllability

The system (2.100) is said to be controllable on the interval $\left[t_{0}, t_{f}\right]$, where $t_{f}>$ $t_{0} \geq 0$, if for any states $x_{0}$ and $x_{f}$, a continuous input $u(t)$ exists that drives the system to the state $x\left(t_{f}\right)=x_{f}$ at time $t=t_{f}$ starting from the state $x\left(t_{0}\right)=x_{0}$ at time $t=t_{0}$. Conditions for controllability of the system (2.100) are given in the following lemma:

Lemma 2.10. Assuming that the matrix $\boldsymbol{A}(t)$ is $n-2$ times differentiable and $\boldsymbol{B}(t)$ is $n-1$ times differentiable, a sufficient condition for the system (2.100) to be controllable on the interval $\left[t_{0}, t_{f}\right]$ is that the following matrix has rank $n$ for at least one $t \in\left[t_{0}, t_{f}\right]$ :

$$
\begin{equation*}
\boldsymbol{K}(t)=\left[\boldsymbol{K}_{0}(t) \boldsymbol{K}_{1}(t) \ldots \boldsymbol{K}_{n-1}(t)\right] \tag{2.101}
\end{equation*}
$$

where $\boldsymbol{K}_{i}, i=1, \ldots, n-1$ are defined by

$$
\begin{align*}
& \boldsymbol{K}_{0}(t)=\boldsymbol{B}(t) \\
& \boldsymbol{K}_{i}(t)=-\boldsymbol{A}(t) \boldsymbol{K}_{i-1}(t)+\dot{\boldsymbol{K}}_{i-1}(t) \tag{2.102}
\end{align*}
$$

Moreover, if the matrix $\boldsymbol{K}(t)$ has rank $n$ for all $t \in\left[t_{0}, t_{f}\right]$, then the system (2.100) is uniformly controllable on the interval $\left[t_{0}, t_{f}\right]$.

Proof. See [20], Chapter 25.

### 2.7.2.2 Observability

Now suppose that the control input is zero, i.e., the system (2.100) is given by

$$
\begin{align*}
\dot{x}(t) & =\boldsymbol{A}(t) x(t), \\
y(t) & =\boldsymbol{C}(t) x(t) . \tag{2.103}
\end{align*}
$$

The system (2.103) is said to be observable on the interval $\left[t_{0}, t_{f}\right]$ if any initial state $x\left(t_{0}\right)=x_{0}$ can be determined from the output $y(t)$ for $t \in\left[t_{0}, t_{f}\right]$. Conditions for observability of the system (2.103) are given in the following lemma.

Lemma 2.11. Assuming that the matrix $\boldsymbol{A}(t)$ is $n-2$ times differentiable and $\boldsymbol{B}(t)$ is $n-1$ times differentiable, a sufficient condition for the system (2.100) to be observable on the interval $\left[t_{0}, t_{f}\right]$ is that the following matrix has rank $n$ for at least one $t \in\left[t_{0}, t_{f}\right]$

$$
L(t)=\left[\begin{array}{c}
\boldsymbol{L}_{0}(t)  \tag{2.104}\\
\boldsymbol{L}_{1}(t) \\
\vdots \\
\boldsymbol{L}_{n-1}(t)
\end{array}\right],
$$

where $\boldsymbol{L}_{i}, i=1, \ldots, n-1$ are defined by

$$
\begin{align*}
& L_{0}(t)=C(t) \\
& L_{i}(t)=L_{i-1}(t) \boldsymbol{A}(t)+\dot{\boldsymbol{L}}_{i-1}(t) \tag{2.105}
\end{align*}
$$

Moreover, if the matrix $\boldsymbol{L}(t)$ has rank $n$ for all $t \in\left[t_{0}, t_{f}\right]$, then the system (2.100) is uniformly observable on the interval $\left[t_{0}, t_{f}\right]$.

Proof. See [20], Chapter 25.

### 2.7.3 Controllability and Observability of Nonlinear Systems

### 2.7.3.1 Controllability

This subsection deals with controllability of the following nonlinear system

$$
\begin{equation*}
\dot{x}=f(x)+\sum_{i=1}^{m} g_{i}(x) u_{i}, \quad x \in \Omega_{x} \subset \mathbb{R}^{n}, \tag{2.106}
\end{equation*}
$$

where $u=\left[u_{1}, u_{2}, \ldots, u_{m}\right]^{T} \in \Omega_{u} \subset \mathbb{R}^{m}$ is the input vector. The system (2.106) is controllable if there exists an admissible input vector $u(t)$ such that the state $x(t)$ can travel from an initial point $x\left(t_{0}\right)=x_{0} \in \Omega_{x}$ to $x\left(t_{f}\right) \in \Omega_{x}$ within a finite time interval $T=t_{f}-t_{0}$. The controllability reveals whether the control system has a set of "healthy" input channels through which the input can excite the states effectively to reach the destination $x_{f}$. Based on this interpretation, the controllability of (2.106) should clearly depend on the function forms of all $f(x)$ and $g_{i}(x)$. The controllability of the nonlinear system (2.106) is based on a useful mathematical concept called Lie algebra, which is defined as follows.

Definition 2.9. A Lie algebra over the real field $\mathbb{R}$ or the complex field $\mathbb{C}$ is a vector space $\mathbb{G}$ for which a bilinear map $(X, Y) \rightarrow[X, Y]$ is defined from $\mathbb{G} \times \mathbb{G} \rightarrow \mathbb{G}$ such that

$$
\begin{aligned}
& \text { 1. }[X, Y]=-[Y, X], X, Y \in \mathbb{G} \\
& \text { 2. }[X,[Y, Z]]+[Y,[Z, X]]+[Z,[X, Y]]=0 \text { for } X, Y, Z \in \mathbb{G} .
\end{aligned}
$$

According to this definition, a Lie algebra is a vector space where an operator [.,.] is installed. Such a general operator, conventionally called a Lie bracket, can be defined arbitrarily as long as it satisfies the preceding two specified conditions simultaneously. The first condition is often called a skew symmetric relation and obviously implies that $[X, X]=0$. The second condition is called the Jacobi identity, which reveals a closed loop cyclic relation among any three elements in a Lie algebra.

Let us now define a special Lie algebra $E$ that collects all $n$-dimensional differentiable vector fields in $\mathbb{R}^{n}$ along with a commutative derivative relation: For any two vector fields $f$ and $g \in \mathbb{R}^{n}$, which are functions of $x \in \mathbb{R}^{n}$,

$$
\begin{equation*}
[f, g]=\frac{\partial g}{\partial x} f-\frac{\partial f}{\partial x} g \tag{2.107}
\end{equation*}
$$

It can be immediately shown that this Lie bracket satisfies the two conditions of a Lie algebra.

To extend the above Lie bracket between two vector fields to higher order derivatives, a more compact notation may be defined based on an adjoint operator, that is, $[f, g]=\operatorname{ad}_{f} g$. This new notation treats the Lie bracket $[f, g]$ as vector field $g$ operated on by an adjoint operator $\operatorname{ad}_{f}=[f,$.$] . Therefore, for an n$-order Lie bracket ( $n>1$ ), one can simply write

$$
[f, \underbrace{\ldots}_{n}[f, g] \underbrace{\ldots}_{n}]=\operatorname{ad}_{f}^{n} g .
$$

For a general control system given by (2.106), we now specifically define a control Lie algebra $E_{c}$, which is spanned by all up to order $(n-1)$ Lie brackets among $f$ and $g_{1}$ through $g_{m}$ as

$$
\begin{equation*}
E_{c}=\operatorname{span}\left\{g_{1}, \ldots, g_{m}, \operatorname{ad}_{f} g_{1}, \ldots, \operatorname{ad}_{f} g_{m}, \ldots, \operatorname{ad}_{f}^{n-1} g_{1}, \ldots, \operatorname{ad}_{f}^{n-1} g_{m}\right\} \tag{2.108}
\end{equation*}
$$

With the control Lie algebra concept, we can show that the following theorem is true and is also a general effective testing criterion for system controllability.

Theorem 2.8. The control system (2.106) is controllable if and only if $\operatorname{dim}\left(E_{c}\right)=$ $\operatorname{dim}\left(\Omega_{x}\right)=n$.

Proof. See [2].
Note that because each element in $E_{c}$ is a function of $x$, the dimension of $E_{c}$ may be different from one point to another. Thus, if the preceding condition of dimension is valid only in a neighborhood of a point in $\Omega_{x} \subset \mathbb{R}^{n}$, we say that the system (2.106) is locally controllable. On the other hand, if the condition of dimension can cover all of region $\Omega_{x}$, then it is globally controllable. Moreover, it is not hard to show that Theorem 2.8 covers the controllability condition for a linear time-invariant system (2.91).

### 2.7.3.2 Observability

We address the observability for the following nonlinear system

$$
\begin{align*}
& \dot{x}=f(x), \\
& y=h(x), \tag{2.109}
\end{align*}
$$

where $y \in \Omega_{y} \subset \mathbb{R}^{p}$ is the output vector. This system is said to be observable if for each pair of distinct states $x_{1}$ and $x_{2}$, the corresponding outputs $y_{1}$ and $y_{2}$ are also distinguishable. Clearly, the observability can be interpreted as a testing criterion to check whether the entire system has sufficient output channels to measure (or
observe) each internal state change. Intuitively, the observability should depend on the function forms of both $f(x)$ and $h(x)$.

We now introduce a Lie derivative, which is virtually a directional derivative for a scalar field $\lambda(x)$, with $x \in \mathbb{R}^{n}$ along the direction of an $n$-dimensional vector field $f(x)$. The mathematical expression is given as

$$
\begin{equation*}
L_{f} \lambda(x)=\frac{\partial \lambda(x)}{\partial x} f(x) \tag{2.110}
\end{equation*}
$$

Since $\frac{\partial \lambda(x)}{\partial x}$ is a $1 \times n$ gradient vector of the scalar $\lambda(x)$ and the norm of a gradient vector represents the maximum rate of function value changes, the product of the gradient and the vector field $f(x)$ in (2.109) becomes the directional derivative of $\lambda(x)$ along $f(x)$. Therefore, the Lie derivative of a scalar field defined by (2.110) is also a scalar field.

If each component of a vector field $h(x) \in \mathbb{R}^{p}$ is considered to take a Lie derivative along $f(x) \in \mathbb{R}^{n}$, then all components can be acted on concurrently and the result is a vector field that has the same dimension as $h(x)$; its $i$ th element is the Lie derivative of the $i$ th component of $h(x)$. Namely, if $h(x)=\left[h_{1}(x), \ldots, h_{p}(x)\right]^{T}$ and each component $h_{i}(x), i=1, \ldots, p$ is a scalar field, then the Lie derivative of the vector field $h(x)$ is defined as

$$
L_{f} h(x)=\left[\begin{array}{c}
L_{f} h_{1}(x)  \tag{2.111}\\
\vdots \\
L_{f} h_{p}(x)
\end{array}\right]
$$

With the Lie derivative concept, we now define an observation space $\Omega_{o}$ over $\mathbb{R}^{n}$ as

$$
\begin{equation*}
\Omega_{o}=\operatorname{span}\left\{h(x), L_{f} h(x), \ldots, L_{f}^{n-1} h(x)\right\} \tag{2.112}
\end{equation*}
$$

In other words, this space is spanned by all up to order $(n-1)$ Lie derivatives of the output function $h(x)$. Then, we further define an observability distribution, denoted by $d \Omega_{o}$, which collects the "gradient" vector of every component in $\Omega_{o}$. Namely,

$$
\begin{equation*}
d \Omega_{o}=\operatorname{span}\left\{\left.\frac{\partial \phi}{\partial x} \right\rvert\, \phi \in \Omega_{o}\right\} . \tag{2.113}
\end{equation*}
$$

With these definitions, we can present the following theorem for testing the observability.

Theorem 2.9. The system (2.109) is observable if and only if $\operatorname{dim}\left(d \Omega_{o}\right)=n$.
Proof. See [2].
Similarly to the controllability case, this testing criterion also has locally observable and globally observable cases, depending on whether the condition of dimension in the theorem is valid only in a neighborhood of a point or over the entire state region.

### 2.7.4 Brockett's Theorem on Feedback Stabilization

The following theorem, which is due to Brockett [21], gives a necessary condition for the existence of a stabilizing control law for the system

$$
\begin{equation*}
\dot{x}=f(x, u) \tag{2.114}
\end{equation*}
$$

at an equilibrium point $x_{0}$ with $x$ being the state and $u$ being the control input.
Theorem 2.10. Let the system (2.114) be given with $f\left(x_{0}, 0\right)=0$ and $f(\cdot, \cdot)$ continuously differentiable in a neighborhood of $\left(x_{0}, 0\right)$. A necessary condition for the existence of a continuously differentiable control law that makes $\left(x_{0}, 0\right)$ asymptotically stable is that

1. the linearized system should have no uncontrollable modes associated with eigenvalues whose real part is positive,
2. there exists a neighborhood $\Omega$ of $\left(x_{0}, 0\right)$ such that for each $\xi \in \Omega$ there exists a control $u_{\xi}(\cdot)$ defined on $[0, \infty)$ such that this control steers the solution of $\dot{x}=f(x, u \xi)$ from $x=\xi$ at $t=0$ to $x=x_{0}$ at $t=\infty$,
3. the mapping

$$
\gamma: A \times \mathbb{R}^{m} \rightarrow \mathbb{R}^{n}
$$

defined by $\gamma:(x, u) \rightarrow f(x, u)$ should be onto an open set containing 0 .
Proof. See [21].
Remark 2.3. If the system (2.114) is of the form

$$
\begin{equation*}
\dot{x}=f(x)+\sum_{i=1}^{m} u_{i} g_{i}(x) ; \quad x(t) \in \Omega \subset \mathbb{R}^{n}, \tag{2.115}
\end{equation*}
$$

then condition 3 of Theorem 2.10 implies that the stabilization problem cannot have a solution if there is a smooth distribution $D$ containing $f(\cdot)$ and $g_{1}(\cdot), \ldots, g_{m}(\cdot)$ with $\operatorname{dim} D<n$. One further special case: If the system (2.114) is of the form

$$
\begin{equation*}
\dot{x}=\sum_{i=1}^{m} u_{i} g_{i}(x) ; \quad x(t) \in \Omega \subset \mathbb{R}^{n} \tag{2.116}
\end{equation*}
$$

with the vectors $g_{i}(x)$ being linearly independent at $x_{0}$, then there exists a solution to the stabilization problem if and only if $m=n$. In this case, we must have as many control parameters as we have dimensions of $x$.

### 2.8 Conclusions

This chapter briefly provides the fundamental tools that will be used for control design and stability analysis of underactuated vehicles in the coming chapters. The
interested reader is referred to [3-6,22-24] for a detailed and comprehensive coverage of the topics discussed in this chapter.

# Part II <br> Modeling and Control Properties of Ocean Vessels 

## Chapter 3 <br> Modeling of Ocean Vessels

In this chapter, we classify the basic motion tasks for ocean vessels and their mathematical models, which will be used for the design of various control systems in the subsequent chapters.

### 3.1 Introduction

In automatic control, feedback improves system performance by allowing the successful completion of a task even in the presence of external disturbances and initial errors, and inaccuracy of the system parameters. To this end, real-time sensor measurements are used to reconstruct the vehicle state. Throughout this study, the latter is assumed to be available at every instant, as provided by local/global position and orientation measurement sensors. In some cases, we also assume that the vehicle velocities are measurable or constructible from position measurements.

We will concentrate on the case of a vessel workspace free of obstacles. In fact, we implicitly consider the vessel controller to be embedded in a hierarchical architecture in which a higher-level planner solves the obstacle avoidance problem and provides a series of motion goals to the lower control layer. In this perspective, the controller deals with the basic issue of converting ideal plans into actual motion execution. The nonholonomic nature of the ocean vessels is related to the fact that the vessel does not usually have independent actuators in the sway and heave axes. This implies the presence of a nonintegrable set of second-order differential constraints on the configuration variables. While these nonholonomic constraints reduce the instantaneous motions that the vessel can perform, they still allow almost global controllability in the configuration space. This feature leads to some challenging problems in the synthesis of feedback controllers, which parallel the new research issues arising in nonholonomic motion planning. Indeed, the ocean vessel application has triggered the search for innovative types of feedback controllers that can be used also for more general nonlinear systems that describe the motion of
more complicated vessel systems such as ocean vessels and air vehicles working in a group.

### 3.2 Basic Motion Tasks

In order to derive the most suitable feedback controllers for each case, it is convenient to classify the possible motion tasks as follows:

- Point-to-point motion: The vessel must reach a desired goal configuration starting from a given initial configuration, see Figure 3.1a.
- Path-following: The vessel must reach and follow a geometric reference path in the Cartesian space starting from a given initial configuration (on or off the path), see Figure 3.1b.
- Trajectory-tracking and path-tracking: The vessel must reach and follow a reference trajectory/path in the Cartesian space (i.e., a geometric path with an associated timing law) starting from a given initial configuration (on or off the trajectory/path), see Figure 3.1c. Trajectory-tracking is referred to as the case where the reference trajectory is generated by a suitable virtual vessel whereas the reference path is not required to be generated by a virtual vessel for the pathtracking.

The above tasks for an ocean vessel are sketched in Figure 3.1. Execution of these tasks can be achieved using either feedforward commands, or feedback control, or a combination of the two. Indeed, feedback solutions exhibit an intrinsic degree of robustness.

Using a more control-oriented terminology, the point-to-point motion task is a stabilization problem for a (equilibrium) point in the vessel state space. When using a feedback strategy, the point-to-point motion task leads to a state regulation control problem for a point in the vessel state space. Posture stabilization is another frequently used term. Without loss of generality, the goal can be taken as the origin of the $n$-dimensional vessel configuration space. Contrary to the usual situation, trajectory-tracking, path-tracking, and path-following are easier than regulation for a nonholonomic vessel. An intuitive explanation of this can be given in terms of a comparison between the number of controlled variables (outputs) and the number of control inputs. For the ship or underwater vehicle moving in a horizontal plane, two input commands are available while three variables (position and orientation) are needed to determine its configuration. Thus, regulation of the surface ship or the underwater vehicle in a horizontal position to a desired configuration implies zeroing three independent configuration errors.

In the path-following task, the controller is given a geometric description of the assigned Cartesian path. This information is usually available in a parameterized form expressing the desired motion in terms of a path parameter, which may be in particular the arc length along the path. For this task, time dependence is not relevant because one is concerned only with the geometric displacement between the vessel


Figure 3.1 Basic motion tasks for an ocean vessel
and the path. In this context, the time evolution of the path parameter is usually free and, accordingly, the command inputs can be arbitrarily scaled with respect to time without changing the resulting vessel path. It is then customary to set the vessel forward velocity (one of the inputs) to an arbitrary constant or time-varying value, leaving the other input variables for control. The path-following problem is thus rephrased as the stabilization to zero of a suitable scalar path error function using only the rest of the control inputs.

In the trajectory-tracking and path-tracking tasks, the vessel must follow the desired Cartesian path with a specified timing law. Although the reference trajectory/path can be split into a parameterized geometric path and a timing law for the parameter, such separation is not strictly necessary. Often, it is simpler to specify the workspace trajectory as the desired time evolution for the position of some representative point of the vessel. The trajectory-tracking and path-tracking problems consist then in the stabilization to zero of the Cartesian errors using all the available control inputs.

The point stabilization problem can be formulated in a local or in a global sense, the latter meaning that we allow for initial configurations that are arbitrarily far from the destination. The same is true also for path-following, trajectory-tracking, and path-tracking, although locality has two different meanings in these tasks. For
path-following, a local solution means that the controller works properly, provided that we start sufficiently close to the path; for trajectory-tracking and path-tracking, closeness should be evaluated with respect to the current position of the reference vessel and a current reference point, which moves on the reference path with a specified time law, on the reference path, respectively. The amount of information that should be provided by a high-level motion planner varies for each control task. In point-to-point motion, information is reduced to a minimum (i.e., the goal configuration only) when a globally stabilizing feedback control solution is available. However, if the initial error is large, such a control may produce erratic behavior and/or large control effort, which are unacceptable in practice. On the other hand, a local feedback solution requires the definition of intermediate subgoals at the task planning level in order to get closer to the final desired configuration. For the other motion tasks, the planner should provide a path that is kinematically feasible (namely, that complies with the nonholonomic constraints of the specific vessel), so as to allow its perfect execution in nominal conditions. While for a fully or overactuated vessel in which any path is feasible, some degree of geometric smoothness is in general required for nonholonomic vessels. Nevertheless, the intrinsic feedback structure of the driving commands enables it to recover transient errors due to isolated path discontinuities. Note also that the infeasibility arising from a lack of continuity in some higher-order derivative of the path may be overcome by appropriate motion timing. For example, paths with discontinuous curvature (like the Reeds and Shepp optimal paths under maximum curvature constraint) can be executed by choosing an appropriate point on the vessel provided that the vessel is allowed to stop, whereas paths with discontinuous tangent are not feasible. In this analysis, the selection of the vessel representative point for path/trajectory planning is critical. The timing profile is the additional item needed in trajectory-tracking and path-tracking control tasks. This information is seldom provided by current motion planners, also because the actual dynamics of the specific vessel are typically neglected at this level. The above example suggests that it may be reasonable at the planning stage to enforce requirements such as "move slower where the path curvature is higher".

### 3.3 Modeling of Ocean Vessels

Modeling of the ocean vessels is usually based on mechanics, principles of statics and dynamics. Statics is concerned with the equilibrium of bodies at rest or moving with a constant velocity. Dynamics deals with bodies having accelerated motion resulting from disturbances or/and control forces. Since we are interested in a mathematical model of the ocean vessels for the purpose of designing the control systems, this section focuses on dynamics of the vessels rather than statics. The following briefly presents the ocean vessel equations of motion based on the results in [11]. The resulting nonlinear model presented in this section is mainly intended for designing control systems in the next chapters. For a detailed and comprehensive derivation of the model, the reader is referred to [ $11,12,25]$. The physical and control
properties of the model are also presented for control design and stability analysis. In this section, we use the notation, see Table 3.1 and Figure 3.2, that complies with the Society of Naval Architects and Marine Engineers (SNAME) [26].


Figure 3.2 Motion variables for an ocean vessel

For an ocean vessel moving in six degrees of freedom, six independent coordinates are required to determine its position and orientation. The first three coordinates $(x, y, z)$ and their first time derivatives correspond to the position and translational motion along the $x$-, $y$ - and $z$-axes, while the last three coordinates $(\phi, \theta, \psi)$ and their first time derivatives describe orientation and rotational motion.

Table 3.1 SNAME Notation for ocean vessels

| Degree of freedom |  | Force and <br> moment | Linear and <br> angular velocity | Position and <br> Euler angles |
| :---: | :---: | :---: | :---: | :---: |
| 1 | Surge | $X$ | $u$ | $x$ |
| 2 | Sway | $Y$ | $v$ | $y$ |
| 3 | Heave | $Z$ | $w$ | $z$ |
| 4 | Roll | $K$ | $p$ | $\phi$ |
| 5 | Pitch | $M$ | $q$ | $\theta$ |
| 6 | Yaw | $N$ | $r$ | $\psi$ |

According to SNAME, the six different motion components are defined as surge, sway, heave, roll, pitch, and yaw. To determine the equations of motion, two reference frames are considered: the inertial or fixed to earth frame $O_{E} X_{E} Y_{E} Z_{E}$ that may be taken to coincide with the vessel fixed coordinates in some initial condition and the body-fixed frame $O_{b} X_{b} Y_{b} Z_{b}$ see Figure 3.2. Since the motion of the Earth hardly affects ocean vessels (different from air vehicles), the earth-fixed frame $O_{E} X_{E} Y_{E} Z_{E}$ can be considered to be inertial. For ocean vessels in general, the
most commonly adopted position for the body-fixed frame is such that it gives hull symmetry about the $O_{b} X_{b} Z_{b}$-plane and approximate symmetry about the $O_{b} Y_{b} Z_{b^{-}}$ plane. In this sense, the body axes $O_{b} X_{b}, O_{b} Y_{b}$, and $O_{b} Z_{b}$ coincide with the principal axes of inertia and are usually defined as follows: $O_{b} X_{b}$ is the longitudinal axis (directed from aft to fore); $O_{b} Y_{b}$ is the transverse axis (directed to starboard); and $O_{b} Z_{b}$ is normal axis (directed from top to bottom. Based on the notion in Table 3.1, the general motion of an ocean vessel can be described by the following vectors:

$$
\begin{array}{lll}
\boldsymbol{\eta}=\left[\begin{array}{ll}
\boldsymbol{\eta}_{\mathbf{1}} & \eta_{\mathbf{2}}
\end{array}\right]^{T}, & \boldsymbol{\eta}_{\mathbf{1}}=\left[\begin{array}{lll}
x & y & z
\end{array}\right]^{T}, & \boldsymbol{\eta}_{\mathbf{2}}=\left[\begin{array}{lll}
\phi & \theta & \psi
\end{array}\right]^{T}, \\
\boldsymbol{v}=\left[\begin{array}{ll}
\boldsymbol{v}_{\mathbf{1}} & \boldsymbol{v}_{\mathbf{2}}
\end{array}\right]^{T}, & \boldsymbol{v}_{\mathbf{1}}=\left[\begin{array}{lll}
l & v & w
\end{array}\right]^{T}, & \boldsymbol{v}_{\mathbf{2}}=\left[\begin{array}{lll}
p & q & r
\end{array}\right]^{T}, \\
\boldsymbol{\tau}=\left[\begin{array}{lll}
\boldsymbol{\tau}_{\mathbf{1}} & \boldsymbol{\tau}_{\mathbf{2}}
\end{array}\right]^{T}, & \boldsymbol{\tau}_{\mathbf{1}}=\left[\begin{array}{lll}
X & Y & Z
\end{array}\right]^{T}, & \boldsymbol{\tau}_{\mathbf{2}}=\left[\begin{array}{lll}
K & M & N
\end{array}\right]^{T},
\end{array}
$$

where $\eta$ denotes the position and orientation vector with coordinates in the earthfixed frame, $\boldsymbol{v}$ denotes the linear and angular velocity vector with coordinates in the body-fixed frame, and $\boldsymbol{\tau}$ denotes the forces and moments acting on the vessel in the body-fixed frame.

In deriving equations of motion of the ocean vessels, we divide the study of vessel dynamics into two parts kinematics, which treats only geometrical aspects of motion, and kinetics, which is the analysis of the forces resulting in the motion.

### 3.3.1 Kinematics

The first time derivative of the position vector $\eta_{1}$ is related to the linear velocity vector $\boldsymbol{v}_{\mathbf{1}}$ via the following transformation:

$$
\begin{equation*}
\dot{\eta}_{1}=J_{1}\left(\eta_{2}\right) v_{1}, \tag{3.1}
\end{equation*}
$$

where $\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ is a transformation matrix, which is related through the functions of the Euler angles: roll $(\phi)$, pitch $(\theta)$, and yaw $(\psi)$. This matrix is given by

$$
\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)=\left[\begin{array}{cc}
\cos (\psi) \cos (\theta) & -\sin (\psi) \cos (\phi)+\sin (\phi) \sin (\theta) \cos (\psi) \\
\sin (\psi) \cos (\theta) & \cos (\psi) \cos (\phi)+\sin (\phi) \sin (\theta) \sin (\psi)  \tag{3.2}\\
-\sin (\theta) & \sin (\phi) \cos (\theta) \\
& \sin (\psi) \sin (\phi)+\sin (\theta) \cos (\psi) \cos (\phi) \\
& -\cos (\psi) \sin (\phi)+\sin (\theta) \sin (\psi) \cos (\phi) \\
& \cos (\phi) \cos (\theta)
\end{array}\right] .
$$

It is noted that the matrix $J_{1}\left(\eta_{2}\right)$ is globally invertible since $J_{1}^{-1}\left(\eta_{2}\right)=J_{1}^{T}\left(\eta_{2}\right)$.
On the other hand, the first time derivative of the Euler angle vector $\boldsymbol{\eta}_{\mathbf{2}}$ is related to the body-fixed velocity vector $\boldsymbol{v}_{\mathbf{2}}$ through the following transformation:

$$
\begin{equation*}
\dot{\eta}_{2}=J_{2}\left(\eta_{2}\right) v_{2}, \tag{3.3}
\end{equation*}
$$

where the transformation matrix $\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{2}\right)$ is given by

$$
\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)=\left[\begin{array}{ccc}
1 & \sin (\phi) \tan (\theta) & \cos (\phi) \tan (\theta)  \tag{3.4}\\
0 & \cos (\phi) & -\sin (\phi) \\
0 & \sin (\phi) / \cos (\theta) & \cos (\phi) / \cos (\theta)
\end{array}\right]
$$

Note that the transformation matrix $\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ is singular at $\theta= \pm \frac{\pi}{2}$. However, during practical operations ocean vessels are not likely to enter the neighborhood of $\theta=$ $\pm \frac{\pi}{2}$ because of the metacentric restoring forces. For the case where it is essential to consider a region containing $\theta= \pm \frac{\pi}{2}$, a four-parameter description based on Euler parameters can be used instead. The interested reader is referred to [12] for more details. Combining (3.1) and (3.3) results in the kinematics of the ocean vessels:

$$
\left[\begin{array}{c}
\dot{\eta}_{1}  \tag{3.5}\\
\dot{\eta}_{2}
\end{array}\right]=\left[\begin{array}{cc}
\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{2}\right) & 0_{3 \times 3} \\
0_{3 \times 3} & \boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{2}\right)
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right] \Leftrightarrow \quad \dot{\boldsymbol{\eta}}=\boldsymbol{J}(\eta) v .
$$

### 3.3.2 Kinetics

### 3.3.2.1 Rigid Body Equations of Motion

Let us define the following vectors:

- $\boldsymbol{f}_{O b}=\left[\begin{array}{lll}X & Y & Z\end{array}\right]^{T}$ : force decomposed in the body-fixed frame.
- $\boldsymbol{m}_{O b}=\left[\begin{array}{lll}K & M & N\end{array}\right]^{T}$ : moment decomposed in the body-fixed frame.
- $\boldsymbol{v}_{O b}=\left[\begin{array}{ll}u & v\end{array}\right]^{T}$ : linear velocity decomposed in the body-fixed frame.
- $\omega_{O b}^{E}=\left[\begin{array}{lll}p & q & r\end{array}\right]^{T}$ : angular velocity of the body-fixed frame relative to the earthfixed frame.
- $\boldsymbol{r}_{O b}=\left[\begin{array}{lll}x_{g} & y_{g} & z_{g}\end{array}\right]^{T}$ : vector from $O_{b}$ to $C G$ (center of gravity of the vessel) decomposed in the body-fixed frame.

By the Newton-Euler formulation for a rigid body with a mass of $m$, we have the following balancing forces and moments:

$$
\begin{align*}
& m\left[\dot{\boldsymbol{v}}_{O b}+\dot{\boldsymbol{\omega}}_{O b}^{E} \times \boldsymbol{r}_{O b}+\boldsymbol{\omega}_{O b}^{E} \times \boldsymbol{v}_{O b}+\boldsymbol{\omega}_{O b}^{E} \times\left(\boldsymbol{\omega}_{O b}^{E} \times \boldsymbol{r}_{O b}\right)\right]=\boldsymbol{f}_{O b}, \\
& \boldsymbol{I}_{o} \dot{\boldsymbol{\omega}}_{O b}^{E}+\boldsymbol{\omega}_{O b}^{E} \times \boldsymbol{I}_{o} \boldsymbol{\omega}_{O b}^{E}+m \boldsymbol{r}_{O b} \times\left(\dot{\boldsymbol{v}}_{O b}+\boldsymbol{\omega}_{O b}^{E} \times \boldsymbol{v}_{O b}\right)=\boldsymbol{m}_{O b}, \tag{3.6}
\end{align*}
$$

where $\boldsymbol{I}_{o}$ is the inertia matrix about $O_{b}$ defined by

$$
\boldsymbol{I}_{\boldsymbol{o}}=\left[\begin{array}{ccc}
I_{x} & -I_{x y} & -I_{x z}  \tag{3.7}\\
-I_{y x} & I_{y} & -I_{y z} \\
-I_{z x} & -I_{z y} & I_{z}
\end{array}\right] .
$$

Here $I_{x}, I_{y}$, and $I_{z}$ are the moments of inertia about the $O_{b} X_{b}, O_{b} Y_{b}$, and $O_{b} Z_{b}$ axes, and $I_{x y}=I_{y x}, I_{x z}=I_{z x}$, and $I_{y z}=I_{z y}$ are the products of inertia. These quantities are defined as

$$
\begin{align*}
& I_{x}=\int_{V}\left(y^{2}+z^{2}\right) \rho_{m} \mathrm{~d} V, I_{x y}=\int_{V} x y \rho_{m} \mathrm{~d} V \\
& I_{y}=\int_{V}\left(x^{2}+z^{2}\right) \rho_{m} \mathrm{~d} V, I_{x z}=\int_{V} x z \rho_{m} \mathrm{~d} V  \tag{3.8}\\
& I_{z}=\int_{V}\left(x^{2}+y^{2}\right) \rho_{m} \mathrm{~d} V, I_{z y}=\int_{V} z y \rho_{m} \mathrm{~d} V
\end{align*}
$$

where $\rho_{m}$ and $V$ are, respectively, the mass density and the volume of the rigid body. Substituting the definitions of $\boldsymbol{f}_{O b}, \boldsymbol{m}_{O b}, \boldsymbol{v}_{O b}, \boldsymbol{\omega}_{O b}^{E}$, and $\boldsymbol{r}_{O b}$ into (3.6) results in the following equations of motion of a rigid body:

$$
\begin{equation*}
M_{R B} \dot{v}+C_{R B}(v) v=\tau_{R B}, \tag{3.9}
\end{equation*}
$$

where $\boldsymbol{v}=[u v w p q r]^{T}$ is the generalized velocity vector decomposed in the body-fixed frame, $\boldsymbol{\tau}_{\boldsymbol{R} \boldsymbol{B}}=\left[\begin{array}{ll}X Y Z K M N\end{array}\right]^{T}$ is the generalized vector of external forces and moments, the rigid body system inertia matrix $\boldsymbol{M}_{\boldsymbol{R} \boldsymbol{B}}$ is given by

$$
\boldsymbol{M}_{\boldsymbol{R} \boldsymbol{B}}=\left[\begin{array}{cccccc}
m & 0 & 0 & 0 & m z_{g} & -m y_{g}  \tag{3.10}\\
0 & m & 0 & -m z_{g} & 0 & m x_{g} \\
0 & 0 & m & m y_{g} & -m x_{g} & 0 \\
0 & -m z_{g} & m y_{g} & I_{x} & -I_{x y} & -I_{x z} \\
m z_{g} & 0 & -m x_{g} & -I_{y x} & I_{y} & -I_{y z} \\
-m y_{g} & m x_{g} & 0 & -I_{z x} & -I_{z y} & I_{z}
\end{array}\right],
$$

and the rigid body Coriolis and centripetal matrix $\boldsymbol{C}_{\boldsymbol{R} \boldsymbol{B}}(\boldsymbol{v})$ is given by

$$
\begin{align*}
\boldsymbol{C}_{\boldsymbol{R} \boldsymbol{B}}(\boldsymbol{v})= & {\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
0 & 0 & 0 \\
-m\left(y_{g} q+z_{g} r\right) & m\left(y_{g} p+w\right) & m\left(z_{g} p-v\right) \\
m\left(x_{g} q-w\right) & -m\left(z_{g} r+x_{g} p\right) & m\left(z_{g} q+u\right) \\
m\left(x_{g} r+v\right) & m\left(y_{g} r-u\right) & -m\left(x_{g} p+y_{g} q\right) \\
& & \\
m\left(y_{g} q+z_{g} r\right) & -m\left(x_{g} q-w\right) & -m\left(x_{g} r+v\right) \\
-m\left(y_{g} p+w\right) & m\left(z_{g} r+x_{g} p\right) & -m\left(y_{g} r-u\right) \\
-m\left(z_{g} p-v\right) & -m\left(z_{g} q+u\right) & m\left(x_{g} p+y_{g} q\right) \\
0 & -I_{y z} q-I_{x z} p+I_{z} r & I_{y z} r+I_{x y} p-I_{y} q \\
& 0 & -I_{x z} r-I_{x y} q+I_{x} p \\
I_{y z} q+I_{x z} p-I_{z} r & 0 & 0
\end{array}\right.} \tag{3.11}
\end{align*}
$$

The generalized external force and moment vector, $\boldsymbol{\tau}_{\boldsymbol{R} \boldsymbol{B}}$, is a sum of hydrodynamic force and moment vector $\boldsymbol{\tau}_{\boldsymbol{H}}$, external disturbance force and moment vector $\boldsymbol{\tau}_{\boldsymbol{E}}$, and propulsion force and moment vector $\boldsymbol{\tau}$. Each of these vectors is detailed in the following sections.

### 3.3.2.2 Hydrodynamic Forces and Moments

In hydrodynamics, it is usually assumed that the hydrodynamic forces and moments on a rigid body can be linearly superimposed, see [27]. The hydrodynamic forces and moments are forces and moments on the body when the body is forced to oscillate with the wave excitation frequency and there are no incident waves. These forces and moments can be identified as the sum of three components: (1) added mass due to the inertia of the surrounding fluid, (2) radiation-induced potential damping due to the energy carried away by generated surface waves, and (3) restoring forces due to Archimedian forces (weight and buoyancy). The hydrodynamic force and moment vector $\boldsymbol{\tau}_{\boldsymbol{H}}$ is given by

$$
\begin{equation*}
\tau_{H}=-M_{A} \dot{v}-C_{A}(v) v-D(v) v-g(\eta) \tag{3.12}
\end{equation*}
$$

where $\boldsymbol{M}_{\boldsymbol{A}}$ is the added mass matrix, $\boldsymbol{C}_{\boldsymbol{A}}(\boldsymbol{v})$ is the hydrodynamic Coriolis and centripetal matrix, $\boldsymbol{D}(\boldsymbol{v})$ is the damping matrix, and $\boldsymbol{g}(\boldsymbol{\eta})$ is the position and orientation depending vector of restoring forces and moments.

The added mass matrix $\boldsymbol{M}_{\boldsymbol{A}}$ is given by

$$
\boldsymbol{M}_{\boldsymbol{A}}=-\left[\begin{array}{cccccc}
X_{\dot{u}} & X_{\dot{v}} & X_{\dot{w}} & X_{\dot{p}} & X_{\dot{q}} & X_{\dot{r}}  \tag{3.13}\\
Y_{\dot{u}} & Y_{\dot{v}} & Y_{\dot{w}} & Y_{\dot{p}} & Y_{\dot{q}} & Y_{\dot{r}} \\
Z_{\dot{u}} & Z_{\dot{v}} & Z_{\dot{w}} & Z_{\dot{p}} & Z_{\dot{q}} & Z_{\dot{r}} \\
K_{\dot{u}} & K_{\dot{v}} & K_{\dot{w}} & K_{\dot{p}} & K_{\dot{q}} & K_{\dot{r}} \\
M_{\dot{u}} & M_{\dot{v}} & M_{\dot{w}} & M_{\dot{p}} & M_{\dot{q}} & M_{\dot{r}} \\
N_{\dot{u}} & N_{\dot{v}} & N_{\dot{w}} & N_{\dot{p}} & N_{\dot{q}} & N_{\dot{r}}
\end{array}\right]
$$

where the SNAME notation has been used. For example, the hydrodynamic added mass force $Y$ along the $y$-axis due to an acceleration $\dot{u}$ in the $x$-direction is written as

$$
\begin{equation*}
Y=-Y_{\dot{u}} \dot{u}, Y_{\dot{u}}:=\frac{\partial Y}{\partial \dot{u}} . \tag{3.14}
\end{equation*}
$$

The hydrodynamic Coriolis and centripetal matrix is given by

$$
\boldsymbol{C}_{\boldsymbol{A}}(\boldsymbol{v})=\left[\begin{array}{cccccc}
0 & 0 & 0 & 0 & -a_{3} & a_{2}  \tag{3.15}\\
0 & 0 & 0 & a_{3} & 0 & -a_{1} \\
0 & 0 & 0 & -a_{2} & a_{1} & 0 \\
0 & -a_{3} & a_{2} & 0 & -b_{3} & b_{2} \\
a_{3} & 0 & -a_{1} & b_{3} & 0 & -b_{1} \\
-a_{2} & a_{1} & 0 & -b_{2} & b_{1} & 0
\end{array}\right],
$$

where

$$
\begin{aligned}
& a_{1}=X_{\dot{u}} u+X_{\dot{v}} v+X_{\dot{w}} w+X_{\dot{p}} p+X_{\dot{q}} q+X_{\dot{r}} r, \\
& a_{2}=Y_{\dot{u}} u+Y_{\dot{v}} v+Y_{\dot{w}} w+Y_{\dot{p}} p+Y_{\dot{q}} q+Y_{\dot{r}} r, \\
& a_{3}=Z_{\dot{u}} u+Z_{\dot{v}} v+Z_{\dot{w}} w+Z_{\dot{p}} p+Z_{\dot{q}} q+Z_{\dot{r}} r, \\
& b_{1}=K_{\dot{u}} u+K_{\dot{v}} v+K_{\dot{w}} w+K_{\dot{p}} p+K_{\dot{q}} q+K_{\dot{r}} r,
\end{aligned}
$$

$$
\begin{align*}
b_{2} & =M_{\dot{u}} u+M_{\dot{v}} v+M_{\dot{w}} w+M_{\dot{p}} p+M_{\dot{q}} q+M_{\dot{r}} r, \\
b_{3} & =N_{\dot{u}} u+N_{\dot{v}} v+N_{\dot{w}} w+N_{\dot{p}} p+N_{\dot{q}} q+N_{\dot{r}} r . \tag{3.16}
\end{align*}
$$

In general, hydrodynamic damping for ocean vessels is mainly caused by potential damping, skin friction, wave drift damping, and damping due to vortex shedding. It is difficult to give a general expression of the hydrodynamic damping matrix $\boldsymbol{D}(\boldsymbol{v})$. However, it is common to write the hydrodynamic damping matrix $\boldsymbol{D}(\boldsymbol{v})$ as

$$
\begin{equation*}
D(v)=D+D_{n}(v) . \tag{3.17}
\end{equation*}
$$

Here the linear damping matrix $\boldsymbol{D}$ is given by

$$
\boldsymbol{D}=-\left[\begin{array}{cccccc}
X_{u} & X_{v} & X_{w} & X_{p} & X_{q} & X_{r}  \tag{3.18}\\
Y_{u} & Y_{v} & Y_{w} & Y_{p} & Y_{q} & Y_{r} \\
Z_{u} & Z_{v} & Z_{w} & Z_{p} & Z_{q} & Z_{r} \\
K_{u} & K_{v} & K_{w} & K_{p} & K_{q} & K_{r} \\
M_{u} & M_{v} & M_{w} & M_{p} & M_{q} & M_{r} \\
N_{u} & N_{v} & N_{w} & N_{p} & N_{q} & N_{r}
\end{array}\right]
$$

The nonlinear damping matrix $\boldsymbol{D}_{\boldsymbol{n}}(\boldsymbol{v})$ is usually modeled by using a third-order Taylor series expansion or modulus functions (quadratic drag). If the $x z$-plane is a plane of symmetry (starboard/port symmetry) an odd Taylor series expansion containing first-order and third-order terms in velocity can be sufficient to describe most manoeuvres. An approximate expression of each of this matrices will be given in the next section when specific vessels are considered.

### 3.3.2.3 Restoring Forces and Moments

In this section, a model for $\boldsymbol{g}(\boldsymbol{\eta})$ is described. Let $\nabla$ be the volume of fluid displaced by the vessel, $g$ the acceleration of gravity (positive downwards), and $\rho$ the water density. The submerged weight of the body and buoyancy force are defined as

$$
\begin{align*}
& W=m g \\
& B=\rho g \nabla \tag{3.19}
\end{align*}
$$

With the above definition, the restoring force and moment vector $\boldsymbol{g}(\boldsymbol{\eta})$ is due to gravity and buoyancy forces, and is given by

$$
\boldsymbol{g}(\boldsymbol{\eta})=\left[\begin{array}{c}
(W-B) \sin (\theta)  \tag{3.20}\\
-(W-B) \cos (\theta) \sin (\phi) \\
-(W-B) \cos (\theta) \cos (\phi) \\
-\left(y_{g} W-y_{b} B\right) \cos (\theta) \cos (\phi)+\left(z_{g} W-z_{b} B\right) \cos (\theta) \sin (\phi) \\
\left(z_{g} W-z_{b} B\right) \sin (\theta)+\left(x_{g} W-x_{b} B\right) \cos (\theta) \cos (\phi) \\
-\left(x_{g} W-x_{b} B\right) \cos (\theta) \sin (\phi)-\left(y_{g} W-y_{b} B\right) \sin (\theta)
\end{array}\right],
$$

where $\left(x_{b}, y_{b}, z_{b}\right)$ denote coordinates of the center of buoyancy.

### 3.3.2.4 Environmental Disturbances

In this section, we detail the vector, $\boldsymbol{\tau}_{\boldsymbol{E}}$, of forces and moments induced by environmental disturbances including ocean currents, waves (wind generated) and wind, i.e., we can write

$$
\begin{equation*}
\tau_{E}=\tau_{E}^{c u}+\tau_{E}^{w a}+\tau_{E}^{w i} \tag{3.21}
\end{equation*}
$$

where $\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{c} \boldsymbol{u}}, \boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{w} \boldsymbol{a}}$, and $\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{w} \boldsymbol{i}}$ are vectors of forces and moments induced by ocean currents, waves and wind, respectively.

## Current-induced Forces and Moments

The vector $\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{c u}}$ of the current-induced forces and moments is given by

$$
\begin{equation*}
\tau_{E}^{c u}=\left(M_{R B}+M_{A}\right) \dot{v}_{c}+C\left(v_{r}\right) v_{r}-C(v) v+D\left(v_{r}\right) v_{r}-D(v) v \tag{3.22}
\end{equation*}
$$

where $\boldsymbol{v}_{\boldsymbol{r}}=\boldsymbol{v}-\boldsymbol{v}_{\boldsymbol{c}}$ and $\boldsymbol{v}_{\boldsymbol{c}}=\left[u_{c}, v_{c}, w_{c}, 0,0,0\right]^{T}$ is a vector of irrotational bodyfixed current velocities. Let the earth-fixed current velocity vector be denoted by $\left[u_{c}^{E}, v_{c}^{E}, w_{c}^{E}\right]^{T}$. Then, the body-fixed components $\left[u_{c}, v_{c}, w_{c}\right]^{T}$ can be computed as

$$
\left[\begin{array}{c}
u_{c}  \tag{3.23}\\
v_{c} \\
w_{c}
\end{array}\right]=\boldsymbol{J}_{\mathbf{1}}^{\boldsymbol{T}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\left[\begin{array}{c}
u_{c}^{E} \\
v_{c}^{E} \\
w_{c}^{E}
\end{array}\right] .
$$

## Wave-induced Forces and Moments

The vector $\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{w} \boldsymbol{a}}$ of the wave-induced forces and moments is given by

$$
\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{w} \boldsymbol{i}}=\left[\begin{array}{c}
\sum_{i=1}^{N} \rho g B L T \cos (\beta) s_{i}(t)  \tag{3.24}\\
\sum_{i=1}^{N}-\rho g B L T \sin (\beta) s_{i}(t) \\
0 \\
0 \\
0 \\
\sum_{i=1}^{N} \frac{1}{24} \rho g B L\left(L^{2}-B^{2}\right) \sin (2 \beta) s_{i}^{2}(t)
\end{array}\right]
$$

where $\beta$ is the vessel's heading (encounter) angle, see Figure 3.3, $\rho$ is the water density, $L$ is the length of the vessel, $B$ is the breadth of the vessel, and $T$ is the draft of the vessel. Ignoring the higher-order terms of the wave amplitude, the wave slope $s_{i}(t)$ for the wave component $i$ is defined as:

$$
\begin{equation*}
s_{i}(t)=A_{i} \frac{2 \pi}{\lambda_{i}} \sin \left(\omega_{e i} t+\phi_{i}\right) \tag{3.25}
\end{equation*}
$$

where $A_{i}$ is the wave amplitude, $\lambda_{i}$ is the wave length, $\omega_{e i}$ is the encounter frequency, and $\phi_{i}$ is a random phase uniformly distributed and constant with time in $[02 \pi)$ corresponding to the wave component $i$.


Figure 3.3 Definition of a vessel's heading (encounter) angle

## Wind-induced Forces and Moments

For the case where the vessel is at rest (zero speed), the vector $\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{w} \boldsymbol{i}}$ of the windinduced forces and moments is given by

$$
\boldsymbol{\tau}_{\boldsymbol{E}}^{w i}=\frac{1}{2} \rho_{a} V_{w}^{2}\left[\begin{array}{c}
C_{X}\left(\gamma_{w}\right) A_{F w}  \tag{3.26}\\
C_{Y}\left(\gamma_{w}\right) A_{L w} \\
C_{Z}\left(\gamma_{w}\right) A_{F w} \\
C_{K}\left(\gamma_{w}\right) A_{L w} H_{L w} \\
C_{M}\left(\gamma_{w}\right) A_{F w} H_{F w} \\
C_{N}\left(\gamma_{w}\right) A_{L w} L_{o a}
\end{array}\right],
$$

where $V_{w}$ is the wind speed, $\rho_{a}$ is the air density, $A_{F w}$ is the frontal projected area, $A_{L w}$ is the lateral projected area, $H_{F w}$ is the centroid of $A_{F w}$ above the water line, $H_{L w}$ is the centroid of $A_{L w}$ above the water line, $L_{o a}$ is the over all length of the vessel, $\gamma_{w}$ is the angle of relative wind of the vessel bow, see Figure 3.4, and is given by

$$
\begin{equation*}
\gamma_{w}=\psi-\beta_{w}-\pi \tag{3.27}
\end{equation*}
$$

with $\beta_{w}$ being the wind direction. All the wind coefficients (look-up tables) $C_{X}\left(\gamma_{w}\right)$, $C_{Y}\left(\gamma_{w}\right), C_{Z}\left(\gamma_{w}\right), C_{K}\left(\gamma_{w}\right), C_{M}\left(\gamma_{w}\right)$, and $C_{N}\left(\gamma_{w}\right)$ are computed numerically or by experiments in a wind tunnel, see [28].


Figure 3.4 Definition of wind speed and direction

For the case where the vessel is moving, the vector $\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{w} \boldsymbol{i}}$ is given by

$$
\boldsymbol{\tau}_{\boldsymbol{E}}^{\boldsymbol{w} \boldsymbol{i}}=\frac{1}{2} \rho_{a} V_{r w}^{2}\left[\begin{array}{c}
C_{X}\left(\gamma_{r w}\right) A_{F w}  \tag{3.28}\\
C_{Y}\left(\gamma_{r w}\right) A_{L w} \\
C_{Z}\left(\gamma_{r w}\right) A_{F w} \\
C_{K}\left(\gamma_{r w}\right) A_{L w} H_{L w} \\
C_{M}\left(\gamma_{r w}\right) A_{F w} H_{F w} \\
C_{N}\left(\gamma_{r w}\right) A_{L w} L_{o a}
\end{array}\right]
$$

where

$$
\begin{align*}
V_{r w} & =\sqrt{u_{r w}^{2}+v_{r w}^{2}} \\
\gamma_{r w} & =-\arctan 2\left(v_{r w}, u_{r w}\right) \tag{3.29}
\end{align*}
$$

with

$$
\begin{align*}
& u_{r w}=u-V_{w} \cos \left(\beta_{w}-\psi\right), \\
& v_{r w}=v-V_{w} \sin \left(\beta_{w}-\psi\right) \tag{3.30}
\end{align*}
$$

### 3.3.2.5 Propulsion Forces and Moments

The vector, $\boldsymbol{\tau}$, of propulsion forces and moments depends on a specific configuration of actuators such as propellers, rudders, and water jets on a particular vessel. In the next section where $\boldsymbol{\tau}$ is specified, we consider some classes of the ocean vessels that are common in practice. In this book, we neglect the dynamics of the actuators that provide the propulsion forces and moments since the response of the actuators such as hydraulic systems and electrical motors is much faster than the response of the vessel.

### 3.3.2.6 Model Summary and its Properties

## Body-fixed Representation

Now substituting $\boldsymbol{\tau}_{\boldsymbol{R} \boldsymbol{B}}=\boldsymbol{\tau}_{\boldsymbol{H}}+\boldsymbol{\tau}_{\boldsymbol{E}}+\boldsymbol{\tau}$ into (3.9) and combining it with (3.5) results in the equations of motion of an ocean vessel in six degrees of freedom as follows:

$$
\begin{align*}
& \dot{\eta}=J(\eta) v, \\
& M \dot{v}=-C(v) v-D(v) v-g(\eta)+\tau+\tau_{E}, \tag{3.31}
\end{align*}
$$

where

$$
\begin{align*}
& M=M_{R B}+M_{A}, \\
& C(v)=C_{R B}(v)+C_{A}(v) . \tag{3.32}
\end{align*}
$$

Under the assumption that the body is at rest (or at most is moving at low speed) in ideal fluid, the matrix $\boldsymbol{M}$ is always symmetric positive definite, i.e.,

$$
\begin{equation*}
\boldsymbol{M}=M^{\boldsymbol{T}}>\mathbf{0} . \tag{3.33}
\end{equation*}
$$

For a rigid body moving in fluid, the Coriolis and centripetal matrix $\boldsymbol{C}(\boldsymbol{v})$ can always be parameterized such that it is skew-symmetric, i.e.,

$$
\begin{equation*}
C(v)=-C^{\boldsymbol{T}}(v), \forall v \in \mathbb{R}^{\mathbf{6}} . \tag{3.34}
\end{equation*}
$$

For a rigid body moving in an ideal fluid, the hydrodynamic damping matrix $\boldsymbol{D}(\boldsymbol{v})$ is real, non-symmetric and strictly positive, i.e.,

$$
\begin{equation*}
D(v)>\mathbf{0}, \forall v \in \mathbb{R}^{\mathbf{6}} . \tag{3.35}
\end{equation*}
$$

## Earth-fixed Representation

The mathematical model (3.31) can also be written using a representation of the earth-fixed coordinates by applying the following kinematic transformations (with
the assumption that $\boldsymbol{J}^{-\mathbf{1}}(\eta)$ exists, i.e., $\left.\theta \neq \pm \frac{\pi}{2}\right)$ :

$$
\begin{align*}
v & =J^{-\mathbf{1}}(\eta) \dot{\eta} \\
\dot{v} & =J^{-\mathbf{1}}(\eta)\left[\ddot{\eta}-\dot{J}(\eta) J^{-\mathbf{1}}(\eta) \dot{\eta}\right] \tag{3.36}
\end{align*}
$$

Now substituting (3.36) into the second equation of (3.31) results in

$$
\begin{equation*}
M^{\star}(\eta) \ddot{\eta}=-C^{\star}(v, \eta) \dot{\eta}-D^{\star}(v, \eta) \dot{\eta}-g^{\star}(\eta)+J^{T}(\eta)\left(\tau+\tau_{E}\right) \tag{3.37}
\end{equation*}
$$

where

$$
\begin{align*}
& M^{\star}(\eta)=J^{-T}(\eta) M J^{-\mathbf{1}}(\eta) \\
& C^{\star}(v, \eta)=J^{-T}(\eta)\left[C(v)-M J^{-1}(\eta) \dot{J}(\eta)\right] J^{-1}(\eta) \\
& D^{\star}(v, \eta)=J^{-T}(\eta) D(v) J^{-1}(\eta) \\
& g^{\star}(\eta)=J^{-T}(\eta) g(\eta) \tag{3.38}
\end{align*}
$$

Under the same assumptions used in the body-fixed representation, the model (3.37) using the earth-fixed representation has the following properties:

$$
\begin{align*}
& M^{\star}(\eta)=M^{\star}(\eta)^{T}, \forall \eta \in \mathbb{R}^{6} \\
& s^{T}\left[\dot{M}^{\star}(\eta)-2 C^{\star}(v, \eta)\right] s=0, \forall \eta \in \mathbb{R}^{6}, v \in \mathbb{R}^{6}, s \in \mathbb{R}^{6}, \\
& D^{\star}(v, \eta)>0, \quad \forall \eta \in \mathbb{R}^{6}, v \in \mathbb{R}^{6} . \tag{3.39}
\end{align*}
$$



Figure 3.5 Motion variables for an ocean vessel moving in a horizontal plane

### 3.4 Standard Models for Ocean Vessels

In this section, the main results of the previous section are simplified to give a set of standard models for surface ships and underwater vehicles. These standard models will be extensively used for the control design in the coming chapters.

### 3.4.1 Three Degrees of Freedom Horizontal Model

### 3.4.1.1 Standard Three Degrees of Freedom Horizontal Model

The horizontal motion of a surface ship or an underwater vehicle moving in a horizontal plane is often described by the motion components in surge, sway, and yaw. Figure 3.5 illustrates the motion variables in this case. Therefore, we choose $\eta=\left[\begin{array}{ll}x & y\end{array} \psi\right]^{T}$ and $\boldsymbol{v}=\left[\begin{array}{lll}u & v & r\end{array}\right]^{T}$. This model is obtained from the general model (3.31) under the following assumption.

## Assumption 3.1.

1. The motion in roll, pitch, and heave is ignored. This means that we ignore the dynamics associated with the motion in heave, roll, and pitch, i.e., $z=0, w=0$, $\phi=0, p=0, \theta=0$, and $q=0$.
2. The vessel has homogeneous mass distribution and $x z$-plane of symmetry so that

$$
\begin{equation*}
I_{x y}=I_{y z}=0 . \tag{3.40}
\end{equation*}
$$

3. The center of gravity $C G$ and the center of buoyancy, $C B$, are located vertically on the $z$-axis.

With Assumption 3.1, the dynamics of a surface ship or an underwater vehicle moving in a horizontal plane is simplified from the general model (3.31) as follows:

$$
\begin{align*}
& \dot{\eta}=J(\eta) v \\
& M \dot{v}=-C(v) v-\left(D+D_{n}(v)\right) v+\tau+\tau_{E} \tag{3.41}
\end{align*}
$$

where the matrices $\boldsymbol{J}(\boldsymbol{\eta}), \boldsymbol{M}, \boldsymbol{C}(\boldsymbol{v}), \boldsymbol{D}$, and $\boldsymbol{D}_{\boldsymbol{n}}(\boldsymbol{v})$ are given by

$$
\boldsymbol{J}(\boldsymbol{\eta})=\left[\begin{array}{ccc}
\cos (\psi) & -\sin (\psi) & 0 \\
\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right], \boldsymbol{M}=\left[\begin{array}{ccc}
m-X_{\dot{u}} & 0 & 0 \\
0 & m-Y_{\dot{v}} & m x_{g}-Y_{\dot{r}} \\
0 & m x_{g}-Y_{\dot{r}} & I_{z}-N_{\dot{r}}
\end{array}\right]
$$

$$
\begin{align*}
& \boldsymbol{C}(\boldsymbol{v})=\left[\begin{array}{cc}
0 & 0 \\
0 & 0 \\
m\left(x_{g} r+v\right)-Y_{\dot{v}} v-Y_{\dot{r}} r & -m u+X_{\dot{u}} u \\
-m\left(x_{g} r+v\right)+Y_{\dot{v}} v+Y_{\dot{r}} r \\
m u-X_{\dot{u}} u \\
0
\end{array}\right], \\
& \boldsymbol{D}=-\left[\begin{array}{ccc}
X_{u} & 0 & 0 \\
0 & Y_{v} & Y_{r} \\
0 & N_{v} & N_{r}
\end{array}\right], \\
& \boldsymbol{D}_{\boldsymbol{n}}(\boldsymbol{v})=-\left[\begin{array}{ccc}
X_{|u| u|u|} & 0 & 0 \\
0 & Y_{|v| v}|v|+Y_{|r| v}|r| & Y_{|v| r}|v| \\
0 & N_{|v| v \mid}|v|+N_{|r| v}|r| N_{|v| r}|v|+N_{|r| r}|r|
\end{array}\right] .
\end{align*}
$$

The propulsion force and moment vector $\boldsymbol{\tau}$ is given by

$$
\boldsymbol{\tau}=\left[\begin{array}{c}
\tau_{u}  \tag{3.43}\\
0 \\
\tau_{r}
\end{array}\right]
$$

The above propulsion force and moment vector $\boldsymbol{\tau}$ implies that we are considering a surface vessel, which does not have an independent actuator in the sway, i.e., an underactuated vessel is under consideration. Such a vessel can be one equipped with a pair of water jets or a pair of propellers.

The environmental disturbance vector $\boldsymbol{\tau}_{\boldsymbol{E}}$ is given by

$$
\boldsymbol{\tau}_{\boldsymbol{E}}=\left[\begin{array}{c}
\tau_{u E}  \tag{3.44}\\
\tau_{v E} \\
\tau_{r E}
\end{array}\right],
$$

where $\tau_{u E}$ and $\tau_{v E}$ are disturbance forces acting in surge and sway respectively, and $\tau_{r E}$ is the disturbance moment acting in yaw.

### 3.4.1.2 Simplified Three Degrees of Freedom Horizontal Model

In some cases, in addition to Assumption 3.1 we ignore the off-diagonal terms of the matrices $\boldsymbol{M}$ and $\boldsymbol{D}$, all elements of the nonlinear damping matrix $\boldsymbol{D}_{\boldsymbol{n}}(\boldsymbol{v})$. These assumptions hold when the vessel has three planes of symmetry, for which the axes of the body-fixed reference frame are chosen to be parallel to the principal axis of the displaced fluid, which are equal to the principal axis of the vessel. Most ships have port/starboard symmetry, and moreover, bottom/top symmetry is not required for horizontal motion. Ship fore/aft nonsymmetry implies that the off-diagonal terms of the inertia and damping matrices are nonzero. However, these terms are small compared to the main diagonal terms. Furthermore, disturbances induced by waves, wind, and ocean currents are ignored. Under the just-mentioned assumptions, the
dynamics of a surface ship or an underwater vehicle moving in a horizontal plane is simplified from the three degrees of freedom model (3.41) as follows:

$$
\begin{align*}
& \dot{\eta}=J(\eta) v, \\
& M \dot{v}=-C(v) v-D v+\tau, \tag{3.45}
\end{align*}
$$

where the matrices $\boldsymbol{J}(\eta), \boldsymbol{M}, \boldsymbol{C}(\boldsymbol{v})$ and $\boldsymbol{D}$ are given by

$$
\begin{align*}
& \boldsymbol{J}(\boldsymbol{\eta})=\left[\begin{array}{ccc}
\cos (\psi) & -\sin (\psi) & 0 \\
\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right], \boldsymbol{M}=\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & 0 \\
0 & 0 & m_{33}
\end{array}\right],  \tag{3.46}\\
& \boldsymbol{C}(\boldsymbol{v})=\left[\begin{array}{ccc}
0 & 0 & -m_{22} v \\
0 & 0 & m_{11} u \\
m_{22} v & -m_{11} u & 0
\end{array}\right], \boldsymbol{D}=\left[\begin{array}{ccc}
d_{11} & 0 & 0 \\
0 & d_{22} & 0 \\
0 & 0 & d_{33}
\end{array}\right],
\end{align*}
$$

with

$$
\begin{align*}
& m_{11}=m-X_{\dot{u}}, m_{22}=m-Y_{\dot{v}}, m_{33}=I_{z}-N_{\dot{r}}, \\
& d_{11}=-X_{u}, d_{22}=-Y_{v}, d_{33}=-N_{r} . \tag{3.47}
\end{align*}
$$

The propulsion force and moment vector $\boldsymbol{\tau}$ is still given by (3.43), i.e., $\boldsymbol{\tau}=$ $\left[\begin{array}{lll}\tau_{u} & 0 & \tau_{r}\end{array}\right]^{T}$.

### 3.4.1.3 Spherical Three Degrees of Freedom Horizontal Model

In addition to the assumptions made in Subsection 3.4.1.2, we assume that the vessel has bottom/top symmetry. An example of this type of vessels is an ODIN moving in a horizontal plane, see Figure 3.6. In this case, the model is further simplified to

$$
\begin{align*}
& \dot{\eta}=J(\eta) v, \\
& M \dot{v}=-C(v) v-D v+\tau \tag{3.48}
\end{align*}
$$

where the matrices $\boldsymbol{J}(\eta), \boldsymbol{M}, \boldsymbol{C}(\boldsymbol{v})$ and $\boldsymbol{D}$ are given by

$$
\begin{gather*}
\boldsymbol{J}(\boldsymbol{\eta})=\left[\begin{array}{ccc}
\cos (\psi) & -\sin (\psi) & 0 \\
\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right], \boldsymbol{M}=\left[\begin{array}{ccc}
m_{x y} & 0 & 0 \\
0 & m_{x y} & 0 \\
0 & 0 & m_{33}
\end{array}\right], \\
\boldsymbol{C}(\boldsymbol{v})=\left[\begin{array}{ccc}
0 & 0 & -m_{x y} v \\
0 & 0 & m_{x y} u \\
m_{x y} v & -m_{x y} u & 0
\end{array}\right], \boldsymbol{D}=\left[\begin{array}{ccc}
d_{x y} & 0 & 0 \\
0 & d_{x y} & 0 \\
0 & 0 & d_{33}
\end{array}\right], \tag{3.49}
\end{gather*}
$$

with

$$
\begin{align*}
& m_{x y}=m-X_{\dot{u}}=m-Y_{\dot{v}}, m_{33}=I_{z}-N_{\dot{r}} \\
& d_{x y}=-X_{u}=-Y_{v}, d_{33}=-N_{r} \tag{3.50}
\end{align*}
$$

The propulsion force and moment vector $\boldsymbol{\tau}$ is still given by (3.43), i.e. $\boldsymbol{\tau}=$ $\left[\begin{array}{lll}\tau_{u} & 0 & \tau_{r}\end{array}\right]^{T}$.


Figure 3.6 An omnidirectional intelligent navigator (ODIN).
Courtesy http://www.math.hawaii.edu/~ryan/STOMP/Photos 20ODIN.html

### 3.4.2 Six Degrees of Freedom Model

### 3.4.2.1 Standard Model

In addition to the assumptions made in Section 3.3, we assume that the center of gravity and the center of buoyancy are located vertically on the $O_{b} Z_{b}$-axis, and that there are no couplings (off-diagonal terms) in the matrices $\boldsymbol{M}, \boldsymbol{D}$, and $\boldsymbol{D}_{\boldsymbol{n}}(\boldsymbol{v})$. In this case, the model presented in Section 3.3 is simplified to

$$
\begin{align*}
& \dot{\eta}_{1}=J_{1}\left(\eta_{2}\right) v_{1}, \\
& M_{1} \dot{v}_{1}=-C_{1}\left(v_{1}\right) v_{2}-D_{1} v_{1}-D_{n 1}\left(v_{1}\right) v_{1}+\tau_{1}+\tau_{1 E}, \\
& \dot{\eta}_{2}=J_{2}\left(\eta_{2}\right) v_{2}, \\
& M_{2} \dot{v}_{2}=-C_{1}\left(v_{1}\right) v_{1}-C_{2}\left(v_{2}\right) v_{2}-D_{2} v_{2}-D_{n 2}\left(v_{2}\right) v_{2}- \\
& \quad g_{2}\left(\eta_{2}\right)+\tau_{2}+\tau_{2 E}, \tag{3.51}
\end{align*}
$$

where $\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ and $\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ are given in (3.2) and (3.4). The matrices $\boldsymbol{M}_{\mathbf{1}}$ and $\boldsymbol{M}_{\mathbf{2}}$ are

$$
\begin{align*}
& \boldsymbol{M}_{\mathbf{1}}=\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & 0 \\
0 & 0 & m_{33}
\end{array}\right], \\
& \boldsymbol{M}_{\mathbf{2}}=\left[\begin{array}{ccc}
m_{44} & 0 & 0 \\
0 & m_{55} & 0 \\
0 & 0 & m_{66}
\end{array}\right], \tag{3.52}
\end{align*}
$$

where

$$
\begin{aligned}
& m_{11}=m-X_{\dot{u}}, m_{22}=m-Y_{\dot{v}} \\
& m_{33}=m-Z_{\dot{w}}, m_{44}=I_{x}-K_{\dot{p}} \\
& m_{55}=I_{y}-M_{\dot{q}}, m_{66}=I_{z}-N_{\dot{r}}
\end{aligned}
$$

The matrices $\boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right)$ and $\boldsymbol{C}_{\mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right)$ are

$$
\begin{align*}
& \boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{1}\right)=\left[\begin{array}{ccc}
0 & m_{33} w & -m_{22} v \\
-m_{33} w & 0 & m_{11} u \\
m_{22} v & -m_{11} u & 0
\end{array}\right], \\
& \boldsymbol{C}_{\mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right)=\left[\begin{array}{ccc}
0 & m_{66} r & -m_{55} q \\
-m_{66} r & 0 & m_{44} p \\
m_{55} q & -m_{44} p & 0
\end{array}\right] . \tag{3.53}
\end{align*}
$$

The linear damping matrices $\boldsymbol{D}_{\mathbf{1}}$ and $\boldsymbol{D}_{\mathbf{2}}$ are

$$
\begin{align*}
& \boldsymbol{D}_{\mathbf{1}}=\left[\begin{array}{ccc}
d_{11} & 0 & 0 \\
0 & d_{22} & 0 \\
0 & 0 & d_{33}
\end{array}\right], \\
& \boldsymbol{D}_{\mathbf{2}}=\left[\begin{array}{ccc}
d_{44} & 0 & 0 \\
0 & d_{55} & 0 \\
0 & 0 & d_{66}
\end{array}\right], \tag{3.54}
\end{align*}
$$

where

$$
\begin{aligned}
& d_{11}=-X_{u}, \\
& d_{22}=-Y_{v}, \\
& d_{33}=-Z_{w}, \\
& d_{44}=-K_{p}, \\
& d_{55}=-M_{q}, \\
& d_{66}=-N_{r} .
\end{aligned}
$$

The nonlinear damping matrices $\boldsymbol{D}_{\boldsymbol{n} 1}\left(\boldsymbol{v}_{\mathbf{1}}\right)$ and $\boldsymbol{D}_{\boldsymbol{n} 2}\left(\boldsymbol{v}_{\mathbf{2}}\right)$ are

$$
\begin{gather*}
\boldsymbol{D}_{\boldsymbol{n} 1}\left(\boldsymbol{v}_{\mathbf{1}}\right)=\left[\begin{array}{ccc}
\sum_{i=2}^{3} d_{u i}|u|^{i-1} & 0 & 0 \\
0 & \sum_{i=2}^{3} d_{v i}|v|^{i-1} & 0 \\
0 & 0 & \sum_{i=2}^{3} d_{w i}|w|^{i-1}
\end{array}\right], \\
\boldsymbol{D}_{\boldsymbol{n} 2}\left(\boldsymbol{v}_{\mathbf{2}}\right)=\left[\begin{array}{ccc}
\sum_{i=2}^{3} d_{p i}|p|^{i-1} & 0 & 0 \\
0 & \sum_{i=2}^{3} d_{q i}|q|^{i-1} & 0 \\
0 & 0 & \sum_{i=2}^{3} d_{r i}|r|^{i-1}
\end{array}\right], \tag{3.55}
\end{gather*}
$$

where $d_{u i}, d_{v i}, d_{w i}, d_{p i}, d_{q i}$, and $d_{r i}$ with $i=2,3$ are the nonlinear hydrodynamic damping coefficients.

The restoring force and moment vector $\boldsymbol{g}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ is given by

$$
\boldsymbol{g}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)=\left[\begin{array}{c}
\rho g \nabla \overline{G M_{T}} \sin (\phi) \cos (\theta)  \tag{3.56}\\
\rho g \nabla \overline{G M_{L}} \sin (\theta) \\
0
\end{array}\right]
$$

where $\rho, g, \nabla, \overline{G M_{T}}$ and $\overline{G M_{L}}$ are the water density, gravity acceleration, displaced volume of water, transverse metacentric height and longitudinal metacentric height, respectively.

The propulsion force and moment vectors $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{\mathbf{2}}$ are

$$
\boldsymbol{\tau}_{\mathbf{1}}=\left[\begin{array}{c}
\tau_{u}  \tag{3.57}\\
0 \\
0
\end{array}\right], \boldsymbol{\tau}_{\mathbf{2}}=\left[\begin{array}{c}
\tau_{p} \\
\tau_{q} \\
\tau_{r}
\end{array}\right]
$$

which imply that the vessel under consideration does not have independent actuators in the sway and heave.

The environmental disturbance vectors $\boldsymbol{\tau}_{\mathbf{1} \boldsymbol{E}}$ and $\boldsymbol{\tau}_{\mathbf{2} \boldsymbol{E}}$ are given by

$$
\boldsymbol{\tau}_{\mathbf{1 E}}=\left[\begin{array}{c}
\tau_{E u}  \tag{3.58}\\
\tau_{E v} \\
\tau_{E w}
\end{array}\right], \boldsymbol{\tau}_{\mathbf{2 E}}=\left[\begin{array}{c}
\tau_{E p} \\
\tau_{E q} \\
\tau_{E r}
\end{array}\right],
$$

where $\tau_{E u}, \tau_{E v}, \tau_{E w}, \tau_{E p}, \tau_{E q}$, and $\tau_{E r}$ are the environmental disturbance forces or moments acting on the surge, sway, heave, roll, pitch, and yaw axes, respectively.

### 3.4.2.2 Ignoring Nonlinear Damping Terms and Roll Model

In addition to the assumptions made in Section 3.4.2.1, it is sometimes reasonable to ignore nonlinear hydrodynamic damping terms and roll, and environmental disturbances. This holds when the vessel is operating at low speed and is equipped with
independent internal/external roll actuators. As such, the model (3.51) is simplified to:

1. Kinematics

$$
\begin{align*}
\dot{x} & =\cos (\psi) \cos (\theta) u-\sin (\psi) v+\sin (\theta) \cos (\psi) w, \\
\dot{y} & =\sin (\psi) \cos (\theta) u+\cos (\psi) v+\sin (\theta) \sin (\psi) w, \\
\dot{z} & =-\sin (\theta) u+\cos (\theta) w, \\
\dot{\theta} & =q, \\
\dot{\psi} & =\frac{r}{\cos (\theta)} . \tag{3.59}
\end{align*}
$$

2. Kinetics

$$
\begin{align*}
& \dot{u}=\frac{m_{22}}{m_{11}} v r-\frac{m_{33}}{m_{11}} w q-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}, \\
& \dot{v}=-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v, \\
& \dot{w}=\frac{m_{11}}{m_{33}} u q-\frac{d_{33}}{m_{33}} w, \\
& \dot{q}=\frac{m_{33}-m_{11}}{m_{55}} u w-\frac{d_{55}}{m_{55}} q-\frac{\rho g \nabla \overline{G M_{L}} \sin (\theta)}{m_{55}}+\frac{1}{m_{55}} \tau_{q}, \\
& \dot{r}=\frac{m_{11}-m_{22}}{m_{66}} u v-\frac{d_{66}}{m_{66}} r+\frac{1}{m_{66}} \tau_{r} . \tag{3.60}
\end{align*}
$$

### 3.5 Conclusions

This chapter sets out material about the basic motion tasks and mathematical models of the ocean vessels that will be used in the subsequent chapters. More details on deriving the mathematical models of the ocean vessels are given in [11, 12,25]. It has also been pointed out that regulation/stabilization is much more difficult than trajectory-tracking, path-tracking, and path-following for underactuated ocean vessels.

## Chapter 4

# Control Properties and Previous Work on Control of Ocean Vessels 

In this chapter, first control properties of ocean vessels are presented. Then, the existing literature on the control of underactuated ocean vessels is reviewed. Through the review of the previous work in the areas of stabilization, trajectory-tracking, path-following, and output feedback control of underactuated ocean vessels, challenging questions are raised. Illustration of the background and process of solutions of those questions, as well as an explanation of the solutions in terms of their physical insights and practical applications are then presented in subsequent chapters.

### 4.1 Controllability Properties

### 4.1.1 Acceleration Constraints

The number $\left(m_{c}\right)$ of independent control inputs (the number of nonzero elements of the propulsion force and moment vector $\boldsymbol{\tau})$ is smaller than the number $\left(n_{c}\right)$ of degrees of freedom to be controlled for a standard model of the ocean vessels. As such, we remove all zero elements of $\boldsymbol{\tau}$ and denote the resulting vector by $\boldsymbol{\tau}_{\boldsymbol{a}}$. Thus, if $\boldsymbol{\tau}_{\boldsymbol{a}} \in \mathbb{R}^{m_{c}}$ and $\boldsymbol{\eta} \in \mathbb{R}^{n_{c}}$, then $m_{c}<n_{c}$. For example, for the case of the vessels with six degrees of freedom to be controlled we have $m_{c}<6$, for the case of the vessels with five degrees of freedom to be controlled $m_{c}<5$, and $m_{c}<3$ for the case of the vessels with three degrees of freedom to be controlled. For clarity, we ignore the environmental disturbance forces and moments to investigate acceleration constraints on the aforementioned ocean vessels. Let $\boldsymbol{M}_{\boldsymbol{u}}, \boldsymbol{C}_{\boldsymbol{u}}(\boldsymbol{v}), \boldsymbol{D}_{\boldsymbol{u}}(\boldsymbol{v})$, and $\boldsymbol{g}_{\boldsymbol{u}}(\boldsymbol{\eta})$ denote the rows of $\boldsymbol{M}, \boldsymbol{C}(\boldsymbol{v}), \boldsymbol{D}(\boldsymbol{v})$, and $\boldsymbol{g}(\boldsymbol{\eta})$ that correspond to those rows without propulsion forces or moments, i.e.,

$$
\begin{equation*}
M_{u} \dot{v}+C_{u}(v) v+D_{u}(v) v+g_{u}(\eta)=0 . \tag{4.1}
\end{equation*}
$$

The above equation describes the acceleration constraints, i.e., second-order constraints. The following results give the conditions whether the constraints given in (4.1) are partially integrable or totally integrable.

Lemma 4.1. The constraints (4.1) are partially integrable if and only if the following conditions hold:

1. $g_{u}(\eta)$ is a constant vector.
2. $\left(\boldsymbol{C}_{\boldsymbol{u}}(\boldsymbol{v})+\boldsymbol{D}_{\boldsymbol{u}}(\boldsymbol{v})\right)$ is a constant matrix.
3. The distribution $\Omega^{\perp}(\eta)=\operatorname{ker}\left(\left(\boldsymbol{C}_{\boldsymbol{u}}(\boldsymbol{v})+\boldsymbol{D}_{\boldsymbol{u}}(\boldsymbol{v})\right) \boldsymbol{J}^{\mathbf{- 1}}(\eta)\right)$ is completely integrable.

Proof. See [29].
Lemma 4.2. The constraints (4.1) are totally integrable if and only if the following conditions hold:

1. The constraints are partially integrable.
2. $\left(\boldsymbol{C}_{\boldsymbol{u}}(\boldsymbol{v})+\boldsymbol{D}_{\boldsymbol{u}}(\boldsymbol{v})\right)=0$.
3. The distribution $\Delta(\eta)=\operatorname{ker}\left(\boldsymbol{M}_{\boldsymbol{u}} \boldsymbol{J}^{\mathbf{1}}(\boldsymbol{\eta})\right)$ is completely integrable.

Proof. See [29].
The following lemma gives a result on the stabilizability of an underactuated ocean vessel.

Lemma 4.3. Consider the system (3.31) with $\boldsymbol{\tau}_{\boldsymbol{E}}=\mathbf{0}^{n_{c} \times 1}$. Assume that the elements of the restoring force and moment vector $\boldsymbol{g}(\boldsymbol{\eta})$ corresponding to the unactuated dynamics are zero, i.e., the vector $\boldsymbol{g}(\eta)$ can be written in the form of

$$
\boldsymbol{g}(\eta)=\left[\begin{array}{c}
g_{\boldsymbol{a}}(\eta)  \tag{4.2}\\
\mathbf{0}^{\left(n_{c}-m_{c}\right) \times 1}
\end{array}\right],
$$

where $\boldsymbol{g}_{\boldsymbol{a}}(\eta) \in \mathbb{R}^{m_{c}}$ is the restoring force and moment vector corresponding to the actuated dynamics. Let $(\boldsymbol{\eta}, \boldsymbol{v})=\left(\eta^{\boldsymbol{e}}, \mathbf{0}^{n_{c}-m_{c}}\right)$ be an equilibrium. There is no $C^{1}$ state feedback law $\boldsymbol{\alpha}(\boldsymbol{\eta}, \boldsymbol{v}): \mathbb{R}^{n_{c}} \times \mathbb{R}^{n_{c}} \rightarrow \mathbb{R}^{m_{c}}$ that makes the equilibrium $\left(\eta^{e}, \boldsymbol{0}^{n_{c}-m_{c}}\right)$ asymptotically stable.

Proof. See [29].

### 4.1.2 Kinematic Constraints

In this section, we address controllability properties of ocean vessels. Since the vessel under consideration has a number of degrees of freedom to be controlled greater than control inputs (e.g., underwater vehicles do not have independent actuators in the heave and sway axes, see Section 3.4.2 and surface ships do not have an independent actuator in the sway axis, see Section 3.4.1), we can address the controllability
issue of the vessel kinematic by considering the kinematics with the linear velocity vector

$$
\boldsymbol{v}_{\mathbf{1}}=\left[\begin{array}{c}
u  \tag{4.3}\\
0 \\
0
\end{array}\right]
$$

We analyze controllability properties of six degrees of freedom vessels. The case of three degrees of freedom vessels can be obtained directly from the results for the six degrees of freedom vessels. With (4.3), we now write the kinematics of the vessel as follows:

$$
\begin{gather*}
\dot{\eta}=\gamma_{\mathbf{1}}(\eta) u+\underset{\mathbf{2}}{ }(\boldsymbol{\eta}) p+\gamma_{\mathbf{3}}(\eta) q+\gamma_{\mathbf{4}}(\eta) r \\
\dot{\Downarrow}  \tag{4.4}\\
\dot{\boldsymbol{\eta}}=\boldsymbol{\Upsilon}(\boldsymbol{\eta}) \boldsymbol{u},
\end{gather*}
$$

where

$$
\begin{gather*}
\boldsymbol{\gamma}_{\mathbf{1}}(\boldsymbol{\eta})=\left[\begin{array}{c}
\cos (\theta) \cos (\psi) \\
\cos (\theta) \sin (\psi) \\
-\sin (\theta) \\
0 \\
0 \\
0
\end{array}\right], \quad \boldsymbol{\gamma}_{\mathbf{2}}(\boldsymbol{\eta})=\left[\begin{array}{l}
0 \\
0 \\
0 \\
1 \\
0 \\
0
\end{array}\right], \\
\boldsymbol{\gamma}_{\mathbf{3}}(\boldsymbol{\eta})=\left[\begin{array}{c}
0 \\
0 \\
0 \\
\sin (\phi) \tan (\theta) \\
\cos (\phi) \\
\sin (\phi) \sec (\theta)
\end{array}\right], \quad \boldsymbol{\gamma} \mathbf{4}(\boldsymbol{\eta})=\left[\begin{array}{c}
0 \\
0 \\
0 \\
\cos (\phi) \tan (\theta) \\
-\sin (\phi) \\
\cos (\phi) \sec (\theta)
\end{array}\right], \tag{4.5}
\end{gather*}
$$

and

$$
\begin{align*}
& \boldsymbol{\Upsilon}(\boldsymbol{\eta})=\left[\boldsymbol{\gamma}_{\mathbf{1}}(\boldsymbol{\eta}) \boldsymbol{\gamma}_{\mathbf{2}}(\boldsymbol{\eta}) \boldsymbol{\gamma}_{\mathbf{3}}(\boldsymbol{\eta}) \boldsymbol{\gamma}_{\mathbf{4}}(\boldsymbol{\eta})\right] \\
& \boldsymbol{u}=\left[\begin{array}{lll}
u & p & q
\end{array}\right]^{T} \tag{4.6}
\end{align*}
$$

From (4.4) and (4.5), a calculation shows that the following nonhonolomic (nonintegrable) constraints are satisfied:

$$
\begin{align*}
& (\cos (\psi) \sin (\theta) \sin (\phi)-\sin (\psi) \cos (\phi)) \dot{x}+ \\
& \quad(\sin (\psi) \sin (\theta) \sin (\phi)+\cos (\psi) \cos (\phi)) \dot{y}+\cos (\theta) \sin (\phi) \dot{z}=0 \\
& (\sin (\psi) \sin (\theta) \cos (\phi)-\sin (\psi) \sin (\phi)) \dot{x}+  \tag{4.7}\\
& \quad(\sin (\psi) \sin (\theta) \cos (\phi)-\cos (\psi) \sin (\phi)) \dot{y}+\cos (\theta) \cos (\phi) \dot{z}=0
\end{align*}
$$

Based on (4.4), we will address the following controllability issues: Controllability about a point (i.e., stabilization) and controllability about a trajectory (i.e., trajectory-tracking).

### 4.1.3 Controllability at a Point

We will first consider a linear approximation of the system (4.4) at an equilibrium point $\boldsymbol{\eta}_{\boldsymbol{e}}$. Let the error associated with the equilibrium point $\boldsymbol{\eta}_{\boldsymbol{e}}$ be as follows:

$$
\begin{equation*}
\tilde{\eta}=\eta-\eta_{e} . \tag{4.8}
\end{equation*}
$$

With (4.8), we can write the tangent linearization of (4.4) at the equilibrium point $\eta_{e}$ as

$$
\begin{equation*}
\dot{\tilde{\eta}}=\Upsilon\left(\eta_{e}\right) u \tag{4.9}
\end{equation*}
$$

which is not controllable because the rank of the matrix $\boldsymbol{\Upsilon}\left(\boldsymbol{\eta}_{\boldsymbol{e}}\right)$ is 4 . This implies that a linear controller will never achieve posture stabilization, not even in a local sense. In order to study the controllability of the vessel in question, we need to use some tools (the Lie algebra rank condition and nilpotent concepts) from nonlinear control theory [4].

Given a set of generators or basis vector fields $\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{2}}, \ldots, \boldsymbol{\gamma}_{m_{c}}$, the length of a Lie product recursively defined as

$$
\begin{align*}
& \ell\left\{\gamma_{i}\right\}=1, \quad i=1,2, \ldots, m_{c} \\
& \ell([A, B])=\ell[A]+\ell[B], \tag{4.10}
\end{align*}
$$

where $A$ and $B$ are themselves Lie products. Alternatively, $\ell[A]$ is the number of generators in the expansion for $A$. A Lie algebra or basis is nilpotent if there exists an integer $k$ such that all Lie products of length greater than $k$ are zero. The integer $k$ is called the order of nilpotency. The use of the nilpotent basis eliminates the need for cumbersome computations as we see that all higher order Lie brackets above some particular order are zero.

The above concepts and conditions imply that Lie algebra $L\left\{\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{2}}, \gamma_{\mathbf{3}}, \gamma_{\mathbf{4}}\right\}$ is nilpotent algebra of order $k=2$, i.e., the vector fields $\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{2}}, \boldsymbol{\gamma}_{\mathbf{3}}$, and $\boldsymbol{\gamma}_{\mathbf{4}}$ are the nilpotent basis. Thus all Lie brackets of order greater than two are zero. The only independent Lie brackets computed from the four basis vector fields are $\left[\gamma_{\mathbf{1}}, \gamma_{\mathbf{3}}\right]$ and $\left[\gamma_{\mathbf{1}}, \gamma_{\mathbf{4}}\right]$. Therefore, for our system the Lie algebra rank condition becomes

$$
\begin{equation*}
\operatorname{rank}\left[\boldsymbol{C}_{\boldsymbol{c}}\right]=6 \Leftrightarrow \operatorname{rank}\left[\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{2}}, \boldsymbol{\gamma}_{\mathbf{3}}, \boldsymbol{\gamma}_{\mathbf{4}},\left[\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{3}}\right],\left[\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{4}}\right]\right]=6, \tag{4.11}
\end{equation*}
$$

where $\left[\gamma_{\mathbf{1}}, \gamma_{\mathbf{3}}\right]$ and $\left[\gamma_{\mathbf{2}}, \gamma_{\mathbf{4}}\right]$ are the two independent Lie brackets computed from the four vector fields $\left(\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{2}}, \boldsymbol{\gamma}_{\mathbf{3}}, \boldsymbol{\gamma}_{\mathbf{4}}\right)$ and $C_{c}$ is called the controllability matrix. For two vector fields $g(x)$ and $h(x)$, a Lie bracket is computed based on the following formula:

$$
\begin{equation*}
[g, h](x)=\frac{\partial h}{\partial x} g-\frac{\partial g}{\partial x} h . \tag{4.12}
\end{equation*}
$$

Using the definition (4.12), Lie brackets $\left[\gamma_{\mathbf{1}}, \gamma_{\mathbf{3}}\right]$ and $\left[\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{4}}\right]$ are given by

$$
\begin{gather*}
{\left[\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{3}}\right]=\frac{\partial \boldsymbol{\gamma}_{\mathbf{3}}}{\partial \boldsymbol{\eta}} \boldsymbol{\gamma}_{\mathbf{1}}-\frac{\partial \boldsymbol{\gamma}_{\mathbf{1}}}{\partial \boldsymbol{\eta}} \boldsymbol{\gamma}_{\mathbf{3}}=\left[\begin{array}{c}
\cos (\psi) \sin (\theta) \cos (\phi)+\sin (\psi) \sin (\phi) \\
\sin (\psi) \sin (\theta) \cos (\phi)-\cos (\psi) \sin (\phi) \\
\cos (\theta) \cos (\phi) \\
0 \\
0 \\
0
\end{array}\right]} \\
{\left[\boldsymbol{\gamma}_{\mathbf{1}}, \boldsymbol{\gamma}_{\mathbf{4}}\right]=\frac{\partial \boldsymbol{\gamma}_{\mathbf{4}}}{\partial \boldsymbol{\eta}} \boldsymbol{\gamma}_{\mathbf{1}}-\frac{\partial \boldsymbol{\gamma}_{\mathbf{1}}}{\partial \boldsymbol{\eta}} \boldsymbol{\gamma}_{\mathbf{4}}=\left[\begin{array}{c}
-\cos (\psi) \sin (\theta) \sin (\phi)+\sin (\psi) \cos (\phi) \\
-\sin (\psi) \sin (\theta) \sin (\phi)-\cos (\psi) \cos (\phi) \\
-\cos (\theta) \sin (\phi) \\
0 \\
0 \\
0
\end{array}\right]} \tag{4.13}
\end{gather*}
$$

Therefore, the controllability matrix $\boldsymbol{C}_{\boldsymbol{c}}$ is given by

$$
\begin{align*}
& \boldsymbol{C}_{\boldsymbol{c}}=\left[\begin{array}{cccc}
\cos (\psi) \cos (\theta) & 0 & 0 & 0 \\
\sin (\psi) \cos (\theta) & 0 & 0 & 0 \\
-\sin (\theta) & 0 & 0 & 0 \\
0 & 1 \sin (\phi) \tan (\theta) & \cos (\phi) \tan (\theta) \\
0 & 0 & \cos (\phi) & -\sin (\phi) \\
0 & 0 \sin (\phi) \sec (\theta) \cos (\phi) \sec (\theta)
\end{array}\right. \\
& \cos (\psi) \sin (\theta) \cos (\phi)+\sin (\psi) \sin (\phi) \\
& \sin (\psi) \sin (\theta) \cos (\phi)-\cos (\psi) \sin (\phi) \\
& \cos (\theta) \cos (\phi)  \tag{4.14}\\
& 0 \\
& 0 \\
& 0 \\
& -\cos (\psi) \sin (\theta) \sin (\phi)+\sin (\psi) \cos (\phi) \\
& -\sin (\psi) \sin (\theta) \sin (\phi)-\cos (\psi) \cos (\phi) \\
& -\cos (\theta) \sin (\phi) \\
& 0 \\
& 0 \\
& 0
\end{align*}
$$

It can be seen that the above matrix $\boldsymbol{C}_{\boldsymbol{c}}$ has one nonzero minor of order 6. Therefore, this matrix is full rank provided that $\theta \neq \frac{\pi}{2}$. This implies that the vessel is locally controllable and also globally controllable as long as the singular condition $\theta \neq \frac{\pi}{2}$ is avoided. As for the stabilizability of system (4.4) to a point, the failure of the previous linear analysis indicates that exponential stability cannot be achieved by
smooth feedback [21]. Things turn out to be even worse: If smooth (in fact, even continuous) time-invariant feedback laws are used, Lyapunov stability cannot be used directly. This negative result is established on the basis of a necessary condition due to Brockett [21], see Section 2.7.4: Smooth stabilizability of a driftless regular system (i.e., such that the input vector fields are well defined and linearly independent at $\eta_{\boldsymbol{e}}$ ) requires that the number of inputs be equal to the number of states. The above difficulty has a deep impact on the control design. In fact, to obtain a posture stabilizing controller it is either necessary to give up the continuity requirement and/or to resort to time-varying control laws.

### 4.1.4 Controllability About a Trajectory

For the system (4.4), let the reference trajectory $\boldsymbol{\eta}_{\boldsymbol{d}}$ and the reference trajectory input $\boldsymbol{u}_{\boldsymbol{d}}$ be

$$
\boldsymbol{\eta}_{\boldsymbol{d}}=\left[\begin{array}{c}
x_{d}(t)  \tag{4.15}\\
y_{d}(t) \\
z_{d}(t) \\
\phi_{d}(t) \\
\theta_{d}(t) \\
\psi_{d}(t)
\end{array}\right], \boldsymbol{u}_{\boldsymbol{d}}=\left[\begin{array}{c}
u_{d}(t) \\
p_{d}(t) \\
q_{d}(t) \\
r_{d}(t)
\end{array}\right] .
$$

Indeed, the reference trajectory $\boldsymbol{\eta}_{\boldsymbol{d}}$ and the reference trajectory input $\boldsymbol{u}_{\boldsymbol{d}}$ should satisfy the nonhonolomic constraints (4.7), i.e.,

$$
\begin{align*}
& \left(\cos \left(\psi_{d}\right) \sin \left(\theta_{d}\right) \sin \left(\phi_{d}\right)-\sin \left(\psi_{d}\right) \cos \left(\phi_{d}\right)\right) \dot{x}_{d}+ \\
& \quad\left(\sin \left(\psi_{d}\right) \sin \left(\theta_{d}\right) \sin \left(\phi_{d}\right)+\cos \left(\psi_{d}\right) \cos \left(\phi_{d}\right)\right) \dot{y}_{d}+\cos \left(\theta_{d}\right) \sin \left(\phi_{d}\right) \dot{z}_{d}=0, \\
& \left(\sin \left(\psi_{d}\right) \sin \left(\theta_{d}\right) \cos \left(\phi_{d}\right)-\sin \left(\psi_{d}\right) \sin \left(\phi_{d}\right)\right) \dot{x}_{d}+ \\
& \quad\left(\sin \left(\psi_{d}\right) \sin \left(\theta_{d}\right) \cos \left(\phi_{d}\right)-\cos \left(\psi_{d}\right) \sin \left(\phi_{d}\right)\right) \dot{y}_{d}+\cos \left(\theta_{d}\right) \cos \left(\phi_{d}\right) \dot{z}_{d}=0 . \tag{4.16}
\end{align*}
$$

Let the errors associated with the reference trajectory and the reference input trajectory be

$$
\begin{align*}
& \eta_{\boldsymbol{e}}=\eta-\eta_{\boldsymbol{d}}, \\
& \boldsymbol{u}_{\boldsymbol{e}}=\boldsymbol{u}-\boldsymbol{u}_{\boldsymbol{d}} . \tag{4.17}
\end{align*}
$$

Using (4.17), we can write (4.4) as

$$
\begin{equation*}
\dot{\eta}=\Upsilon\left(\eta_{d}+\eta_{e}\right)\left(u_{d}+u_{e}\right) . \tag{4.18}
\end{equation*}
$$

The Taylor series expansion of $\boldsymbol{\Upsilon}\left(\boldsymbol{\eta}_{\boldsymbol{d}}+\boldsymbol{\eta}_{\boldsymbol{e}}\right)$ about the nominal solution $\boldsymbol{\eta}_{\boldsymbol{d}}$ is given by

$$
\begin{equation*}
\dot{\eta}=\left(\boldsymbol{\Upsilon}\left(\boldsymbol{\eta}_{\boldsymbol{d}}, t\right)+\left.\frac{\partial \boldsymbol{\Upsilon}(\boldsymbol{\eta})}{\partial \boldsymbol{\eta}}\right|_{\boldsymbol{\eta}=\boldsymbol{\eta}_{\boldsymbol{d}}} \eta_{\boldsymbol{e}}(t)+\operatorname{HOT}\right)\left(\boldsymbol{u}_{\boldsymbol{d}}(t)+\boldsymbol{u}_{\boldsymbol{e}}(t)\right) . \tag{4.19}
\end{equation*}
$$

Since the reference trajectory $\boldsymbol{\eta}$ and the reference input trajectory $\boldsymbol{u}_{\boldsymbol{d}}$ satisfy the nonholonomic constraints (4.16), we have

$$
\begin{equation*}
\dot{\boldsymbol{\eta}}_{\boldsymbol{d}}=\boldsymbol{\Upsilon}\left(\boldsymbol{\eta}_{\boldsymbol{d}}, t\right) \boldsymbol{u}_{\boldsymbol{d}}(t) \tag{4.20}
\end{equation*}
$$

Subtracting (4.19) by (4.20) and ignoring the high-order terms (HOT) gives

$$
\begin{align*}
\dot{\eta}_{\boldsymbol{e}} & =\left(\left.\frac{\partial \boldsymbol{\Upsilon}(\eta)}{\partial \boldsymbol{\eta}}\right|_{\boldsymbol{\eta}=\boldsymbol{\eta}_{\boldsymbol{d}}} \eta_{\boldsymbol{e}}(t)\right) \boldsymbol{u}_{\boldsymbol{d}}(t)+\boldsymbol{\Upsilon}\left(\boldsymbol{\eta}_{\boldsymbol{d}}, t\right) \boldsymbol{u}_{\boldsymbol{e}}(t) \\
& :=\boldsymbol{A}(t) \boldsymbol{\eta}_{\boldsymbol{e}}(t)+\boldsymbol{B}(t) \boldsymbol{u}_{\boldsymbol{e}}(t) \tag{4.21}
\end{align*}
$$

where

$$
\boldsymbol{A}(t)=\left[\begin{array}{ll}
0_{3 \times 3} & \boldsymbol{A}_{\mathbf{1}}(t)  \tag{4.22}\\
0_{3 \times 3} & \boldsymbol{A}_{\mathbf{2}}(t)
\end{array}\right], \boldsymbol{B}(t)=\left[\begin{array}{cc}
\boldsymbol{J}_{\boldsymbol{d} \mathbf{1}}(t) & 0_{3 \times 3} \\
0_{3 \times 1} & \boldsymbol{J}_{\boldsymbol{d} \mathbf{2}}(t)
\end{array}\right],
$$

with $\boldsymbol{A}_{\mathbf{1}}$ and $\boldsymbol{A}_{\mathbf{2}}$ given by

$$
\begin{align*}
& \boldsymbol{A}_{\mathbf{1}}(t)=\left[\begin{array}{cc}
0-\cos \left(\psi_{d}\right) \sin \left(\theta_{d}\right) u_{d} & -\sin \left(\psi_{d}\right) \cos \left(\theta_{d}\right) u_{d} \\
0-\sin \left(\psi_{d}\right) \sin \left(\theta_{d}\right) u_{d} & \cos \left(\psi_{d}\right) \cos \left(\theta_{d}\right) u_{d} \\
0 & -\cos \left(\theta_{d}\right) u_{d}
\end{array}\right], \\
& \boldsymbol{A}_{\mathbf{2}}(t)=\left[\begin{array}{c}
\cos \left(\phi_{d}\right) \tan \left(\phi_{d}\right) q_{d}-\sin \left(\phi_{d}\right) \tan \left(\phi_{d}\right) r_{d} \\
-\sin \left(\phi_{d}\right) q_{d}-\cos \left(\phi_{d}\right) r_{d} \\
\cos \left(\phi_{d}\right) \sec \left(\theta_{d}\right) q_{d}-\sin \left(\phi_{d}\right) \sec \left(\theta_{d}\right) r_{d}
\end{array}\right. \\
& \sin \left(\phi_{d}\right) \sec ^{2}\left(\theta_{d}\right) q_{d}+\cos \left(\phi_{d}\right) \sec ^{2}\left(\theta_{d}\right) r_{d}  \tag{4.23}\\
& 0 \\
& \left.\sin \left(\phi_{d}\right) \sec \left(\theta_{d}\right) \tan \left(\theta_{d}\right) q_{d}+\cos \left(\phi_{d}\right) \sec \left(\theta_{d}\right) \tan \left(\theta_{d}\right) r_{d} 0\right],
\end{align*}
$$

and $\boldsymbol{J}_{\boldsymbol{d} \mathbf{1}}(t)$ and $\boldsymbol{J}_{\boldsymbol{d} \mathbf{2}}(t)$ given by

$$
\begin{align*}
& \boldsymbol{J}_{\boldsymbol{d} \mathbf{1}}(t)=\left[\begin{array}{c}
\cos \left(\theta_{d}\right) \cos \left(\psi_{d}\right) \\
\cos \left(\theta_{d}\right) \sin \left(\psi_{d}\right) \\
-\sin \left(\theta_{d}\right)
\end{array}\right], \\
& \boldsymbol{J}_{\boldsymbol{d} \mathbf{2}}(t)=\left[\begin{array}{ccc}
1 & \sin \left(\phi_{d}\right) \tan \left(\theta_{d}\right) & \cos \left(\phi_{d}\right) \tan \left(\theta_{d}\right) \\
0 & \cos \left(\phi_{d}\right) & -\sin \left(\phi_{d}\right) \\
0 & \sin \left(\phi_{d}\right) \sec \left(\theta_{d}\right) & \cos \left(\phi_{d}\right) \sec \left(\theta_{d}\right)
\end{array}\right] . \tag{4.24}
\end{align*}
$$

In (4.23) and (4.24), the argument $t$ of $\phi_{d}, \theta_{d}, \psi_{d}, u_{d}, p_{d}$, and $r_{d}$ is omitted for simplicity.

The system (4.21) is linear time-varying. The controllability condition becomes

$$
\begin{equation*}
\operatorname{rank}\left\{\boldsymbol{B}, \boldsymbol{A} \boldsymbol{B}, \boldsymbol{A}^{\mathbf{2}} \boldsymbol{B}, \boldsymbol{A}^{\mathbf{3}} \boldsymbol{B}, \boldsymbol{A}^{\mathbf{4}} \boldsymbol{B}, \boldsymbol{A}^{\mathbf{5}} \boldsymbol{B}\right\}=6 . \tag{4.25}
\end{equation*}
$$

A calculation shows that the above matrix has a nonzero minor of order 6 provided that $\left(u_{d} \neq 0, p_{d} \neq 0, q_{d} \neq 0, r_{d} \neq 0\right)$ and $\left(\theta_{d} \neq \frac{\pi}{2}\right)$. Therefore, we conclude that the kinematic system (4.21) can be locally stabilized by linear feedback about trajectories consisting of linear or circular or helix paths, which do not collapse to a point.

### 4.2 Previous Work on Control of Underactuated Ocean Vessels

This section starts with a brief review on the control of nonholonomic systems, due to their relevance to the control of underactuated ocean vessels. Next, the existing methods on control of underactuated ocean vessels are reviewed. Limitations of the existing methods are then pointed out and hence motivate the contributions of the book.

### 4.2.1 Control of Nonholonomic Systems

The term "nonholonomic system" originates from classical mechanics and has its widely accepted meaning as a "Lagrange system with linear constraints being nonintegrable". A mechanical system is said to be nonholonomic if its generalized velocity satisfies an equality condition that cannot be written as an equivalent condition on the generalized position, see [30]. Control of nonholonomic dynamic systems has formed an active area in the control community - see surveys by Kolmanovsky and McClamroch in [31], Canudas de Wit et al. in [15], Murray and Sastry in [32], and references therein for an overview and interesting introductory examples in this expanding area.

Nonholonomic systems have inherent difficulties in feedback stabilization at the origin or at a given equilibrium point since the tangent linearization of these systems is uncontrollable. In fact, a direct application of Brockett's necessary condition, see Section 2.7.4 for more details, for feedback stabilization implies that nonholonomic systems cannot be stabilized by any stationary continuous state feedback although they are open loop controllable. As a consequence, the classical smooth control theory cannot be applied. This motivates researchers to seek novel approaches. These approaches can be roughly classified into discontinuous feedback, see for example [33-45] and time-varying feedback, see for example, [15, 32, 46-48]. The discontinuous feedback approach often uses the state scaling originated from the $\sigma$ process [49] and a switching control strategy to overcome the difficulty due to the loss of controllability. This approach results in a fast transient response and usually an exponential convergence can be achieved. The drawback is discontinuity in the control input. On the other hand, the time-varying feedback approach provides
a smooth/continuous controller, i.e., no switching is required, however the price is slow convergence. The stability analysis is often based on linear time-varying system theory and Barbalat's lemma. The backstepping technique [3] is usually used for high-order chained form systems in both discontinuous and time-varying approaches. Those aforementioned systems are either driftless or have weak nonlinear drifts. When nonholonomic systems are perturbed by drifts with uncertainties, robust and adaptive control approaches are often applied. The robust control design schemes are based on the size domination concept [50]. The control is conservative when a priori knowledge of uncertainties is poor. A class of nonholonomic systems with strong nonlinear uncertainties was recently considered in [51]. Discontinuous state feedback and output feedback controllers were designed to achieve global exponential stability. However the $x_{0}$-subsystem is required to be Lipschitz since a constant control input $u_{0}$ is used to get around the difficulty due to the loss of controllability. The adaptive approach [38,40, 46] provides less conservative control input but increases the dynamics of the closed loop system. The systems studied in these papers do not allow drifts in the $x_{0}$-subsystem. A difficulty in designing adaptive stabilization controllers for chained systems with drifts is that the state of the $x_{0}$-subsystem can have several zero crossings due to transient behavior of the unknown parameter estimate. This phenomenon causes difficulties in applying the state scaling. For a solution of the stabilization of nonholonomic systems in a chained form with strong nonlinear drifts and unknown parameters, the reader is referred to [52].

### 4.2.2 Control of Underactuated Ships and Underwater Vehicles

Control of underactuated ships and autonomous underwater vehicles (AUVs) is an active field due to its important applications such as passenger and goods transportation, environmental surveying, undersea cable inspection, and offshore oil installations.

Based on its practical requirement, motion control of underactuated ocean vessels has been divided into three areas: Stabilization, trajectory-tracking, and pathfollowing. These control problems are challenging due to the fact that the motion of underactuated surface ships and AUVs possesses more degrees of freedom to be controlled than the number of the independent controls under some nonintegrable second-order nonholonomic constraints [29,53,54]. In particular, underactuated ships do not usually have an actuator in the sway axis while in the case of AUVs there are no actuators in the sway and heave directions. This configuration is by far the most common among marine vessels. Therefore, Brockett's condition indicates that any continuous time-invariant feedback control law does not make a null solution of the underactuated surface ship and AUV dynamics asymptotically stable in the sense of Lyapunov. Furthermore as observed in [22, 54], the underactuated ship and AUV system is not transformable into a standard chain system. Consequently, existing control schemes [15, 32-48] developed for chained systems
cannot be applied directly. Nevertheless, in the past decade, stabilization, trajectorytracking control, and path-following of underactuated ocean vessels have been studied separately from different viewpoints.

### 4.2.2.1 Stabilization

An underactuated ocean vessel belongs to a class of underactuated mechanical systems subject to some nonintegrable second-order nonholonomic constraints, see [29,53,54]. Therefore, design of a feedback stabilizer using linear and classical nonlinear control theories is not possible. There are two main approaches to deal with stabilization of an underactuated ocean vessel. They are (time-invariant and timevarying) discontinuous feedback and time-varying continuous/smooth feedback. We here mention some typical results of both approaches.

## Time-invariant and Time-varying Discontinuous Approach

A discontinuous state feedback control law was proposed in [55] using the $\sigma$-process to exponentially stabilize an underactuated ship at the origin where the ship model is discontinuously transformed to an extended chained form system. The dynamics of an underactuated ship is considered in [55], see also Section 3.4.1:

$$
\begin{align*}
& \dot{\boldsymbol{\eta}}=\boldsymbol{J}(\boldsymbol{\eta}) \boldsymbol{v}, \\
& \boldsymbol{M} \dot{\boldsymbol{v}}=-\boldsymbol{C}(\boldsymbol{v}) \boldsymbol{v}-\boldsymbol{D} \boldsymbol{v}+\boldsymbol{\tau}, \\
& \boldsymbol{\eta}=[x, y, \psi]^{T}, \boldsymbol{v}=[u, v, r]^{T}, \boldsymbol{\tau}=\left[\tau_{u}, 0, \tau_{r}\right]^{T}, \tag{4.26}
\end{align*}
$$

where $(x, y)$ denotes the earth-fixed position of the center of mass of the ship, $\psi$ denotes the orientation angle, $(u, v)$ and $r$ are the linear and angular velocities in the body-fixed frame, and $\left(\tau_{u}, \tau_{r}\right)$ are the surge force and yaw moment. The matrices $\boldsymbol{J}(\boldsymbol{\eta}), \boldsymbol{M}, \boldsymbol{C}(\boldsymbol{v})$, and $\boldsymbol{D}$ are given by

$$
\begin{align*}
& \boldsymbol{J}(\boldsymbol{\eta})=\left[\begin{array}{ccc}
\cos (\psi) & -\sin (\psi) & 0 \\
\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right], \boldsymbol{M}=\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & 0 \\
0 & 0 & m_{33}
\end{array}\right], \\
& \boldsymbol{C}(\boldsymbol{v})=\left[\begin{array}{ccc}
0 & 0 & -m_{22} v \\
0 & 0 & m_{11} u \\
m_{22} v & -m_{11} u & 0
\end{array}\right], \boldsymbol{D}=\left[\begin{array}{ccc}
d_{11} & 0 & 0 \\
0 & d_{22} & 0 \\
0 & 0 & d_{33}
\end{array}\right], \tag{4.27}
\end{align*}
$$

where $m_{11}, m_{22}$, and $m_{33}$ denote the ship inertia including added mass, and $d_{11}$, $d_{22}$, and $d_{33}$ are hydrodynamic damping constants, see Chapter 3 for more details. The control objective is to design the control inputs $\tau_{u}$ and $\tau_{r}$ to stabilize (4.26) asymptotically at the origin. In [55], the coordinate transformation

$$
\left[\begin{array}{c}
x_{1}  \tag{4.28}\\
x_{2} \\
x_{3} \\
x_{4} \\
x_{5} \\
x_{6}
\end{array}\right]=\left[\begin{array}{c}
\psi \\
x \cos (\psi)+y \sin (\psi) \\
-x \sin (\psi)+y \cos (\psi) \\
v \\
r \\
u
\end{array}\right]
$$

is used to transform the ship model (4.26) to the following system

$$
\begin{align*}
& \dot{x}_{1}=x_{5}, \\
& \dot{x}_{2}=x_{6}+x_{3} x_{5}, \\
& \dot{x}_{3}=x_{4}-x_{2} x_{5}, \\
& \dot{x}_{4}=-\alpha x_{4}-\beta x_{5} x_{6},  \tag{4.29}\\
& \dot{x}_{5}=\Omega_{1}, \\
& \dot{x}_{6}=\Omega_{2},
\end{align*}
$$

where $\alpha=d_{22} / m_{22}, \beta=m_{11} / m_{22}$, and

$$
\begin{equation*}
\Omega_{1}=\frac{\tau_{r}-d_{33} r+\left(m_{11}-m_{22}\right) u v}{m_{33}}, \Omega_{2}=\frac{\tau_{u}+m_{22} v r-d_{11} u}{m_{11}} \tag{4.30}
\end{equation*}
$$

It can be seen that the system (4.29) consists of two subsystems, namely ( $x_{1}, x_{2}, x_{3}, x_{4}$ ) and ( $x_{5}, x_{6}$ ), connected to each other in a strict feedback form [3]. The control design can be simply carried out in two steps as follows.

## Step 1

In this step, the author of [55] considers the first four equations of (4.29), and $\left(x_{5}, x_{6}\right)$ as controls $\left(v_{1}, v_{2}\right)$. With the assumption of $x_{1} \neq 0$, the coordinate transformation ( $\sigma$-process)

$$
\begin{equation*}
y=x_{1}, z_{1}=x_{2}, z_{2}=\frac{x_{3}}{x_{1}}, z_{3}=\frac{x_{4}}{x_{1}} \tag{4.31}
\end{equation*}
$$

results in

$$
\begin{align*}
& \dot{y}=v_{1}, \\
& \dot{z}_{1}=v_{2}+y z_{2} v_{1} \\
& \dot{z}_{2}=z_{3}-\frac{z_{1}+z_{2}}{y} v_{1},  \tag{4.32}\\
& \dot{z}_{3}=-\alpha z_{3}-\frac{z_{3}+\beta v_{2}}{y} v_{1} .
\end{align*}
$$

The feedback control law is designed as

$$
\begin{align*}
& v_{1}=-k_{1} y \\
& v_{2}=-k_{21} z_{1}-k_{22} z_{2}-k_{23} z_{3} \tag{4.33}
\end{align*}
$$

where $k_{1}, k_{21}, k_{22}$, and $k_{23}$ are the control gains chosen such that the matrix

$$
A_{1}=\left[\begin{array}{ccc}
-k_{21} & -k_{22} & -k_{23}  \tag{4.34}\\
k_{1} & k_{1} & 1 \\
-k_{1} \beta k_{21} & -k_{1} \beta k_{22} & k_{1}-\alpha-k_{1} \beta k_{23}
\end{array}\right]
$$

is Hurwitz.

## Step 2

At this step, the last two equations of (4.29) are considered. Using the standard backstepping technique results in the following control law

$$
\begin{align*}
\Omega_{1} & =-k_{3}\left(x_{5}-v_{1}\right)-k_{1} x_{5} \\
\Omega_{2} & =-k_{4}\left(x_{6}-v_{2}\right)-k_{21}\left(x_{6}+x_{3} x_{5}\right)  \tag{4.35}\\
& -k_{22} \frac{x_{4}-x_{2} x_{5}}{x_{1}}+k_{23} \frac{\alpha x_{4}+\beta x_{5} x_{6}}{x_{1}}+k_{22} \frac{x_{3} x_{5}}{x_{1}^{2}}+k_{23} \frac{x_{4} x_{5}}{x_{1}^{2}},
\end{align*}
$$

where $k_{3}>k_{1}$ and $k_{4}$ are positive constants. The actual controls $\tau_{u}$ and $\tau_{r}$ can be found from (4.35) and (4.30). In [55] it is shown that if the initial conditions $x_{1}\left(t_{0}\right) \neq 0$ and $x_{1}\left(t_{0}\right)\left(x_{5}\left(t_{0}\right)+k_{1} x_{1}\left(t_{0}\right)\right) \geq 0$ then $\left(x_{1}(t), x_{2}(t), x_{3}(t), x_{4}(t), x_{5}(t)\right.$, $\left.x_{6}(t)\right)$ is bounded for all $t \geq t_{0} \geq 0$, and exponentially converges to zero. If the above conditions do not hold, the controls

$$
\begin{align*}
& \Omega_{1}=-\left|x_{1}-\epsilon\right|^{a} \operatorname{sign}\left(x_{1}-\epsilon\right)-\left|x_{5}\right|^{b} \operatorname{sign}\left(x_{5}\right), \\
& \Omega_{2}=0, \tag{4.36}
\end{align*}
$$

with $\epsilon \neq 0, b \in(0,1)$, and $a>b /(2-b)$ being constants, can be used to make the above conditions hold in finite time. For more details, the reader is referred to [55].

Remark 4.1. The aforementioned discontinuous stabilizer provides a fast convergence of the stabilizing errors to zero. However, the control inputs $\tau_{u}$ and $\tau_{r}$ are discontinuous. Moreover, under arbitrarily small nonvanishing environmental disturbances induced by waves, wind, and ocean currents, the closed loop system consisting of (4.35) and (4.29) can be unstable in the sense that the states $\left(x_{1}(t), x_{2}(t)\right.$, $\left.x_{3}(t), x_{4}(t), x_{5}(t), x_{6}(t)\right)$ can go to infinity exponentially fast.

The work mentioned in $[22,53,54,56-58]$ can also be grouped in the discontinuous approach. The authors of [22] developed a discontinuous time-varying feedback stabilizer for a nonholonomic system and applied it to underactuated ships. Some local exponential stabilization results were reported in $[53,54]$ based on the time-varying homogeneous control approach. An application of averaging and backstepping tech-
niques was proposed in [59] to design a global practical controller for stabilization and tracking control of surface ships. Experimental results on dynamic positioning of underactuated ships were reported in [56]. By transforming the underactuated ship kinematics and dynamics into the so-called skew form, some dynamic feedback results on stabilization were given in [57]. In [58], the authors proposed a discontinuous solution to the problem of steering an underactuated AUV to a point with desired orientation using the polar coordinate transformation motivated from the work in [60].

## Time-varying Continuous/Smooth Approach

A typical result on stabilization of the underactuated ship (4.26) in the time-varying continuous/smooth approach is given in [61]. In [61], the coordinate transformations (similar to the ones given in (4.30) and (4.28))

$$
\begin{align*}
& z_{1}=\cos (\psi) x+\sin (\psi) y, \\
& z_{2}=-\sin (\psi) x+\cos (\psi) y, \\
& z_{3}=\psi \\
& \Omega_{1}=\frac{\tau_{r}-d_{33} r+\left(m_{11}-m_{22}\right) u v}{m_{33}},  \tag{4.37}\\
& \Omega_{2}=\frac{\tau_{u}+m_{22} v r-d_{11} u}{m_{11}}
\end{align*}
$$

are first used to transform the ship system (4.26) to

$$
\begin{align*}
& \dot{z}_{1}=u+z_{2} r, \\
& \dot{z}_{2}=v-z_{1} r, \\
& \dot{z}_{3}=r, \\
& \dot{u}=\Omega_{2},  \tag{4.38}\\
& \dot{v}=-c u r-d v, \\
& \dot{r}=\Omega_{1},
\end{align*}
$$

where $c=m_{11} / m_{22}$ and $d=d_{22} / m_{22}$. Then the following nontrivial coordinate transformations

$$
\begin{align*}
& Z_{2}=z_{2}+\frac{v}{d}, \\
& u=-\frac{d}{c} z_{1}-\frac{d}{c} \mu, \\
& \Omega_{2 \mu}=\frac{d}{c} z_{1}+\frac{d}{c} \mu-Z_{2} r+\frac{v}{d} r-\frac{c}{d} \Omega_{2} \tag{4.39}
\end{align*}
$$

are applied to (4.38) to obtain the system

$$
\begin{align*}
& \dot{z}_{1}=-\frac{d}{c} z_{1}-\frac{d}{c} \mu+Z_{2} r-\frac{v}{d} r, \\
& \dot{Z}_{2}=\mu r \\
& \dot{z}_{3}=r  \tag{4.40}\\
& \dot{v}=-d v+d\left(z_{1}+\mu\right) r, \\
& \dot{\mu}=\Omega_{2 \mu} \\
& \dot{r}=\Omega_{1}
\end{align*}
$$

Let $k_{2}, k_{3}, k_{\mu}$, and $k_{r}$ be strictly positive constants such that $1 \geq k_{2} \geq k_{3}$. The controls $\Omega_{1}$ and $\Omega_{2 \mu}$ are designed in [61] as

$$
\begin{align*}
& \Omega_{1}=-k_{r}\left(r-r_{f}\right)+\dot{r}_{f}-\lambda\left(Z_{2} \mu_{f}+2 Z_{3}+2 Z_{2} k_{2} \cos (t) \mu_{f}\right) \\
& \Omega_{2 \mu}=-k_{\mu}\left(\mu-\mu_{f}\right)+\dot{\mu}_{f}-\lambda\left(Z_{2}+2 Z_{3} k_{2} \cos (t)\right) r \tag{4.41}
\end{align*}
$$

where

$$
\begin{align*}
& \lambda=2+\frac{k_{3}}{3}-\frac{k_{3} \sin (2 t)}{6} \frac{2 V_{1}+V_{1}^{2}}{\left(1+V_{1}\right)^{2}}, \\
& Z_{3}=z_{3}+k_{2} \cos (t) Z_{2} \\
& V_{1}=Z_{2}^{2}+2 Z_{3}^{2} \\
& \mu_{f}=-\frac{\sin (t) Z_{2}^{2}}{2\left(0.001+Z_{2}^{2}\right)}  \tag{4.42}\\
& r_{f}=\frac{-k_{3} Z_{3}+k_{2} \sin (t) Z_{2}}{1+k_{2} \cos (t) \mu_{f}}
\end{align*}
$$

In [61], it is proven that the closed loop system consisting of (4.41), (4.39), (4.37), and (4.26) is GAS at the origin.

Remark 4.2. The design of the feedback given in (4.41) is nontrivial. Overall, convergence of the stabilizing errors to zero is slow. This is a well-known phenomenon of the continuous/smooth time-varying approach applying not to only underactuated ships but also to mobile robots. Moreover, since the stabilizer design mentioned above is nontrivial, it is difficult to extend the control design scheme to solve a trajectory-tracking problem, see the next section. In addition, the physical meaning of the feedback is not clear.

In addition to the aforementioned results on stabilization of underactuated vessels, the following results are also related to the topic under discussion. In [62], several control configurations were considered, and a technique for synthesizing open loop controls was given. The first control scheme with the dynamic AUV model taken into account was proposed in [63]. A kinematic drift free model of the underwater vehicles with four control inputs was used to design a regulation controller in [64]. The authors of [65] proposed a controller that is able to stabilize an AUV to some equilibria based on the interconnection and damping assignment passivitybased control approach, which has been successfully applied to many other mechan-
ical systems [66]. See also [67] for stabilization results of underactuated mechanical systems on Riemannian manifolds.

### 4.2.2.2 Trajectory-tracking

Trajectory-tracking is here defined as a control problem of forcing an underactuated surface ship or AUV to track a reference trajectory generated by a suitable vessel model, i.e., the vessel model that has the same parameters as the real one. There are two main approaches to solve the trajectory-tracking control problems. The first approach is based on linear time-varying control system theory while the second approach relies on the Lyapunov direct method. We here briefly describe typical results of the two approaches.

## Linear Time-varying Approach

A typical work belonging to this approach is given in [68] on a global $K$-exponential tracking result for the underactuated ship (4.26). In [68], the authors consider a problem of designing the control $\tau_{u}$ and $\tau_{r}$ to force the position $(x, y)$ and orientation $\psi$ of the ship (4.26) to track the reference position $\left(x_{d}, y_{d}\right)$ and orientation $\psi_{d}$ generated by the reference ship model

$$
\begin{align*}
& \dot{\boldsymbol{\eta}_{\boldsymbol{d}}}=\boldsymbol{J}\left(\boldsymbol{\eta}_{\boldsymbol{d}}\right) \boldsymbol{v}_{\boldsymbol{d}}, \\
& \boldsymbol{M} \dot{\boldsymbol{v}}_{\boldsymbol{d}}=-\boldsymbol{C}\left(\boldsymbol{v}_{\boldsymbol{d}}\right) \boldsymbol{v}_{\boldsymbol{d}}-\boldsymbol{D} \boldsymbol{v}_{\boldsymbol{d}}+\boldsymbol{\tau}_{\boldsymbol{d}}, \\
& \boldsymbol{\eta}_{\boldsymbol{d}}=\left[\begin{array}{l}
x_{d} \\
y_{d} \\
\psi_{d}
\end{array}\right], \boldsymbol{v}_{\boldsymbol{d}}=\left[\begin{array}{l}
u_{d} \\
v_{d} \\
r_{d}
\end{array}\right], \boldsymbol{\tau}_{\boldsymbol{d}}=\left[\begin{array}{c}
\tau_{u d} \\
0 \\
\tau_{r d}
\end{array}\right] . \tag{4.43}
\end{align*}
$$

In [68], the coordinate transformations

$$
\left\{\begin{array} { l } 
{ z _ { 1 } = \operatorname { c o s } ( \psi ) x + \operatorname { s i n } ( \psi ) y , }  \tag{4.44}\\
{ z _ { 2 } = - \operatorname { s i n } ( \psi ) x + \operatorname { c o s } ( \psi ) y , } \\
{ z _ { 3 } = \psi , }
\end{array} \quad \left\{\begin{array}{l}
z_{1 d}=\cos \left(\psi_{d}\right) x_{d}+\sin \left(\psi_{d}\right) y_{d} \\
z_{2 d}=-\sin \left(\psi_{d}\right) x_{d}+\cos \left(\psi_{d}\right) y_{d} \\
z_{3 d}=\psi_{d}
\end{array}\right.\right.
$$

and the tracking errors

$$
\begin{align*}
& u_{e}=u-u_{d}, v_{e}=v-v_{d}, r_{e}=r-r_{d} \\
& z_{1 e}=z_{1}-z_{1 d}, z_{2 e}=z_{2}-z_{2 d}, z_{3 e}=z_{3}-z_{3 d} \tag{4.45}
\end{align*}
$$

are used to obtain the tracking error dynamics of a chained form

$$
\dot{u}_{e}=\frac{m_{22}}{m_{11}}\left(v_{e} r_{e}+v_{e} r_{d}+v_{d} r_{e}\right)-\frac{d_{11}}{m_{11}} u_{e}+\frac{1}{m_{11}}\left(\tau_{u}-\tau_{u d}\right),
$$

$$
\begin{align*}
& \dot{v}_{e}=-\frac{m_{11}}{m_{22}}\left(u_{e} r_{e}+u_{e} r_{d}+u_{d} r_{e}\right)-\frac{d_{22}}{m_{22}} v_{e}, \\
& \dot{r}_{e}=\frac{m_{11}-m_{22}}{m_{33}}\left(u_{e} v_{e}+u_{e} v_{d}+u_{d} v_{e}\right)-\frac{d_{33}}{m_{33}} r_{e}+\frac{1}{m_{33}}\left(\tau_{r}-\tau_{r d}\right),  \tag{4.46}\\
& \dot{z}_{1 e}=u_{e}+z_{2 e} r_{e}+z_{2 e} r_{d}+z_{2 d} r_{e}, \\
& \dot{z}_{2 e}=v_{e}-z_{1 e} r_{e}-z_{1 e} r_{d}-z_{1 d} r_{e}, \\
& \dot{z}_{3 e}=r_{e} .
\end{align*}
$$

Assuming that $u_{d}, v_{d}, z_{1 d}$, and $z_{2 d}$ are bounded, and that $r_{d}(t)$ is persistently exciting, the controls

$$
\begin{align*}
\tau_{u}= & \tau_{u d}-k_{1} u_{e}+k_{2} r_{d} v_{e}-k_{3} z_{1 e}+k_{4} r_{d} z_{2 e}, \\
\tau_{r}= & \tau_{r d}-\left(m_{11}-m_{22}\right)\left(u_{e} v_{e}+v_{d} u_{e}+u_{d} v_{e}\right)-  \tag{4.47}\\
& k_{5} r_{e}-k_{6} z_{3 e}
\end{align*}
$$

where the control gains $k_{i}, i=1, \ldots, 6$ satisfy

$$
\begin{align*}
& k_{1}>d_{22}-d_{11}, \\
& k_{2}=\frac{m_{22} k_{4}\left(k_{4}+k_{1}+d_{11}-d_{22}\right)}{d_{22} k_{4}+m_{11} k_{3}}, \\
& 0<k_{3}<\left(k_{1}+d_{11}-d_{22}\right) \frac{d_{22}}{m_{11}},  \tag{4.48}\\
& k_{4}>0 \\
& k_{5}>-d_{33}, \\
& k_{6}>0,
\end{align*}
$$

make the closed loop system consisting of (4.47), (4.43), and (4.26), that is,

$$
\begin{aligned}
{\left[\begin{array}{c}
\dot{u}_{e} \\
\dot{v}_{e} \\
\dot{z}_{1 e} \\
\dot{z}_{2 e}
\end{array}\right]=} & {\left[\begin{array}{cccc}
-\frac{k_{1}+d_{11}}{m_{11}} & \frac{k_{2}+m_{22}}{m_{11}} r_{d}(t) & -\frac{k_{3}}{m_{11}} & \frac{k_{4}}{m_{11}} r_{d}(t) \\
-\frac{m_{11}}{m_{22}} r_{d}(t) & -\frac{d_{22}}{m_{22}} & 0 & 0 \\
1 & 0 & 0 & r_{d}(t) \\
0 & 1 & -r_{d}(t) & 0
\end{array}\right]\left[\begin{array}{c}
u_{e} \\
v_{e} \\
z_{1 e} \\
z_{2 e}
\end{array}\right]+} \\
& {\left[\begin{array}{cc}
\frac{m_{22}}{m_{11}}\left(v_{e}+v_{d}\right) & 0 \\
-\frac{m_{11}}{m_{22}}\left(u_{e}+u_{d}\right) & 0 \\
z_{2 e}+z_{2 d} & 0 \\
-\left(z_{1 e}+z_{1 d}\right) & 0
\end{array}\right]\left[\begin{array}{c}
r_{e} \\
z_{3 e}
\end{array}\right], }
\end{aligned}
$$

$$
\left[\begin{array}{c}
\dot{r}_{e}  \tag{4.49}\\
\dot{z}_{3 e}
\end{array}\right]=\left[\begin{array}{cc}
-\frac{d_{33}+k_{5}}{m_{33}} & -\frac{k_{6}}{m_{33}} \\
1 & 0
\end{array}\right]\left[\begin{array}{c}
r_{e} \\
z_{3 e}
\end{array}\right]
$$

globally $K$-exponentially stable at the origin. Proof of $K$-exponential stability of the above closed loop system is straightforward using the results on stability of cascade systems in [17] and [69], and those of linear time-varying system theory in [6], see [68] for details. It should be noted that the persistently exciting condition on the yaw reference velocity, $r_{d}$, is required to prove $K$-exponential stability of the closed loop system (4.49).

Remark 4.3. The persistent exciting condition on the yaw reference velocity $r_{d}$ excludes a straight-line reference trajectory. In comparison with the direct Lyapunov approach summarized below, it is difficult to deal with any external disturbances and/or actuator dynamics using the linear time-varying approach.

## Direct Lyapunov Approach

The direct Lyapunov method has been widely used in designing controllers for underactuated ocean vessels. However, the use of the Lyapunov direct method for designing control systems for underactuated ocean vessels is not straightforward due to the underactuated nature of ocean vessels. We here describe typical results of trajectory-tracking control based on the Lyapunov direct method.

## Local Trajectory-tracking Results

An application of the recursive technique proposed in [70] for the standard chainform systems yields a high-gain based local tracking result in [59] for surface ships. The experimental results of this proposed controller were reported in [71]. The control design in $[59,71]$ starts from $(4.46)$ as follows. First of all, the following restrictive assumption is made on the yaw reference velocity:

$$
\begin{equation*}
0<r_{d \min }<\left|r_{d}(t)\right|<r_{d \max } \tag{4.50}
\end{equation*}
$$

where $r_{d \text { min }}$ and $r_{d \text { max }}$ are strictly positive constants. Motivated by the work in [70], the authors define new error variables as

$$
\begin{align*}
& \omega_{1}=z_{1 e}-z_{2} z_{3 e} \\
& \omega_{2}=z_{2 e}+z_{1} z_{3 e} \\
& y_{1}=v_{e}+c u z_{3 e}+k_{2} \omega_{2}  \tag{4.51}\\
& y_{2}=u_{e}+k_{1} \omega_{1}-\frac{k_{2}\left(d-k_{2}\right)}{c r_{d}} \omega_{2} \\
& y_{3}=z_{3 e}
\end{align*}
$$

where $c=m_{11} / m_{22}, d=d_{22} / m_{22}$, and $k_{2}$ is a parameter to be determined later. It should be stressed that the condition (4.50) on the yaw reference velocity, $r_{d}$ is required so that the error transformations (4.51) are valid. With (4.51), the tracking error system (4.46) is written in a triangular-like structure as follows:

$$
\begin{align*}
\dot{\omega}_{1}= & y_{2}-k_{1} \omega_{1}+\frac{k_{2}\left(d-k_{2}\right)}{c r_{d}} \omega_{2}+\omega_{2} r_{d}-\left(v-z_{1} r_{e}\right) y_{3}, \\
\dot{\omega}_{2}= & y_{1}-k_{2} \omega_{2}-\omega_{1} r_{d}+\left((1-c) u+z_{2} r_{e}\right) y_{3}, \\
\dot{y}_{1}= & -c y_{2} r_{d}+\left(c k_{1}-k_{2}\right) r_{d} \omega_{1}-\left(d-k_{2}\right) y_{1}+\left(c \Omega_{2}+c d u+\right. \\
& \left.k_{2}\left((1-c) u+z_{2} r_{e}\right)\right) y_{3}, \\
\dot{y}_{2}= & \Omega_{2}-\Omega_{2 d}+k_{1} y_{2}-k_{1}^{2} \omega_{1}+\frac{1}{c r_{d}} k_{1} k_{2}\left(d-k_{2}\right) \omega_{2}+k_{1} r_{d} \omega_{2}+ \\
& \frac{\dot{r}_{d}}{c r_{d}^{2}} k_{2}\left(d-k_{2}\right) \omega_{2}-\frac{1}{c r_{d}} k_{2}\left(d-k_{2}\right)\left(y_{1}-k_{2} \omega_{2}-r_{d} \omega_{1}\right)- \\
& \left(k_{1}\left(v-z_{1} r_{e}\right)+\frac{1}{c r_{d}} k_{2}\left(d-k_{2}\right)\left((1-c) u+z_{2} r_{e}\right)\right) y_{3}, \\
\dot{y}_{3}= & r_{e}, \\
\dot{r}_{e}= & \Omega_{1}-\Omega_{1 d}, \tag{4.52}
\end{align*}
$$

where $\Omega_{1}$ and $\Omega_{2}$ are given in (4.30), and

$$
\begin{align*}
& \Omega_{1 d}=\frac{\tau_{r d}-d_{33} r_{d}+\left(m_{11}-m_{22}\right) u_{d} v_{d}}{m_{33}} \\
& \Omega_{2 d}=\frac{\tau_{u d}+m_{22} v_{d} r_{d}-d_{11} u_{d}}{m_{11}} \tag{4.53}
\end{align*}
$$

The triangular structure (4.52) allows us to use the backstepping technique [3] to design the controls $\Omega_{1}$ and $\Omega_{2}$. In [71], the the controls $\Omega_{1}$ and $\Omega_{2}$ are designed as

$$
\begin{align*}
\Omega_{1}= & -a_{3}\left(r-\alpha_{r}\right)+\dot{\alpha}_{r}-\kappa y_{3}, \\
\Omega_{2}= & -a_{1} y_{2}-\gamma \omega_{1}+c a r_{d} y_{1}-\left(-\Omega_{2 d}+k_{1} y_{2}-k_{1}^{2} \omega_{1}+\right. \\
& \frac{1}{c r_{d}} k_{1} k_{2}\left(d-k_{2}\right) \omega_{2}+k_{1} r_{d} \omega_{2}+\frac{\dot{r}_{d}}{c r_{d}^{2}} k_{2}\left(d-k_{2}\right) \omega_{2}-\frac{1}{c r_{d}} \times \\
& \left.k_{2}\left(d-k_{2}\right)\left(y_{1}-k_{2} \omega_{2}-r_{d} \omega_{1}\right)\right), \tag{4.54}
\end{align*}
$$

where

$$
\begin{aligned}
\alpha_{r}= & \left(\lambda+\gamma\left(\omega_{1} z_{1}+\omega_{2} z_{2}\right)+a k_{2} y_{1} z_{2}+k_{1} y_{2} z_{1}-\frac{1}{c r_{d}} k_{2}\left(d-k_{2}\right) y_{2} z_{2}\right)^{-1} \times \\
& \left(-a_{2} y_{3}+\gamma \omega_{1} v-\gamma \omega_{2}(1-c) u-a y_{1}\left(c\left(\Omega_{2}+d u\right)+k_{2}(1-c) u\right)+\right.
\end{aligned}
$$

$$
\begin{align*}
& \left.k_{1} y_{2} v+\frac{1}{c r_{d}} k_{2}\left(d-k_{2}\right)(1-c) u y_{2}\right)+r_{d}, \\
\kappa= & \lambda+\gamma \omega_{1} z_{1}+\gamma \omega_{2} z_{2}+a k_{2} y_{1} z_{2}+k_{1} y_{2} z_{1}-\frac{1}{c r_{d}} k_{2}\left(d-k_{2}\right) y_{2} z_{2} . \tag{4.55}
\end{align*}
$$

In (4.54) and (4.55), the control parameters $k_{1}, k_{2}, a, a_{1}, a_{2}, a_{3}$, and $\lambda$ are positive constants, and are chosen such that

$$
\begin{align*}
& k_{2}<d, k_{1}>\frac{k_{2}\left(d-k_{2}\right)^{2}}{c^{2} r_{d \min }^{2}} \\
& \frac{1}{k_{2}\left(d-k_{2}\right)}<\frac{a}{\gamma}<\frac{k_{1}\left(d-k_{2}\right)}{\left(c k_{1}-k_{2}\right)^{2} r_{d \max }^{2}} \tag{4.56}
\end{align*}
$$

It is noted that the virtual control $\alpha_{r}$ given in (4.55) is solvable if and only if

$$
\begin{equation*}
\lambda>-\left(\gamma \omega_{1} z_{1}+\gamma \omega_{2} z_{2}+a k_{2} y_{1} z_{2}+k_{1} y_{2} z_{1}-\frac{1}{c r_{d}} k_{2}\left(d-k_{2}\right) y_{2} z_{2}\right) \tag{4.57}
\end{equation*}
$$

Proof of local exponential stability of the closed loop system consisting of (4.54), (4.46), and (4.30) can be carried out by using the Lyapunov function

$$
\begin{equation*}
V=\frac{1}{2} \gamma \omega_{1}^{2}+\frac{1}{2} \gamma \omega_{2}^{2}+\frac{1}{2} a y_{1}^{2}+\frac{1}{2} y_{2}^{2}+\frac{\lambda}{2} y_{3}^{2}+\frac{1}{2}\left(r-\alpha_{r}\right)^{2} . \tag{4.58}
\end{equation*}
$$

Remark 4.4. There are two limitations of the aforementioned tracking controllers. These limitations are described in conditions (4.50) and (4.57). The condition (4.50) implies that the reference yaw velocity $r_{d}$ cannot be zero at any time. This restrictive condition excludes a straight-line reference trajectory. The condition (4.57) implies that the aforementioned trajectory-tracking result is inherently local. One can argue that by the control parameters $k_{1}, k_{2}, a, a_{1}, a_{2}, a_{3}$, and $\lambda$ may increase the size of the attraction region. However, it is very hard to ensure this property since the control parameters must satisfy various conditions specified in (4.56). In fact, this is true, as said in [71].

## Global Trajectory-tracking Results

Based on Lyapunov's direct method and the passivity approach [72], two restricted tracking solutions of an underactuated surface ship were proposed in [19]. We here briefly summarize the result based on the passivity approach in [19]. The result based on the standard backstepping technique [3] is discussed later. In [19], the starting point is the tracking error system (4.46). The passivity based method consists of two steps as follows.

## Step 1

Design of the surge force $\tau_{u}$ : This force is designed based on the Lyapunov function

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right)^{2}+\frac{1}{2} z_{2 e}^{2}+\frac{\lambda_{0}}{2} v_{e}^{2}+\frac{1}{2} \bar{u}_{e}^{2} \tag{4.59}
\end{equation*}
$$

where $\bar{u}_{e}=u_{e}-\alpha_{0}$ with

$$
\begin{equation*}
\alpha_{0}=-\lambda_{2}\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right) . \tag{4.60}
\end{equation*}
$$

In (4.59) and (4.60), the control parameters $\lambda_{0}, \lambda_{1}$ and $\lambda_{2}$ are chosen such that

$$
\begin{equation*}
c(t)=\min \left(2 \epsilon\left(\lambda_{2}-\lambda_{1} r_{d}^{2}(t)\right), 2\left(\lambda_{1} r_{d}^{2}(t)-\frac{m_{22}}{(1-\epsilon) \lambda_{0} d_{22}}\right), 2 \epsilon \frac{d_{22}}{m_{22}}, 2 c_{1}\right) \geq c^{*} \tag{4.61}
\end{equation*}
$$

where $c_{1}$ is a positive constant, $0<\epsilon<1$, and $c^{*}$ is strictly positive. In (4.59) and (4.60), $\alpha_{0}$ is understood as a virtual control of $u_{e}$. From the first time derivative of the Lyapunov function $V_{1}$ given in (4.59) along the solutions of (4.46), a choice of the surge force $\tau_{u}$

$$
\begin{align*}
\tau_{u}= & \tau_{u d}+m_{11}\left[-\frac{m_{22}}{m_{11}}\left(v r-v_{d} r_{d}\right)+\frac{d_{11}}{m_{11}} u_{e}-c_{1}\left(u_{e}+\lambda_{2}\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right)\right)-\right. \\
& \left(\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right)-\frac{\lambda_{0} m_{11}}{m_{22}} r_{d} v_{e}\right)-\lambda_{2}\left(u_{e}+z_{2 e} r_{d}+z_{2} r_{e}\right)+ \\
& \left.\lambda_{1} \lambda_{2} \dot{r}_{d} z_{2 e}+\lambda_{1} \lambda_{2} r_{d}\left(v_{e}-z_{1 e} r_{d}-z_{1} r_{e}\right)\right] \tag{4.62}
\end{align*}
$$

gives

$$
\begin{equation*}
\dot{V}_{1} \leq-c(t) V_{1}+\left[\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right)\left(z_{2}+\lambda_{1} r_{d} z_{1}\right)-z_{2 e} z_{1}-\frac{\lambda_{0} m_{11}}{m_{22}} v_{e} u\right] r_{e}, \tag{4.63}
\end{equation*}
$$

where $c(t)$ is given in (4.61).

## Step 2

Design of the yaw moment $\tau_{r}$ : This moment is designed based on the Lyapunov function

$$
\begin{equation*}
V_{2}=V_{1}+\frac{1}{2} z_{3 e}^{2}+\frac{1}{2} \bar{r}_{e}^{2} \tag{4.64}
\end{equation*}
$$

where $\bar{r}_{e}=r_{e}-\alpha_{1}$ and

$$
\begin{equation*}
\alpha_{1}=-c_{2}\left[\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right)\left(z_{2}+\lambda_{1} r_{d} z_{1}\right)-z_{2 e} z_{1}-\frac{\lambda_{0} m_{11}}{m_{22}} v_{e} u+z_{3 e}\right] \tag{4.65}
\end{equation*}
$$

with $c_{2}>0$. From the first time derivative of the Lyapunov function $V_{2}$ given in (4.64) along the solutions of (4.46), a choice of the yaw moment

$$
\begin{align*}
\tau_{r}= & \tau_{r d}+m_{33}\left[-\frac{m_{11}-m_{22}}{m_{33}}\left(u v-u_{d} v_{d}\right)+\frac{d_{33}}{m_{33}} r_{e}-c_{3} \bar{r}_{e}+\dot{\alpha}_{1}-\right. \\
& \left.\left(\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right)\left(z_{2}+\lambda_{1} r_{d} z_{1}\right)-z_{2 e} z_{1}-\frac{\lambda_{0} m_{11}}{m_{22}} v_{e} u+z_{3 e}\right)\right] \tag{4.66}
\end{align*}
$$

where $c_{3}>0$ results in

$$
\begin{align*}
\dot{V}_{2} \leq & -c(t) V_{1}-c_{2}\left(\left(z_{1 e}-\lambda_{1} z_{2 e} r_{d}\right)\left(z_{2}+\lambda_{1} r_{d} z_{1}\right)-z_{2 e} z_{1}-\right. \\
& \left.\frac{\lambda_{0} m_{11}}{m_{22}} v_{e} u+z_{3 e}\right)^{2}-c_{3} \bar{r}_{e}^{2} \tag{4.67}
\end{align*}
$$

This implies global asymptotic stability of the closed loop system at the origin as long as the control parameters $\lambda_{0}, \lambda_{1}$, and $\lambda_{2}$ are chosen such that (4.61) holds. Note that this condition is feasible only when the reference yaw velocity $r_{d}$ satisfies the following restrictive condition

$$
\begin{equation*}
0<r_{\star} \leq\left|r_{d}(t)\right| \leq r^{\star} \tag{4.68}
\end{equation*}
$$

where $r_{\star}$ and $r^{\star}$ are positive constants. In [19], the result based on the standard backstepping technique also consists of two steps. The first step is to design the surge force $\tau_{u}$. This step is the same as Step 1 mentioned above. The second step is to design the yaw moment $\tau_{r}$. This step is slightly different from Step 2. In this step, a simple controller to stabilize the $\left(z_{3 e}, r_{e}\right)$-subsystem, that is the third and last equations of (4.46), is designed as

$$
\begin{equation*}
\tau_{r}=\tau_{r d}+m_{33}\left[-\frac{m_{11}-m_{22}}{m_{33}}\left(u v-u_{d} v_{d}\right)+\frac{d_{33}}{m_{33}} r_{e}-k_{1} z_{3 e}-k_{2} r_{e}\right] \tag{4.69}
\end{equation*}
$$

With the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$ designed as in (4.62) and (4.70), it is proven in [19] that the tracking errors $\left(z_{1 e}, z_{2 e}, z_{3 e}, u_{e}, v_{e}, r_{e}\right)$ exponentially converge (not exponential stability of the closed loop system) to zero as long as the following restrictive condition on the reference yaw velocity $r_{d}$ holds

$$
\begin{equation*}
\int_{t_{0}}^{t} r_{d}^{2}(\tau) \mathrm{d} \tau \geq \sigma_{r}\left(t-t_{0}\right), \quad \forall 0 \leq t_{0} \leq t<\infty \tag{4.70}
\end{equation*}
$$

where $\sigma_{r}$ is a strictly positive constant.
Remark 4.5. In comparison with the trajectory-tracking results in [71], we see that the control design in [19] is much simpler and gives global solutions. However, the restrictive conditions on the yaw reference velocity cannot be relaxed, see (4.68) and (4.70). Moreover, it is only possible to find the control parameters such that $c(t)$ given in (4.61) is strictly positive for vessels with a small ratio $\frac{m_{22}}{d_{22}}$, i.e., the vessels with large damping in the sway axis.

Remark 4.6. A common restriction of the above results on trajectory-tracking control of underactuated ships is that the reference yaw velocity has to satisfy various kinds of persistently exciting conditions. This implies that the reference trajectory must be curved, and indeed excludes a straight-line reference trajectory, hence, it substantially limits the practical use of the aforementioned control systems. A curious question is why all the above controllers suffer from the must-be-curved reference trajectory restriction. An answer is that the design of the above controllers starts from the chained form (4.46). The reader will find that this book provides various solutions for trajectory-tracking control of underactuated ships without imposing a persistent exciting condition on the yaw reference velocity. As such, we will not use the chained form (4.46) but will project the tracking errors, $x-x_{d}$, $y-y_{d}$, and $\psi-\psi_{d}$, on the body-fixed frame.

Apart from the aforementioned results on trajectory-tracking control of underactuated ships there are a few more results that are worth reviewing. Using sliding mode control, output redefinition and results on tracking of nonlinear nonminimum phase system [73], a path controller for surface ships was proposed in [74]. However, the convergence of the combined output does not guarantee convergence of its components. A continuous time-invariant state feedback controller was developed in [75] to achieve global exponential position tracking under the assumption that the reference surge velocity is always positive. Unfortunately, the orientation of the ship was not controlled. In [76,77], (see also [78]), the authors developed a highgain dynamic feedback control law to achieve global ultimate regulation and tracking of underactuated ships. The dynamics of the closed loop system is increased due to the controller designed to make the state of the transformed system track the auxiliary signals generated by some oscillator. The same approach was extended to the case of adaptive tracking control in [77]. It is worth mentioning that in [47], a time-varying velocity feedback controller was proposed to achieve both stabilization and tracking of unicycle mobile robots at the kinematics level motivated by the work in [18]. However this controller cannot be extended directly to the case of underactuated ships or AUVs due to the nonintegrable second-order constraint. Some related independent work includes [79,80] on local $H_{\infty}$ tracking control and output redefinition, and the trajectory planning approach, see $[81,82]$.

### 4.2.2.3 Path-following

Path-following is here defined as a control problem of forcing an underactuated ship or AUV to follow a specified path at a desired forward speed. Due to the high dependence on the reference model and complicated control laws of the trajectorytracking approach, several researchers have studied the path-following problem, which is more suitable for practical implementation. The problem of path-following for air and underwater vehicles was introduced in [83] where some local results were obtained using linearization techniques. In [84], a feedforward cancelation of simplified vessel dynamics scheme followed by a linear quadratic regulator design was proposed to obtain local results on "track-keeping". A fourth-order ship model
in Serret-Frenet frame was used in [85] to develop a control strategy to track both a straight line and a circumference under constant ocean current disturbance. The ocean-current direction was assumed to be known. A path-following controller was proposed in [86] by using a kinematic model written in polar coordinates, which is inspired by the solution for mobile robots in [60]. However, the controller was designed at the kinematic level with an assumption of constant ocean current and its direction known to be to achieve an adjustable boundedness of the path-following error. Since ocean vessels do not have direct control over velocities, a static mapping implementation might result in an unstable closed loop system due to nonvanishing environmental disturbances. Recently, a path-following controller based on a transformation of the ship kinematics to the Serret-Frenet frame, which was used for mobile robot control [44], on the path was proposed in [87], where an acceleration feedback and linearization of ship dynamics were used. It is worth mentioning that in $[88,89]$, a simple control scheme was proposed to make mobile robots follow a specified path using a polar coordinate transformation. Since underactuated surface ships have fewer numbers of actuators than the to-be-controlled degrees of freedom and are subject to nonintegrable acceleration constraints, their dynamic models are not transformable into a system without drifts. Therefore, the above control scheme is not directly applicable. In [54], a continuous, periodic time-varying feedback control law was proposed to locally exponentially stabilize an underactuated underwater vehicle at the origin. When the hydrodynamic restoring force in roll is large enough, this controller can be used without a roll control torque. However the closed loop system exhibits undesired oscillatory motions.

In [90], a linearization technique with an assumption of reference trajectories of underwater vehicles, which are helices parameterized by the vehicles' linear speed, yaw rate, and path angle, was introduced to develop the so-called time-invariant generalized vehicle error dynamics and kinematics. Various controllers were then designed based on the gain-scheduling technique to yield some local stability result about the trimming trajectories.

### 4.2.2.4 Output Feedback

Output feedback control of an underactuated ocean vessel is here defined as a control problem of forcing the vessel to achieve the aforementioned tasks (stabilization, trajectory-tracking, and path-following) without using measurements of the vessel's velocities for feedback. For ocean vessels, output feedback control usually consists of two stages. The first stage is to design an observer to reconstruct unmeasured states. Using the reconstructed states, a controller is designed to achieve control objectives in the second stage. In the literature, there are two main approaches to designing an observer for ocean vessels.

The first approach is based on the output-injection method applied directly to the vessel's equations of motion. This approach is simple and usually results in a semiglobal observer due to the quadratic terms of the vessel's velocities. Belonging to this approach are the results presented in $[11,14,91-94]$ on output feedback
control of fully actuated ocean vessels or Lagrange systems. In addition, the reader is referred to [93] for an exponential observer and output feedback controller for a special class of multi-degree of freedom Lagrange systems without cross terms of quadratic velocities, and to [14,95-98] for output feedback control of robot manipulators and rigid body without measurements of angular velocities. Another method to design an observer is the use of contraction theory, see for example, $[99,100]$. This method has been applied to Lagrange systems with monotonic velocity terms but without any quadratic velocity terms. Below, we summarize the aforementioned results on an observer design. We will show that a standard observer design cannot be used to obtain a global exponential/asymptotical observer for the ocean vessel system (1.1). Assume that the vessel velocity vector $\boldsymbol{v}$ is not measurable for feedback. We would then design an output injection observer to estimate $\boldsymbol{v}$ as follows:

$$
\begin{align*}
& \dot{\hat{\eta}}=J(\eta) \hat{v}+K_{01}(\eta-\hat{\eta}), \\
& M \dot{\hat{v}}=-C(\hat{v}) \hat{v}-D(\hat{v}) \hat{v}-g(\eta)+\tau+K_{02}(\eta-\hat{\eta}), \tag{4.71}
\end{align*}
$$

where $\hat{\boldsymbol{\eta}}$ and $\hat{\boldsymbol{v}}$ are estimates of $\boldsymbol{\eta}$ and $\boldsymbol{v}$, respectively, and the positive definite symmetric matrices $K_{\mathbf{0 1}} \in \mathbb{R}^{6 \times 6}$ and $\boldsymbol{K}_{\mathbf{0 2}} \in \mathbb{R}^{6 \times 6}$ are the observer gain matrices. It is noted that in some of the aforementioned work, the observer gain matrices $\boldsymbol{K}_{\mathbf{0 1}}$ and $\boldsymbol{K}_{\mathbf{0 2}}$ depend on the measurable state $\boldsymbol{\eta}$. Letting the observer errors be

$$
\begin{align*}
& \tilde{\boldsymbol{\eta}}=\boldsymbol{\eta}-\hat{\boldsymbol{\eta}}, \\
& \tilde{\boldsymbol{v}}=\boldsymbol{v}-\hat{\boldsymbol{v}} \tag{4.72}
\end{align*}
$$

and differentiating (4.72) along the solutions of (1.1) and (4.71) results in

$$
\begin{align*}
& \dot{\tilde{\eta}}=-K_{01} \tilde{\eta}+J(\eta) \tilde{v} \\
& M \dot{\tilde{v}}=-K_{02} \tilde{\eta}-(C(v) v-C(\hat{v}) \hat{v})-(D(v) v-D(\hat{v}) \hat{v}) . \tag{4.73}
\end{align*}
$$

The term $(\boldsymbol{D}(\boldsymbol{v}) \boldsymbol{v}-\boldsymbol{D}(\hat{\boldsymbol{v}}) \hat{\boldsymbol{v}})$ does not cause a problem if the damping matrix $\boldsymbol{D}(\boldsymbol{v})$ is monotonic, i.e., $\left((\boldsymbol{v}-\hat{\boldsymbol{v}})^{T}(\boldsymbol{D}(\boldsymbol{v}) \boldsymbol{v}-\boldsymbol{D}(\hat{\boldsymbol{v}}) \hat{\boldsymbol{v}})\right.$ is nonnegative for all $\boldsymbol{v} \in \mathbb{R}^{6}$ and $\hat{v} \in \mathbb{R}^{6}$. However, we can see a serious problem with (4.73) because of the Coriolis matrix, i.e., $(\boldsymbol{v}-\hat{\boldsymbol{v}})^{\boldsymbol{T}}(\boldsymbol{C}(\boldsymbol{v}) \boldsymbol{v}-\boldsymbol{C}(\hat{\boldsymbol{v}}) \hat{\boldsymbol{v}})$ is not nonnegative for all $\boldsymbol{v} \in \mathbb{R}^{6}$ and $\hat{\boldsymbol{v}} \in \mathbb{R}^{6}$. Therefore, only a local or semiglobal observer can be obtained.

The second approach involves a nontrivial coordinate transformation to transform the vessel's equations of motion to a new set of differential equations that are linear in unmeasured states. Then the output-injection method is used to design an observer. This approach usually results in a global observer if the nontrivial coordinate transformation can be found. Unfortunately, this coordinate transformation depends heavily on a solution of a set of partial differential equations, which in general are hard to solve. The main idea of this approach is to find a coordinate transformation

$$
\begin{equation*}
X=Q(\eta) v \tag{4.74}
\end{equation*}
$$

where $\boldsymbol{Q}(\eta)$ is an invertible. This matrix is to be determined later. Substituting (4.74) into (4.26) results in

$$
\begin{align*}
& \dot{\eta}=J(\eta) Q^{-1}(\eta) X \\
& \dot{X}=\left[\dot{Q}(\eta) v-Q(\eta) M^{-1} C(v) v\right]-Q(\eta) M^{-1} D Q^{-1} X+Q(\eta) M^{-1} \tau \tag{4.75}
\end{align*}
$$

The goal is to determine the matrix $\boldsymbol{Q}(\boldsymbol{\eta})$ such that

$$
\begin{equation*}
\dot{Q}(\eta) v-Q(\eta) M^{-1} C(v) v=0 \tag{4.76}
\end{equation*}
$$

for all $\eta \in \mathbb{R}^{3}$ and $\boldsymbol{v} \in \mathbb{R}^{3}$. With (4.76), we can write (4.75) as

$$
\begin{align*}
& \dot{\eta}=J(\eta) Q^{-1}(\eta) X \\
& \dot{X}=-Q(\eta) M^{-1} D Q^{-1} X+Q(\eta) M^{-1} \tau \tag{4.77}
\end{align*}
$$

It is seen that the transformed system (4.77) is linear in the unmeasured state $\boldsymbol{X}$. This allows us to design an exponential/asymptotical observer to estimate $\boldsymbol{X}$. After that an estimate, $\hat{\boldsymbol{v}}$, of $\boldsymbol{v}$ can be found from (4.74), i.e.,

$$
\begin{equation*}
\hat{v}=Q^{-1}(\eta) \hat{X} \tag{4.78}
\end{equation*}
$$

where $\hat{\boldsymbol{X}}$ denotes an estimate of $\boldsymbol{X}$. It is noted that combining the first equation of (4.26) and (4.76) results in a set of partial differential equations. Finding a solution to this set of partial differential equations is a hard task. A simple application of the above idea gives the results in [101-104] for some single degree of freedom Lagrange systems. It is noted that the method of solving the set of partial differential equations in [101-104] is not applicable for systems of more than one degree of freedom. For more complicated Lagrange systems, it is hard to find a result in this approach. However, the reader is referred to [105] where an output eedback control solution for simultaneous stabilization and tracking control of an underactuated ODIN is given.

Remark 4.7. The main difficulty in designing an observer-based output feedback for surface ships and Lagrange systems in general is because of the Coriolis matrix, which results in cross terms of unmeasured velocities. In addition, the underactuation of surface ships makes the output feedback problem much more challenging. For example, many solutions proposed for robot control, see [14] and references therein, cannot directly be applied. The reader will find that a set of special coordinate transformations is derived in this book to transform the ship dynamics to a system that is linear in unmeasured velocities, and another set of coordinate transformations that makes it possible to design global output feedback control controllers for underactuated ships.

### 4.3 Conclusions

This chapter presented the main control properties of ocean vessels. The literature on the control of underactuated ocean vessels including ships and underwater vehicles was then reviewed. Through this review, several challenging questions were raised. These questions motivate contributions of the coming chapters of this book.

## Part III <br> Control of Underactuated Ships

## Chapter 5 <br> Trajectory-tracking Control of Underactuated Ships

This chapter addresses the problem of trajectory-tracking control of underactuated surface ships. In particular, we present a method to design a controller for underactuated surface ships with only surge force and yaw moment available to globally asymptotically track a reference trajectory generated by a suitable virtual ship. The reference yaw velocity does not have to satisfy a persistently exciting condition as was often required in previous literature. Hence, the reference trajectory is allowed to be a curve including a straight line and a circle. In addition, a new solution to global $K$-exponential tracking as given in previous work is obtained. The control development is based on Lyapunov's direct method and the backstepping technique, and utilizes passive properties of ship dynamics and their interconnected structure.

### 5.1 Control Objective

We consider an underactuated ship with simplified dynamics meaning that all offdiagonal terms of the linear and nonlinear damping matrices, and environmental disturbances induced by waves, wind and ocean currents are ignored. For the reader's convenience the mathematical model of an underactuated ship moving in surge, sway, and yaw directions, see Section 3.4.1.2, is rewritten as

$$
\begin{align*}
\dot{x} & =u \cos (\psi)-v \sin (\psi), \\
\dot{y} & =u \sin (\psi)+v \cos (\psi), \\
\dot{\psi} & =r, \\
\dot{u} & =\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u},  \tag{5.1}\\
\dot{v} & =-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v, \\
\dot{r} & =\frac{m_{11}-m_{22}}{m_{33}} u v-\frac{d_{33}}{m_{33}} r+\frac{1}{m_{33}} \tau_{r},
\end{align*}
$$

where all symbols in (5.1) are defined as in Section 3.4.1.2. The available control inputs are the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$. Since the sway control force is not available, the ship model (5.1) is underactuated. The tracking control problem is to force the underactuated ship to track a reference trajectory generated by a virtual ship as

$$
\begin{align*}
\dot{x}_{d} & =u_{d} \cos \left(\psi_{d}\right)-v_{d} \sin \left(\psi_{d}\right), \\
\dot{y}_{d} & =u_{d} \sin \left(\psi_{d}\right)+v_{d} \cos \left(\psi_{d}\right), \\
\dot{\psi}_{d} & =r_{d} \\
\dot{u}_{d} & =\frac{m_{22}}{m_{11}} v_{d} r_{d}-\frac{d_{11}}{m_{11}} u_{d}+\frac{1}{m_{11}} \tau_{u d},  \tag{5.2}\\
\dot{v}_{d} & =-\frac{m_{11}}{m_{22}} u_{d} r_{d}-\frac{d_{22}}{m_{22}} v_{d}, \\
\dot{r}_{d} & =\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u_{d} v_{d}-\frac{d_{33}}{m_{33}} r_{d}+\frac{1}{m_{33}} \tau_{r d},
\end{align*}
$$

where all variables have similar meanings as in system (5.1) for the virtual reference ship. In this chapter, we propose a novel method to design a controller such that it forces the ship model (5.1) to globally asymptotically track a reference trajectory generated by a virtual ship described by (5.2) under the following assumption.

## Assumption 5.1.

1. The reference signals $u_{d}, r_{d}, \dot{u}_{d}, \ddot{u}_{d}$ and $\dot{r}_{d}$ are bounded.
2. One of the following conditions holds:
(1) There exists a positive constant $\sigma_{r}$ such that, for any pair of $\left(t_{0}, t\right), 0 \leq t_{0} \leq$ $t<\infty$,

$$
\begin{equation*}
\int_{t_{0}}^{t} r_{d}^{2}(\tau) \mathrm{d} \tau \geq \sigma_{r}\left(t-t_{0}\right) \tag{5.3}
\end{equation*}
$$

(2) There exist constants $\sigma_{u}^{\min }>0, \chi_{1} \geq 0$ and $\chi_{2}>0$ such that

$$
\begin{equation*}
\sigma_{u}^{\min } \leq\left|u_{d}(t)\right|,\left|r_{d}(t)\right| \leq \chi_{1} e^{-\chi_{2}\left(t-t_{0}\right)}, \quad \forall t \geq t_{0} \geq 0 \tag{5.4}
\end{equation*}
$$

Remark 5.1. A persistently exciting condition similar to (5.3) is also required in [19, 71, 106], and (5.4) implies that the sign of $u_{d}(t)$ remains unchanged. However, we only require either (5.3) or (5.4) to be satisfied. Assumption 5.1 is quite realistic from a practical point of view. Roughly speaking, either $r_{d}$ or $u_{d}$ is allowed to approach zero. When both $r_{d}$ and $u_{d}$ are equal to zero, the tracking problem becomes one of stabilizing (5.1) and cannot be solved by any time-invariant smooth feedback. It is also noted that the exponential vanishing of $r_{d}$ in (5.4) can be easily replaced by $\int_{0}^{\infty}\left|r_{d}(t)\right| \mathrm{d} t \leq \vartheta, 0 \leq \vartheta<\infty$.

In $[19,71,106]$, a coordinate transformation of the form

$$
\begin{align*}
& z_{1}=\cos (\psi) x+\sin (\psi) y \\
& z_{2}=-\sin (\psi) x+\cos (\psi) y  \tag{5.5}\\
& z_{3}=\psi
\end{align*}
$$

was used for both (5.1) and (5.2), then different controllers were designed for the resulting tracking error system. Consequently, the reference yaw velocity $r_{d}$ must satisfy persistently exciting conditions of various kinds. Hence, the way-point tracking is excluded from consideration. In addition, the physical meaning of the tracking errors is less clear. To overcome the above-mentioned drawbacks, we introduce the position and orientation errors

$$
\begin{align*}
& x-x_{d}, \\
& y-y_{d},  \tag{5.6}\\
& \psi-\psi_{d}
\end{align*}
$$

in a frame attached to the ship body, see Figure 5.1. In this figure, $O_{E} X_{E} Y_{E}$ is the earth-fixed frame, $O_{b} X_{b} Y_{b}$ is the body-fixed frame, and $C G$ is the center of gravity of the ship. With the above definition in mind, we have the error coordinates


Figure 5.1 Definition of tracking errors

$$
\left[\begin{array}{l}
x_{e}  \tag{5.7}\\
y_{e} \\
\psi_{e}
\end{array}\right]=\left[\begin{array}{lll}
\cos (\psi) & \sin (\psi) & 0 \\
-\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{l}
x-x_{d} \\
y-y_{d} \\
\psi-\psi_{d}
\end{array}\right] .
$$

Indeed, convergence of $\left[\begin{array}{lll}x_{e} & y_{e} & \psi_{e}\end{array}\right]^{T}$ to zero implies that of $\left[x-x_{d} y-y_{d} \psi-\psi_{d}\right]^{T}$ since the matrix

$$
\left[\begin{array}{lll}
\cos (\psi) & \sin (\psi) & 0 \\
-\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right]
$$

is nonsingular for all $\psi \in \mathbb{R}$. We also define the velocity tracking errors as

$$
\begin{align*}
u_{e} & =u-u_{d} \\
v_{e} & =v-v_{d}  \tag{5.8}\\
r_{e} & =r-r_{d}
\end{align*}
$$

Then, differentiating both sides of (5.7) and (5.8) along the solutions of (5.1) and (5.2) results in

$$
\begin{align*}
& \dot{x}_{e}=u_{e}-u_{d}\left(\cos \left(\psi_{e}\right)-1\right)-v_{d} \sin \left(\psi_{e}\right)+r_{e} y_{e}+r_{d} y_{e}, \\
& \dot{y}_{e}=v_{e}-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)+u_{d} \sin \left(\psi_{e}\right)-r_{e} x_{e}-r_{d} x_{e}, \\
& \dot{\psi}_{e}=r_{e} \\
& \dot{u}_{e}=\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d},  \tag{5.9}\\
& \dot{v}_{e}=-\frac{m_{11}}{m_{22}} u_{e} r_{d}-\frac{m_{11}}{m_{22}}\left(u_{e}+u_{d}\right) r_{e}-\frac{d_{22}}{m_{22}} v_{e}, \\
& \dot{r}_{e}=\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r+\frac{1}{m_{33}} \tau_{r}-\dot{r}_{d} .
\end{align*}
$$

It is now clear that the tracking problem of underactuated surface ships becomes one of stabilizing (5.9) at the origin since, as has already been mentioned convergence of $x_{e}, y_{e}$, and $\psi_{e}$ to zero implies that of $x-x_{d}, y-y_{d}$, and $\psi-\psi_{d}$. In particular, we will design explicit expressions for $\tau_{u}$ and $\tau_{r}$ such that $\left\|X_{e}(t)\right\| \leq$ $\gamma\left(\left\|X_{e}\left(t_{0}\right)\right\|\right) e^{-\mu\left(t-t_{0}\right)}, t_{0} \geq 0, X_{e}\left(t_{0}\right) \in \mathbb{R}^{6}$, with $X_{e}=\left[x_{e}, y_{e}, \psi_{e}, u_{e}, v_{e}, r_{e}\right]^{T}, \gamma$ being a class- $K$ function, and $\mu$ being a nonnegative continuous decreasing function of $\|\left(X_{e}\left(t_{0}\right) \|\right.$. When Assumption 5.1 satisfies (5.3), it can be shown that $\mu$ is a positive constant, i.e., global $K$-exponential tracking is achieved.

### 5.2 Control Design

In this section, a procedure to design a global asymptotic stabilizer for the tracking error system (5.9) is presented. The triangular structure of (5.9) suggests that we design the actual controls $\tau_{u}$ and $\tau_{r}$ in two stages. First, we design the virtual velocity controls $u_{e}$ and $r_{e}$ to globally asymptotically stabilize $x_{e}, y_{e}, \psi_{e}$, and $v_{e}$ at the origin. Based on the backstepping technique, the controls $\tau_{u}$ and $\tau_{r}$ will then be designed.

By looking at (5.9), we can directly see that $x_{e}$ and $\psi_{e}$ can be stabilized at the origin by $u_{e}$ and $r_{e}$. There are several options to stabilize $y_{e}$ at the origin. We can either use $r_{e}, v_{e}$ or $\psi_{e}$. If $r_{e}$ is directly used, the control design will be extremely complicated since $r_{e}$ enters all of the three equations of (5.9). In addition, $r_{e}$ couples with $x_{e}$ and $y_{e}$ in the first and second equations of (5.9), respectively. This causes difficulties in obtaining a global solution. On the other hand, the use of $v_{e}$ to stabilize $y_{e}$ at the origin will result in an undesired feature of the vehicle control practice, namely the vehicle will slide in the sway direction. Hence we will choose $\psi_{e}$ to stabilize the sway error $y_{e}$ at the origin. This also coincides with the ship control practice as described in [107]: A good helmsman will use the ship course angle rather than use the ship sway velocity to steer the ship. Toward this end, we define the following coordinate

$$
\begin{equation*}
z_{e}=\psi_{e}+\arcsin \left(\frac{k u_{d} y_{e}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}\right) \tag{5.10}
\end{equation*}
$$

where $k$ is a positive constant satisfying

$$
\begin{equation*}
k u_{d}^{\max } \leq k_{*} \tag{5.11}
\end{equation*}
$$

for some positive constant $k_{*}<1$ to be chosen later in the stability analysis. $u_{d}^{\max }$ denotes the maximum value of $\left|u_{d}\right|$. It is seen that (5.10) is well defined and convergence of $x_{e}, z_{e}$, and $y_{e}$ to zero implies that of $\psi_{e}$. Using (5.10) instead of $z_{e}=\psi_{e}+k y_{e}$, we avoid the ship whirling around for large $y_{e}$. This equation is also different from the one in [108] and [18], which resulted in a local tracking result for mobile robots. With (5.10), the ship error dynamics (5.9) are rewritten as

$$
\begin{align*}
& \dot{x}_{e}=u_{e}+u_{d}\left(1-\varpi_{e}\right)+k u_{d} v_{d} y_{e} \varpi_{e 1}+r_{e} y_{e}+r_{d} y_{e}+p_{x}, \\
& \dot{y}_{e}=v_{e}+v_{d}\left(1-\varpi_{e}\right)-k u_{d}^{2} y_{e} \varpi_{e 1}-r_{e} x_{e}-r_{d} x_{e}+p_{y}, \\
& \dot{v}_{e}=-\alpha v_{e}-\beta u_{e} r_{d}-\beta\left(u_{e}+u_{d}\right) r_{e}, \\
& \dot{z}_{e}=\left(1-k u_{d} \varpi_{e 2}^{-1} x_{e}\right) r_{e}+\varpi_{e 2}^{-1}\left(p_{z}+p_{u}\right)-k u_{d} x_{e} y_{e} \varpi_{e 2}^{-1} \varpi_{e 1}^{2} \tilde{u}_{e},  \tag{5.12}\\
& \dot{u}_{e}=\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}}\left(u_{e}+u_{d}\right)+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d}, \\
& \dot{r}_{e}=\frac{m_{11}-m_{22}}{m_{33}} u v-\frac{d_{33}}{m_{33}}\left(r_{e}+r_{d}\right)+\frac{1}{m_{33}} \tau_{r}-\dot{r}_{d},
\end{align*}
$$

where, for notational simplicity, we have defined

$$
\begin{aligned}
\alpha= & \frac{d_{22}}{m_{22}}, \beta=\frac{m_{11}}{m_{22}}, \\
p_{x}= & -\left(\left(\cos \left(z_{e}\right)-1\right) u_{d}+\sin \left(z_{e}\right) v_{d}\right) \varpi_{e}- \\
& \left(\sin \left(z_{e}\right) u_{d}-\left(\cos \left(z_{e}\right)-1\right) v_{d}\right) k u_{d} y_{e} \varpi_{e 1}, \\
p_{y}= & -\left(\left(\cos \left(z_{e}\right)-1\right) v_{d}-\sin \left(z_{e}\right) u_{d}\right) \varpi_{e}- \\
& \left(\sin \left(z_{e}\right) v_{d}+\left(\cos \left(z_{e}\right)-1\right) u_{d}\right) k u_{d} y_{e} \varpi_{e 1},
\end{aligned}
$$

$$
\begin{align*}
p_{z}= & k u_{d} p_{y}-k u_{d} y_{e}\left(x_{e} p_{x}+y_{e} p_{y}\right) \varpi_{e 1}^{2}, \\
p_{u}= & k u_{d} v_{e}+k u_{d} v_{d}\left(1-\varpi_{e}\right)-k^{2} u_{d}^{3} y_{e} \varpi_{e 1}-k u_{d} r_{d} x_{e}+k \dot{u}_{d} y_{e}- \\
& k u_{d} \varpi_{e 1}^{2} y_{e}\left(x_{e} u_{e}^{d}+x_{e} u_{d}\left(1-\varpi_{e}\right)+k u_{d} v_{d} \varpi_{e 1} x_{e} y_{e}+y_{e} v_{e}+\right. \\
& \left.y_{e} v_{d}\left(1-\varpi_{e}\right)-k u_{d}^{2} \varpi_{e 1} y_{e}^{2}\right), \\
\varpi_{e 1}= & \sqrt{\frac{1}{1+x_{e}^{2}+y_{e}^{2}}}, \\
\varpi_{e 2}= & \sqrt{1+x_{e}^{2}+\left(1-\left(k u_{d}\right)^{2}\right) y_{e}^{2}}, \\
\varpi_{e}= & \varpi_{e 1} \varpi_{e 2}, \tag{5.13}
\end{align*}
$$

with $u_{e}^{d}$ and $\tilde{u}_{e}$ being defined in (5.15). Before designing the control laws, $\tau_{u}$ and $\tau_{r}$, we note from (5.11) that

$$
\begin{equation*}
0<1-k_{*} \leq 1-\frac{x_{e}}{\varpi_{e 2}} k u_{d} . \tag{5.14}
\end{equation*}
$$

This allows us to design global control laws $\tau_{u}$ and $\tau_{r}$ to asymptotically stabilize (5.12) at the origin. The control design consists of two steps as follows.

## Step 1

We define the following virtual control errors

$$
\begin{align*}
& \tilde{u}_{e}=u_{e}-u_{e}^{d} \\
& \tilde{r}_{e}=r_{e}-r_{e}^{d} \tag{5.15}
\end{align*}
$$

where $u_{e}^{d}$ and $r_{e}^{d}$ are the virtual velocity controls of $u_{e}$ and $r_{e}$, respectively.
The virtual surge and yaw velocity controls are chosen as

$$
\begin{align*}
& u_{e}^{d}=-k_{1} x_{e}+k_{2} r_{d} y_{e} \\
& r_{e}^{d}=r_{1 e}^{d}+r_{2 e}^{d} \tag{5.16}
\end{align*}
$$

where

$$
\begin{gather*}
r_{1 e}^{d}=-\frac{1}{\varpi_{e 2}-k u_{d} x_{e}} p_{u}  \tag{5.17}\\
r_{2 e}^{d}=-\frac{1}{\varpi_{e 2}-k u_{d} x_{e}}\left(k_{3} \varpi_{e 2} z_{e}+p_{z}\right), \tag{5.18}
\end{gather*}
$$

and $k_{i}, i=1,2,3$, are positive constants to be selected later. We have written $r_{e}^{d}$ as a sum of $r_{1 e}^{d}$ and $r_{2 e}^{d}$ to simplify notation in the stability analysis later.

Remark 5.2. Unlike the standard application of the backstepping technique, in order to reduce complexity of the controller expressions, we have chosen a simple virtual control law $u_{e}^{d}$ without canceling the known terms. The first term in $u_{e}^{d}$ is used to
stabilize the $x_{e}$-dynamics while the second term plays the role of stabilizing the $y_{e}$-dynamics when (5.3) holds. From (5.17) and (5.13), we observe that $r_{1 e}^{d}$ is Lipschitz in $\left(x_{e}, y_{e}, v_{e}\right)$ and with (5.18), $r_{2 e}^{d}$ exponentially vanishes when $z_{e}$ does. This observation plays a crucial role in the stability analysis of the closed loop system.

## Step 2

Differentiating (5.15) along the solutions of (5.12) and (5.16) yields

$$
\begin{align*}
& \dot{\tilde{u}}_{e}=\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d}-\dot{u}_{e}^{d}, \\
& \dot{\tilde{r}}_{e}=\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r+\frac{1}{m_{33}} \tau_{r}-\dot{r}_{d}-\dot{r}_{e}^{d}, \tag{5.19}
\end{align*}
$$

where $\dot{u}_{e}^{d}$ and $\dot{r}_{e}^{d}$ are the first time derivatives of $u_{e}^{d}$ and $r_{e}^{d}$ along the solutions of (5.12).

From (5.19) we choose the actual controls $\tau_{u}$ and $\tau_{r}$ without canceling the useful damping terms as

$$
\begin{align*}
& \tau_{u}=m_{11}\left(-\delta_{1} \tilde{u}_{e}-\frac{m_{22}}{m_{11}} v r+\frac{d_{11}}{m_{11}}\left(u_{e}^{d}+u_{d}\right)+\dot{u}_{d}+\dot{u}_{e}^{d}\right) \\
& \tau_{r}=m_{33}\left(-\delta_{2} \tilde{r}_{e}-\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v+\frac{d_{33}}{m_{33}}\left(r_{e}^{d}+r_{d}\right)+\dot{r}_{d}+\dot{r}_{e}^{d}\right) \tag{5.20}
\end{align*}
$$

where $\delta_{1}$ and $\delta_{2}$ are positive constants to be chosen later. Substituting (5.16) and (5.20) into (5.12) yields the closed loop system

$$
\begin{align*}
\dot{X}_{1 e} & =f_{1}\left(t, X_{1 e}\right)+g_{1}\left(t, X_{1 e}, X_{2 e}\right)  \tag{5.21}\\
\dot{X}_{2 e} & =f_{2}\left(t, X_{1 e}, X_{2 e}\right)
\end{align*}
$$

where $X_{1 e}=\left[\begin{array}{lll}x_{e} & y_{e} & v_{e}\end{array}\right]^{T}, X_{2 e}=\left[\begin{array}{lll}z_{e} & \tilde{u}_{e} & \tilde{r}_{e}\end{array}\right]^{T}$,

$$
\begin{align*}
& f_{1}\left(t, X_{1 e}\right)=\left[\begin{array}{l}
-k_{1} x_{e}+k_{2} r_{d} y_{e}+u_{d}\left(1-\varpi_{e}\right)+\left(k u_{d} v_{d} \varpi_{e 1}+r_{1 e}^{d}+r_{d}\right) y_{e} \\
v_{e}+v_{d}\left(1-\varpi_{e}\right)-k u_{d}^{2} \varpi_{e 1} y_{e}-r_{1 e}^{d} x_{e}-r_{d} x_{e} \\
-\alpha v_{e}-\beta u_{e}^{d} r_{d}-\beta\left(u_{e}^{d}+u_{d}\right) r_{1 e}^{d}
\end{array}\right],  \tag{5.22}\\
& g_{1}\left(t, X_{1 e}, X_{2 e}\right)=\left[\begin{array}{l}
r_{2 e}^{d} y_{e}+\tilde{r}_{e} y_{e}+p_{x}+\tilde{u}_{e} \\
-r_{2 e}^{d} x_{e}-\tilde{r}_{e} x_{e}+p_{y} \\
-\beta\left(\tilde{u}_{e} r_{d}+\tilde{u}_{e}\left(r_{1 e}^{d}+r_{2 e}^{d}+\tilde{r}_{e}\right)+\left(u_{e}^{d}+u_{d}\right)\left(r_{2 e}^{d}+\tilde{r}_{e}\right)\right)
\end{array}\right], \\
& f_{2}\left(t, X_{1 e}, X_{2 e}\right)=\left[\begin{array}{l}
-k_{3} z_{e}+\left(1-k u_{d} x_{e} / \varpi_{e 2}\right) \tilde{r}_{e}-k u_{d} x_{e} y_{e} \varpi_{e 1}^{2} \tilde{u}_{e} / \varpi_{e 2} \\
-\left(\delta_{1}+d_{11} / m_{11}\right) \tilde{u}_{e} \\
-\left(\delta_{2}+d_{33} / m_{33}\right) \tilde{r}_{e}
\end{array}\right] . \tag{5.24}
\end{align*}
$$

The time dependence in $f_{1}\left(t, X_{1 e}\right), g_{1}\left(t, X_{1 e}, X_{2 e}\right)$ and $f_{2}\left(t, X_{1 e}, X_{2 e}\right)$ results from the time-varying reference velocities. We now state our main result, the proof of which is given in the next section.

Theorem 5.1. Assume that the reference signals $\left(x_{d}, y_{d}, \psi_{d}, v_{d}\right)$ are generated by the virtual ship model (5.2) and that the reference velocities $\left(u_{d}, r_{d}\right)$ satisfy Assumption 5.1. If the state feedback control law (5.20) is applied to the ship system (5.1) then the tracking errors $x(t)-x_{d}(t), y(t)-y_{d}(t), \psi(t)-\psi_{d}(t)$, and $v(t)-v_{d}(t)$ converge to zero asymptotically from arbitrary initial values with an appropriate choice of the design constants $k_{i}, 1 \leq i \leq 3$ and $\delta_{j}, j=1,2$, i.e., the closed loop system (5.21) is GAS at the origin. Furthermore if Assumption 5.1 holds with (5.3), the closed loop system (5.21) is globally $K$-exponentially stable at the origin.

### 5.3 Stability Analysis

To prove the global asymptotic stability at the origin of the closed loop system (5.21), we note that (5.21) consists of two subsystems, $X_{1 e}$ and $X_{2 e}$, in an interconnected structure. If $X_{1 e}$ does not appear in the second equation of (5.21), we can use the stability analysis approaches for a cascaded system in [17]. One might claim that if the system

$$
\begin{align*}
\dot{X}_{1 e} & =f_{1}\left(t, X_{1 e}\right) \\
\dot{X}_{2 e} & =f_{2}\left(t, X_{1 e}, X_{2 e}\right) \tag{5.25}
\end{align*}
$$

is uniformly globally asymptotically stable (UGAS), and the term $g_{1}\left(t, X_{1 e}, X_{2 e}\right)$ somehow is of linear growth in $X_{1 e}$, then the closed loop system (5.21) is UGAS. However, we are interested in proving stronger stability properties of (5.21). Hence we present the following lemma.

Lemma 5.1. Consider the following nonlinear system:

$$
\begin{equation*}
\dot{x}=f(t, x)+g(t, x, \xi(t)) \tag{5.26}
\end{equation*}
$$

where $x \in \mathbb{R}^{n}, \xi(t) \in \mathbb{R}^{m}, f(t, x)$ is piecewise continuous in t and locally Lipschitz in $x$. Assume the following:

C1. There exists a proper function $V(t, x)$ satisfying:

$$
\begin{align*}
& c_{1}\|x\|^{2} \leq V(t, x) \leq c_{2}\|x\|^{2}, \\
& \left\|\frac{\partial V}{\partial x}(t, x)\right\| \leq c_{3}\|x\| \\
& \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x) \leq-l_{1}(t)\|x\|^{2}-l_{2}(t) \frac{\|x\|^{2}}{\sqrt{1+\|x\|^{2}}} \tag{5.27}
\end{align*}
$$

where $c_{i}>0, i=1,2,3$, and the pair $l_{1}(t)$ and $l_{2}(t)$ satisfy one of the conditions:

1. $\int_{t_{0}}^{t} l_{1}(\tau) \mathrm{d} \tau \geq \eta_{1}\left(t-t_{0}\right), \eta_{1}>0$ and $l_{2}(t) \geq 0, \forall t \geq t_{0} \geq 0$,
2. $0<\eta_{2} \leq l_{2}(t)<\infty$, and $\left|l_{1}(t)\right| \leq \rho_{1} e^{-\rho_{2}\left(t-t_{0}\right)}, \rho_{1} \geq 0, \rho_{2}>0, \forall t \geq t_{0} \geq 0$.

C2. $g(t, x, \xi(t))$ grows linearly in $x$ :

$$
\|g(t, x, \xi(t))\| \leq\left(\lambda_{1}+\lambda_{2}\|x\|\right)\|\xi(t)\|, \lambda_{i} \geq 0, i=1,2
$$

C3. $\xi(t)$ vanishes exponentially: $\|\xi(t)\| \leq \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}, \sigma_{0}>0, \forall t \geq$ $t_{0} \geq 0$, with $\gamma_{0}$ being a class-K function.

Then, the solution $x(t)$ of (5.26) is GAS in the sense that

$$
\begin{equation*}
\|x(t)\| \leq \gamma\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right) e^{-\left(\sigma_{1}+\sigma_{2}\right)\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0, \sigma_{1} \geq 0 \tag{5.28}
\end{equation*}
$$

where $\gamma$ is a class- $K$ function, $\sigma_{2}$ is a nonnegative continuous decreasing function of $\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|$, and $\sigma_{1}>0$ if condition C1 holds with 1. It can be understood that $\xi(t)$ is generated from a globally $K$-exponentially stable dynamic system.

Proof. We first consider the case of condition C1 satisfying 1. From conditions C1, C2, and C3, we have

$$
\begin{equation*}
\dot{V} \leq-\frac{l_{1}(t)}{c_{2}} V+c_{3}\left(\frac{\lambda_{1}}{\sqrt{c_{1}}} \sqrt{V}+\frac{\lambda_{2}}{c_{1}} V\right) \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)} \tag{5.29}
\end{equation*}
$$

By defining $\kappa(t)=\sqrt{V(t)}$, we can rewrite (5.29) as

$$
\begin{align*}
\dot{\kappa}(t) \leq & -\left(\frac{l_{1}(t)}{2 c_{2}}-\frac{c_{3} \lambda_{2}}{2 c_{1}} \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}\right) \kappa(t)+\frac{c_{3} \lambda_{1}}{2 \sqrt{c_{1}}} \times \\
& \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)} \tag{5.30}
\end{align*}
$$

Solving the differential inequality (5.30) readily yields (5.28). It is of interest to note that in this case, the system (5.26) is GES due to independence of $\sigma_{1}$ in (5.28) on $\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|$. We now move to the case of condition C1 satisfying 2. We first show that $x(t)$ is bounded for all $t \geq t_{0} \geq 0$. From conditions $\mathrm{C} 1, \mathrm{C} 2$, and C 3 , we have

$$
\begin{align*}
\dot{V} \leq & \rho_{1} e^{-\rho_{2}\left(t-t_{0}\right)}\|x\|^{2}-l_{2}(t) \frac{\|x\|^{2}}{\sqrt{1+\|x\|^{2}}}+ \\
& c_{3}\|x\|\left(\lambda_{1}+\lambda_{2}\|x\|\right) \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}  \tag{5.31}\\
\leq & c_{3}\left(\frac{\lambda_{1}}{\sqrt{c_{1}}} \sqrt{V}+\frac{\lambda_{2}}{c_{1}} V\right) \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+\frac{\rho_{1}}{c_{1}} e^{-\rho_{2}\left(t-t_{0}\right)} V .
\end{align*}
$$

From (5.31), it is not hard to show that the solution $x(t)$ is bounded by $\|x(t)\| \leq$ $b\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right)$ with $b$ being a class- $K$ function. By utilizing this bound and
condition C 1 , we can rewrite the first line in (5.31) as

$$
\begin{align*}
\dot{V} \leq & -\frac{\eta_{2}}{c_{2} \sqrt{1+b^{2}}} V+c_{3}\left(\frac{\lambda_{1}}{\sqrt{c_{1}}} \sqrt{V}+\frac{\lambda_{2}}{c_{1}} V\right) \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+\frac{\rho_{1}}{c_{1}} \times \\
& e^{-\rho_{2}\left(t-t_{0}\right)} V . \tag{5.32}
\end{align*}
$$

By defining $\kappa(t)=\sqrt{V(t)}$, we rewrite (5.32) as follows:

$$
\begin{align*}
\dot{\kappa}(t) \leq & \left(\frac{\eta_{2}}{2 c_{2} \sqrt{1+b^{2}}}-\frac{c_{3} \lambda_{2}}{2 c_{1}} \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}-\frac{\rho_{1}}{2 c_{1}} e^{-\rho_{2}\left(t-t_{0}\right)}\right) \times \\
& \kappa(t)+\frac{c_{3} \lambda_{1}}{2 \sqrt{c_{1}}} \gamma_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)} . \tag{5.33}
\end{align*}
$$

Again, solving the differential inequality (5.33) readily yields (5.28).
We now view our closed loop system (5.21) as a form of the system studied in Lemma 5.1 with $X_{1 e}(t)$ as $x(t)$, and $X_{2 e}(t)$ as $\xi(t)$. In the following, we will show that the closed loop system (5.21) satisfies all the conditions C1, C2, and C3.

Verifying Condition C1. To verify condition C1 of Lemma 5.1 for the system

$$
\begin{equation*}
\dot{X}_{1 e}=f_{1}\left(t, X_{1 e}\right), \tag{5.34}
\end{equation*}
$$

the first time derivative of the following Lyapunov function

$$
\begin{equation*}
V_{1}=\frac{1}{2} x_{e}^{2}+\frac{1}{2} y_{e}^{2}+\frac{k_{4}}{2}\left(v_{e}+k_{5} y_{e}\right)^{2} \tag{5.35}
\end{equation*}
$$

where $k_{4}$ and $k_{5}$ are positive constants, along the solutions of (5.34) satisfies

$$
\begin{equation*}
\dot{V}_{1}=-k_{1} x_{e}^{2}-\beta k_{2} k_{4} k_{5} r_{d}^{2} y_{e}^{2}-k u_{d}^{2}\left(1+k_{4} k_{5}^{2}\right) \varpi_{e 1} y_{e}^{2}-k_{4}\left(\alpha-k_{5}\right) v_{e}^{2}+M, \tag{5.36}
\end{equation*}
$$

where

$$
\begin{align*}
M= & k_{2} r_{d} x_{e} y_{e}+\left(1-\varpi_{e}\right)\left(u_{d} x_{e}+k_{4} k_{5} v_{d}\left(v_{e}+k_{5} y_{e}\right)+v_{d} y_{e}\right)+k u_{d} v_{d} \varpi_{e 1} \times \\
& x_{e} y_{e}-k_{2} k_{4} \beta r_{d}^{2} y_{e} v_{e}-k_{4} k_{5} k u_{d}^{2} \varpi_{e 1} y_{e} v_{e}-\beta k_{4}\left(v_{e}+k_{5} y_{e}\right)\left(k_{2} r_{d} y_{e}+u_{d}\right) r_{1 e}^{d}, \tag{5.37}
\end{align*}
$$

and we have chosen

$$
\begin{equation*}
\beta k_{1}-k_{5}=0,1-k_{4}\left(\alpha-k_{5}\right) k_{5}=0 \tag{5.38}
\end{equation*}
$$

to cancel some common terms. In addition, we note that

$$
\begin{equation*}
1-\varpi_{e}=\frac{\left(k u_{d}\right)^{2} y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}\left(\sqrt{1+x_{e}^{2}+y_{e}^{2}}+\sqrt{1+x_{e}^{2}+\left(1-\left(k u_{d}\right)^{2}\right) y_{e}^{2}}\right)} \tag{5.39}
\end{equation*}
$$

We now write $M=M_{1}+\sum_{i=1}^{4} N_{i}$, where

$$
\begin{align*}
& M_{1}=k_{2} r_{d} x_{e} y_{e}+\left(1-\varpi_{e}\right)\left(u_{d} x_{e}+k_{4} k_{5} v_{d}\left(v_{e}+k_{5} y_{e}\right)+v_{d} y_{e}\right)+ \\
& \quad k u_{d} v_{d} \varpi_{e 1} x_{e} y_{e}-k_{2} k_{4} \beta r_{d}^{2} y_{e} v_{e}-k_{4} k_{5} k u_{d}^{2} \varpi_{e 1} y_{e} v_{e} \\
& N_{1}=-\beta k_{2} k_{4} r_{d} y_{e} v_{e} r_{1 e}^{d}, N_{2}=-\beta k_{4} u_{d} v_{e} r_{1 e}^{d} \\
& N_{3}=-\beta k_{2} k_{4} k_{5} r_{d} y_{e}^{2} r_{1 e}^{d}, N_{4}=-\beta k_{4} k_{5} r_{d} y_{e} r_{1 e}^{d} \tag{5.40}
\end{align*}
$$

Applying (5.39) and the completing square to (5.40) yields

$$
\begin{gather*}
M_{1} \leq \frac{1}{4 \varepsilon_{1}}\left(k_{2}+k u_{d}^{2}+k v_{d}^{2}\right) x_{e}^{2}+\left(k_{2} \varepsilon_{1}+\beta k_{2} k_{4} \varepsilon_{2}\right) r_{d}^{2} y_{e}^{2}+ \\
\left(2 k \varepsilon_{1}+2 k k_{4} k_{5} \varepsilon_{3}\right) \frac{u_{d}^{2} y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}+k\left(1+k_{4} k_{5}^{2}\right)\left|u_{d} v_{d}\right| \times  \tag{5.41}\\
\frac{y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}+\left(\frac{k k_{4} k_{5} v_{d}^{2}}{4 \varepsilon_{3}}+\frac{\beta k_{2} k_{4} r_{d}^{2}}{4 \varepsilon_{2}}+\frac{k k_{4} k_{5} u_{d}^{2}}{4 \varepsilon_{3}}\right) v_{e}^{2}, \\
N_{1} \leq \frac{\beta k_{2} k_{4} k}{1-k_{*}}\left(\frac{\left|u_{d} r_{d}\right|}{\sqrt{1-k_{*}^{2}}}+\left|u_{d} r_{d}\right|+\frac{1}{4 \varepsilon_{2}}\left(\left(u_{d}^{2}\left(v_{d}^{2}+r_{d}^{2}\right)+2 k\left|u_{d}^{3}\right|+\right.\right.\right. \\
\left.\left.\left.\frac{\dot{u}_{d}^{2}}{\sqrt{1-k_{*}^{2}}}+\left(k_{1}+k_{2} r_{d}^{2}\right) u_{d}^{2}+k_{*}^{2} u_{d}^{4}+k_{*}^{2} v_{d}^{2}+\left|u_{d} v_{d}\right|\right)\right)\right) v_{e}^{2}+\frac{\beta k_{2} k_{4}}{1-k_{*}}  \tag{5.42}\\
\times\left(k \varepsilon_{2}\left(3+k_{1}+k_{2}+2 k\left|u_{d}^{3}\right|+\frac{1}{\sqrt{1-k_{*}^{2}}}+\left|u_{d} v_{d}\right|\right)+k_{*}^{2} \varepsilon_{2}\right) r_{d}^{2} y_{e}^{2}, \\
N_{2} \leq \frac{\beta k_{4} k}{1-k_{*}}\left(2 u_{d}^{2}+u_{d}^{2}\left|r_{d}\right| \varepsilon_{1}+\frac{1}{4 \varepsilon_{3}}\left(k_{*}^{2} u_{d}^{2} v_{d}^{2}+2 k u_{d}^{4}+\dot{u}_{d}^{2}+k_{1} u_{d}^{2}+\right.\right. \\
\left.\left.\quad k_{2} u_{d}^{2} r_{d}^{2}+k_{*}^{2} u_{d}^{4}+k u_{d}^{4} v_{d}^{2}+k_{*}^{2} u_{d}^{2} v_{d}^{2}\right)\right) v_{e}^{2}+\frac{\beta k k_{4} u_{d}^{2}\left|r_{d}\right|}{\left(1-k_{*}\right) 4 \varepsilon_{1}} x_{e}^{2}+  \tag{5.43}\\
\frac{\beta k_{4} \varepsilon_{3} k}{1-k_{*}}\left(3 k_{*}^{2}+2 k u_{d}^{2}+\frac{1}{\left.\sqrt{1-k_{*}^{2}}+k_{1}+k_{2}+k\right) \frac{u_{d}^{2} y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}}\right. \\
N_{3} \leq \frac{\beta k_{2} k_{4} k_{5} k}{1-k_{*}}\left(\frac{2 \varepsilon_{2}}{\left.\sqrt{1-k_{*}^{2}}+\left|u_{d}\right|+k_{2}\left|u_{d}\right|\right)} r_{d}^{2} y_{e}^{2}+\frac{\beta k_{2} k_{4} k_{5} k}{1-k_{*}} \times\right. \\
\left(\frac{\left|\dot{u}_{d}\right|}{\sqrt{1-k_{*}^{2}}}+k_{1}\left|u_{d}\right|\right)\left|r_{d}\right| y_{e}^{2}+\frac{\beta k_{4} k_{5} k_{2} k}{1-k_{*}}\left(\frac{2 k\left|u_{d} r_{d}\right|}{\sqrt{1-k_{*}^{2}}+k^{2}\left|r_{d} v_{d}\right|+}\right.  \tag{5.44}\\
\left.k_{*}^{2}+k\left|r_{d} v_{d}\right|+k^{2}\left|u_{d} r_{d} v_{d}\right|\right) \frac{u_{d}^{2} y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}+\frac{\beta k_{2} k_{4} k_{5} k u_{d}^{2} v_{e}^{2}}{\left(1-k_{*}\right) 2 \varepsilon_{2} \sqrt{1-k_{*}^{2}}}
\end{gather*}
$$

$$
\begin{align*}
& N_{4} \leq \frac{\beta k_{4} k_{5}}{1-k_{*}} \frac{k u_{d}^{2}}{2 \varepsilon_{3}} v_{e}^{2}+\frac{\beta k_{4} k_{5} k u_{d}^{2}}{1-k_{*}} \frac{1}{4 \varepsilon_{1}} x_{e}^{2}+\frac{\beta k_{4} k_{5} k u_{d}^{2} \varepsilon_{1}}{1-k_{*}} r_{d}^{2} y_{e}^{2}+ \\
& \frac{\beta k_{4} k_{5} k\left|\dot{u}_{d}\right|}{\left(1-k_{*}\right) \sqrt{1-k_{*}^{2}}} \frac{\left|u_{d}\right| y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}+\frac{\beta k_{4} k_{5} k}{1-k_{*}}\left(\frac{\varepsilon_{3}}{\sqrt{1-k_{*}^{2}}}+2 k_{*}^{2}\left|v_{d}\right|+\right.  \tag{5.45}\\
& \left.\quad 2 k u_{d}^{2}+k_{1}+k_{2}\left|r_{d}\right|+k_{*}^{2}\left|u_{d}\right|+k\left|u_{d} v_{d}\right|+1\right) \frac{u_{d}^{2} y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}} .
\end{align*}
$$

Substituting (5.37) with (5.41)-(5.45) into (5.36) yields

$$
\begin{equation*}
\dot{V}_{1} \leq-\mu_{1}(t) x_{e}^{2}-\mu_{21}(t) y_{e}^{2}-\frac{\mu_{22}(t) y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}-\mu_{3}(t) v_{e}^{2} \tag{5.46}
\end{equation*}
$$

where

$$
\begin{aligned}
& \mu_{1}(t)= k_{1}-\frac{1}{4 \varepsilon_{1}}\left(k_{2}+k u_{d}^{2}+k v_{d}^{2}+\frac{\beta k_{4} k u_{d}^{2}\left|r_{d}\right|}{1-k_{*}}+\frac{\beta k_{4} k_{5} k u_{d}^{2}}{1-k_{*}}\right), \\
& \mu_{21}(t)= \mu_{211}(t) r_{d}^{2}-\mu_{212}(t)\left|r_{d}\right|, \\
& \mu_{211}(t)= {\left[\beta k_{2} k_{4}\left(k_{5}-\varepsilon_{2}\right)-k_{2} \varepsilon_{1}-\frac{\beta k_{2} k_{4}}{1-k_{*}}\left(k \varepsilon _ { 2 } \left(3+k_{1}+k_{2}+\right.\right.\right.} \\
&\left.\left.2 k\left|u_{d}^{3}\right|+\frac{1}{\sqrt{1-k_{*}^{2}}}+\left|u_{d} v_{d}\right|\right)+k_{*}^{2} \varepsilon_{2}\right)-\frac{\beta k_{2} k_{4} k_{5}}{1-k_{*}} \times \\
&\left.\left(\frac{2 k \varepsilon_{2}}{\sqrt{1-k_{*}^{2}}}+k\left|u_{d}\right|+k k_{2}\left|u_{d}\right|\right)-\frac{\beta k_{4} k_{5} k u_{d}^{2} \varepsilon_{1}}{1-k_{*}}\right], \\
& \mu_{212}(t)= \frac{\beta k_{2} k_{4} k_{5}}{1-k_{*}}\left(\frac{k\left|\dot{u}_{d}\right|}{\sqrt{1-k_{*}^{2}}}+k k_{1}\left|u_{d}\right|\right), \\
& \mu_{221}(t)= \mu_{221}(t) u_{d}^{2}-\mu_{222}(t)\left|u_{d}\right|, \\
& k\left(1+k_{4} k_{5}^{2}\right)-2 k_{1}-2 k_{4} k_{5} k \varepsilon_{3}-\frac{\beta k_{4} \varepsilon_{3} k}{1-k_{*}} \times \\
&\left(3 k_{*}^{2}+2 k u_{d}^{2}+\frac{1}{\sqrt{1-k_{*}^{2}}}+k_{1}+k_{2}+k\right)-\frac{\beta k_{4} k_{5} k_{2} k}{1-k_{*}} \times \\
&\left(\frac{2 k\left|u_{d} r_{d}\right|}{\sqrt{1-k_{*}^{2}}}+k^{2}\left|r_{d} v_{d}\right|+k_{*}^{2}+k\left|r_{d} v_{d}\right|+k^{2}\left|u_{d} r_{d} v_{d}\right|\right)- \\
& \frac{\beta k_{4} k_{5} k}{1-k_{*}}\left(\frac{\varepsilon_{3}}{\sqrt{1-k_{*}^{2}}}+2 k_{*}^{2}\left|v_{d}\right|+2 k u_{d}^{2}+k_{1}+k_{2}\left|r_{d}\right|+\right. \\
&\left.k_{*}^{2}\left|u_{d}\right|+k\left|u_{d} v_{d}\right|+1\right), \\
& \mu_{222}(t)= k\left(1+k_{4} k_{5}^{2}\right)\left|v_{d}\right|+\frac{\beta k_{4} k_{5} k\left|\dot{u}_{d}\right|}{\left(1-k_{*}\right) \sqrt{1-k_{*}^{2}}},
\end{aligned}
$$

$$
\begin{align*}
\mu_{3}(t)= & k_{4}\left(\alpha-k_{5}\right)-\frac{k k_{4} k_{5} v_{d}^{2}}{4 \varepsilon_{3}}-\frac{\beta k_{2} k_{4} r_{d}^{2}}{4 \varepsilon_{2}}-\frac{k k_{4} k_{5} u_{d}^{2}}{4 \varepsilon_{3}}- \\
& \frac{\beta k_{2} k_{4} k_{5} k u_{d}^{2}}{\left(1-k_{*}\right) 2 \varepsilon_{2} \sqrt{1-k_{*}^{2}}}-\frac{\beta k_{4} k_{5}}{1-k_{*}} \frac{k u_{d}^{2}}{2 \varepsilon_{3}}-\frac{\beta k_{2} k_{4} k}{1-k_{*}} \times \\
& \left(\frac{\left|u_{d} r_{d}\right|}{\sqrt{1-k_{*}^{2}}}+\left|u_{d} r_{d}\right|+\frac{1}{4 \varepsilon_{2}}\left(u_{d}^{2}\left(v_{d}^{2}+r_{d}^{2}\right)+2 k\left|u_{d}^{3}\right|+\right.\right. \\
& \left.\left.\frac{\dot{u}_{d}^{2}}{\sqrt{1-k_{*}^{2}}}+\left(k_{1}+k_{2} r_{d}^{2}\right) u_{d}^{2}+k_{*}^{2} u_{d}^{4}+k_{*}^{2} v_{d}^{2}+\left|u_{d} v_{d}\right|\right)\right)- \\
& \frac{\beta k_{4} k}{1-k_{*}}\left(2 u_{d}^{2}+u_{d}^{2}\left|r_{d}\right| \varepsilon_{1} \frac{1}{4 \varepsilon_{3}}\left(k_{*}^{2} u_{d}^{2} v_{d}^{2}+2 k u_{d}^{4}+\dot{u}_{d}^{2}+\right.\right. \\
& \left.\left.k_{1} u_{d}^{2}+k_{2} u_{d}^{2} r_{d}^{2}+k_{*}^{2} u_{d}^{4}+k u_{d}^{4} v_{d}^{2}+k_{*}^{2} u_{d}^{2} v_{d}^{2}\right)\right) \tag{5.47}
\end{align*}
$$

for some positive constants $\varepsilon_{i}, i=1,2,3$. The time dependence of $\mu_{1}, \mu_{21}, \mu_{22}$, and $\mu_{3}$ is due to the time-varying reference velocities. Hence condition C1 is satisfied if there exist positive constants $k_{4}$ and $k_{5}$, and the design constants $k, k_{1}$, and $k_{2}$ such that the following conditions are satisfied:

1. $k u_{d}^{\max } \leq k_{*}<1, \beta k_{1}-k_{5}=0,1-k_{4}\left(\alpha-k_{5}\right) k_{5}=0, \mu_{1}(t) \geq \mu_{1}^{*}, \mu_{3}(t) \geq$ $\mu_{3}^{*}, \forall t \geq 0$,
2. $\mu_{221}(t) \geq 0, \mu_{211}^{\sigma_{r}}-\mu_{212}^{m} r_{d}^{\max }-\mu_{222}^{m} u_{d}^{\max } \geq \mu_{21}^{*}$, if (5.3) holds, see Assumption 5.1,
3. $\mu_{22}(t) \geq \mu_{22}^{*}$, if (5.4) holds, see Assumption 5.1,
where $\mu_{1}^{*}, \mu_{21}^{*}, \mu_{22}^{*}$, and $\mu_{3}^{*}$ are some positive constants, $\mu_{212}^{m}$ and $\mu_{222}^{m}$ are the values of $\mu_{211}(t)$ and $\mu_{212}(t)$ with the maximum values of $\left|u_{d}(t)\right|,\left|\dot{u}_{d}(t)\right|$ and $\left|r_{d}(t)\right|$ substituted in, and $r_{d}^{\text {max }}$ denotes the maximum value of $\left|r_{d}\right|$. We now need to show that the above constants always exist. From (5.38), we have

$$
\begin{equation*}
k_{1}=\frac{k_{5}}{\beta}, k_{5}=\frac{\alpha}{2}-\sqrt{\frac{\alpha^{2}}{4}-\frac{1}{k_{4}}} . \tag{5.48}
\end{equation*}
$$

It is noted that we take $k_{5}=\frac{\alpha}{2}-\sqrt{\frac{\alpha^{2}}{4}-\frac{1}{k_{4}}}$ instead of $k_{5}=\frac{\alpha}{2}+\sqrt{\frac{\alpha^{2}}{4}-\frac{1}{k_{4}}}$ since we want to have a small $k_{5}$ for the condition $\mu_{3}(t) \geq \mu_{3}^{*}>0$. We choose $k_{4} \geq 4 / \alpha^{2}$ to make $k_{5}$ real and positive, and set $k_{2}=\delta k$ with $\delta$ being a small positive constant, say $\delta \leq \max \left(u_{d}^{\max }, \dot{u}_{d}^{\max }, r_{d}^{\max }\right)$, with $\dot{u}_{d}^{\max }$ being the maximum of $\left|\dot{u}_{d}\right|$. By observing from all functions $\mu_{i}(t), i=1,21,22,3$ that their negative parts have $k$ as a factor, and that the mass including added mass in the sway dynamics, $m_{22}$, is always larger than that in the surge dynamics, $m_{11}$, for surface ships, i.e., $\beta<1$, therefore we can always find a positive constant $k$ or even $k=0$ such that the conditions 1 and 2 hold for some small $\varepsilon_{i}>0, i=1,2,3$. For simplicity, one can take $\varepsilon_{i}=\beta$. In fact, $k=0$ results in the controllers proposed in $[19,106]$ (with nonzero $k_{2}$ selected). However if $k=0$ the condition 3 cannot be satisfied. Hence we should choose $k$ to be a small positive constant. Note that a small $k$ automatically implies a small $k_{*}$. The value
of this constant should be reduced if the reference yaw, $r_{d}$, and surge, $u_{d}$, velocities and surge acceleration, $\dot{u}_{d}$, are large. The physical meaning of this interpretation is that the distance from the ship to the point it aims to track should be increased if the velocities are large, otherwise the ship will miss that point due to high velocities. Furthermore if $\alpha$ is small, $k_{5}$ is automatically small. This physically means that if the damping in the sway is small, the control gain in the surge dynamics should also be small, otherwise the ship will slide in the sway direction. This can be mathematically seen from the expression of $\mu_{3}(t)$. Due to complicated expressions of $\mu_{i}(t)$, we give a simple procedure to choose the design constants $k_{1}$ and $k_{2}$ rather than present their complex explicit expressions.

1. Pick a small positive constant $k$, say $k \ll \max \left(u_{d}^{\max }, r_{d}^{\max }, \dot{u}_{d}^{\max }\right), \delta$ such that $\delta \leq \max \left(u_{d}^{\max }, r_{d}^{\max }\right)$ and set $k_{2}=\delta k, k_{4}=4 / \alpha^{2}$.
2. Substitute (5.48) into $\mu_{i}(t)$ and slowly increase $k$ and $k_{4}$ until all conditions 1, 2, and 3 hold with arbitrarily positive constants $\mu_{1}^{*}, \mu_{21}^{*}, \mu_{22}^{*}$, and $\mu_{3}^{*}$.

Once conditions 1,2 , and 3 hold, the design constant $k_{1}$ is calculated from (5.48) and $k_{2}=\delta k$. Based on the above choice of design constants, it directly follows that

- if (5.3) holds:

$$
\begin{equation*}
\dot{V}_{1} \leq-\mu_{1}^{*} x_{e}^{2}-\left(\mu_{211}^{m} r_{d}^{2}-\mu_{212}^{m} r_{d}^{\max }-\mu_{222}^{m} u_{d}^{\max }\right) y_{e}^{2}-\mu_{3}^{*} v_{e}^{2} \tag{5.49}
\end{equation*}
$$

- if (5.4) holds:

$$
\begin{align*}
\dot{V}_{1} \leq & -\mu_{1}^{*} x_{e}^{2}-\frac{\mu_{22}^{*} y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}-\mu_{3}^{*} v_{e}^{2}+ \\
& \left(\left|\mu_{211}(t)\right| \chi_{1}+\left|\mu_{212}(t)\right|\right) \chi_{1} e^{-\chi_{2}\left(t-t_{0}\right)} y_{e}^{2} \tag{5.50}
\end{align*}
$$

It is now seen that condition C 1 of Lemma 5.1 follows readily from (5.50) since the last term in the second line of (5.50) is linear in $y_{e}^{2}$ and its coefficient exponentially vanishes, see also proof of Lemma 5.1.

Remark 5.3. When the reference surge and sway velocities and surge acceleration are very large, the above procedure will result in very small control gains $k, k_{1}$, and $k_{2}$, which will give very slow convergence. Hence it is suggested that for large $u_{d}, r_{d}$, and $\dot{u}_{d}$, the control gains $k, k_{1}$, and $k_{2}$ should be chosen such that either conditions 1 and 2 or 1 and 3 hold to improve the convergence. The trade-off is that some "switches" in the control gains for different reference trajectories will occur.

Verifying Condition C2. We need to show that $g_{1}\left(t, X_{1 e}, X_{2 e}\right)$ satisfies condition C2. It can be seen that

$$
\begin{align*}
& \varpi_{e} \leq 1 \\
& k u_{d} \varpi_{e 1} y_{e} \leq k^{*} \\
& \left|\sin \left(z_{e}\right) / z_{e}\right| \leq 1  \tag{5.51}\\
& \left|\left(\cos \left(z_{e}\right)-1\right) / z_{e}\right| \leq 1 .
\end{align*}
$$

From (5.13) and (5.51), it is not hard to show that:

$$
\begin{align*}
& \left|p_{x}\right| \leq\left(1+k^{*}\right)\left(\left|u_{d}\right|+\left|v_{d}\right|\right)\left|z_{e}\right|, \\
& \left|p_{y}\right| \leq\left(1+k^{*}\right)\left(\left|u_{d}\right|+\left|v_{d}\right|\right)\left|z_{e}\right|,  \tag{5.52}\\
& \left|p_{z}\right| \leq k^{*}\left(2\left|p_{y}\right|+\left|p_{x}\right|\right) .
\end{align*}
$$

Applying (5.52) to (5.21) and noting Remark 5.3 yields the condition C2 readily.
Verifying Condition C3. This condition is verified by showing that the $X_{2 e^{-}}$ subsystem is GES. Take the following quadratic function:

$$
\begin{equation*}
V_{2}=\frac{1}{2}\left(z_{e}^{2}+\tilde{u}_{e}^{2}+\tilde{r}_{e}^{2}\right) \tag{5.53}
\end{equation*}
$$

Differentiating (5.53) along the solutions of the $X_{2 e}$-subsystem satisfies

$$
\begin{align*}
\dot{V}_{2} & \leq-\left(k_{3}-1.5\right) z_{e}^{2}-\left(\delta_{1}+\frac{d_{11}}{m_{11}}-0.5\right) \tilde{u}_{e}^{2}-\left(\delta_{2}+\frac{d_{33}}{m_{33}}-1\right) \tilde{r}_{e}^{2} \\
& \leq-\rho_{2} V_{2} \tag{5.54}
\end{align*}
$$

where

$$
\rho_{2}=\min \left(2\left(k_{3}-1.5\right), 2\left(\delta_{1}+\frac{d_{11}}{m_{11}}-0.5\right), 2\left(\delta_{2}+\frac{d_{33}}{m_{33}}-1\right)\right)
$$

Thus it suffices to choose the design constants $k_{3}, \delta_{1}$, and $\delta_{2}$ such that

$$
\begin{align*}
& k_{3}>1.5, \\
& \delta_{1} \geq 0.5,  \tag{5.55}\\
& \delta_{2} \geq 1 .
\end{align*}
$$

Hence condition C3 follows from (5.54). We have thus verified all the conditions of Lemma 5.1. Therefore the closed loop system (5.21) is GAS at the origin.

Remark 5.4. From the proof of Theorem 5.1 and Lemma 5.1, it is seen that when the yaw reference velocity $r_{d}$ satisfies the persistently exciting condition as stated in Assumption 5.1, the tracking error system is GES. This is the case reported in [ $19,71,106]$ based on different approaches. Therefore the results in those papers are a special case of the result in Theorem 5.1.

### 5.4 Simulations

This section validates the control laws (5.20) by simulating them on a monohull ship with the length of 32 m , mass of $118 \times 10^{3} \mathrm{~kg}$, and other parameters are calculated by using MARINTEK Ship Motion program version 3.18, a program for calculating
the added mass and damping matrices of the ship, as

$$
\begin{aligned}
& m_{11}=120 \times 10^{3} \mathrm{~kg}, m_{22}=177.9 \times 10^{3} \mathrm{~kg}, m_{33}=636 \times 10^{5} \mathrm{kgm}^{2}, \\
& d_{11}=215 \times 10^{2} \mathrm{kgs}^{-1}, d_{u 2}=43 \times 10^{2} \mathrm{kgm}^{-1}, d_{u 3}=21.5 \times 10^{2} \mathrm{kgsm}^{-2}, \\
& d_{22}=117 \times 10^{3} \mathrm{kgs}^{-1}, d_{v 2}=23.4 \times 10^{3} \mathrm{kgm}^{-1}, d_{v 3}=11.7 \times 10^{3} \mathrm{kgsm}^{-2}, \\
& d_{33}=802 \times 10^{4} \mathrm{kgm}^{2} s^{-1}, d_{r 2}=160.4 \times 10^{4} \mathrm{kgm}^{2}, d_{r 3}=80.2 \times 10^{4} \mathrm{kgm}^{2} s, \\
& d_{u i}=0, d_{v i}=0, d_{r i}=0, \forall i>3 .
\end{aligned}
$$

This ship has a minimum turning circle with the radius of 150 m , a maximum surge force of $5.2 \times 10^{9} \mathrm{~N}$, and a maximum yaw moment of $8.5 \times 10^{8} \mathrm{Nm}$. From the ship parameters, we have $\alpha=0.54$ and $\beta=0.55$. The reference trajectory is generated by the virtual ship (5.2) where the reference surge force $\tau_{u d}$ is taken as $\tau_{u d}=d_{11}-$ $m_{22} v_{d} r_{d}$, and the reference yaw moment $\tau_{r d}$ is taken as $\tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} r_{d}$ for the first 800 seconds and $\tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} r_{d}+0.01 d_{33}$ for the last 1000 seconds. This choice means that the reference surge velocity $u_{d}$ is $1 \mathrm{~ms}^{-1}$ over the entire simulation time, and that the reference trajectory is a straight line for the first 800 seconds and a circle with a radius of 200 m for the last 1000 seconds. We first pick $k=0.02$ and $k_{4}=13.9$, then increase these constants until conditions $\mathrm{a}, \mathrm{b}$, and c hold. We have

$$
\begin{aligned}
& k=0.1, k_{1}=0.34, k_{2}=0.1, k_{3}=2 \\
& k_{4}=15.26, k_{5}=0.19, \delta_{1}=5, \delta_{2}=5
\end{aligned}
$$

The initial conditions are chosen as

$$
\begin{aligned}
& {[x(0), y(0), \psi(0), u(0), v(0), r(0)]} \\
& =\left[15 \mathrm{~m}, 50 \mathrm{~m}, 0.7 \mathrm{rad}, 0 \mathrm{~ms}^{-1}, 0.5 \mathrm{~ms}^{-1}, 0 \mathrm{rads}^{-1}\right]
\end{aligned}
$$

The reference trajectory is generated by a virtual ship with the initial conditions of

$$
\begin{aligned}
& {\left[x_{d}(0), y_{d}(0), \psi_{d}(0), u_{d}(0), v_{d}(0), r_{d}(0)\right]} \\
& =\left[0.5 \mathrm{~m}, 0.5 \mathrm{~m}, 0.2 \mathrm{rad}, 1 \mathrm{~ms}^{-1}, 0 \mathrm{~ms}^{-1}, 0 \mathrm{rads}^{-1}\right] .
\end{aligned}
$$

The tracking errors are plotted in Figure 5.2a and control inputs are plotted in Figure 5.2b. In addition, the real and reference trajectories in the $(x, y)$-plane are plotted in Figure 5.3. It can be seen from Figure 5.2a that the tracking errors asymptotically converge to zero as proven in Theorem 5.1. From Figure 5.2b, we can see that the control inputs $\tau_{u}$ and $\tau_{r}$ are quite large in the first few seconds because of a short transient response time. Indeed, we can reduce the control effort by tuning the control design gains such as reducing the control gains $\delta_{1}$ and $\delta_{2}$. However, this will increase the transient response time.

For a comparison, we also simulate the controller proposed in [19] with the same initial conditions and reference trajectory in Figures 5.4 and 5.5. It is clearly seen that the controller proposed in [19] cannot track the straight line as discussed before. It should be noted that the controllers proposed in $[71,106]$ also result in similar nonzero errors although they were designed based on different approaches.


Figure 5.2 Simulation results of the controller proposed in this chapter: a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dashed line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; and $\mathbf{b}$. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line)

### 5.5 Conclusions

The key to control development is the coordinate change (5.10) to transform the tracking error system, in which the tracking errors are interpreted in the frame attached to the ship's body, to a triangular form, to which the popular backstepping technique can be applied. This coordinate change made it possible to relax the severe restrictions in $[19,71,106]$ on the reference trajectories generated by the virtual ships in the sense that the developed controllers can force the vehicles to track a straight line (way-point tracking), a curve and a combination of both. Several assumptions on the ship dynamics have been made such as all off-diagonal terms of the damping matrix, nonlinear damping matrix, and environmental disturbances induced by waves, wind, and ocean currents being ignored. For a justification, the reader is referred to Section 3.4.1.2. It is seen from Assumption 5.1 that the surge and yaw reference velocities cannot be zero, i.e., the problem of stabilization and/or parking is not considered in this chapter. In the next chapter, this limitation will be removed when a simultaneous stabilization and tracking control is addressed. This chapter is based on [109-111].


Figure 5.3 Simulation results of the controller proposed in this chapter: Tracking trajectory in the $(x, y)$-plane $\left((x, y)\right.$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$


Figure 5.4 Simulation results of the controller proposed in [19]: a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dashed line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; $\mathbf{b}$. Control inputs ( $\tau_{u}[\mathrm{~N}]$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line)


Figure 5.5 Simulation results of the controller proposed in [19]: Tracking trajectory in the $(x, y)$ plane $\left((x, y)\right.$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$

## Chapter 6

## Simultaneous Stabilization and Trajectory-tracking Control of Underactuated Ships

This chapter examines the problem of designing a single controller that achieves both stabilization and tracking simultaneously for underactuated surface ships without an independent sway actuator and with simplified dynamics. The proposed controller guarantees that stabilization and tracking errors converge to zero asymptotically from any initial values. In comparison with the preceding chapter, a path approaching the origin and a set-point can also be included in the reference trajectory, i.e., stabilization/regulation is also considered. The control development is based on some special coordinate transformations, Lyapunov's direct method and the backstepping technique, and utilizes passive properties of ship dynamics and their interconnected structure.

### 6.1 Control Objective

For the reader's convenience, the mathematical model of the underactuated ship moving in surge, sway, and yaw, see Section 3.4.1.2, is recaptured as

$$
\begin{align*}
\dot{x} & =u \cos (\psi)-v \sin (\psi), \\
\dot{y} & =u \sin (\psi)+v \cos (\psi), \\
\dot{\psi} & =r, \\
\dot{u} & =\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u},  \tag{6.1}\\
\dot{v} & =-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v, \\
\dot{r} & =\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r+\frac{1}{m_{33}} \tau_{r},
\end{align*}
$$

where all symbols in (6.1) are defined as in Section 3.4.1.2. The available controls are the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$. Since the sway control force is not
available, the ship model (6.1) is underactuated. Similar to the preceding chapter, it is assumed that the reference trajectory is generated by a virtual ship as

$$
\begin{align*}
\dot{x}_{d} & =u_{d} \cos \left(\psi_{d}\right)-v_{d} \sin \left(\psi_{d}\right) \\
\dot{y}_{d} & =u_{d} \sin \left(\psi_{d}\right)+v_{d} \cos \left(\psi_{d}\right) \\
\dot{\psi}_{d} & =r_{d} \\
\dot{u}_{d} & =\frac{m_{22}}{m_{11}} v_{d} r_{d}-\frac{d_{11}}{m_{11}} u_{d}+\frac{1}{m_{11}} \tau_{u d},  \tag{6.2}\\
\dot{v}_{d} & =-\frac{m_{11}}{m_{22}} u_{d} r_{d}-\frac{d_{22}}{m_{22}} v_{d} \\
\dot{r}_{d} & =\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u_{d} v_{d}-\frac{d_{33}}{m_{33}} r_{d}+\frac{1}{m_{33}} \tau_{r d},
\end{align*}
$$

where all variables have similar meanings as in the system (6.1) for the virtual reference ship. In this chapter, a single controller that simultaneously solves stabilization and tracking problems of underactuated surface ships is proposed under the following assumptions.

Assumption 6.1. The reference velocities, $u_{d}$ and $r_{d}$, are bounded and differentiable with bounded derivatives $\dot{u}_{d}, \ddot{u}_{d}$ and $\dot{r}_{d}$.

Assumption 6.2. One of the following conditions holds:

$$
\begin{aligned}
& \text { C1. } u_{d}=r_{d}=0 \text {. } \\
& \text { C2. } \int_{t_{0}} r_{d}^{2}(\tau) \mathrm{d} \tau \geq \sigma_{r}\left(t-t_{0}\right), \sigma_{r}>0, \forall 0 \leq t_{0} \leq t<\infty \text {. } \\
& \text { C3. }\left|u_{d}(t)\right| \geq \sigma_{u}>0 \text { and } \int_{0}^{\infty}\left|r_{d}(t)\right| \mathrm{d} t \leq \mu_{1}, 0 \leq \mu_{1}<\infty \text {. } \\
& \text { C4. } \int_{0}^{\infty}\left(\left|r_{d}(t)\right|+\left|u_{d}(t)\right|+\left|\dot{u}_{d}(t)\right|\right) \mathrm{d} t \leq \mu_{2}, 0 \leq \mu_{2}<\infty \text {. }
\end{aligned}
$$

Remark 6.1. The problems of regulation, stabilization, or dynamic positioning are included in condition C1. Tracking a circular path belongs to the case when condition C2 holds. Condition C3 covers the case of straight-line or way-point tracking. By Barbalat's lemma together with Assumption 6.2, Condition C4 implies that $\lim _{t \rightarrow \infty} r_{d}(t)=\lim _{t \rightarrow \infty} u_{d}(t)=0$. Hence the parking problem is captured by condition C4. Indeed, one can see that Condition C4 covers Condition C1. However we study these two conditions separately because Condition C 1 has been studied in the literature but Condition C 4 has not.

Remark 6.2. The work in $[19,53-56,71,106]$ can deal with either Condition C1 or Condition C2. In the preceding chapter of this book, we allowed for either Condition C2 or Condition C3.

Similar to the preceding chapter, we introduce position and orientation errors $x-x_{d}, y-y_{d}$, and $\psi-\psi_{d}$ in a frame attached to the ship body. This results in the error coordinates as

$$
\left[\begin{array}{l}
x_{e}  \tag{6.3}\\
y_{e} \\
\psi_{e}
\end{array}\right]=\left[\begin{array}{lll}
\cos (\psi) & \sin (\psi) & 0 \\
-\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{l}
x-x_{d} \\
y-y_{d} \\
\psi-\psi_{d}
\end{array}\right] .
$$

One can see that the convergence of $x_{e}, y_{e}$, and $\psi_{e}$ to the origin implies that of $x-x_{d}, y-y_{d}$ and $\psi-\psi_{d}$. We also define the velocity tracking errors as

$$
\begin{equation*}
u_{e}=u-u_{d}, v_{e}=v-v_{d}, r_{e}=r-r_{d} . \tag{6.4}
\end{equation*}
$$

Differentiating both sides of (6.3) along the solutions of (6.1) and (6.2) yields:

$$
\begin{align*}
& \dot{x}_{e}=u_{e}-u_{d}\left(\cos \left(\psi_{e}\right)-1\right)-v_{d} \sin \left(\psi_{e}\right)+r_{e} y_{e}+r_{d} y_{e}, \\
& \dot{y}_{e}=v_{e}-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)+u_{d} \sin \left(\psi_{e}\right)-r_{e} x_{e}-r_{d} x_{e}, \\
& \dot{\psi}_{e}=r_{e}, \\
& \dot{u}_{e}=\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d},  \tag{6.5}\\
& \dot{v}_{e}=-\frac{m_{11}}{m_{22}} u_{e} r_{d}-\frac{m_{11}}{m_{22}}\left(u_{e}+u_{d}\right) r_{e}-\frac{d_{22}}{m_{22}} v_{e}, \\
& \dot{r}_{e}=\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r+\frac{1}{m_{33}} \tau_{r}-\dot{r}_{d} .
\end{align*}
$$

The tracking and regulation problem of underactuated surface ships is therefore equivalent to stabilizing system (6.5) at the origin. In particular, we are interested in designing explicit expressions for $\tau_{u}$ and $\tau_{r}$ such that $\lim _{t \rightarrow \infty}\left\|X_{e}(t)\right\|=0$, for all $t \geq t_{0} \geq 0$ and $X_{e}\left(t_{0}\right) \in \mathbb{R}^{6}$ with $X_{e}=\left[x_{e}, y_{e}, \psi_{e}, u_{e}, v_{e}, r_{e}\right]^{T}$. In case of Condition C 2 , it will be shown that $\left\|X_{e}(t)\right\| \leq \gamma\left(\left\|X_{e}\left(t_{0}\right)\right\|\right) e^{-\mu\left(t-t_{0}\right)}$ with $\gamma$ being a class- $K$ function, and $\mu$ a positive constant, i.e., global $K$-exponential tracking is achieved.

### 6.2 Control Design

Similar to the preceding chapter, we observe from (6.5) that $x_{e}$ and $\psi_{e}$ can be stabilized by $u_{e}$ and $r_{e}$. On the other hand, $\psi_{e}$ should be chosen as a virtual control to stabilize the sway error $y_{e}$. However when the surge reference velocity is zero or approaches zero, i.e., the case of stabilization and parking, we cannot use $\psi_{e}$ to stabilize $y_{e}$ since $u_{d} \sin \left(\psi_{e}\right)$ is zero or approaches zero. In this case, we need some persistently exciting signal in $r_{e}$ to stabilize $y_{e}$. As such, we choose the coordinate transformation

$$
\begin{equation*}
z_{e}=\psi_{e}+\arcsin \left(\frac{k(t) y_{e}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}\right) \tag{6.6}
\end{equation*}
$$

with

$$
k(t)=\lambda_{1}\left(u_{d}+\lambda_{2} \cos \left(\lambda_{3} t\right)\right)
$$

where the constants $\lambda_{i}, 1 \leq i \leq 3$ are such that

$$
\begin{equation*}
\sup _{t \geq 0}|k(t)|<1 \tag{6.7}
\end{equation*}
$$

They will be specified in the stability analysis later. Notice that (6.6) is well defined and that convergence of $z_{e}$ and $y_{e}$ to the origin implies that of $\psi_{e}$.

Remark 6.3. Using the nonlinear coordinate transformation (6.6) instead of $z_{e}=$ $\psi_{e}+k(t) y_{e}$, we avoid the ship moving left and right largely when $y_{e}$ is large. If one uses $z_{e}=\psi_{e}+\arcsin \left(k(t) y_{e} / \sqrt{1+y_{e}^{2}}\right)$, the problem of the ship whirling around is avoided. Indeed using these coordinate transformations will result in a much simpler tracking error system than using (6.6). However, it is readily shown that these coordinate transformations result in a local control solution because of the term $r_{e} x_{e}$ in the $y_{e}$-dynamics. The term $\lambda_{1} u_{d}$ plays the role of stabilizing $y_{e}$ when $u_{d}$ is not zero and does not approach to zero while the term $\lambda_{1} \lambda_{2} \cos \left(\lambda_{3} t\right)$ is used to introduce the persistently exciting signal in $r_{e}$ to stabilize $y_{e}$ when $u_{d}$ is zero or approaches zero.

Remark 6.4. The function arcsin in (6.6) can be replaced by several other smooth functions such as arctan and tanh, or the identity function. We have chosen the function arcsin because of its simplicity.

Using the new coordinate (6.6), the ship error dynamics (6.5) are rewritten as

$$
\begin{align*}
\dot{x}_{e}= & u_{e}+u_{d} \frac{\left(\varpi_{2}-\varpi_{1}\right)}{\varpi_{2}}+\frac{k(t) v_{d} y_{e}}{\varpi_{2}}+r_{e} y_{e}+r_{d} y_{e}+p_{x}, \\
\dot{y}_{e}= & v_{e}+v_{d} \frac{\left(\varpi_{2}-\varpi_{1}\right)}{\varpi_{2}}-\frac{k(t) u_{d} y_{e}}{\varpi_{2}}-r_{e} x_{e}-r_{d} x_{e}+p_{y}, \\
\dot{z}_{e}= & \left(1-\frac{k(t) x_{e}}{\varpi_{1}}\right) r_{e}+\frac{1}{\varpi_{1}}\left(\dot{k}(t) y_{e}+k(t)\left(v_{e}+v_{d} \frac{\left(\varpi_{2}-\varpi_{1}\right)}{\varpi_{2}}-\right.\right. \\
& \left.\frac{k(t) u_{d} y_{e}}{\varpi_{2}}-r_{d} x_{e}\right)-\frac{k(t) y_{e}}{\varpi_{2}^{2}}\left(x_{e} u_{e}+y_{e} v_{e}+\left(x_{e} u_{d}+y_{e} v_{d}\right) \times\right. \\
& \left.\frac{\left(\varpi_{2}-\varpi_{1}\right)}{\varpi_{2}}+\left(x_{e} v_{d}-y_{e} u_{d}\right) \frac{k(t) y_{e}}{\varpi_{2}}\right)+p_{z}, \\
\dot{u}_{e}= & \frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}}\left(u_{e}+u_{d}\right)+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d}, \\
\dot{v}_{e}= & -\frac{d_{22}}{m_{22}} v_{e}-\frac{m_{11}}{m_{22}} u_{e} r_{d}-\frac{m_{11}}{m_{22}}\left(u_{e}+u_{d}\right) r_{e}, \\
\dot{r}_{e}= & \frac{m_{11}-m_{22}}{m_{33}} u v-\frac{d_{33}}{m_{33}}\left(r_{e}+r_{d}\right)+\frac{1}{m_{33}} \tau_{r}-\dot{r}_{d}, \tag{6.8}
\end{align*}
$$

where, for notational simplicity, we have used $k(t)$ instead of using its explicit expression given in (6.6), and defined as

$$
\begin{align*}
& \varpi_{1}=\sqrt{1+x_{e}^{2}+\left(1-k^{2}(t)\right) y_{e}^{2}}, \varpi_{2}=\sqrt{1+x_{e}^{2}+y_{e}^{2}} \\
& p_{x}=-u_{d}\left(\left(\cos \left(z_{e}\right)-1\right) \frac{\varpi_{1}}{\varpi_{2}}+\sin \left(z_{e}\right) \frac{k(t) y_{e}}{\varpi_{2}}\right)- \\
& v_{d}\left(\sin \left(z_{e}\right) \frac{\varpi_{1}}{\varpi_{2}}-\left(\cos \left(z_{e}\right)-1\right) \frac{k(t) y_{e}}{\varpi_{2}}\right) \\
& p_{y}=-v_{d}\left(\left(\cos \left(z_{e}\right)-1\right) \frac{\varpi_{1}}{\varpi_{2}}+\sin \left(z_{e}\right) k(t) \varpi_{2}^{-1} y_{e}\right)+ \\
& u_{d}\left(\sin \left(z_{e}\right) \frac{\varpi_{1}}{\varpi_{2}}-\left(\cos \left(z_{e}\right)-1\right) \frac{k(t) y_{e}}{\varpi_{2}}\right) \\
& p_{z}=\frac{1}{\varpi_{1}}\left(k(t) p_{y}-\frac{k(t) y_{e}}{\varpi_{2}}\left(x_{e} p_{x}+y_{e} p_{y}\right)\right) \tag{6.9}
\end{align*}
$$

The effort that we have made so far is to put the tracking error dynamics in the triangular form of (6.8), and to have the term $-k(t) u_{d} y_{e} / \varpi_{2}$ in the $y_{e}$-dynamics. This term plays an important role in stabilizing the $y_{e}$-dynamics. In this section, a procedure to design a universal controller for the tracking error system (6.8) is presented in detail. The triangular structure of (6.8) suggests that we design the actual controls $\tau_{u}$ and $\tau_{r}$ in two stages. First, we design virtual velocity controls $u_{e}$ and $r_{e}$ to globally asymptotically stabilize $x_{e}, y_{e}, z_{e}$ and $v_{e}$ at the origin. Based on the backstepping technique, the controls $\tau_{u}$ and $\tau_{r}$ are then designed to force the error between the virtual velocity controls and their actual values exponentially to zero. Before designing the control laws, $\tau_{u}$ and $\tau_{r}$, we note if the constants $\lambda_{i}, 1 \leq$ $i \leq 3$ are chosen such that $\sup _{t \geq 0}|k(t)|<1$, then (6.6) implies that

$$
\begin{equation*}
0<1-\sup _{t \geq 0}|k(t)| \leq 1-\frac{k(t) x_{e}}{\varpi_{1}}<2 \tag{6.10}
\end{equation*}
$$

This allows us to design control laws for $\tau_{u}$ and $\tau_{r}$ to globally asymptotically stabilize (6.8). The control design consists of two steps as follows.

## Step 1

In this step, we define the following virtual control errors

$$
\begin{align*}
& \tilde{u}_{e}=u_{e}-u_{e}^{d} \\
& \tilde{r}_{e}=r_{e}-r_{e}^{d} \tag{6.11}
\end{align*}
$$

where $u_{e}^{d}$ and $\mathrm{r}_{e}^{d}$ are the virtual velocity controls of $u_{e}$ and $r_{e}$. Unlike the standard application of backstepping, in order to reduce complexity of the controller expressions, we will choose a simpler virtual control law $u_{e}^{d}$ without canceling the known terms. The virtual surge and yaw velocity controls are chosen as

$$
\begin{align*}
& u_{e}^{d}=-k_{1} x_{e}+k_{2} r_{d} y_{e}, \\
& r_{e}^{d}=r_{1 e}^{d}+r_{2 e}^{d} \tag{6.12}
\end{align*}
$$

where

$$
\begin{align*}
& r_{1 e}^{d}=-\frac{1}{\left(1-k(t) x_{e} / \varpi_{1}\right)} \frac{1}{\varpi_{1}}\left(\dot{k}(t) y_{e}+k(t)\left(v_{e}+v_{d} \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}}-\right.\right. \\
&\left.\frac{k(t) u_{d} y_{e}}{\varpi_{2}}-r_{d} x_{e}\right)-\frac{k(t) y_{e}}{\varpi_{2}^{2}}\left(x_{e}\left(-k_{1} x_{e}+k_{2} r_{d} y_{e}\right)+y_{e} v_{e}+\right. \\
&\left.\left.\left(x_{e} u_{d}+y_{e} v_{d}\right) \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}}+\left(x_{e} v_{d}-y_{e} u_{d}\right) \frac{k(t) y_{e}}{\varpi_{2}}\right)\right),  \tag{6.13}\\
& r_{2 e}^{d}=-\left[\left(1-k(t) x_{e} / \varpi_{1}\right)\right]^{-1}\left(k_{3} z_{e}+p_{z}\right), \tag{6.14}
\end{align*}
$$

and $k_{i}, i=1,2,3$, are positive constants to be selected later.

## Step 2

Differentiating (6.11) along the solutions of (6.8) and (6.12) yields

$$
\begin{align*}
& \dot{\tilde{u}}_{e}=\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d}-\dot{u}_{e}^{d}, \\
& \dot{\tilde{r}}_{e}=\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r+\frac{1}{m_{33}} \tau_{r}-\dot{r}_{d}-\dot{r}_{e}^{d}, \tag{6.15}
\end{align*}
$$

where $\dot{u}_{e}^{d}$ and $\dot{r}_{e}^{d}$ are the first time derivatives of $u_{e}^{d}$ and $r_{e}^{d}$ along the solutions of (6.8). From (6.15) we choose the actual controls $\tau_{u}$ and $\tau_{r}$ as

$$
\begin{align*}
\tau_{u}= & m_{11}\left(-c_{1} \tilde{u}_{e}-\frac{m_{22}}{m_{11}} v r+\frac{d_{11}}{m_{11}}\left(u_{e}^{d}+u_{d}\right)+\dot{u}_{d}+\dot{u}_{e}^{d}+\frac{k(t) x_{e} y_{e} z_{e}}{\varpi_{1} \varpi_{2}^{2}}\right) \\
\tau_{r}= & m_{33}\left(-c_{2} \tilde{r}_{e}-\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v+\frac{d_{33}}{m_{33}}\left(r_{e}^{d}+r_{d}\right)+\right. \\
& \left.\dot{r}_{d}+\dot{r}_{e}^{d}-\left(1-\frac{k(t) x_{e}}{\varpi_{1}}\right) z_{e}\right) \tag{6.16}
\end{align*}
$$

where $c_{1}$ and $c_{2}$ are positive constants. Substituting (6.16) and (6.12) into (6.8) yields the closed loop system

$$
\begin{aligned}
\dot{x}_{e}= & -k_{1} x_{e}+k_{2} r_{d} y_{e}+u_{d}\left(\varpi_{2}-\varpi_{1}\right) \varpi_{2}^{-1}+k(t) v_{d} \varpi_{2}^{-1} y_{e}+ \\
& r_{1 e}^{d} y_{e}+r_{d} y_{e}+r_{2 e}^{d} y_{e}+\tilde{u}_{e}+\tilde{r}_{e} y_{e}+p_{x}, \\
\dot{y}_{e}= & v_{e}+v_{d}\left(\varpi_{2}-\varpi_{1}\right) \varpi_{2}^{-1}-k(t) u_{d} \varpi_{2}^{-1} y_{e}-r_{1 e}^{d} x_{e}-r_{d} x_{e}- \\
& r_{2 e}^{d} x_{e}-\tilde{r}_{e} x_{e}+p_{y},
\end{aligned}
$$

$$
\begin{align*}
\dot{v}_{e}= & -\frac{d_{22}}{m_{22}} v_{e}-\frac{m_{11}}{m_{22}} u_{e}^{d} r_{d}-\frac{m_{11}}{m_{22}}\left(u_{e}+u_{d}\right) r_{1 e}^{d}-\frac{m_{11}}{m_{22}}\left(r_{d}+r_{1 e}^{d}\right) \tilde{u}_{e}- \\
& \frac{m_{11}}{m_{22}}\left(\tilde{u}_{e}+u_{e}^{d}+u_{d}\right)\left(r_{2 e}^{d}+\tilde{r}_{e}\right) \\
\dot{z}_{e}= & -k_{3} z_{e}+\left(1-\frac{k(t) x_{e}}{\varpi_{1}}\right) \tilde{r}_{e}-\frac{k(t) x_{e} y_{e} \tilde{u}_{e}}{\varpi_{1} \varpi_{2}^{2}}, \\
\dot{\tilde{u}}_{e}= & -\left(c_{1}+\frac{d_{11}}{m_{11}}\right) \tilde{u}_{e}+\frac{k(t) x_{e} y_{e} z_{e}}{\varpi_{1} \varpi_{2}^{2}} \\
\dot{\tilde{r}}_{e}= & -\left(c_{2}+\frac{d_{33}}{m_{33}}\right) \tilde{r}_{e}-\left(1-\frac{k(t) x_{e}}{\varpi_{1}}\right) z_{e} . \tag{6.17}
\end{align*}
$$

We now state our main result, the proof of which is given in the next section.
Theorem 6.1. Assume that the reference signals $\left(x_{d}, y_{d}, \psi_{d}, v_{d}\right)$ are generated by the virtual ship model (6.2) and that the reference velocities $\left(u_{d}, r_{d}\right)$ satisfy Assumptions 6.1 and 6.2. If the universal state feedback control law (6.16) is applied to the ship system (6.1) then the tracking errors $\left(x(t)-x_{d}(t), y(t)-y_{d}(t), \psi(t)-\psi_{d}(t)\right.$, and $\left.v(t)-v_{d}(t)\right)$ converge to zero asymptotically from any initial values, i.e., the closed loop system (6.17) is GAS at the origin, with an appropriate choice of the design constants $\lambda_{i}$ and $k_{i}, 1 \leq i \leq 3$. Furthermore, if Assumption 6.2 holds with $C 2$, the closed loop system (6.17) is globally $K$-exponentially stable at the origin.

### 6.3 Stability Analysis

We only need to show that the transformed tracking errors, $\left(x_{e}, y_{e}, z_{e}, v_{e}\right)$, converge to zero from any initial values. Under Assumption 6.1, boundedness of the controls $\tau_{u}$ and $\tau_{r}$ follows readily. We note that the closed loop system (6.17) consists of two subsystems, $\left(z_{e}, \tilde{u}_{e}, \tilde{r}_{e}\right)$ and $\left(x_{e}, y_{e}, v_{e}\right)$, in an interconnected structure. To prove Theorem 6.1, we take the following Lyapunov function:

$$
\begin{equation*}
V_{1}\left(X_{1 e}\right)=\frac{1}{2}\left(z_{e}^{2}+\tilde{u}_{e}^{2}+\tilde{r}_{e}^{2}\right) \tag{6.18}
\end{equation*}
$$

for the $\left(z_{e}, \tilde{u}_{e}, \tilde{r}_{e}\right)$-subsystem, and

$$
\begin{equation*}
V_{2}\left(X_{2 e}\right)=\frac{1}{2}\left(x_{e}^{2}+y_{e}^{2}+k_{4}\left(y_{e}+k_{5} v_{e}\right)^{2}\right) \tag{6.19}
\end{equation*}
$$

for the $\left(x_{e}, y_{e}, v_{e}\right)$-subsystem, where $k_{4}$ and $k_{5}$ are positive constants to be selected later and

$$
X_{1 e}=\left[\begin{array}{c}
\tilde{z}_{e} \\
\tilde{u}_{e} \\
\tilde{r}_{e}
\end{array}\right], X_{2 e}=\left[\begin{array}{c}
x_{e} \\
y_{e} \\
v_{e}
\end{array}\right] .
$$

From (6.18) and (6.19), if we consider $V=V_{1}+V_{2}$ as a Lyapunov function for the closed loop system (6.17) it can then be seen that this function has a connection with the energy of the ship system as follows. We write the function $V$ as

$$
\begin{equation*}
V=X^{T} A X \tag{6.20}
\end{equation*}
$$

where $X=\left[x_{e}, y_{e}, z_{e}, u-\left(u_{d}+u_{e}^{d}\right), v-v_{d}, r-\left(r_{d}+r_{e}^{d}\right)\right]^{T}$ and $A$ is a diagonal positive definite matrix. It is now seen that the function $V$ consists of two positive definite parts. The first part related to $\left(x_{e}, y_{e}, z_{e}\right)$ can be regarded as the "potential" energy. The second part related to $\left(u-\left(u_{d}+u_{e}^{d}\right), v-v_{d}, r-\left(r_{d}+r_{e}^{d}\right)\right)$ is referred to as the "kinetic" energy of the ship, which is defined with respect to a coordinate that moves at a speed of $\left(\left(u_{d}+u_{e}^{d}\right), v_{d},\left(r_{d}+r_{e}^{d}\right)\right)$. Hence we wish to bring the energy function $V$ to zero. However considering this function together with the closed loop (6.17) is difficult, we consider the subsystems $\left(z_{e}, \tilde{u}_{e}, \tilde{r}_{e}\right)$ and $\left(x_{e}, y_{e}, v_{e}\right)$ separately.
( $\left.z_{e}, \tilde{u}_{e}, \tilde{r}_{e}\right)$-subsystem. From the last three equations of (6.17), it can be seen that the ( $z_{e}, \tilde{u}_{e}, \tilde{r}_{e}$ )-subsystem is GES at the origin by considering the Lyapunov function $V_{1}\left(X_{1 e}\right)$, see (6.18), whose time derivative along the solutions of the last three equations of (6.17) is

$$
\begin{equation*}
\dot{V}_{1}=-k_{3} z_{e}^{2}-\left(c_{1}+\frac{d_{11}}{m_{11}}\right) \tilde{u}_{e}^{2}-\left(c_{2}+\frac{d_{33}}{m_{33}}\right) \tilde{r}_{e}^{2} \tag{6.21}
\end{equation*}
$$

Therefore we have

$$
\begin{equation*}
\left\|X_{1 e}(t)\right\| \leq\left\|X_{1 e}\left(t_{0}\right)\right\| e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.22}
\end{equation*}
$$

where $\sigma_{1}=\min \left\{k_{3},\left(c_{1}+d_{11} / m_{11}\right),\left(c_{2}+d_{33} / m_{33}\right)\right\}$.
$\left(x_{e}, y_{e}, v_{e}\right)$-subsystem. To investigate stability of this subsystem, we take the Lyapunov function $V_{2}\left(X_{2 e}\right)$, see (6.19), whose first time derivative, after some manipulation, along the solutions of the first three equations of (6.17) is

$$
\begin{align*}
\dot{V}_{2}= & -k_{1} x_{e}^{2}-\frac{k(t) u_{d}\left(1+k_{4}\right) y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}-\beta k_{2} k_{4} k_{5} r_{d}^{2} y_{e}^{2}- \\
& k_{4} k_{5}\left(\alpha k_{5}-1\right) v_{e}^{2}+M+\Omega+\Phi \tag{6.23}
\end{align*}
$$

where for simplicity, we have defined $\alpha=d_{22} / m_{22}, \beta=m_{11} / m_{22}$,

$$
\begin{align*}
M= & k_{2} r_{d} x_{e} y_{e}+\left(u_{d} x_{e}+v_{d}\left(k_{4}+1\right) y_{e}\right) \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}}+ \\
& \frac{k(t) y_{e}\left(v_{d} x_{e}-k_{4} k_{5} u_{d} v_{e}\right)}{\varpi_{2}}+k_{4} k_{5} v_{d} \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}} v_{e}- \\
& \beta k_{2} k_{4} k_{5}^{2} r_{d}^{2} y_{e} v_{e}-\beta k_{4} k_{5}\left(y_{e}+k_{5} v_{e}\right)\left(k_{2} r_{d} y_{e}+u_{d}\right) r_{1 e}^{d} \tag{6.24}
\end{align*}
$$

$$
\begin{gather*}
\Omega=x_{e}\left(p_{x}+\tilde{u}_{e}\right)+y_{e} p_{y}+k_{4}\left(y_{e}+k_{5} v_{e}\right)\left(p_{y}-\tilde{r}_{e} x_{e}-\right. \\
\left.r_{2 e}^{d} x_{e}-\beta k_{5}\left(\left(r_{d}+r_{1 e}^{d}\right) \tilde{u}_{e}+\left(\tilde{u}_{e}+u_{e}^{d}+u_{d}\right)\left(\tilde{r}_{e}+r_{2 e}^{d}\right)\right)\right)  \tag{6.25}\\
\Phi=k_{4}\left(y_{e}+k_{5} v_{e}\right)\left(-1+\beta k_{1} k_{5}\right)\left(r_{1 e}^{d}+r_{d}\right) x_{e}+\left(1-k_{4}\left(\alpha k_{5}-1\right)\right) y_{e} v_{e} \tag{6.26}
\end{gather*}
$$

We first choose the design constants such that

$$
\begin{align*}
& 1-k_{4}\left(\alpha k_{5}-1\right)=0, \\
& \beta k_{1} k_{5}-1=0 \tag{6.27}
\end{align*}
$$

to cancel the common term $\Phi$. Noticing that

$$
\begin{align*}
& \left|\frac{\sin \left(z_{e}\right)}{z_{e}}\right| \leq 1,\left|\frac{\cos \left(z_{e}\right)-1}{z_{e}}\right| \leq 1, \\
& \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}}=\frac{k^{2}(t) y_{e}^{2}}{\varpi_{2}\left(\varpi_{2}+\varpi_{1}\right)},  \tag{6.28}\\
& \frac{1}{\varpi_{1}} \leq \frac{1}{\varpi_{2} \sqrt{1-k^{2}(t)}}
\end{align*}
$$

after a lengthy but simple calculation of upper bounds of $M$ and $\Omega$ by completing the squares, we arrive at

$$
\begin{align*}
\dot{V}_{2} \leq & -p_{1}(t) x_{e}^{2}-p_{21}(t) y_{e}^{2}-\frac{p_{22}(t) y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}+\frac{p_{23}(t) y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}- \\
& p_{3}(t) v_{e}^{2}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)}, \tag{6.29}
\end{align*}
$$

where

$$
\begin{align*}
p_{1}(t)= & k_{1}-\frac{1}{4 \rho_{1}}\left(k_{2}+|k(t)|\left|v_{d}\right|+\frac{\beta k_{4} k_{5}|k(t)|\left|u_{d}\right|\left|r_{d}\right|}{1-|k(t)|}+\right. \\
& \left.\frac{\beta k_{1} k_{2} k_{4} k_{5}|k(t)|}{1-|k(t)|}\right),  \tag{6.30}\\
p_{21}(t)= & p_{211}(t) r_{d}^{2}-p_{212}(t)\left|r_{d}\right|, \\
p_{211}(t)= & \beta k_{2} k_{4} k_{5}-k_{2} \rho_{1}-\frac{\beta k_{2} k_{4} k_{5}^{2}}{4 \rho_{2}}-\frac{\beta k_{2} k_{4} k_{5}}{1-|k(t)|}\left(\frac{|k(t)|}{2 \rho_{2} \sqrt{1-k^{2}(t)}}+\right. \\
& \left.\left(k_{2}+1+k_{1} \rho_{1}\right)|k(t)|\right)-\frac{\beta k_{2} k_{4} k_{5}^{2}}{4 \rho_{2}(1-|k(t)|)}\left(\frac{|\dot{k}(t)|+|k(t)| k^{2}(t)}{\sqrt{1-k^{2}(t)}}+\right. \\
& \left(k^{2}(t)\left|u_{d}\right|+\left(1+k_{1}+k_{2}\right)|k(t)|+|k(t)| k^{2}(t)\left|u_{d}\right|+\right.
\end{align*}
$$

$$
\begin{align*}
& \left.|k(t)| k^{2}(t)\left|v_{d}\right|+k^{2}(t)\left|u_{d}\right|+k^{2}(t)\left|v_{d}\right|\right)-\frac{\beta k_{4} k_{5}^{2}|k(t)|\left|u_{d}\right|}{4 \rho_{2}(1-|k(t)|)}, \\
& p_{212}(t)=\frac{\beta k_{2} k_{4} k_{5}|\dot{k}(t)|}{(1-|k(t)|) \sqrt{1-k^{2}(t)}}, \\
& p_{22}(t)=\left(1+k_{4}\right) k(t) u_{d}-p_{221}(t), \\
& p_{221}(t)=k_{4} k_{5} \rho_{3}|k(t)|\left|u_{d}\right|+\frac{\beta k_{2} k_{4} k_{5}}{1-|k(t)|}\left(|k(t)|+1+\frac{1}{\sqrt{1-k^{2}(t)}}\right) \times \\
& k^{2}(t)\left|u_{d}\right|\left|r_{d}\right|+\frac{\beta k_{4} k_{5}}{1-|k(t)|}\left(\frac{|\dot{k}(t)|\left|u_{d}\right|}{\sqrt{1-k^{2}(t)}}+\frac{|k(t)|\left|u_{d}\right|}{4 \rho_{2} \sqrt{1-k^{2}(t)}}+k^{2}(t) \times\right. \\
& \left(|k(t)|\left|u_{d}\right|\left|v_{d}\right|+u_{d}^{2}\right)+\frac{\rho_{1}|k(t)|\left|u_{d}\right|\left|r_{d}\right|}{\sqrt{1-k^{2}(t)}}+k_{1}|k(t)|\left|u_{d}\right|+k_{2}|k(t)| \times \\
& \left|u_{d}\right|\left|r_{d}\right|+\frac{|k(t)|\left|u_{d}\right|}{4 \rho_{2} \sqrt{1-k^{2}(t)}}+|k(t)| k^{2}(t) u_{d}^{2}+|k(t)| k^{2}(t)\left|u_{d}\right|\left|v_{d}\right|+ \\
& \left.k^{2}(t)\left(u_{d}^{2}+\left|u_{d}\right|\left|v_{d}\right|\right)\right)+\frac{\beta k_{4} k_{5}^{2}}{4 \rho_{2}(1-|k(t)|)}\left(\frac{|\dot{k}(t)|\left|u_{d}\right|+\left|k^{3}(t)\right|\left|u_{d} v_{d}\right|}{\sqrt{1-k^{2}(t)}}+\right. \\
& k^{2}(t) u_{d}^{2}+\left(k_{1}+k_{2}\left|r_{d}\right|\right)|k(t)|\left|u_{d}\right|+|k(t)| k^{2}(t) u_{d}^{2}+k^{2}(t)(|k(t)|+1) \times \\
& \left.\left|u_{d}\right|\left|v_{d}\right|+k^{2}(t) u_{d}^{2}\right)+k^{2}(t)\left|u_{d}\right| \text {, }  \tag{6.32}\\
& p_{23}(t)=\left(k^{2}(t)\left(1+k_{4}\right)+\rho_{1}|k(t)|+\rho_{3} k^{2}(t) k_{4} k_{5}\right)\left|v_{d}\right|+\frac{\beta k_{4} k_{2} k_{5}}{1-|k(t)|} \times \\
& \left(\frac{|k(t)| k^{2}(t)\left|v_{d}\right|\left|r_{d}\right|}{\sqrt{1-k^{2}(t)}}+k^{2}(t)(|k(t)|+1)\left|v_{d}\right|\left|r_{d}\right|\right),  \tag{6.33}\\
& p_{3}(t)=k_{4} k_{5}\left(k_{5} \alpha-1\right)-\left(\frac{k_{4} k_{5}|k(t)|\left|u_{d}\right|}{4 \rho_{3}}+\frac{k_{4} k_{5} k^{2}(t)\left|v_{d}\right|}{4 \rho_{3}}+\right. \\
& \left.\rho_{2} \beta k_{2} k_{4} k_{5}^{2} r_{d}^{2}\right)-\frac{2 \rho_{2} \beta k_{2} k_{4} k_{5}|k(t)|}{(1-|k(t)|) \sqrt{1-k^{2}(t)}}- \\
& \frac{2 \beta \rho_{2} k_{4} k_{5}|k(t)|\left|u_{d}\right|}{1-|k(t)|}-\frac{\beta k_{2} k_{4} k_{5}^{2}}{1-|k(t)|}\left(\rho _ { 2 } \left(\frac{|\dot{k}(t)|}{\sqrt{1-k^{2}(t)}}+\right.\right. \\
& |k(t)| k^{2}(t) v_{d}^{2}+k^{2}(t)\left|u_{d}\right|+|k(t)| r_{d}^{2}+\left(k_{1}+k_{2} r_{d}^{2}\right) \times \\
& \left.\left.|k(t)|+(|k(t)|+1) k^{2}(t)\left(\left|v_{d}\right|+\left|u_{d}\right|\right)\right)+\frac{2|k(t)|\left|r_{d}\right|}{\sqrt{1-k^{2}(t)}}\right)-
\end{align*}
$$

$$
\begin{align*}
& \frac{\beta k_{4} k_{5}^{2}}{1-|k(t)|}\left(\rho _ { 2 } \left(\left|\dot{k}(t) u_{d}\right|+\left|k(t) u_{d}\right|\left(\left|v_{d}\right|+3\right)+\left(k_{1}+k_{2}\left|r_{d}\right|\right) \times\right.\right. \\
& \left.\left.\left|k(t) u_{d}\right|+(|k(t)|+2) k^{2}(t) u_{d}^{2}+k^{2}(t)\left|u_{d} v_{d}\right|(|k(t)|+1)\right)\right) \tag{6.34}
\end{align*}
$$

where $\chi_{1}, \chi_{2}$ are nondecreasing functions of $\left\|X_{1 e}\left(t_{0}\right)\right\|, \sigma_{1}$ is given in (6.22), and $\rho_{i}>0, i=1,2,3$.

We are now in a position to select the design parameters, $\lambda_{i}, 1 \leq i \leq 3$ and $k_{j}, j=$ $1,2,4,5$ so that the closed loop system (6.17) is GAS at the origin. We proceed to choose these parameters by investigating all the cases of Assumption 6.2. For each case, we choose a subset of $\lambda_{i}$ and $k_{j}$. As discussed before, we firstly choose

$$
\begin{equation*}
0<\lambda_{1}<1 /\left(u_{d}^{\max }+\lambda_{2}\right), \lambda_{2}>0, \lambda_{3}>0 \tag{6.35}
\end{equation*}
$$

so that $\sup _{t \geq 0}|k(t)| \leq k_{*}<1$, where $u_{d}^{\max }$ is the maximum value of $\left|u_{d}(t)\right|$. Note that this primary choice guarantees $\sup _{t \geq 0}|k(t)|<1$.

Secondly, we choose $\lambda_{i}, 1 \leq i \leq 3$ and $k_{j}$ such that

$$
\begin{equation*}
p_{1}(t) \geq p_{1}^{*}>0, p_{3}(t) \geq p_{3}^{*}>0 \tag{6.36}
\end{equation*}
$$

Before going further to choose $\lambda_{i}$ and $k_{j}$, let us discuss each case of Assumption 6.2.
Case C1. From the fourth equation of (6.2), we have

$$
v_{d}(t)=v_{d}\left(t_{0}\right) e^{-\alpha_{1}\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0, \alpha_{1}>0
$$

Therefore there exists a nondecreasing function $\pi_{1}$ of $\left|v_{d}\left(t_{0}\right)\right|$ such that

$$
\begin{equation*}
p_{23}(t) \leq \pi_{1} e^{-\alpha_{1}\left(t-t_{0}\right)} \tag{6.37}
\end{equation*}
$$

It is noted that in this case

$$
\begin{equation*}
p_{21}(t)=p_{22}(t)=0 \tag{6.38}
\end{equation*}
$$

From (6.35)-(6.38), we can write (6.29) as

$$
\begin{equation*}
\dot{V}_{2} \leq-p_{1}^{*} x_{e}^{2}-p_{3}^{*} v_{e}^{2}+\pi_{1} e^{-\alpha_{1}\left(t-t_{0}\right)} \frac{y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.39}
\end{equation*}
$$

which yields

$$
\begin{equation*}
\dot{V}_{2} \leq-p_{1}^{*} x_{e}^{2}-p_{3}^{*} v_{e}^{2}+\left(\eta_{1} V_{2}+\eta_{2}\right) e^{-\sigma_{2}\left(t-t_{0}\right)}, \sigma_{2}=\min \left(\alpha_{1}, \sigma_{1}\right) \tag{6.40}
\end{equation*}
$$

for some nondecreasing functions $\eta_{1}$ and $\eta_{2}$ of $\left\|\left(X_{1 e}\left(t_{0}\right), v_{d}\left(t_{0}\right)\right)\right\|$. We now present a technical lemma to investigate stability of the $\left(x_{e}, y_{e}, v_{e}\right)$-subsystem, in this case based on (6.40).

Lemma 6.1. Consider the following first-order scalar differential equation

$$
\begin{equation*}
\dot{x}=(a x+b) e^{-c\left(t-t_{0}\right)}+\delta(t) x, \tag{6.41}
\end{equation*}
$$

where $a \geq 0, b \geq 0$ and $c>0$ are constants. If $\delta(t)$ enjoys the property that there exists a constant $\sigma \geq 0$ such that

$$
\begin{equation*}
\int_{0}^{\infty}|\delta(t)| \mathrm{d} t \leq \sigma<\infty \tag{6.42}
\end{equation*}
$$

then the solution of (6.41) satisfies

$$
\begin{equation*}
|x(t)| \leq\left|x\left(t_{0}\right)\right| e^{\sigma} e^{a / c}+\frac{b}{a} e^{\sigma}\left(e^{a / c}-1\right):=\pi\left(\left|x\left(t_{0}\right)\right|\right) \tag{6.43}
\end{equation*}
$$

Proof. Proof of this lemma is followed directly by solving the differential equation (6.41).

By applying Lemma 6.1 and the comparison principle found in [6] to (6.40), we have

$$
\begin{equation*}
V_{2}(t) \leq V_{2}\left(t_{0}\right) e^{\chi_{1} / \sigma_{1}}+\frac{\chi_{2}}{\chi_{1}}\left(e^{\chi_{1} / \sigma_{1}}-1\right):=\eta_{3}\left(\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|\right) \tag{6.44}
\end{equation*}
$$

Therefore we now rewrite (6.40) as

$$
\begin{equation*}
\dot{V}_{2} \leq-W_{2}\left(x_{e}, v_{e}\right)+\left(\eta_{1} \eta_{3}+\eta_{2}\right) e^{-\sigma_{2}\left(t-t_{0}\right)} \tag{6.45}
\end{equation*}
$$

where $W_{2}\left(x_{e}, v_{e}\right)=p_{1}^{*} x_{e}^{2}+p_{3}^{*} v_{e}^{2}$. Integrating both sides of (6.45) yields

$$
\begin{equation*}
\int_{t_{0}}^{t} W_{2}\left(x_{e}(\tau), v_{e}(\tau)\right) \mathrm{d} \tau \leq V_{2}\left(t_{0}\right)+\int_{t_{0}}^{t}\left(\eta_{1} \eta_{3}+\eta_{2}\right) e^{-\sigma_{2}\left(\tau-t_{0}\right)} \mathrm{d} \tau \tag{6.46}
\end{equation*}
$$

It is seen that the right-hand side of (6.46) exists and is bounded. On the other hand, $W_{2}\left(x_{e}(t), v_{e}(t)\right)$ is uniformly continuous because its time derivative is bounded. Hence from Barbalat's lemma, we have $W_{2}\left(x_{e}(t), v_{e}(t)\right) \rightarrow 0$ as $t \rightarrow \infty$. Therefore $\left(x_{e}(t), v_{e}(t)\right) \rightarrow 0$ as $t \rightarrow \infty$. To prove that $y_{e}(t) \rightarrow 0$ as $t \rightarrow \infty$, applying Lemma 2.6 to the $x_{e}$-dynamic equation of (6.17) yields

$$
\begin{gather*}
\lim _{t \rightarrow \infty}\left(k_{2} r_{d} y_{e}+u_{d} \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}}+\frac{k(t) v_{d} y_{e}}{\varpi_{2}}+r_{1 e}^{d} y_{e}+r_{d} y_{e}+\right. \\
\left.r_{2 e}^{d} y_{e}+\tilde{u}_{e}+\tilde{r}_{e} y_{e}+p_{x}\right)=0 \tag{6.47}
\end{gather*}
$$

Since $u_{d}=0$ and $\left(X_{1 e}(t), x_{e}(t), v_{e}(t), v_{d}(t)\right) \rightarrow 0$ as $t \rightarrow \infty$, (6.47) is equivalent to

$$
\begin{equation*}
\lim _{t \rightarrow \infty}\left(r_{1 e}^{d} y_{e}\right)=0 \tag{6.48}
\end{equation*}
$$

From the expression of $r_{1 e}^{d}$, it is directly shown that (6.48) is equivalent to

$$
\begin{equation*}
\lim _{t \rightarrow \infty}\left(\frac{y_{e}^{2}}{\sqrt{1+\left(1-k^{2}(t)\right) y_{e}^{2}}} \dot{k}(t)\right)=0 \tag{6.49}
\end{equation*}
$$

On the other hand from (6.45), we have

$$
\frac{\mathrm{d}}{\mathrm{~d} t}\left(V_{2}+\sigma_{2}^{-1}\left(\eta_{1} \eta_{3}+\eta_{2}\right) e^{-\sigma_{2}\left(t-t_{0}\right)}\right) \leq 0
$$

which implies that

$$
V_{2}+\sigma_{2}^{-1}\left(\eta_{1} \eta_{3}+\eta_{2}\right) e^{-\sigma_{2}\left(t-t_{0}\right)}
$$

is nonincreasing. Since $V_{2}$ is bounded from below by zero, $V_{2}$ tends to a finite nonnegative constant depending on $\left\|X_{e}\left(t_{0}\right)\right\|$. This implies that the limit of $\left|y_{e}(t)\right|$ exists and is finite, say $l_{y_{e}}$. If $l_{y_{e}}$ were not zero, there would exist a sequence of increasing time instants $\left\{t_{i}\right\}_{i=1}^{\infty}$ with $t_{i} \rightarrow \infty$, such that both of the limits of $\dot{k}\left(t_{i}\right)$ and $\dot{k}\left(t_{i}\right) y_{e}^{2}\left(t_{i}\right)$ were not zero, which is impossible because of (6.49). Hence $l_{y_{e}}$ must be zero. Therefore we conclude from (6.49) that $y_{e}(t) \rightarrow 0$ as $t \rightarrow \infty$. Hence we have proven $\lim _{t \rightarrow \infty} X_{2 e}(t)=0$ for this case. We define the subset of the design parameters that satisfy the conditions (6.27), (6.35), and (6.36) as $\Xi_{1}^{\lambda k}$.

Case C2. In this case, we choose the design parameters $\lambda_{i}$ and $k_{j}$ such that

$$
\begin{equation*}
p_{211}(t) \geq p_{211}^{*}>0 \tag{6.50}
\end{equation*}
$$

Under Assumption 6.1, there exists a choice of $\lambda_{i}$ and $k_{j}$ such that (6.27), (6.35), (6.36), and (6.50) hold. Substituting (6.35), (6.36), and (6.50) into (6.29) yields

$$
\begin{align*}
\dot{V}_{2} \leq & -p_{1}^{*} x_{e}^{2}-\left(p_{211}^{*} r_{d}^{2}-\left|p_{22}(t)\right|-p_{212}(t)\left|r_{d}\right|-p_{23}(t)\right) y_{e}^{2}- \\
& p_{3}^{*} v_{e}^{2}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.51}
\end{align*}
$$

On the other hand, it is noted that

$$
\begin{equation*}
V_{2} \leq \frac{1}{2}\left(x_{e}^{2}+\left(1+k_{4}+k_{4} k_{5}\right) y_{e}^{2}+k_{4} k_{5}\left(1+k_{5}\right) v_{e}^{2}\right) . \tag{6.52}
\end{equation*}
$$

From (6.52) and (6.51), we have

$$
\begin{align*}
\dot{V}_{2} \leq & -2 \min \left(p_{1}^{*}, \frac{p_{211}^{*} r_{d}^{2}-\left|p_{22}(t)\right|-p_{212}(t)\left|r_{d}\right|-p_{23}(t)}{1+k_{4}+k_{4} k_{5}}, \frac{p_{3}^{*}}{k_{4} k_{5}\left(1+k_{5}\right)}\right) \times \\
& V_{2}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.53}
\end{align*}
$$

By applying Lemma 2.6 to (6.53), there exist a positive constant $\sigma_{3}$ independent of initial conditions and a nondecreasing function $\pi_{3}$ of $\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|$ such that

$$
\begin{equation*}
\left\|X_{2 e}(t)\right\| \leq \pi_{3}\left(\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{3}\left(t-t_{0}\right)} \tag{6.54}
\end{equation*}
$$

as long as

$$
\begin{equation*}
\int_{t_{0}}^{t}\left(p_{211}^{*} r_{d}^{2}(\tau)-p_{212}(\tau)\left|r_{d}(\tau)\right|-\left|p_{22}(\tau)\right|-p_{23}(\tau)\right) \mathrm{d} \tau \geq p_{21}^{*}\left(t-t_{0}\right) \tag{6.55}
\end{equation*}
$$

where $p_{21}^{*}>0$. For simplicity of calculation, we can replace the last three terms in the left-hand side of (6.55) by their maximum values. Note that (6.54) implies that $X_{2 e}(t)$ is globally $K$-exponentially stable at the origin but not GES in the sense of Lyapunov, i.e., $\pi_{3}$ in (6.54) must be a linear function of $\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|$. However, it can be shown that the local exponential stability (in the sense of Lyapunov) of the closed loop system (6.17) is guaranteed.

We define the subset of $\lambda_{i}$ and $k_{j}$ satisfying the conditions (6.27), (6.35), (6.36), (6.50), and (6.55) as $\Xi_{2}^{\lambda k}$.

Case C3. Under Assumption 6.2 and $\int_{0}^{\infty}\left|r_{d}(t)\right| \mathrm{d} t \leq \mu_{1}<\infty$, there exists

$$
0 \leq \sigma_{p 21}<\infty
$$

such that

$$
\begin{equation*}
\int_{0}^{\infty}\left|p_{21}(t)\right| \mathrm{d} t \leq \sigma_{p 21} \tag{6.56}
\end{equation*}
$$

In this case, we choose $\lambda_{i}$ and $k_{j}$ such that

$$
\begin{equation*}
p_{22}(t)-p_{23}(t) \geq p_{22}^{*}>0 . \tag{6.57}
\end{equation*}
$$

With (6.36) and (6.57), we rewrite (6.29) as

$$
\begin{gather*}
\dot{V}_{2} \leq-p_{1}^{*} x_{e}^{2}-p_{22}^{*} \frac{y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}-p_{3}^{*} v_{e}^{2}+ \\
\left|p_{21}(t)\right| y_{e}^{2}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} . \tag{6.58}
\end{gather*}
$$

Again by applying Lemma 6.1 to (6.58), we have

$$
\begin{align*}
V_{2}(t) \leq & V_{2}\left(t_{0}\right) e^{2 \sigma_{p 21}} e^{\chi_{1} / \sigma_{1}}+\frac{\chi_{2}}{\chi_{1}} e^{2 \sigma_{p 21}}\left(e^{\chi_{1} / \sigma_{1}}-1\right) \\
& :=\eta_{3}\left(\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|\right) \tag{6.59}
\end{align*}
$$

Substituting (6.59) and (6.52) into (6.58) yields

$$
\begin{equation*}
\dot{V}_{2} \leq-\rho_{4} V_{2}+2\left|p_{21}(t)\right| V_{2}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.60}
\end{equation*}
$$

where

$$
\begin{equation*}
\rho_{4}=2 \min \left(p_{1}^{*}, \frac{p_{22}^{*}}{\sqrt{1+2 \eta_{3}}\left(1+k_{4}+k_{4} k_{5}\right)}, \frac{p_{3}^{*}}{k_{4} k_{5}\left(1+k_{5}\right)}\right) \tag{6.61}
\end{equation*}
$$

From (6.52) and (6.60), it can be shown that

$$
\begin{equation*}
\left\|X_{2 e}(t)\right\| \leq \pi_{4}\left(\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{4}\left(t-t_{0}\right)} \tag{6.62}
\end{equation*}
$$

where

$$
\begin{align*}
& \text { if } \rho_{4} \neq \sigma_{1} \text { then } \\
& \quad \sigma_{4}=0.5 \min \left(\rho_{4}, \sigma_{1}\right), \\
& \quad \pi_{4}(\bullet)=\sqrt{4 \lambda_{4}^{-1} e^{2 \sigma_{p 21}} e^{\chi_{1} / \sigma_{1}}\left(V_{2}\left(t_{0}\right)+\chi_{2}\left|\sigma_{1}-\rho_{4}\right|^{-1}\right)}, \\
& \text { if } \rho_{4}=\sigma_{1} \text { then } \\
& \quad \sigma_{4}=0.5\left(\sigma_{1}-d_{4}\right), \\
& \quad \pi_{4}(\bullet)=\sqrt{4 \lambda_{4}^{-1} e^{2 \sigma_{p 21}} e^{\chi_{1} / \sigma_{1}}\left(V_{2}\left(t_{0}\right)+\chi_{2} \phi_{4}\right),} \tag{6.63}
\end{align*}
$$

with

$$
\begin{gathered}
\lambda_{4}=\min \left(2,1+k_{4}+k_{4} k_{5}^{2}-\sqrt{\left(1+k_{4}-k_{4} k_{5}^{2}\right)^{2}+4\left(k_{4} k_{5}\right)^{2}}\right) \\
0<d_{4}<\sigma_{1} \\
\phi_{4} \geq\left(t-t_{0}\right) e^{-d_{4}\left(t-t_{0}\right)} .
\end{gathered}
$$

Note that $\phi_{4}$ is finite for an arbitrarily small $d_{4}$. It can be seen from (6.59), (6.61) and (6.63) that the rate $\sigma_{4}>0$ depends on the initial conditions $\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|$. Hence the closed loop system (6.17) is GAS but not exponentially stable at the origin. In other words, global asymptotic tracking is achieved in this case. We define the subset of $\lambda_{i}$ and $k_{j}$ satisfying the conditions (6.27), (6.35), (6.36), and (6.57) as $\Xi_{3}^{\lambda k}$.

Case C4. We first show that $V_{2}$ is bounded. Substituting (6.35) and (6.36) into (6.29) yields

$$
\begin{equation*}
\dot{V}_{2} \leq-p_{1}^{*} x_{e}^{2}-p_{3}^{*} v_{e}^{2}+p_{4}(t) y_{e}^{2}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.64}
\end{equation*}
$$

where

$$
\begin{equation*}
p_{4}(t)=\left|p_{21}(t)\right|+\left|p_{22}(t)\right|+p_{23}(t) \tag{6.65}
\end{equation*}
$$

Note that in this case, there exists $0 \leq \sigma_{p 4}<\infty$ such that $\int_{0}^{\infty} p_{4}(t) \mathrm{d} t \leq \sigma_{p 4}$. From (6.19) and (6.65), we have

$$
\begin{equation*}
\dot{V}_{2} \leq 2 p_{4}(t) V_{2}+\left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.66}
\end{equation*}
$$

Applying Lemma 6.1 to (6.66) yields

$$
\begin{align*}
V_{2}(t) \leq & V_{2}\left(t_{0}\right) e^{2 \sigma_{p 4}} e^{\chi_{1} / \sigma_{1}}+\frac{\chi_{2}}{\chi_{1}} e^{2 \sigma_{p 4}}\left(e^{\chi_{1} / \sigma_{1}}-1\right) \\
& :=\eta_{4}\left(\left\|\left(X_{1 e}\left(t_{0}\right), X_{2 e}\left(t_{0}\right)\right)\right\|\right) \tag{6.67}
\end{align*}
$$

Hence $y_{e}$ is also bounded. Substituting this upper bound into (6.64) yields

$$
\begin{equation*}
\dot{V}_{2} \leq-p_{1}^{*} x_{e}^{2}-p_{3}^{*} v_{e}^{2}+2 \eta_{4} p_{4}(t)+\left(\chi_{1} \eta_{4}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{6.68}
\end{equation*}
$$

From (6.68), we have

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t}\left(V_{2}-2 \eta_{4} \int_{0}^{t} p_{4}(\tau) \mathrm{d} \tau+\sigma_{1}^{-1}\left(\chi_{1} \eta_{4}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)}\right) \leq 0 \tag{6.69}
\end{equation*}
$$

which implies that $\left(x_{e}, v_{e}\right) \rightarrow 0$ as $t \rightarrow \infty$ and $\lim _{t \rightarrow \infty} y_{e}=0$. We define the subset of $\lambda_{i}$ and $k_{j}$ satisfying the conditions (6.27), (6.35), and (6.36) as $\Xi_{4}^{\lambda k}$.

### 6.4 Selection of Design Constants

In this section, we show that $\bigcap_{1 \leq l \leq 4} \Xi_{l}^{\lambda k} \neq \emptyset$. From the above section, it can be seen that these design constants must satisfy (6.27), (6.35), (6.36), (6.50), (6.55), and (6.57). From (6.27), we have

$$
\begin{equation*}
k_{1}=\frac{1}{\beta k_{5}}, k_{4}=\frac{1}{\alpha k_{5}-1} . \tag{6.70}
\end{equation*}
$$

Hence we first choose $k_{5}>1 / \alpha$ and set $k_{2}=\delta k_{*}$ with $\delta$ being a small positive constant, say $\delta \leq \max \left(u_{d}^{\max }, \dot{u}_{d}^{\max }, r_{d}^{\max }\right)$, with $\dot{u}_{d}^{\max }$ being the maximum value of $\left|\dot{u}_{d}\right|$. Next we replace $k$ in all negative terms of $p_{i}(t), i=1,21,22,23,3$ by $k_{*}$. It is now observed that the negative parts of $p_{i}(t)$ have $k_{*}$ as a factor. The mass including added mass in the sway dynamics, $m_{22}$, is always larger than that in the surge dynamics, $m_{11}$, for surface ships, i.e., $\beta<1$. Therefore we can always find positive constants $k_{*}$ and $k_{5}$ such that the conditions (6.36), (6.50), (6.55), and (6.57) hold with some small $\rho_{i}>0, i=1,2,3$, and $\lambda_{i}$ being chosen based on (6.35). The value of constant $k_{*}$ should be reduced if the reference yaw, $r_{d}$, and surge, $u_{d}$, velocities and surge acceleration, $\dot{u}_{d}$, are large. This physically means that the distance from the ship to the point it aims to track should be increased if the velocities and surge acceleration are large, otherwise the ship will miss that point. In fact, setting $\lambda_{i}=0$ and picking $k_{2}>0$ results in the tracking controllers proposed in [19,71,106] with a restrictive assumption of the yaw reference velocity satisfying a persistently exciting condition. It can also be seen that setting $\lambda_{2}=0$ yields the tracking controller in the preceding chapter, which does not require the yaw reference velocity to be persistently exciting. Furthermore, small $\alpha$ automatically results in large $k_{5}$ and small $k_{1}$, see (6.70). This can be physically interpreted as follows: If the damping in the sway dynamics is small, the control gain in the surge dynamics should be small, otherwise
the ship will slide in the sway direction. Due to complicated expressions of $p_{i}(t)$, we provide some guidelines to choose the design constants rather than present their extremely complex explicit expressions.

1. Pick positive constants $\lambda_{i}$ such that the constant $k_{*}$ is small, say $k_{*} \ll \max \left(u_{d}^{\max }\right.$, $\left.r_{d}^{\max }, \dot{u}_{d}^{\max }\right), \delta$ such that $\delta \leq \max \left(u_{d}^{\max }, r_{d}^{\max }\right)$ and set $k_{2}=\delta k_{*}, k_{5}=1 / \alpha+\varepsilon$ with $\varepsilon$ being a small positive constant.
2. Substitute (6.70) into $p_{i}(t)$ and slowly increase $\lambda_{i}$ and $k_{5}$ until all the conditions (6.36), (6.50), (6.55), and (6.57) hold with some positive constants $p_{1}^{*}, p_{211}^{*}, p_{21}^{*}$, $p_{22}^{*}$, and $p_{3}^{*}$.

Remark 6.5. When the reference surge and sway velocities and surge acceleration are very large, the above procedure will result in very small control gains $k, k_{1}$, and $k_{2}$, which will give slow convergence. Hence it is suggested that for large $u_{d}, r_{d}$, and $\dot{u}_{d}$, the control gains $k, k_{1}$, and $k_{2}$ should be chosen such that either conditions ((6.27), (6.35), (6.36)) or ((6.27), (6.35), (6.36), (6.50), (6.55)), or ((6.27), (6.35), (6.36), (6.57)) hold to improve the convergence. The trade-off is that the control gains will be different for different reference trajectories.

### 6.5 Sensitivity Analysis

The control law (6.16) has been designed under the assumptions that the system parameters are precisely known and there are no environmental disturbances. Indeed, these assumptions are unrealistic in practice. The aim of this section is to discuss the sensitivity of our proposed controller in relation to the inaccurate knowledge of the ship parameters. A discussion related to environmental disturbances can be carried out similarly.

The control law (6.16) can be easily modified to account for the inaccurate knowledge of the ship parameters entering (6.15), by adding some robustifying terms or can be changed into an adaptive version such that the $\left(z_{e}, \tilde{u}_{e}, \tilde{r}_{e}\right)$-subsystem is globally exponentially/asymptotically stable at the origin. However, this control law cannot overcome the imprecise knowledge of the ship parameters, which enter the last equation of (6.2) due to the fact that the reference trajectory is generated by the virtual ship model (6.2). Therefore, we only focus on this issue to simplify the analysis, i.e., we are interested in the question: Assuming that the control law (6.16) has been modified as above, how do the ship parameters affect the tracking errors when the virtual ship model (6.2) is formed by the nominal values of the ship parameters?

Defining

$$
\begin{equation*}
\Delta \alpha=\alpha-\alpha_{c}, \Delta \beta=\beta-\beta_{c}, \tag{6.71}
\end{equation*}
$$

where $\alpha_{c}$ and $\beta_{c}$ are the known parts of $\alpha$ and $\beta$, we rewrite the $v_{e}$-dynamics of (6.8) as

$$
\begin{equation*}
\dot{v}_{e}=-\alpha v_{e}-\beta u_{e} r_{d}-\beta\left(u_{e}+u_{d}\right) r_{e}+\Delta \alpha v_{d}+\Delta \beta u_{d} r_{d} \tag{6.72}
\end{equation*}
$$

Notice that we cannot choose the design constants such that $1-k_{4}\left(\alpha k_{5}-1\right)=0$ as in the previous sections to cancel the term $y_{e} v_{e}\left(1-k_{4}\left(\alpha k_{5}-1\right)\right)$, see (6.26), due to the inaccurately known ship parameters. If one replaces $\alpha$ in the above expression by $\alpha_{c}+\Delta \alpha$, the term $y_{e} v_{e}\left(1-k_{4}\left(\alpha_{c} k_{5}-1\right)\right)$ will be canceled. However the remaining term $\Delta \alpha k_{4} k_{5} y_{e} v_{e}$ cannot be canceled in the case of C3 because this term does not contain any reference velocities. Therefore, we replace $k_{5}$ in the Lyapunov function (6.19) by $k_{5}^{*} / \alpha$ where $k_{5}^{*}$ is a positive constant to be determined, i.e.,

$$
\begin{equation*}
V_{2}^{*}\left(X_{2 e}\right)=\frac{1}{2}\left(x_{e}^{2}+y_{e}^{2}+k_{4}\left(y_{e}+k_{5}^{*} \alpha^{-1} v_{e}\right)^{2}\right) . \tag{6.73}
\end{equation*}
$$

The condition (6.27) is replaced by

$$
\begin{equation*}
1-k_{4}\left(k_{5}^{*}-1\right)=0, \beta_{c} \alpha_{c}^{-1} k_{1} k_{5}^{*}-1=0 \tag{6.74}
\end{equation*}
$$

Note that the term $\Phi$ is now written as

$$
\begin{equation*}
\Phi=\Delta \beta^{*} k_{1} k_{4} \frac{k_{5}^{*}}{\alpha}\left(y_{e}+\frac{k_{5}^{*}}{\alpha} v_{e}\right)\left(r_{1 e}^{d}+r_{d}\right) x_{e}, \tag{6.75}
\end{equation*}
$$

where $\Delta \beta^{*}=\beta / \alpha-\beta_{c} / \alpha_{c}$. With (6.72) and (6.75), we have

$$
\begin{align*}
\dot{V}_{2}^{*} \leq & -q_{1}(t) x_{e}^{2}-q_{21}(t) y_{e}^{2}-\frac{q_{22}(t) y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}+\frac{q_{23}(t) y_{e}^{2}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}-q_{3}(t) v_{e}^{2}+ \\
& \left(\chi_{1} V_{2}+\chi_{2}\right) e^{-\sigma_{1}\left(t-t_{0}\right)}+\left(\frac{1}{4 \phi_{1}} \frac{k_{4} k_{5}^{*}}{\alpha}+\frac{1}{4 \phi_{2}} \frac{k_{4} k_{5}^{* 2}}{\alpha^{2}}\right)\left|\Delta \alpha v_{d}+\Delta \beta u_{d} r_{d}\right| \tag{6.76}
\end{align*}
$$

where

$$
\begin{align*}
q_{1}(t)= & \hat{p}_{1}(t)-\frac{k_{1} k_{4} k_{5}^{*}\left|\Delta \beta^{*}\right|}{\alpha}\left(\frac{1}{4 \rho_{1}}+\frac{k_{5}^{*}}{\alpha}+\frac{|k(t)|}{1-|k(t)|}\left(\frac{1+4 \rho_{1}\left|r_{d}\right|}{4 \rho_{1} \sqrt{1-k^{2}(t)}}+\right.\right. \\
& \left.k_{1}+k_{2}\left|r_{d}\right|+k^{2}(t)\left|u_{d}\right|+\left|k(t) v_{d}\right|\right)+\frac{|k(t)| k_{5}^{*}}{4 \rho_{1} \alpha(1-|k(t)|)} \times \\
& \left.\left(k_{1}+\left|r_{d}\right|+k_{2}\left|r_{d}\right|+k^{2}(t)\left|u_{d}\right|+\left|k(t) v_{d}\right|+\left|k(t) u_{d}\right|\right)\right)  \tag{6.77}\\
q_{21}(t)= & \hat{p}_{21}(t)-\left(k_{4} \rho_{1}\left|\Delta \beta^{*}\right| r_{d}^{2}+\frac{k_{4} k_{5}^{*}}{\alpha}\left|\Delta \alpha v_{d}+\Delta \beta u_{d} r_{d}\right|\right)  \tag{6.78}\\
q_{22}(t)= & \hat{p}_{22}(t)-\frac{k_{1} k_{4} k_{5}^{*}\left|\Delta \beta^{*}\right|}{\alpha(1-|k(t)|)}\left(\frac{|\dot{k}(t)|+|k(t)| \rho_{1}}{\sqrt{1-k^{2}(t)}}+\frac{k_{5}^{*}}{4 \rho_{2} \alpha} \times\right. \\
& \left.\frac{|\dot{k}(t)|}{\sqrt{1-k^{2}(t)}}+2\left(k^{2}(t)\left|u_{d}\right|+\left|k^{3}(t) v_{d}\right|\right)\left(1+\frac{k_{5}^{*}}{4 \rho_{2} \alpha}\right)\right)  \tag{6.79}\\
q_{23}(t)= & \hat{p}_{23}(t) \tag{6.80}
\end{align*}
$$

$$
\begin{align*}
q_{3}(t)= & \hat{p}_{3}(t)-\left(\frac{k_{4} k_{5}^{* 2}}{\alpha^{2}}\left|\Delta \alpha v_{d}+\Delta \beta u_{d} r_{d}\right| \phi_{2}+\frac{k_{4} k_{5}^{*}\left|\Delta \beta^{*}\right| \rho_{1}}{\alpha}+\right. \\
& \frac{2|k(t)| k_{4}\left|\Delta \beta^{*}\right| \rho_{1}}{(1-|k(t)|) \sqrt{1-k^{2}(t)}}+\frac{k_{1} k_{4} k_{5}^{* 2}\left|\Delta \beta^{*}\right|}{\alpha^{2}(1-|k(t)|)}\left(\frac{|\dot{k}(t)| \rho_{2}}{\sqrt{1-k^{2}(t)}}+\right. \\
& 2|k(t)|+\left(\rho_{1}+\rho_{2}\right)\left|k^{3}(t)\right|\left(\left|v_{d}\right|+\left|u_{d}\right|\right)+\rho_{1}|k(t)|\left(k_{1}+\right. \\
& \left.\left.\left.\left.k_{2}+\left|r_{d}\right|\right)+\rho_{1} k^{2}(t)\left(\left|v_{d}\right|+\left|u_{d}\right|\right)\right)\right)\right) \tag{6.81}
\end{align*}
$$

$\hat{p}_{1}, \hat{p}_{21}, \hat{p}_{22}, \hat{p}_{23}$, and $\hat{p}_{3}$ are $p_{1}, p_{21}, p_{22}, p_{23}$, and $p_{3}$ respectively, with $k_{5}$ being replaced by $k_{5}^{*} / \alpha$, and $\phi_{1}$ and $\phi_{2}$ are some positive constants. From (6.76), the stability analysis of the ( $x_{e}, y_{e}, v_{e}$ )-subsystem can be carried out using the same arguments as in the previous section. It can be shown that there always exist the design constants $k_{1}, k_{3}, k_{4}$, and $k_{5}^{*}$ such that the required conditions hold if $\Delta \alpha$ and $\Delta \beta$ are not too large, say $\Delta \alpha \ll \alpha, \Delta \beta \ll \beta$ and $\Delta \beta^{*} \ll \min (\beta, \alpha)$. It is noted that when Assumption 6.2 holds with either cases $\mathrm{C} 1, \mathrm{C} 3$, or C 4 , the $\left(x_{e}, y_{e}, v_{e}\right)$ subsystem is still GAS since the term $\left|\Delta \alpha v_{d}+\Delta \beta u_{d} r_{d}\right|$ in (6.76) globally exponentially/asymptotically vanishes at the origin. However, only global practical stability can be achieved in the case of C 2 . This coincides with the ship control practice in the sense that a big ship, in general, cannot accurately track the reference trajectory generated by a small ship because of underactuated configuration in the sway dynamics.

Remark 6.6. Since the inaccurate knowledge of the ship parameters directly affects the tracking errors when Assumption 6.2 holds with case C 2 and the $v_{e}$-dynamics is globally ISS with respect to $u_{e}$ and $r_{e}$ as inputs, see (6.8), we can also treat the $v_{e}$-dynamics as unmodeled dynamics. In this case, one can apply the methodology proposed in this chapter and the work on controlling nonlinear systems with unmodeled dynamics [112].

### 6.6 Simulations

The same monohull ship for the trajectory tracking simulations in Section 5.4 is used in this section to validate the control law (6.16) where all the four conditions in Assumption 6.2 are considered. The reference trajectory is generated by the virtual ship with $\left[x_{d}(0), y_{d}(0), \psi_{d}(0), u_{d}(0), v_{d}(0), r_{d}(0)\right]=\left[0 \mathrm{~m}, 0 \mathrm{~m}, 0 \mathrm{~ms}^{-1}, 0 \mathrm{rads}^{-1}\right]$ and

- for the case $\mathrm{C} 1: \tau_{u d}=-m_{22} v_{d} r_{d}, \tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} r_{d}$, i.e., the reference trajectory is a point at the origin,
- for the case C2: $\tau_{u d}=-m_{22} v_{d} r_{d}+d_{11}, \tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} r_{d}+0.01 d_{33}$, i.e., the reference trajectory is a circle with a radius of 200 m and the reference velocity $u_{d}$ is $1 \mathrm{~ms}^{-1}$,
- for the case C3: $\tau_{u d}=-m_{22} v_{d} r_{d}+d_{11}, \tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} r_{d}$, i.e., the reference trajectory is a straight line and the reference velocity $u_{d}$ is $1 \mathrm{~ms}^{-1}$,
- for the case C4: $\tau_{u d}=-m_{22} v_{d} r_{d}+0.1 d_{11} e^{-0.3 t}, \tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} r_{d}+$ $0.1 d_{33} e^{-5 t}$, i.e., the reference trajectory is a path approaching to the origin.

We first choose $k_{3}=0.8, \lambda_{3}=3.5$, and $c_{1}=c_{2}=5$. Note that these constants can be any arbitrarily positive real numbers. The bigger these constants, the faster the transient response will be, but the control effort becomes larger. Next, we initialize $k_{*}=0.01$ and $k_{5}=1.8$, then increase these values until all of the conditions (6.36), (6.50), (6.55), and (6.57) hold. We get $k_{*}=0.02$ and $k_{5}=2.14$, which result in $k_{1}=0.85, k_{2}=0.54, k_{4}=6.4, \lambda_{1}=0.015$, and $\lambda_{2}=0.005$, see ( 6.35 ) and (6.70). With the above design constant selection, one can verify that all conditions (6.27), (6.35), (6.36), (6.50), (6.55), and (6.57) hold with $p_{1}^{*}=0.47, p_{211}^{*}=0.15$, $p_{21}^{*}=0.07, p_{22}^{*}=0.13, p_{3}^{*}=0.16, \rho_{1}=0.2, \rho_{2}=2$, and $\rho_{3}=0.5$. The initial conditions of the ship are chosen as $[x(0), y(0), \psi(0), u(0), v(0), r(0)]=$ $\left[5 \mathrm{~m},-40 \mathrm{~m}, 0.7 \mathrm{rad}, 0 \mathrm{~ms}^{-1}, 0.5 \mathrm{~ms}^{-1}, 0 \mathrm{rads}^{-1}\right]$.

Figures 6.1-6.4 plot the tracking errors, control inputs, and tracking trajectory in the $(x, y)$-plane for cases $\mathrm{C} 1-\mathrm{C} 4$. It can be seen from Figures 6.1a-6.4a that the tracking errors asymptotically ( $K$-exponentially for the case of C 2 ) converge to zero as proven in Theorem 6.1. It is noted that the transient response for cases C1 and C4 is much slower and is more oscillatory than cases C2 and C3 because the reference velocities $u_{d}$ and $r_{d}$ are zero or approach zero. Hence the term $\lambda_{1} u_{d}$ does not help in the control gain $k(t)$, see (6.6). The observation about the transient response for cases C 1 and C 4 should not be surprising due to the sinusoidal signal introduced in the control gain $k(t)$, and is also true for any time-varying controllers that stabilize an underactuated system at the origin [15].

To test the robustness of our proposed controller with respect to the ship parameters and small environmental disturbances induced by wave, wind, and ocean current, we simulate the control law (6.16) with the same design constants selected as above. Figures 6.5-6.8 plot the tracking errors, control inputs, and tracking trajectory in ( $x, y$ )-plane for cases C1-C4 with $8 \%$ variation in all of the system parameters in the sense that the ship parameters are taken as $0.92 m_{11}, 1.08 m_{22}, 0.92 m_{33}$, $0.92 d_{11}, 1.08 d_{22}, 1.08 d_{33}$, and with the environmental disturbances acting on the surge, sway and yaw dynamics as $\tau_{w u}(t)=0.1 m_{11} \operatorname{rand}(\cdot), \tau_{w v}(t)=0.1 m_{22} \operatorname{rand}(\cdot)$, $\tau_{w v}(t)=0.1 m_{33} \operatorname{rand}(\cdot)$, where $\operatorname{rand}(\cdot)$ is the zero-mean random noise with the uniform distribution on the interval $[-0.50 .5]$. The above choice of the variation in the ship parameters is only for illustrating robustness properties of the controller designed in this chapter with respect to the system parameter uncertainties. Note that this variation directly affects the tracking errors in Case C2 since the reference sway velocity has to be generated by the virtual ship. It is worth mentioning that under arbitrarily small nonvanishing environmental disturbances, the discontinuous stabilization controller proposed in [55] results in an unstable closed loop system in the sense that the closed loop trajectory goes to infinity exponentially fast.

With the above disturbances, the last three equations of (6.1) are of the form

$$
\dot{u}=\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}+\frac{1}{m_{11}} \tau_{w u}(t),
$$

$$
\begin{align*}
& \dot{v}=-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v+\frac{1}{m_{22}} \tau_{w v}(t), \\
& \dot{r}=\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t) . \tag{6.82}
\end{align*}
$$



Figure 6.1 Simulation results without disturbances for case $\mathrm{C} 1:$ a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]:$ dash-dotted line); b. Control inputs ( $\tau_{u}[\mathrm{~N}]$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line); $\mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line)

### 6.7 Conclusions

Restrictive assumptions on reference velocities required in the literature and Chapter 5 have been removed thanks to the coordinate transformation (6.6). The proposed methodology in this chapter can be readily extended to design a single controller for simultaneous stabilization and tracking of underactuated underwater vehicles. The work presented in this chapter is based on [105, 113-115].


Figure 6.2 Simulation results without disturbances for case C 2 : a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; $\mathbf{b}$. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line $)$; $\mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$


Figure 6.3 Simulation results without disturbances for case C 3 : a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; $\mathbf{b}$. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line $)$; $\mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$


Figure 6.4 Simulation results without disturbances for case C 4 : a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $) ; \mathbf{b}$. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line $) ; \mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$


Figure 6.5 Simulation results with disturbances for case C 1 : a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; b. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line $) ; \mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$


Figure 6.6 Simulation results with disturbances for case C 2 : a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; $\mathbf{b}$. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line $)$; $\mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$


Figure 6.7 Simulation results with disturbances for case C3: a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; $\mathbf{b}$. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line $)$; $\mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$


Figure 6.8 Simulation results with disturbances for case C 4 : a. Convergence of tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]$ : dot-dotted line, $\psi-\psi_{d}[\mathrm{rad}]$ : dash-dotted line $)$; $\mathbf{b}$. Control inputs $\left(\tau_{u}[\mathrm{~N}]\right.$ : solid line, $\tau_{r}[\mathrm{Nm}]$ : dashed line $) ; \mathbf{c}$. Tracking trajectory in the $(x, y)$-plane $((x, y)$ : solid line, $\left(x_{d}, y_{d}\right)$ : dashed line $)$

# Chapter 7 <br> Partial-state and Output Feedback Trajectory-tracking Control of Underactuated Ships 

Global partial-state feedback and output feedback control schemes are discussed in this chapter for tracking control of an underactuated surface ship without sway force. For the case of partial-state feedback, we do not require measurements of the ship sway and surge velocities, while for the case of output feedback, none of the ship velocities are required for feedback. The reference trajectory to be tracked can be a curve including a straight line. Global nonlinear coordinate changes are introduced to transform the ship dynamics to a system affine in the ship velocities to design observers to globally exponentially estimate unmeasured velocities. These observers plus the techniques in the previous chapter facilitate the development of controllers in the following sections.

### 7.1 Control Objective

For the convenience of the reader, the mathematical model of the underactuated ship moving in surge, sway and yaw, see Section 3.4.1.2 (i.e., (3.45) and (3.46)), is once again presented:

$$
\begin{align*}
& \dot{\eta}=J(\eta) v, \\
& M \dot{v}=-C(v) v-D v+\tau \tag{7.1}
\end{align*}
$$

where the matrices $\boldsymbol{J}(\boldsymbol{\eta}), \boldsymbol{M}, \boldsymbol{C}(\boldsymbol{v})$, and $\boldsymbol{D}$ are given by

$$
\begin{align*}
& \boldsymbol{J}(\boldsymbol{\eta})=\left[\begin{array}{ccc}
\cos (\psi) & -\sin (\psi) & 0 \\
\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right], \boldsymbol{M}=\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & 0 \\
0 & 0 & m_{33}
\end{array}\right],  \tag{7.2}\\
& \boldsymbol{C}(\boldsymbol{v})=\left[\begin{array}{ccc}
0 & 0 & -m_{22} v \\
0 & 0 & m_{11} u \\
m_{22} v & -m_{11} u & 0
\end{array}\right], \boldsymbol{D}=\left[\begin{array}{ccc}
d_{11} & 0 & 0 \\
0 & d_{22} & 0 \\
0 & 0 & d_{33}
\end{array}\right],
\end{align*}
$$

with

$$
\begin{align*}
& m_{11}=m-X_{\dot{u}}, m_{22}=m-Y_{\dot{v}}, m_{33}=I_{z}-N_{\dot{r}} \\
& d_{11}=-X_{u}, d_{22}=-Y_{v}, d_{33}=-N_{r} . \tag{7.3}
\end{align*}
$$

The propulsion force and moment vector $\boldsymbol{\tau}$ is still given by (3.43), i.e.,

$$
\boldsymbol{\tau}=\left[\begin{array}{c}
\tau_{u}  \tag{7.4}\\
0 \\
\tau_{r}
\end{array}\right]
$$

We assume that the reference trajectory is generated by a virtual ship as follows:

$$
\begin{align*}
& \dot{\boldsymbol{\eta}}_{\boldsymbol{d}}=\boldsymbol{J}\left(\boldsymbol{\eta}_{\boldsymbol{d}}\right) \boldsymbol{v}_{\boldsymbol{d}} \\
& \dot{v}_{d}=-\frac{m_{11}}{m_{22}} u_{d} r_{d}-\frac{d_{22}}{m_{22}} v_{d} \tag{7.5}
\end{align*}
$$

where all the variables have similar meanings as in system (7.1). It is noted that we do not require the reference surge and yaw velocities to be generated by the virtual ship. In this chapter we impose the following assumptions on the reference model (7.5):

Assumption 7.1. The reference signals $x_{d}, y_{d}, u_{d}, r_{d}, \dot{u}_{d}, \ddot{u}_{d}$ and $\dot{r}_{d}$ are bounded. There exists a strictly positive constant $u_{d \min }$, such that $\left|u_{d}(t)\right| \geq u_{d \min }, \forall t \geq 0$. The reference sway velocity satisfies $\left|v_{d}(t)\right|<\left|u_{d}(t)\right|, \forall t \geq 0$.

Assumption 7.2. One of the following conditions holds:
C1. The surge and sway displacements $(x, y)$, yaw angle, $\psi$, and yaw velocity, $r$, are measurable but the surge and sway velocities, $u$ and $v$, are not.
C2. The surge and sway displacements $(x, y)$ and yaw angle $\psi$ are measurable but none of the velocities $u$, $v$, and $r$ are measurable.

Remark 7.1. Condition $\left|u_{d}(t)\right| \geq u_{d \text { min }}, \forall t \geq 0$ implies that the reference surge velocity is always nonzero but can be either positive or negative. This means that we consider both forward and backward tracking. From a practical control viewpoint of surface ships, the condition $\left|u_{d}(t)\right| \geq u_{d \text { min }}, \forall t \geq 0$ is much less restrictive than a persistently exciting condition on the yaw reference velocity in the references $[19,71,106]$ in the sense that tracking of a straight line is included. Surface ships are often equipped with a rudder or a pair of propellers or water jets. The yaw moment to steer the ship is generated by changing the rudder angle or the speed of each propeller or water jet. These facts imply that the tracking control is carried out only when the surge speed is nonzero. The condition $\left|v_{d}(t)\right|<\left|u_{d}(t)\right|, \forall t \geq 0$ implies that the ship cannot track a circle with arbitrarily small radius due to the ship's high inertia and underactuation in the sway direction.

Remark 7.2. Assumption 7.2.C1 means that we need to solve a partial-state feedback control problem. Although the yaw velocity is measurable, there is still a cross
term $u v$ in the yaw velocity dynamics, see the last equation of (7.1). Assumption 7.2. C2 implies that we need to solve an output feedback control problem. Indeed, Assumption 7.2.C2 covers Assumption 7.2.C1. We will, however, show later that design of an output feedback tracking controller is much more involved than that of the partial-state feedback controller.

### 7.2 Partial-state Feedback

### 7.2.1 Observer Design

As discussed above, since the term $\boldsymbol{C}(\boldsymbol{v}) \boldsymbol{v}$ in (7.1) causes difficulties in observer design, we first remove this term by proposing the following coordinate transformation:

$$
\begin{equation*}
\boldsymbol{X}=e^{\boldsymbol{Q}(t, \eta) \boldsymbol{v}} \tag{7.6}
\end{equation*}
$$

where $\boldsymbol{Q}(t, \boldsymbol{\eta})$ is a matrix to be determined. Differentiating both sides of (7.6) along the solutions of the second equation of (7.1) yields

$$
\begin{equation*}
\dot{\boldsymbol{X}}=e^{\boldsymbol{Q}(t, \boldsymbol{\eta})}\left[\dot{\boldsymbol{Q}}(t, \boldsymbol{\eta})-\boldsymbol{M}^{-\mathbf{1}} \boldsymbol{C}(\boldsymbol{v})\right] \boldsymbol{v}+e^{\boldsymbol{Q}(t, \boldsymbol{\eta})}(-\boldsymbol{D} \boldsymbol{v}+\boldsymbol{\tau}) . \tag{7.7}
\end{equation*}
$$

It can be seen that the square bracket on the right-hand side of (7.7) is zero, if the matrix $\boldsymbol{Q}(t, \boldsymbol{\eta})$ is chosen such that

$$
\begin{equation*}
\dot{\boldsymbol{Q}}(t, \eta)-M^{-1} C(v)=0 . \tag{7.8}
\end{equation*}
$$

By using the first equation of (7.1), a particular solution of (7.8) is

$$
\boldsymbol{Q}(t, \boldsymbol{\eta})=\boldsymbol{M}^{-1}\left[\begin{array}{lll}
0 & 0 & -m_{22} q_{13}  \tag{7.9}\\
0 & 0 & m_{11} q_{23} \\
m_{22} q_{13} & -m_{11} q_{23} & 0
\end{array}\right]
$$

with

$$
\begin{align*}
q_{13} & =y \cos (\psi)-x \sin (\psi)+p_{13}(t), \\
q_{23} & =y \sin (\psi)+x \cos (\psi)+p_{23}(t), \\
\dot{p}_{13} & =y \sin (\psi) r+x \cos (\psi) r,  \tag{7.10}\\
\dot{p}_{23} & =-y \cos (\psi) r+x \sin (\psi) r .
\end{align*}
$$

Note that the matrix $\boldsymbol{Q}(t, \boldsymbol{\eta})$ contains only the available signals since we assume that $x, y, \psi$, and $r$ are measurable. Using the Taylor expansion the matrix $e^{\boldsymbol{Q}(t, \boldsymbol{\eta})}$ can be expanded as

$$
\left.\begin{array}{rl}
e^{\mathbf{Q}(t, \boldsymbol{\eta})}= & {\left[\begin{array}{c}
\frac{1}{2 a_{5}}\left(2 a_{2} a_{4}+a_{1} a_{3}\left(e^{-\sqrt{a_{5}}}+e^{\sqrt{a_{5}}}\right)\right) \\
\frac{1}{2 a_{5}}\left(a_{2} a_{3}\left(e^{-\sqrt{a_{5}}}+e^{\sqrt{a_{5}}}-2\right)\right) \\
\frac{1}{2 \sqrt{a_{5}}}\left(a_{3}\left(-e^{-\sqrt{a_{5}}}+e^{\sqrt{a_{5}}}\right)\right)
\end{array}\right.} \\
& \frac{1}{2 a_{5}}\left(a_{1} a_{4}\left(e^{-\sqrt{a_{5}}}+e^{\sqrt{a_{5}}}-2\right)\right) \\
& \frac{1}{2 a_{5}}\left(2 a_{1} a_{3}+a_{2} a_{4}\left(e^{-\sqrt{a_{5}}}+e^{\sqrt{a_{5}}}\right)\right) \frac{1}{2 \sqrt{a_{5}}}\left(a_{1}\left(-e^{-\sqrt{a_{5}}}+e^{\sqrt{a_{5}}}\right)\right) \\
\left.\frac{1}{2 \sqrt{a_{5}}}\left(a_{4}^{-\sqrt{a_{5}}}+e^{\sqrt{a_{5}}}\right)\right)
\end{array}\right],
$$

where

$$
\begin{align*}
& a_{1}=-\frac{m_{22}}{m_{11}} q_{13}, \\
& a_{2}=\frac{m_{11}}{m_{22}} q_{23}, \\
& a_{3}=\frac{m_{22}}{m_{33}} q_{13},  \tag{7.12}\\
& a_{4}=-\frac{m_{11}}{m_{33}} q_{23}, \\
& a_{5}=a_{1} a_{3}+a_{2} a_{4} .
\end{align*}
$$

Similarly, $e^{-\boldsymbol{Q}(t, \eta)}$ has the same form as (7.11), but all of the terms $a_{i}, 1 \leq i \leq$ 4 have an opposite sign to those defined in (7.12). From (7.11) and noticing that $a_{5} \leq 0, \forall\left(q_{13}, q_{23}\right) \in \mathbb{R}^{2}$, it is easily seen that all elements of $e^{\boldsymbol{Q}(t, \boldsymbol{\eta})}$ or $e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}$ are bounded by some constants, which depend only on the ship parameters, $m_{11}$, $m_{22}$, and $m_{33}$. Using the coordinate change (7.6), the ship system (7.1) is written in $(\boldsymbol{\eta}, \boldsymbol{X})$ coordinates as

$$
\begin{align*}
& \dot{\eta}=J(\eta) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{X} \\
& \dot{\boldsymbol{X}}=-e^{\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{M}^{-\mathbf{1}} D e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{X}+e^{\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{M}^{-\mathbf{1}} \tau \tag{7.13}
\end{align*}
$$

The system (7.13) has a very nice structure, namely linear in the unmeasured states. Of course, a reduced-order observer can be designed but it is often noise-sensitive. Here we use the following nonlinear observer to construct the unmeasured surge and sway velocities:

$$
\begin{align*}
\dot{\hat{\boldsymbol{\eta}}}= & \boldsymbol{J}(\boldsymbol{\eta}) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \hat{\boldsymbol{X}}+\boldsymbol{K}_{\mathbf{0}}(\boldsymbol{\eta}-\hat{\boldsymbol{\eta}}), \\
\dot{\hat{\boldsymbol{X}}}= & -e^{\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{D} e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \hat{\boldsymbol{X}}+e^{\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{\tau}+ \\
& \left(\boldsymbol{J}(\boldsymbol{\eta}) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}\right)^{T}(\boldsymbol{\eta}-\hat{\boldsymbol{\eta}}), \tag{7.14}
\end{align*}
$$

where $\hat{\boldsymbol{\eta}}$ and $\hat{\boldsymbol{X}}$ are the estimates of $\boldsymbol{\eta}$ and $\boldsymbol{X}$, respectively; $\boldsymbol{K}_{\mathbf{0}}=\boldsymbol{K}_{\mathbf{0}}^{\boldsymbol{T}}$ is the positive diagonal observer gain matrix. From (7.13) and (7.14), we have

$$
\begin{align*}
& \dot{\tilde{\boldsymbol{\eta}}}=\boldsymbol{J}(\boldsymbol{\eta}) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \tilde{\boldsymbol{X}}-\boldsymbol{K}_{\mathbf{0}} \tilde{\boldsymbol{\eta}} \\
& \dot{\tilde{\boldsymbol{X}}}=-e^{\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{D} e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \tilde{\boldsymbol{X}}-\left(\boldsymbol{J}(\boldsymbol{\eta}) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}\right)^{T} \tilde{\boldsymbol{\eta}} \tag{7.15}
\end{align*}
$$

where $\tilde{\eta}=\eta-\hat{\eta}$ and $\tilde{X}=\boldsymbol{X}-\hat{X}$. From (7.15), one can show that

$$
\begin{equation*}
\|(\tilde{\boldsymbol{\eta}}(t), \tilde{\boldsymbol{X}}(t))\| \leq\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\| e^{-\sigma_{0}\left(t-t_{0}\right)}, \forall 0 \leq t_{0} \leq t<\infty \tag{7.16}
\end{equation*}
$$

with $\sigma_{0}=\min \left(\lambda_{\min }\left(\boldsymbol{K}_{\mathbf{0}}\right), \lambda_{\min }\left(\boldsymbol{M}^{\mathbf{1}} \boldsymbol{D}\right)\right)$, which in turn implies that (7.14) is a global exponential observer of (7.13). We define $\hat{\boldsymbol{v}}=[\hat{u}, \hat{v}, \hat{r}]^{T}$ being an estimate of the velocity vector $v$ as

$$
\begin{equation*}
\hat{\boldsymbol{v}}=e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \hat{\boldsymbol{X}} . \tag{7.17}
\end{equation*}
$$

Using (7.17) and (7.14), we rewrite (7.1) in ( $\boldsymbol{\eta}, \hat{\boldsymbol{v}})$ coordinates as

$$
\left[\begin{array}{l}
\dot{\eta}  \tag{7.18}\\
\dot{\hat{\boldsymbol{v}}}
\end{array}\right]=\left[\begin{array}{l}
\boldsymbol{J}(\boldsymbol{\eta}) \hat{\boldsymbol{v}} \\
-\boldsymbol{M}^{-\mathbf{1}} \boldsymbol{C}(\hat{\boldsymbol{v}}) \hat{\boldsymbol{v}}-\boldsymbol{M}^{-\mathbf{1}} \boldsymbol{D} \hat{\boldsymbol{v}}+\boldsymbol{M}^{-\mathbf{1}} \boldsymbol{\tau}
\end{array}\right]+\left[\begin{array}{ll}
0_{3 \times 3} & H_{12} \\
H_{21} & H_{22}
\end{array}\right]\left[\begin{array}{l}
\tilde{\boldsymbol{\eta}} \\
\tilde{\boldsymbol{X}}
\end{array}\right],
$$

where

$$
\begin{align*}
H_{12}= & \boldsymbol{J}(\boldsymbol{\eta}) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}, \\
H_{21}= & e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}\left(\boldsymbol{J}(\boldsymbol{\eta}) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}\right)^{T}, \\
H_{22}= & \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{C}\left(\hat{\boldsymbol{v}}+e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \tilde{\boldsymbol{X}}\right) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}-\boldsymbol{M}^{-\mathbf{1}} \boldsymbol{C}^{*}(\hat{\boldsymbol{v}}) e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})}+ \\
& e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{C}\left(\hat{\boldsymbol{v}}+e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \tilde{\boldsymbol{X}}\right), \tag{7.19}
\end{align*}
$$

with $\boldsymbol{C}^{*}(\hat{\boldsymbol{v}})$ being defined such that $\boldsymbol{C}^{*}(\hat{\boldsymbol{v}}) \tilde{\boldsymbol{v}}=\boldsymbol{C}(\tilde{\boldsymbol{v}}) \hat{\boldsymbol{v}}$. It is now observed that the systems (7.15) and (7.18) are in a cascaded structure. It is also observed that the system $(\tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}})$ is GES at the origin and that the connected terms $H_{12}, H_{21}$, and $H_{22}$ are Lipschitz in $\eta$ and $\hat{\boldsymbol{v}}$. Furthermore from (7.17) and (7.6), the velocity estimate error vector, $\tilde{\boldsymbol{v}}=[\tilde{u}, \tilde{v}, \tilde{r}]^{T}=\boldsymbol{v}-\hat{\boldsymbol{v}}$, satisfies

$$
\begin{equation*}
\tilde{\boldsymbol{v}}=e^{-\boldsymbol{Q}(t, \boldsymbol{\eta})} \tilde{\boldsymbol{X}} \tag{7.20}
\end{equation*}
$$

Since all elements of $e^{-\boldsymbol{Q}(t, \eta)}$ are bounded, (7.16) and (7.20) imply that there exists a positive constant $\gamma_{0}$ such that

$$
\begin{equation*}
\|(\tilde{\boldsymbol{\eta}}(t), \tilde{\boldsymbol{v}}(t))\| \leq \gamma_{0}\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{v}}\left(t_{0}\right)\right)\right\| e^{-\sigma_{0}\left(t-t_{0}\right)}, \forall 0 \leq t_{0} \leq t<\infty \tag{7.21}
\end{equation*}
$$

which means that the estimation errors $\tilde{\boldsymbol{\eta}}(t)$ and $\tilde{\boldsymbol{v}}(t)$ globally exponentially converge to the origin.

### 7.2.2 Coordinate Transformations

We now interpret the position and orientation errors $x-x_{d}, y-y_{d}$, and $\psi-\psi_{d}$ in a frame attached to the ship body. That is, we consider the error coordinates

$$
\begin{align*}
& {\left[\begin{array}{l}
x_{e} \\
y_{e} \\
\psi_{e}
\end{array}\right]=\boldsymbol{J}^{-\mathbf{1}}(\boldsymbol{\eta})\left[\begin{array}{l}
x-x_{d} \\
y-y_{d} \\
\psi-\psi_{d}
\end{array}\right],} \\
& {\left[\begin{array}{l}
u_{e} \\
v_{e} \\
r_{e}
\end{array}\right]=\left[\begin{array}{l}
\hat{u}-u_{d} \\
\hat{v}-v_{d} \\
\hat{r}-r_{d}
\end{array}\right] .} \tag{7.22}
\end{align*}
$$

Differentiating both sides of (7.22) along the solutions of (7.18) and (7.5) yields the error dynamics of the "kinematic part" in the transformed coordinates:

$$
\begin{align*}
& \dot{x}_{e}=u_{e}-u_{d}\left(\cos \left(\psi_{e}\right)-1\right)-v_{d} \sin \left(\psi_{e}\right)+r_{e} y_{e}+r_{d} y_{e}+h_{x}, \\
& \dot{y}_{e}=v_{e}-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)+u_{d} \sin \left(\psi_{e}\right)-r_{e} x_{e}-r_{d} x_{e}+h_{y},  \tag{7.23}\\
& \dot{\psi}_{e}=r_{e}+h_{\psi}
\end{align*}
$$

where $h_{x}, h_{y}$, and $h_{\psi}$ are the first, second, and third rows of $\boldsymbol{J}^{-\mathbf{1}}(\boldsymbol{\eta}) H_{12} \tilde{\boldsymbol{X}}+$ $\tilde{r}\left[\begin{array}{c}y_{e} \\ -x_{e} \\ 0\end{array}\right]$, respectively.

By looking at (7.23), we see that $x_{e}$ and $\psi_{e}$ can be stabilized by $u_{e}$ and $r_{e}$. There are several options to stabilize $y_{e}$, namely $r_{e}, v_{e}$, or $\psi_{e}$. If $r_{e}$ is used, the control design will be extremely complicated since $r_{e}$ enters all of the three equations of (7.23). On the other hand, the use of $v_{e}$ to stabilize $y_{e}$ will result in an undesired feature of ship control practice, namely the ship will slide in the sway direction. Hence we will use $\psi_{e}$ to stabilize the sway error $y_{e}$. As such, we define the following coordinate transformation

$$
\begin{equation*}
z_{e}=\psi_{e}+\arcsin \left(\frac{k y_{e}}{\sqrt{c^{2}+x_{e}^{2}+y_{e}^{2}+v_{e}^{2}}}\right) \tag{7.24}
\end{equation*}
$$

where the constants $k$ and $c$ are such that $|k|<1$ and $c \geq 1$ and will be specified later. It is seen that (7.24) is well defined and that convergence of $z_{e}$ and $y_{e}$ implies that of $\psi_{e}$. By using the nonlinear coordinate transformation (7.24) instead of a linear one like $z_{e}=\psi_{e}+k y_{e}$, we avoid the ship whirling around when $y_{e}$ is large. The coordinate (7.24) is slightly different from the one in the preceding chapter. This will result in bounded virtual velocity controls. Using the nonlinear coordinate (7.24) together with (7.23), the ship error dynamics are rewritten as

$$
\begin{aligned}
& \dot{x}_{e}=u_{e}+u_{d} \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)+k v_{d} \varpi_{2}^{-1} y_{e}+r_{e} y_{e}+r_{d} y_{e}+p_{x}+h_{x}, \\
& \dot{y}_{e}=v_{e}+v_{d} \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)-k u_{d} \varpi_{2}^{-1} y_{e}-r_{e} x_{e}-r_{d} x_{e}+p_{y}+h_{y},
\end{aligned}
$$

$$
\begin{align*}
\dot{z}_{e}= & \left(1-k \varpi_{1}^{-1} x_{e}+k \beta \varpi_{1}^{-1} \varpi_{2}^{-2} y_{e} v_{e}\left(u_{e}+u_{d}\right)\right) r_{e}+ \\
& k \varpi_{1}^{-1}\left(v_{e}+v_{d} \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)-k u_{d} \varpi_{2}^{-1} y_{e}-r_{d} x_{e}-y_{e} \varpi_{2}^{-2} \times\right. \\
& \left.x_{e} u_{e}+y_{e} v_{e}+\left(x_{e} u_{d}+y_{e} v_{d}\right) \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)+\left(x_{e} v_{d}-y_{e} u_{d}\right)\right) \\
& \left.\left(k \varpi_{2}^{-1} y_{e}-v_{e}\left(\alpha v_{e}+\beta u_{e} r_{d}\right)\right)\right)+p_{z}+h_{z}, \\
\dot{u}_{e}= & \frac{m_{22}}{m_{11}} \hat{v} \hat{r}-\frac{d_{11}}{m_{11}} \hat{u}+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d}+h_{u}, \\
\dot{v}_{e}= & -\frac{m_{11}}{m_{22}} u_{e} r_{d}-\frac{m_{11}}{m_{22}}\left(u_{e}+u_{d}\right) r_{e}-\frac{d_{22}}{m_{22}} v_{e}+h_{v}, \\
\dot{r}_{e}= & \frac{\left(m_{11}-m_{22}\right)}{m_{33}} \hat{u} \hat{v}-\frac{d_{33}}{m_{33}} \hat{r}+\frac{1}{m_{33}} \tau_{r}-\dot{r}_{d}+h_{r}, \tag{7.25}
\end{align*}
$$

where $h_{u}, h_{v}$, and $h_{r}$ are the first, second, and third rows of $H_{21} \tilde{\boldsymbol{\eta}}+H_{22} \tilde{\boldsymbol{X}}$, respectively. Also, for notational simplicity, we have defined

$$
\begin{align*}
\varpi_{1}= & \sqrt{c^{2}+x_{e}^{2}+\left(1-k^{2}\right) y_{e}^{2}+v_{e}^{2}}, \\
\varpi_{2}= & \sqrt{c^{2}+x_{e}^{2}+y_{e}^{2}+v_{e}^{2}}, \\
\alpha= & \frac{d_{22}}{m_{22}}, \quad \beta=\frac{m_{11}}{m_{22}}, \\
p_{x}= & -u_{d}\left(\left(\cos \left(z_{e}\right)-1\right) \varpi_{1} \varpi_{2}^{-1}+\sin \left(z_{e}\right) k \varpi_{2}^{-1} y_{e}\right)- \\
& v_{d}\left(\sin \left(z_{e}\right) \varpi_{1} \varpi_{2}^{-1}-\left(\cos \left(z_{e}\right)-1\right) k \varpi_{2}^{-1} y_{e}\right), \\
p_{y}= & -v_{d}\left(\left(\cos \left(z_{e}\right)-1\right) \varpi_{1} \varpi_{2}^{-1}+\sin \left(z_{e}\right) k \varpi_{2}^{-1} y_{e}\right)+  \tag{7.26}\\
& u_{d}\left(\sin \left(z_{e}\right) \varpi_{1} \varpi_{2}^{-1}-\left(\cos \left(z_{e}\right)-1\right) k \varpi_{2}^{-1} y_{e}\right), \\
p_{z}= & k \varpi_{1}^{-1}\left(p_{y}-y_{e} \varpi_{2}^{-2}\left(x_{e} p_{x}+y_{e} p_{y}\right)\right) \\
h_{z}= & k \varpi_{1}^{-1}\left(h_{y}-y_{e} \varpi_{2}^{-1}\left(x_{e} h_{x}+y_{e} h_{y}+v_{e} h_{v}\right)\right)+h_{\psi} .
\end{align*}
$$

It is now clear that the problem of forcing the underactuated ship (7.1) to track the virtual ship (7.5) becomes one of stabilizing the system (7.25) at the origin.

### 7.2.3 Control Design

The triangular structure of (7.25) suggests that we design the actual controls $\tau_{u}$ and $\tau_{r}$ in two stages. First, we design the virtual velocity controls for $u_{e}$ and $r_{e}$ to globally asymptotically stabilize $x_{e}, y_{e}, z_{e}$ and $v_{e}$ at the origin. Based on the backstepping technique, the controls $\tau_{u}$ and $\tau_{r}$ will be then designed. It is noted that the term $\left(1-k x_{e} / \varpi_{1}+k \beta y_{e} v_{e}\left(u_{e}+u_{d}\right) /\left(\varpi_{1} \varpi_{2}^{2}\right)\right)$ in the $z_{e}$-dynamics may vanish and therefore might prevent a global design. This problem can be fixed by decomposing $u_{e}$ and $r_{e}$ as

$$
\begin{align*}
& u_{e}=u_{e}^{d}+\tilde{u}_{e} \\
& r_{e}=r_{e}^{d}+\tilde{r}_{e} \tag{7.27}
\end{align*}
$$

where $u_{e}^{d}$ and $r_{e}^{d}$ are the virtual velocity controls of $u_{e}$ and $r_{e} ; \tilde{u}_{e}$ and $\tilde{r}_{e}$ are the virtual control errors.

## Step 1

In this step, the virtual surge and yaw velocity controls are chosen as

$$
\begin{equation*}
u_{e}^{d}=-\frac{k_{1} x_{e}}{\varpi_{2}}, r_{e}^{d}=r_{1 e}^{d}+r_{2 e}^{d} \tag{7.28}
\end{equation*}
$$

where

$$
\begin{align*}
& r_{1 e}^{d}=-\frac{1}{1-k x_{e} / \varpi_{1}+k \beta y_{e} v_{e}\left(u_{e}^{d}+u_{d}\right) /\left(\varpi_{1} \varpi_{2}^{2}\right)}\left(\frac { k } { \varpi _ { 1 } } \left(v_{e}+\frac{v_{d}}{\varpi_{2}} \times\right.\right. \\
&\left.\left(\varpi_{2}-\varpi_{1}\right)\right)-\frac{k u_{d} y_{e}}{\varpi_{2}}-r_{d} x_{e}-\left(\frac { y _ { e } } { \varpi _ { 2 } } \left(x_{e} u_{e}^{d}+y_{e} v_{e}+\left(x_{e} u_{d}+\right.\right.\right. \\
&\left.\left.\left.\left.y_{e} v_{d}\right) \frac{\left(\varpi_{2}-\varpi_{1}\right)}{\varpi_{2}^{-1}}+\left(x_{e} v_{d}-y_{e} u_{d}\right) \frac{k y_{e}}{\varpi_{2}}-v_{e}\left(\alpha v_{e}+\beta u_{e}^{d} r_{d}\right)\right)\right)\right), \\
& r_{2 e}^{d}=-\frac{1}{1-k x_{e} / \varpi_{1}+k \beta y_{e} v_{e}\left(u_{e}^{d}+u_{d}\right) /\left(\varpi_{1} \varpi_{2}^{2}\right)}\left(\frac{k_{2} z_{e}}{\sqrt{1+z_{e}^{2}}}+p_{z}\right), \tag{7.29}
\end{align*}
$$

and $k_{i}, i=1,2$, are positive constants to be selected later. We have written $r_{e}^{d}$ as a sum of $r_{1 e}^{d}$ and $r_{2 e}^{d}$ to simplify notation in the stability analysis later. Notice that

$$
\begin{equation*}
1-\frac{k x_{e}}{\varpi_{1}}+\frac{k \beta y_{e} v_{e}\left(u_{e}^{d}+u_{d}\right)}{\varpi_{1} \varpi_{2}^{2}} \geq 1-|k|\left(1+\frac{0.5 \beta\left(k_{1}+\left|u_{d}\right|\right)}{c}\right) \tag{7.30}
\end{equation*}
$$

therefore $r_{1 e}^{d}$ and $r_{2 e}^{d}$ are well defined if the design constants $k, c$ and $k_{1}$ are chosen such that

$$
\begin{equation*}
1-|k|\left(1+0.5 \beta\left(k_{1}+\left|u_{d}\right|\right) / c\right) \geq k^{*}>0 \tag{7.31}
\end{equation*}
$$

By noting

$$
\begin{align*}
& \left|p_{x}\right| \leq\left|u_{d}\right|(2+|k|)+\left|v_{d}\right|(1+2|k|), \\
& \left|p_{y}\right| \leq\left|v_{d}\right|(2+|k|)+\left|u_{d}\right|(1+2|k|), \\
& \left|p_{z}\right| \leq|k|\left(2\left|p_{y}\right|+\left|p_{x}\right|\right), \tag{7.32}
\end{align*}
$$

we can show from (7.28) and (7.29) that $u_{e}^{d}$ and $r_{e}^{d}$ are bounded by some constants. Remark 7.3. Unlike the standard application of the backstepping technique, in order to reduce complexity of the controller expressions, we have chosen a simple virtual
control law $u_{e}^{d}$ without canceling the known terms. From (7.29) and (7.28), we observe that $r_{1 e}^{d}$ is Lipschitz in $\left(x_{e}, y_{e}, v_{e}\right)$ and $r_{2 e}^{d}$ vanishes when $z_{e}$ does. This observation plays a crucial role in the stability analysis of the closed loop system.

## Step 2

By differentiating (7.27) along the solutions of (7.25) and (7.29), the actual controls $\tau_{u}$ and $\tau_{r}$ without canceling the useful damping terms are chosen as

$$
\begin{align*}
\tau_{u} & =-m_{11}\left(c_{1} \tilde{u}_{e}+m_{22} m_{11}^{-1} \hat{v} \hat{r}-d_{11} m_{11}^{-1}\left(u_{e}^{d}+u_{d}\right)-\dot{u}_{d}-\right. \\
& \frac{\partial u_{e}^{d}}{\partial x_{e}}\left(u_{e}+u_{d} \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)+k v_{d} \varpi_{2}^{-1} y_{e}+r_{e} y_{e}+r_{d} y_{e}\right)- \\
& \frac{\partial u_{e}^{d}}{\partial y_{e}}\left(v_{e}+v_{d} \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)-k u_{d} \varpi_{2}^{-1} y_{e}-r_{e} x_{e}-r_{d} x_{e}\right)- \\
& \frac{\partial u_{e}^{d}}{\partial v_{e}} m_{22}^{-1}\left(-m_{11} u_{e} r_{d}-m_{11}\left(u_{e}+u_{d}\right) r_{e}-d_{22} v_{e}\right)+ \\
& \left.\varpi_{1}^{-1} \varpi_{2}^{-2}\left(k \beta y_{e} v_{e} r-k y_{e}\left(x_{e}-\beta r_{d} v_{e}\right)\right) z_{e}\right), \\
\tau_{r} & =-m_{33}\left(c_{2} \tilde{r}_{e}+\left(m_{11}-m_{22}\right) m_{33}^{-1} \hat{u} \hat{v}-d_{33} m_{33}^{-1}\left(r_{e}^{d}+r_{d}\right)-\dot{r}_{d}-\right. \\
& \frac{\partial r_{e}^{d}}{\partial u_{d}} \dot{u}_{d}-\frac{\partial r_{e}^{d}}{\partial v_{d}} \dot{v}_{d}-\frac{\partial r_{e}^{d}}{\partial r_{d}} \dot{r}_{d}-\frac{\partial r_{e}^{d}}{\partial x_{e}}\left(u_{e}+u_{d} \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)+\right. \\
& \left.k v_{d} \varpi_{2}^{-1} y_{e}+r_{e} y_{e}+r_{d} y_{e}\right)-\frac{\partial r_{e}^{d}}{\partial y_{e}}\left(v_{e}+v_{d} \varpi_{2}^{-1}\left(\varpi_{2}-\varpi_{1}\right)-\right. \\
& \left.k u_{d} \varpi_{2}^{-1} y_{e}-r_{e} x_{e}-r_{d} x_{e}\right)-\frac{\partial r_{e}^{d}}{\partial v_{e}} m_{22}^{-1}\left(-m_{11} u_{e} r_{d}-m_{11}\left(u_{e}+\right.\right. \\
& \left.\left.u_{d}\right) r_{e}-d_{22} v_{e}\right)-\frac{\partial r_{e}^{d}}{\partial \psi_{e}} r_{e}+\left(1-k \varpi_{1}^{-1} x_{e}+k \beta \varpi_{1}^{-1} \varpi_{2}^{-2} \times\right. \\
& \left.\left.y_{e} v_{e}\left(u_{e}^{d}+u_{d}\right)\right) z_{e}\right), \tag{7.33}
\end{align*}
$$

where $c_{i}, i=1,2$, are positive constants. We now state the first main result of this chapter, the proof of which is given in the next section.

Theorem 7.1. Under Assumption 7.1 assume the following:

1. There are no environmental disturbances
2. The ship parameters are known
3. The reference signals are generated by the virtual ship model (7.5) and the reference velocities satisfy Assumption 7.1.

If the partial state feedback control law (7.33) together with the observer (7.18) are applied to the ship system (7.1), then the tracking errors $x(t)-x_{d}(t), y(t)-y_{d}(t)$, $\psi(t)-\psi_{d}(t)$, and $v(t)-v_{d}(t)$ globally asymptotically and locally exponentially
converge to zero with an appropriate choice of the design constants $c, k$ and $k_{i}, i=$ 1,2. Furthermore, the virtual velocity controls, $u_{e}^{d}$ and $r_{e}^{d}$, are bounded by some computable positive constants.

### 7.2.4 Stability Analysis

Substituting (7.33) and (7.28) into (7.25) results in the following closed loop system:

$$
\begin{align*}
& \dot{X}_{1 e}=f_{1}\left(t, X_{e}\right)+g_{1}\left(t, X_{e}\right)+\phi_{1}\left(t, \hat{v}, X_{e}, \eta, \tilde{\eta}, \tilde{X}\right), \\
& \dot{X}_{2 e}=f_{2}\left(t, X_{e}\right)+\phi_{2}\left(t, \hat{v}, X_{e}, \eta, \tilde{\eta}, \tilde{X}\right) \tag{7.34}
\end{align*}
$$

where

$$
\begin{aligned}
& \boldsymbol{X}_{\mathbf{1} \boldsymbol{e}}=\left[\begin{array}{l}
x_{e} \\
y_{e} \\
v_{e}
\end{array}\right], \boldsymbol{X}_{\mathbf{2} \boldsymbol{e}}=\left[\begin{array}{l}
z_{e} \\
\tilde{u}_{e} \\
\tilde{r}_{e}
\end{array}\right], \boldsymbol{X}_{\boldsymbol{e}}=\left[\begin{array}{c}
\boldsymbol{X}_{\mathbf{1}} \\
\boldsymbol{X}_{\mathbf{2} e}
\end{array}\right], \\
& \boldsymbol{f}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)=\left[\begin{array}{c}
f_{11} \\
f_{12} \\
f_{13}
\end{array}\right], \boldsymbol{f}_{\mathbf{2}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)=\left[\begin{array}{c}
f_{21} \\
f_{22} \\
f_{23}
\end{array}\right], \boldsymbol{g}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)=\left[\begin{array}{c}
p_{x}+\tilde{u}_{e} \\
p_{y} \\
g_{13}
\end{array}\right], \\
& \boldsymbol{\phi}_{\mathbf{1}}\left(t, \hat{\boldsymbol{v}}, \boldsymbol{X}_{\boldsymbol{e}}, \boldsymbol{\eta}, \tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}}\right)=\left[\begin{array}{c}
h_{x}^{*} \\
h_{y}^{*} \\
h_{v}
\end{array}\right], \boldsymbol{\phi}_{\mathbf{2}}\left(t, \hat{\boldsymbol{v}}, \boldsymbol{X}_{\boldsymbol{e}}, \boldsymbol{\eta}, \tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}}\right)=\left[\begin{array}{c}
h_{z} \\
\varphi_{22} \\
\varphi_{23}
\end{array}\right], \\
& f_{11}=-\frac{k_{1} x_{e}}{\varpi_{2}}+u_{d} \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}}+\frac{k v_{d} y_{e}}{\varpi_{2}}+y_{e}\left(r_{e}+r_{d}+\tilde{r}\right), \\
& f_{12}=v_{e}+v_{d} \frac{\varpi_{2}-\varpi_{1}}{\varpi_{2}}-\frac{k u_{d} y_{e}}{\varpi_{2}}-x_{e}\left(r_{e}+r_{d}+\tilde{r}\right), \\
& f_{13}=-\alpha v_{e}-\beta u_{e}^{d} r_{d}-\beta\left(u_{e}^{d}+u_{d}\right) r_{1 e}^{d}, \\
& g_{13}=-\beta\left(\tilde{u}_{e}\left(r_{e}^{d}+\tilde{r}_{e}+r_{d}\right)+\left(u_{e}^{d}+u_{d}\right)\left(r_{2 e}^{d}+\tilde{r}_{e}\right)\right), \\
& f_{21}=-\frac{k_{2} z_{e}}{\sqrt{1+z_{e}^{2}}+\left(1-\frac{k x_{e}}{\varpi_{1}}+\frac{k \beta y_{e} v_{e}\left(u_{e}^{d}+u_{d}\right)}{\varpi_{1} \varpi_{2}^{2}}\right) \tilde{r}_{e}+} \\
& \left(\frac{k \beta y_{e} v_{e} r_{e}}{\varpi_{1} \varpi_{2}^{2}}-\frac{k y_{e}}{\varpi_{1} \varpi_{2}^{2}}\left(x_{e}-\beta r_{d} v_{e}\right)\right) \tilde{u}_{e}, \\
& f_{22}=-\left(c_{1}+\frac{d_{11}}{m_{11}}\right) \tilde{u}_{e}-\left(\frac{k \beta y_{e} v_{e} r_{e}}{\varpi_{1} \varpi_{2}^{2}}-\frac{k y_{e}}{\varpi_{1} \varpi_{2}^{2}}\left(x_{e}-\beta r_{d} v_{e}\right)\right) z_{e}, \\
& f_{23}=-\left(c_{2}+\frac{d_{33}}{m_{33}}\right) \tilde{r}_{e}-\left(1-\frac{k x_{e}}{\varpi_{1}}+\frac{k \beta y_{e} v_{e}\left(u_{e}^{d}+u_{d}\right)}{\varpi_{1} \varpi_{2}^{2}}\right) z_{e},
\end{aligned}
$$

$$
\begin{align*}
\varphi_{22} & =h_{u}-\frac{\partial u_{e}^{d}}{\partial x_{e}}\left(p_{x}+h_{x}\right)-\frac{\partial u_{e}^{d}}{\partial y_{e}}\left(p_{y}+h_{y}\right)-\frac{\partial u_{e}^{d}}{\partial v_{e}} h_{v} \\
\varphi_{23} & =h_{r}-\frac{\partial r_{e}^{d}}{\partial x_{e}}\left(p_{x}+h_{x}\right)-\frac{\partial r_{e}^{d}}{\partial y_{e}}\left(p_{y}+h_{y}\right)-\frac{\partial r_{e}^{d}}{\partial v_{e}} h_{v}-\frac{\partial r_{e}^{d}}{\partial \psi_{e}} h_{\psi}, \tag{7.35}
\end{align*}
$$

with $h_{x}^{*}$ and $h_{y}^{*}$ being the first and second rows of $\boldsymbol{J}^{-\mathbf{1}}(\boldsymbol{\eta}) H_{12} \tilde{\boldsymbol{X}}$. The time dependence of $\boldsymbol{f}_{\boldsymbol{i}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right), \boldsymbol{g}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$, and $\boldsymbol{\phi}_{\boldsymbol{i}}\left(t, \hat{\boldsymbol{v}}, \boldsymbol{X}_{\boldsymbol{e}}, \boldsymbol{\eta}, \tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}}\right), i=1,2$, is due to the time-varying reference velocities. Observe that the closed loop system (7.34) consists of the $\left(\boldsymbol{X}_{\mathbf{1} e}, \boldsymbol{X}_{\mathbf{2} \boldsymbol{e}}\right)$-subsystem and ( $\left.\tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}}\right)$-subsystem (see (7.15)) in a cascaded structure. From (7.35) it can be readily shown that the connected terms $\boldsymbol{\phi}_{\boldsymbol{i}}\left(t, \hat{\boldsymbol{v}}, \boldsymbol{X}_{\boldsymbol{e}}, \boldsymbol{\eta}, \tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}}\right), i=1,2$, satisfy

$$
\begin{equation*}
\left\|\phi_{i}\left(t, \hat{\boldsymbol{v}}, \boldsymbol{X}_{\boldsymbol{e}}, \boldsymbol{\eta}, \tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}}\right)\right\| \leq \phi_{i}(\hat{\boldsymbol{v}}, \boldsymbol{\eta})\|(\tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}})\|, \tag{7.36}
\end{equation*}
$$

where the functions $\phi_{i}(\hat{\boldsymbol{v}}, \boldsymbol{\eta})$ are Lipschitz in $\hat{\boldsymbol{v}}$ and are bounded with respect to any $\eta$. Also we note that by definition $\hat{\mathbf{v}}=[\hat{u}, \hat{v}, \hat{r}]^{T}=\left[u_{e}^{d}+u_{d}+\tilde{u}_{e}, v_{e}+v_{d}, r_{e}^{d}+\right.$ $\left.r_{d}+\tilde{r}_{e}\right]^{T}$, the $(\tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}})$-subsystem is GES at the origin, and the reference velocities $u_{d}, v_{d}$ and $r_{d}$ are bounded. On the other hand, the virtual velocity controls, $u_{e}^{d}$ and $r_{e}^{d}$ are bounded. Hence, using the recent stability results for cascade systems given in [17,69], we need to show that there exist the design constants $c, k, k_{1}$, and $k_{2}$ such that the $\left(\boldsymbol{X}_{\mathbf{1} \boldsymbol{e}}, \boldsymbol{X}_{\mathbf{2} \boldsymbol{e}}\right)$-subsystem without the connected terms $\boldsymbol{\phi}_{\boldsymbol{i}}(t, \hat{\boldsymbol{v}}, \boldsymbol{\eta}, \tilde{\boldsymbol{\eta}}, \tilde{\boldsymbol{X}}), i=$ 1,2 , is GAS at the origin. That is why we did not include some nonlinear damping terms in the control law (7.33). From the above discussions, we will study the system given by

$$
\begin{align*}
& \dot{X}_{1 e}=f_{1}\left(t, X_{e}\right)+g_{1}\left(t, X_{e}\right) \\
& \dot{X}_{2 e}=f_{2}\left(t, X_{e}\right) \tag{7.37}
\end{align*}
$$

To further simplify the investigation of stability of the system (7.37), we note that this system also consists of the $\boldsymbol{X}_{\mathbf{1} \boldsymbol{e}}$-subsystem and the $\boldsymbol{X}_{\mathbf{2} \boldsymbol{e}}$-subsystem in a cascaded structure. From (7.35), it is not hard to show that the connected term $\boldsymbol{g}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$ satisfies $\left\|\boldsymbol{g}_{\mathbf{1}}\left(\boldsymbol{t}, \boldsymbol{X}_{\boldsymbol{e}}\right)\right\| \leq \kappa_{1}\left\|\boldsymbol{X}_{\mathbf{2} e}\right\|$ with $\kappa_{1}$ being some positive constant. Therefore global stability of $\dot{\boldsymbol{X}}_{\boldsymbol{1} \boldsymbol{e}}=\boldsymbol{f}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$ and $\dot{\boldsymbol{X}}_{\mathbf{2} \boldsymbol{e}}=\boldsymbol{f}_{\mathbf{2}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$ implies that of (7.37). We will first study stability of the subsystem $\dot{\boldsymbol{X}}_{\mathbf{2} \boldsymbol{e}}=\boldsymbol{f}_{\mathbf{2}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$ then move to the subsystem $\dot{\boldsymbol{X}}_{\mathbf{1} \boldsymbol{e}}=\boldsymbol{f}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$.

Subsystem $\dot{\boldsymbol{X}}_{\mathbf{2} e}=\boldsymbol{f}_{\mathbf{2}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$. By differentiating $V_{1}=0.5\left(z_{e}^{2}+\tilde{u}_{e}^{2}+\tilde{r}_{e}^{2}\right)$ along the solutions of $\dot{X}_{\mathbf{2} e}=f_{\mathbf{2}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$, one can show that this system is globally asymptotically and locally exponentially stable at the origin for any constants $k_{2}>0, c_{1} \geq 0$, and $c_{2} \geq 0$.

Subsystem $\dot{\boldsymbol{X}}_{\mathbf{1} \boldsymbol{e}}=\boldsymbol{f}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$. To investigate the stability of this subsystem, we take the Lyapunov function

$$
\begin{equation*}
V_{2}=\sqrt{c^{2}+x_{e}^{2}+y_{e}^{2}+v_{e}^{2}}+\frac{1}{2} k_{3} v_{e}^{2}-c \tag{7.38}
\end{equation*}
$$

where $k_{3}$ is a positive constant to be selected later. After some lengthy but simple calculation using the completing squares, the time derivative of (7.38) along the solutions of $\dot{\boldsymbol{X}}_{\mathbf{1} \boldsymbol{e}}=\boldsymbol{f}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$ satisfies

$$
\begin{equation*}
\dot{V}_{2} \leq-p_{1}(t) x_{e}^{2} \varpi_{2}^{-2}-p_{2}(t) y_{e}^{2} \varpi_{2}^{-2}-p_{3}(t) v_{e}^{2} \tag{7.39}
\end{equation*}
$$

with

$$
\begin{align*}
p_{1}(t)= & k_{1}-\beta\left(k_{3}+1 / c\right) k_{1}\left|r_{d}\right| /\left(4 \varepsilon_{3}\right)-\left[\left|k v_{d}\right| /\left(4 \varepsilon_{1}\right)+\beta\left(k_{3}+1 / c\right) \times\right. \\
& \left.\left(k_{1}+\left|u_{d}\right|\right)|k|\left(\left|r_{d}\right| /\left(4 \varepsilon_{3}\left(1-k^{2}\right)\right)+k_{1} /\left(2 c \sqrt{1-k^{2}}\right)\right) / k^{*}\right], \\
p_{2}(t)= & k u_{d}-1 /\left(4 \varepsilon_{2}\right)-\left[k^{2}\left(0.5\left|u_{d}\right|+\left|v_{d}\right|\right) / c+\varepsilon_{1}\left|k v_{d}\right|+\beta\left(k_{3}+1 / c\right) \times\right. \\
& \left.\left(k_{1}+\left|u_{d}\right|\right)|k|\left(k^{2}\left|v_{d}\right|+\left(k^{2}+|k|\right)\left(\left|u_{d}\right|+\left|v_{d}\right|\right)+\left|k u_{d}\right| /\left(c 4 \varepsilon_{3}\right)\right) / k^{*}\right], \\
p_{3}(t)= & \alpha k_{3}-\varepsilon_{2}-\beta\left(k_{3}+1 / c\right) k_{1}\left|r_{d}\right| \varepsilon_{3}-\left[\beta\left(k_{3}+1 / c\right)\left(k_{1}+\left|u_{d}\right|\right)|k| \times\right. \\
& \left.\left(2 / c+\left|k u_{d}\right| \varepsilon_{3} / c+\varepsilon_{3}\left|r_{d}\right|+0.5 \alpha / c+\beta k_{1}\left|r_{d}\right| / c^{2}\right) / k^{*}\right], \tag{7.40}
\end{align*}
$$

where $\varepsilon_{i}, 1 \leq i \leq 3$ are positive constants. Hence the subsystem $\dot{\boldsymbol{X}}_{\mathbf{1} \boldsymbol{e}}=\boldsymbol{f}_{\mathbf{1}}\left(t, \boldsymbol{X}_{\boldsymbol{e}}\right)$ is GAS at the origin if the design constants are chosen such that

$$
\begin{equation*}
p_{i}(t) \geq p_{i}^{*} \tag{7.41}
\end{equation*}
$$

for some positive constants $p_{i}^{*}, i=1,2,3$. In summary, we need to choose the constants $c, k, k_{1}$, and $k_{3}$ such that they satisfy (7.31) and (7.41). Note that the condition (7.31) automatically implies that $|k|<1$ is required in (7.24). In the next section, we will show that under Assumption 7.1, there always exist the constants $c$, $k, k_{1}$, and $k_{3}$ such that (7.31) and (7.41) hold.

### 7.2.5 Selection of Design Constants

To choose the design constants $c, k, k_{1}$, and $k_{3}$, we observe the following: First, it is noted from the expression of $p_{2}(t)$ that the sign of constant $k$ must have the same sign as that of the reference surge velocity, $u_{d}$ (this sign does not change under Assumption 7.1). Second, it is observed that the condition (7.31) can be rewritten in the form of

$$
\begin{equation*}
1-|k|\left(1+0.5 \beta\left(k_{1}+u_{d}^{\max }\right) / c\right) \geq k^{*}>0 \tag{7.42}
\end{equation*}
$$

which implies that for each fixed $k^{*}<1, k_{1}>0, u_{d}^{\max },|k|<k^{*}$, we can always pick a large enough constant $c$ such that (7.42) holds. Third, under Assumption 7.1, the magnitude of the reference sway velocity is always less than that of the reference surge velocity. Fourth, the mass including added mass in the sway dynamics, $m_{22}$, is always larger than that in the surge dynamics, $m_{11}$, for surface ships, i.e., $\beta<1$.

Finally, all of the terms in the square brackets in $p_{i}, i=1,2,3$ have the constant $k$ as a factor. These terms also decrease when the constant $c$ increases. Looking closely at $p_{i}$ with the above observations, if the constant $k$ is chosen small enough and the constant $c$ is selected large enough, then we can pick a positive constant $k_{1}$ such that (7.31) and (7.41) hold with some large enough $k_{3}$. It is noted that $|k|$ should be decreased and $c$ should be increased if $u_{d}$ is large. This physically means that the distance from the ship to the point it aims to track should be increased if the velocities and surge acceleration are large, otherwise the ship will miss that point. Furthermore when $\alpha$ is small, $|k|$ and $k_{1}$ should be decreased, and $c$ should be increased. This can be physically interpreted as follows: If the damping in the sway dynamics is small, the control gain in the surge dynamics should also be small otherwise the ship will slide in the sway direction. Due to complicated expressions of $p_{i}(t)$, we provide some general guidelines to choose the design constants rather than present their extremely complex explicit expressions: Pick $k^{*}<1$, small values for $|k|$ and $k_{1}$, larger values for $c$ and $k_{3}$. Then increase $c$ and $k_{3}$ until (7.31) and (7.41) hold.

### 7.3 Output Feedback

### 7.3.1 Observer Design

We now introduce a more general coordinate change than (7.6) to cancel the term $\boldsymbol{C}(\boldsymbol{v}) \boldsymbol{v}$ in (7.1) as follows:

$$
\begin{equation*}
X=P(\eta) v \tag{7.43}
\end{equation*}
$$

where $\boldsymbol{P}(\eta) \in \mathbb{R}^{3 \times 3}$ is a global invertible matrix to be determined. With (7.43), the second equation of (7.1) is written as

$$
\begin{equation*}
\dot{X}=\left[\dot{P}(\eta) v-P(\eta) M^{-1} C(v) v\right]-P(\eta) M^{-1} D P^{-1}(\eta) X+P(\eta) M^{-1} \tau \tag{7.44}
\end{equation*}
$$

Our goal is to cancel the terms in the square bracket on the right-hand side of (7.44). Assuming that the elements of $\boldsymbol{P}(\boldsymbol{\eta})$ are $p_{i j}(\boldsymbol{\eta}), i=1,2,3, j=1,2,3$, the first bracket in the right-hand side of (7.44) is zero if

$$
\begin{align*}
& \dot{p}_{i 1} u+\dot{p}_{i 2} v+\dot{p}_{i 3} r+\frac{m_{22}}{m_{11}} p_{i 1} v r-\frac{m_{11}}{m_{22}} p_{i 2} u r+\frac{m_{11}-m_{22}}{m_{33}} p_{i 3} u v=0, \\
& i=1,2,3, \forall(\boldsymbol{\eta}, u, v, r) \in \mathbb{R}^{6}, \tag{7.45}
\end{align*}
$$

where for brevity, we omit the argument $\eta$ of $p_{i j}(\eta)$. With the first equation of (7.1), we expand (7.45) as

$$
\begin{align*}
& \left(\frac{\partial p_{i 1}}{\partial x} \cos (\psi)+\frac{\partial p_{i 1}}{\partial y} \sin (\psi)\right) u^{2}+\left(-\frac{\partial p_{i 2}}{\partial x} \sin (\psi)+\frac{\partial p_{i 2}}{\partial y} \cos (\psi)\right) v^{2}+ \\
& \frac{\partial p_{i 3}}{\partial \psi} r^{2}+\left(-\frac{\partial p_{i 1}}{\partial x} \sin (\psi)+\frac{\partial p_{i 1}}{\partial y} \cos (\psi)+\frac{\partial p_{i 2}}{\partial x} \cos (\psi)+\frac{\partial p_{i 2}}{\partial y} \sin (\psi)+\right. \\
& \left.\frac{m_{11}-m_{22}}{m_{33}} p_{i 3}\right) u v+\left(\frac{\partial p_{i 1}}{\partial \psi}+\frac{\partial p_{i 3}}{\partial x} \cos (\psi)+\frac{\partial p_{i 3}}{\partial y} \sin (\psi)-\frac{m_{11}}{m_{22}} p_{i 2}\right) u r+ \\
& \left(\frac{\partial p_{i 2}}{\partial \psi}-\frac{\partial p_{i 3}}{\partial x} \sin (\psi)+\frac{\partial p_{i 3}}{\partial y} \cos (\psi)+\frac{m_{22}}{m_{11}} p_{i 1}\right) v r=0 . \tag{7.46}
\end{align*}
$$

Therefore (7.46) holds for all $(\eta, u, v, r) \in \mathbb{R}^{6}$ if

$$
\begin{align*}
& \frac{\partial p_{i 1}}{\partial x} \cos (\psi)+\frac{\partial p_{i 1}}{\partial y} \sin (\psi)=0 \\
& -\frac{\partial p_{i 2}}{\partial x} \sin (\psi)+\frac{\partial p_{i 2}}{\partial y} \cos (\psi)=0 \\
& \frac{\partial p_{i 3}}{\partial \psi}=0 \\
& \left(\frac{\partial p_{i 2}}{\partial y}-\frac{\partial p_{i 1}}{\partial x}\right) \sin (\psi)+\left(\frac{\partial p_{i 1}}{\partial y}+\frac{\partial p_{i 2}}{\partial x}\right) \cos (\psi)+\frac{m_{11}-m_{22}}{m_{33}} p_{i 3}=0 \\
& \frac{\partial p_{i 1}}{\partial \psi}+\frac{\partial p_{i 3}}{\partial x} \cos (\psi)+\frac{\partial p_{i 3}}{\partial y} \sin (\psi)-\frac{m_{11}}{m_{22}} p_{i 2}=0 \\
& \frac{\partial p_{i 2}}{\partial \psi}-\frac{\partial p_{i 3}}{\partial x} \sin (\psi)+\frac{\partial p_{i 3}}{\partial y} \cos (\psi)+\frac{m_{22}}{m_{11}} p_{i 1}=0 \tag{7.47}
\end{align*}
$$

A family of solutions of the above set of six partial differential equations is

$$
\begin{align*}
& p_{i 1}=\frac{\left(\left(m_{11} C_{i 3} x+m_{33} C_{i 1}\right) \sin (\psi)-\left(m_{11} C_{i 3} y-m_{33} C_{i 2}\right) \cos (\psi)\right)}{m_{33}}, \\
& p_{i 2}=\frac{m_{22}\left(\left(m_{11} C_{i 3} x+m_{33} C_{i 1}\right) \cos (\psi)+\left(m_{11} C_{i 3} y-m_{33} C_{i 2}\right) \sin (\psi)\right)}{m_{11} m_{33}}, \\
& p_{i 3}=C_{i 3}, \tag{7.48}
\end{align*}
$$

where $C_{i 1}, C_{i 2}$ and $C_{i 3}$ are arbitrary constants. It is noted that the above solutions can be obtained by the following Maple ${ }^{\mathrm{TM}}$ code:

```
>PDE1:=diff(p11(x,y,\psi),x) *cos(\psi)+
diff(p11(x,y,\psi),y) *sin(\psi)=0,
-diff(p12(x,y,\psi),x) *sin(\psi)+
diff(p12(x,y,\psi),y) *cos(\psi)=0,
diff(p13(x,y,\psi),\psi)=0,
    (diff(p12(x,y,\psi),y)-diff(p11(x,y,\psi),x)) *
    sin(\psi)+(diff(p11(x,y,\psi),y)+diff(p12(x,y,
    \psi),x))*cos(\psi)+(m11-m22)/m33*p13(x,y,\psi)=0,
    diff(p11(x,y,\psi),\psi)+diff(p13(x,y,\psi),x)*
```

```
        cos(\psi)+diff(p13(x,y,\psi),y)*sin(\psi) -
        m11/m22*p12(x,y,\psi)=0,
        diff(p12(x,y,\psi),\psi)-diff(p13(x,y,\psi),x)*
        sin(\psi)+diff(p13(x,y,\psi),y) *cos(\psi) +
        m22/m11*p11(x,y,\psi)=0;
>PDE2:=diff(p21(x,y,\psi),x) *cos(\psi) +
    diff(p21(x,y,\psi),y) *sin(\psi)=0,
    -diff(p22(x,y,\psi),x) *sin(\psi)+
    diff(p22(x,y,\psi),y)*cos(\psi)=0,
    diff(p23(x,y,\psi),\psi)=0,
    (diff(p22(x,y,\psi),y)-diff(p21(x,y,\psi),x)) *
    sin(\psi)+(diff(p21(x,y,\psi),y)+diff(p22(x,y,
    \psi),x))*cos(\psi)+(m11-m22)/m33*p23(x,y,\psi)=0,
    diff(p21(x,y,\psi),\psi)+diff(p23(x,y,\psi),x)*
    cos(\psi)+diff(p23(x,y,\psi),y)*sin(\psi) -
    m11/m22*p22(x,y,\psi)=0,
    diff(p22(x,y,\psi),\psi)-diff(p23(x,y,\psi),x)*
    sin(\psi)+diff(p23(x,y,\psi),y) *cos(\psi) +
    m22/m11*p21(x,y,\psi)=0;
>PDE3:=diff(p31(x,y,\psi),x) *cos(\psi) +
    diff(p31(x,y,\psi),y)*sin(\psi)=0,
    -diff(p32(x,y,\psi),x) *sin(\psi)+
    diff(p32(x,y,\psi),y)*Cos(\psi)=0,
    diff(p33(x,y,\psi),\psi)=0,
    (diff(p32(x,y,\psi),y)-diff(p31(x,y,\psi),x)) *
    sin(\psi)+(diff(p31(x,y,\psi),y)+diff(p32(x,y,
    \psi),x))*cos(\psi)+(m11-m22)/m33*p33(x,y,\psi)=0,
    diff(p31(x,y,\psi),\psi)+diff(p33(x,y,\psi),x)*
    cos(\psi)+diff(p33(x,y,\psi),y)*sin(\psi) -
    m11/m22*p32(x,y,\psi)=0,
    diff(p32(x,y,\psi),\psi)-diff(p33(x,y,\psi),x) *
    sin(\psi)+diff(p33(x,y,\psi),y) *cos(\psi) +
    m22/m11*p31(x,y,\psi)=0;
>solutions:=pdsolve([PDE1, PDE2, PDE3]);
```

We now choose the constants $C_{i 1}, C_{i 2}$, and $C_{i 3}$ such that the matrix $\boldsymbol{P}(\boldsymbol{\eta})$ is invertible. A choice of $C_{13}=C_{11}=0, C_{12}=1, C_{23}=C_{22}=0, C_{21}=1, C_{31}=$ $C_{32}=0$, and $C_{33}=1$ results in

$$
\boldsymbol{P}(\boldsymbol{\eta})=\left[\begin{array}{ccc}
\cos (\psi) & -\frac{m_{22} \sin (\psi)}{m_{11}} & 0  \tag{7.49}\\
\sin (\psi) & \frac{m_{22} \cos (\psi)}{m_{11}} & 0 \\
\frac{m_{11}(\sin (\psi) x-\cos (\psi) y)}{m_{33}} & \frac{m_{22}(\cos (\psi) x+\sin (\psi) y)}{m_{33}} & 1
\end{array}\right]
$$

Substituting (7.49) into (7.44) and using (7.43) and the first equation of (7.1), we have

$$
\begin{align*}
& \dot{\eta}=J(\eta) P^{-\mathbf{1}}(\eta) X, \\
& \dot{X}=-D_{\eta}(\eta) X+P(\eta) M^{-1} \tau \tag{7.50}
\end{align*}
$$

with $\boldsymbol{D}_{\boldsymbol{\eta}}(\boldsymbol{\eta})=\boldsymbol{P}(\boldsymbol{\eta}) \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{D} \boldsymbol{P}^{-\mathbf{1}}(\boldsymbol{\eta})$. It can be seen that the matrix $\boldsymbol{P}(\boldsymbol{\eta})$ given in (7.49) does not use any ship velocities. This feature results in the main difference between the partial-state feedback design in Section 7.2 and the output feedback design in this section. From (7.50), we use the following full-order nonlinear observer to construct the unmeasured ship velocities:

$$
\begin{align*}
& \dot{\hat{\eta}}=J(\eta) P^{-1}(\eta) \hat{X}+K_{01}(\eta-\hat{\eta}), \\
& \dot{\hat{X}}=-D_{\eta}(\eta) \hat{X}+P(\eta) M^{-1} \tau+K_{02}(\eta-\hat{\eta}), \tag{7.51}
\end{align*}
$$

where $\hat{\boldsymbol{\eta}}$ and $\hat{\boldsymbol{X}}$ are the estimates of $\boldsymbol{\eta}$ and $\boldsymbol{X}$, respectively. The observer gain matrices $\boldsymbol{K}_{\mathbf{0 1}}$ and $\boldsymbol{K}_{\mathbf{0 2}}$ are chosen such that

$$
\begin{aligned}
& Q_{01}=K_{01}^{T} P_{01}+P_{01} K_{01} \\
& Q_{02}=D_{\eta}^{T}(\eta) P_{02}+P_{02} D_{\eta}(\eta)
\end{aligned}
$$

are positive definite, and that

$$
\begin{equation*}
\left(J(\eta) \boldsymbol{P}^{-1}(\eta)\right)^{\boldsymbol{T}} \boldsymbol{P}_{01}-\boldsymbol{P}_{02} \boldsymbol{K}_{02}=0 \tag{7.52}
\end{equation*}
$$

with $\boldsymbol{P}_{\mathbf{0 1}}$ and $\boldsymbol{P}_{\mathbf{0 2}}$ being positive definite matrices. It is straightforward to show that $\boldsymbol{K}_{\mathbf{0 1}}$ and $\boldsymbol{K}_{\mathbf{0 2}}$ always exist since $\boldsymbol{D}_{\boldsymbol{\eta}}(\boldsymbol{\eta})$ is positive definite. From (7.50) and (7.51), we have

$$
\begin{align*}
& \dot{\tilde{\eta}}=J(\eta) P^{-1}(\eta) \tilde{X}-K_{01} \tilde{\eta}, \\
& \dot{\tilde{X}}=-D_{\eta}(\eta) \tilde{X}-K_{02} \tilde{\eta}, \tag{7.53}
\end{align*}
$$

where $\tilde{\boldsymbol{\eta}}:=(\tilde{x}, \tilde{y}, \tilde{\psi})^{T}=\boldsymbol{\eta}-\hat{\boldsymbol{\eta}}$ and $\tilde{\boldsymbol{X}}:=\left(\tilde{x}_{1}, \tilde{x}_{2}, \tilde{x}_{3}\right)^{T}=\boldsymbol{X}-\hat{\boldsymbol{X}}$. From (7.53) and (7.52), we can show that there exist strictly positive constants $\gamma_{0}$ and $\sigma_{0}$ such that

$$
\begin{equation*}
\|(\tilde{\boldsymbol{\eta}}(t), \tilde{\boldsymbol{X}}(t))\| \leq \gamma_{0}\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\| e^{-\sigma_{0}\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0 . \tag{7.54}
\end{equation*}
$$

Define $\hat{\boldsymbol{v}}=[\hat{u}, \hat{v}, \hat{r}]^{T}$ being an estimate of the velocity vector $\boldsymbol{v}$ as

$$
\begin{equation*}
\hat{v}=P^{-1}(\eta) \hat{X} \tag{7.55}
\end{equation*}
$$

Then the velocity estimate error vector, $\tilde{\boldsymbol{v}}:=[\tilde{u}, \tilde{v}, \tilde{r}]^{T}=\boldsymbol{v}-\hat{\boldsymbol{v}}$, satisfies

$$
\begin{equation*}
\tilde{v}=P^{-1}(\eta) \tilde{X} \tag{7.56}
\end{equation*}
$$

Based on (7.56), we cannot conclude anything about the convergence of the velocity estimate errors since some elements of the matrix $\boldsymbol{P}(\boldsymbol{\eta})$, see (7.49), depend linearly on $x$ and $y$. However, our controller will guarantee that $(x, y)$ are globally bounded. Then (7.56) implies that the velocity estimate errors globally exponentially converge to zero. Indeed, the linear dependence of $\boldsymbol{P}(\boldsymbol{\eta})$ on $x$ and $y$ will result in a challenging problem, which the control design will have to take care of. To prepare for the control design, using (7.55), we rewrite (7.51) as

$$
\begin{align*}
\dot{x}= & \hat{u} \cos (\psi)-\hat{v} \sin (\psi)+\left(\cos ^{2}(\psi)+m_{11} m_{22}^{-1} \sin ^{2}(\psi)\right) \tilde{x}_{1}+ \\
& 0.5\left(m_{22}-m_{11}\right) m_{22}^{-1} \sin (2 \psi) \tilde{x}_{2}, \\
\dot{y} & =\hat{u} \sin (\psi)+\hat{v} \cos (\psi)+0.5\left(m_{22}-m_{11}\right) m_{22}^{-1} \sin (2 \psi) \tilde{x}_{1}+ \\
& \left(\sin ^{2}(\psi)+m_{11} m_{22}^{-1} \cos ^{2}(\psi)\right) \tilde{x}_{2}, \\
\dot{\psi}= & \hat{r}+m_{11} m_{33}^{-1}\left(y \tilde{x}_{1}-x \tilde{x}_{2}\right)+\tilde{x}_{3}, \\
\dot{\hat{u}}= & m_{22} m_{11}^{-1} \hat{v} \hat{r}-d_{11} m_{11}^{-1} \hat{u}+m_{11}^{-1} \tau_{u}+\hat{v}\left(m_{22} m_{33}^{-1}\left(y \tilde{x}_{1}-x \tilde{x}_{2}\right)+\right. \\
& \left.m_{22} m_{11}^{-1} \tilde{x}_{3}\right)+\cos (\psi) \tilde{x}+\sin (\psi) \tilde{y}-m_{11} m_{33}^{-1}(\sin (\psi) x-\cos (\psi) y) \tilde{\psi}, \\
\dot{\hat{v}}= & -m_{11} m_{22}^{-1} \hat{u} \hat{r}-d_{22} m_{22}^{-1} \hat{v}-\hat{u} m_{11} m_{22}^{-1}\left(m_{11} m_{33}^{-1}\left(y \tilde{x}_{1}-x \tilde{x}_{2}\right)+\tilde{x}_{3}\right)- \\
& m_{11}^{2} m_{22}^{-2}(\sin (\psi) \tilde{x}-\cos (\psi) \tilde{y})-m_{11}^{2} m_{22}^{-1} m_{33}^{-1}(\cos (\psi) x+\sin (\psi) y) \tilde{\psi}, \\
\dot{\hat{r}}= & \left(m_{11}-m_{22}\right) m_{33}^{-1} \hat{u} \hat{v}-d_{33} m_{33}^{-1} \hat{r}+m_{33}^{-1} \tau_{r}+m_{11}^{2} m_{22}^{-1} m_{33}^{-1} \hat{u} \times \\
& \left(-\sin (\psi) \tilde{x}_{1}+\cos (\psi) \tilde{x}_{2}\right)-m_{22} m_{33}^{-1} \hat{v} \times\left(\cos (\psi) \tilde{x}_{1}+\right. \\
& \left.\sin (\psi) \tilde{x}_{2}\right)+\left(m_{11}^{2}+m_{33}^{2}\right) m_{33}^{-2}\left(x^{2}+y^{2}\right) \tilde{\psi}+m_{22}^{-1} m_{33}^{-1} \times \\
& \left(\left(0.5\left(m_{11}^{2}-m_{11} m_{22}\right) \sin (2 \psi) x+\left(m_{11} m_{22} \cos ^{2}(\psi)+\right.\right.\right. \\
& \left.\left.m_{11}^{2} \sin ^{2}(\psi)\right) y\right) \tilde{x}+\left(0.5\left(m_{11}^{2}-m_{11} m_{22}\right) \sin (2 \psi) y-\right. \\
& \left.\left.\left(m_{11} m_{22} \sin ^{2}(\psi)+m_{11}^{2} \cos ^{2}(\psi)\right) x\right) \tilde{y}\right), \tag{7.57}
\end{align*}
$$

where for simplicity, we have taken $\boldsymbol{K}_{\mathbf{0 2}}=\left(\boldsymbol{J}(\eta) \boldsymbol{P}^{\mathbf{- 1}}(\eta)\right)^{\boldsymbol{T}}$.

### 7.3.2 Coordinate Transformations

If one applies the coordinate change (7.22) to (7.57), it will result in a very complicated system, namely some quadratic terms of $\left(x_{e}, y_{e}, \psi_{e}\right)$ multiplied by the observer errors appearing in the kinematic part of the transformed system due to linear
dependence of $x$ and $y$ on some elements of the matrix $\boldsymbol{P}(\eta)$. This makes the control design extremely difficult and might result in a finite escape. To avoid the said difficulty, we propose the following coordinate transformation:

$$
\left[\begin{array}{c}
x_{e}  \tag{7.58}\\
y_{e} \\
\psi_{e}
\end{array}\right]=\boldsymbol{J}^{-\mathbf{1}}\left(\eta_{\boldsymbol{d}}\right)\left[\begin{array}{c}
x-x_{d} \\
y-y_{d} \\
\psi-\psi_{d}
\end{array}\right],\left\{\begin{array}{l}
u_{e}=\hat{u}-u_{d} \\
v_{e}=\hat{v}-v_{d} \\
r_{e}=\hat{r}-r_{d}
\end{array}\right.
$$

Indeed, convergence to zero of $\left(x_{e}, y_{e}, \psi_{e}\right)$ implies that of $\left(x-x_{d}, y-y_{d}, \psi-\psi_{d}\right)$. Differentiating both sides of (7.58) along the solutions of (7.57) and (7.5) yields the error dynamics of the "kinematic part":

$$
\begin{align*}
& \dot{x}_{e}=u_{e}+\left(u_{e}+u_{d}\right)\left(\cos \left(\psi_{e}\right)-1\right)-\left(v_{e}+v_{d}\right) \sin \left(\psi_{e}\right)+r_{d} y_{e}+h_{x}, \\
& \dot{y}_{e}=v_{e}+\left(v_{e}+v_{d}\right)\left(\cos \left(\psi_{e}\right)-1\right)+\left(u_{e}+u_{d}\right) \sin \left(\psi_{e}\right)-r_{d} x_{e}+h_{y},  \tag{7.59}\\
& \dot{\psi}_{e}=r_{e}+\Omega_{\psi}+h_{\psi},
\end{align*}
$$

where, for notational simplicity, we have defined

$$
\begin{align*}
& h_{x}=\cos \left(\psi_{d}\right) \Delta_{x}+\sin \left(\psi_{d}\right) \Delta_{y}, \\
& h_{y}=-\sin \left(\psi_{d}\right) \Delta_{x}+\cos \left(\psi_{d}\right) \Delta_{y}, \\
& h_{\psi}=m_{11} m_{33}^{-1}\left(y_{d} \tilde{x}_{1}-x_{d} \tilde{x}_{2}\right)+\tilde{x}_{3}, \\
& \Omega_{\psi}=m_{11} m_{33}^{-1}\left(\left(\sin \left(\psi_{d}\right) x_{e}+\cos \left(\psi_{d}\right) y_{e}\right) \tilde{x}_{1}-\left(\cos \left(\psi_{d}\right) x_{e}-\operatorname{sins}\left(\psi_{d}\right) y_{e}\right) \tilde{x}_{2}\right), \\
& \Delta_{x}=\left(\cos ^{2}(\psi)+\sin ^{2}(\psi) m_{11} m_{22}^{-1}\right) \tilde{x}_{1}+0.5 \sin (2 \psi) \tilde{x}_{2}\left(m_{11}-m_{22}\right) m_{22}^{-1}, \\
& \Delta_{y}=0.5 \sin (2 \psi) \tilde{x}_{1}\left(m_{11}-m_{22}\right) m_{22}^{-1}+\left(\sin ^{2}(\psi)+\cos ^{2}(\psi) m_{11} m_{22}^{-1}\right) \tilde{x}_{2} . \tag{7.60}
\end{align*}
$$

We define the following coordinate transformation, which is slightly different from (7.24):

$$
\begin{equation*}
z_{e}=\psi_{e}+\arcsin \left(\frac{k u_{d} y_{e}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}}}\right) \tag{7.61}
\end{equation*}
$$

where the constant $k$ is such that $\left|k u_{d}(t)\right|<1, \forall t \geq 0$. This constant will be specified later. Using the nonlinear coordinate transformation (7.61) together with (7.59), the ship error dynamics are rewritten as

$$
\begin{align*}
\dot{x}_{e} & =u_{e}+\left(u_{e}+u_{d}\right) p_{x}-\left(v_{e}+v_{d}\right) p_{y}+\left(v_{e}+v_{d}\right) k u_{d} \varpi_{2}^{-1} y_{e}+ \\
& r_{d} y_{e}+h_{x}, \\
\dot{y}_{e} & =v_{e}+\left(v_{e}+v_{d}\right) p_{x}+\left(u_{e}+u_{d}\right) p_{y}-k u_{d}^{2} \varpi_{2}^{-1} y_{e}- \\
& k u_{d} \varpi_{2}^{-1} u_{e} y_{e}-r_{d} x_{e}+h_{y}, \\
\dot{z}_{e} & =r_{e}+f_{z}+g_{z} u_{e}+\Omega_{z}+h_{z}, \\
\dot{u}_{e} & =m_{22} m_{11}^{-1} \hat{v} \hat{r}-d_{11} m_{11}^{-1} \hat{u}+m_{11}^{-1} \tau_{u}-\dot{u}_{d}+\Omega_{u}+h_{u}, \\
\dot{v}_{e} & =-m_{11} m_{22}^{-1}\left(u_{e} r_{e}+u_{d} r_{e}+u_{e} r_{d}\right)-d_{22} m_{22}^{-1} v_{e}+\Omega_{v}+h_{v}, \\
\dot{r}_{e} & =\left(m_{11}-m_{22}\right) m_{33}^{-1} \hat{u} \hat{v}-d_{33} m_{33}^{-1} \hat{r}+m_{33}^{-1} \tau_{r}-\dot{r}_{d}+\Omega_{r}+h_{r}, \tag{7.62}
\end{align*}
$$

where, for notational simplicity and convenience of the control design, we have defined the following:

The terms $\omega_{1}, \omega_{2}, p_{x}, p_{y}, f_{z}$, and $g_{z}$ are defined as

$$
\begin{align*}
\varpi_{1} & =\sqrt{1+x_{e}^{2}+\left(1-k^{2} u_{d}^{2}\right) y_{e}^{2}}, \varpi_{2}=\sqrt{1+x_{e}^{2}+y_{e}^{2}}, \\
p_{x} & =\varpi_{2}^{-1}\left(\left(\cos \left(z_{e}\right)-1\right) \varpi_{1}+\left(\varpi_{1}-\varpi_{2}\right)+\sin \left(z_{e}\right) k u_{d} y_{e}\right), \\
p_{y} & =\varpi_{2}^{-1}\left(\sin \left(z_{e}\right) \varpi_{1}-\left(\cos \left(z_{e}\right)-1\right) k u_{d} y_{e}\right),  \tag{7.63}\\
f_{z} & =k \varpi_{1}^{-1}\left(\dot{u}_{d}+u_{d}\left(( 1 + x _ { e } ^ { 2 } ) \varpi _ { 2 } ^ { - 2 } \left(v_{e}-k u_{d}^{2} \varpi_{2}^{-1} y_{e}+\left(v_{e}+v_{d}\right) \times\right.\right.\right. \\
& \left.p_{x}+u_{d} p_{y}\right)-r_{d} x_{e}-\varpi_{2}^{-2} x_{e} y_{e}\left(u_{d} p_{x}-\left(v_{e}+v_{d}\right)\left(p_{y}-\right.\right. \\
& \left.\left.\left.\left.k u_{d} \varpi_{2}^{-1} y_{e}\right)\right)\right)\right), \\
g_{z} & =k u_{d} \varpi_{1}^{-1}\left(\left(1+x_{e}^{2}\right) \varpi_{2}^{-2}\left(-k u_{d} \varpi_{2}^{-1} y_{e}+p_{y}\right)-\varpi_{2}^{-2} y_{e} x_{e}\left(1+p_{x}\right)\right) .
\end{align*}
$$

The terms $\Omega_{z}, \Omega_{u}, \Omega_{v}$, and $\Omega_{r}$ containing states multiplied by the observer errors are defined as

$$
\begin{align*}
\Omega_{z}= & \Omega_{\psi}, \\
\Omega_{u}= & \hat{v}\left(m_{22} m_{33}^{-1}\left(y \tilde{x}_{1}-x \tilde{x}_{2}\right)+m_{22} m_{11}^{-1} \tilde{x}_{3}\right)- \\
& m_{11} m_{33}^{-1}\left(\sin \left(\psi_{e}\right) x_{e}-\cos \left(\psi_{e}\right) y_{e}\right) \tilde{\psi}, \\
\Omega_{v}= & -\left(u_{e}+u_{d}\right) m_{11} m_{22}^{-1}\left(m_{11} m_{33}^{-1}\left(y \tilde{x}_{1}-x \tilde{x}_{2}\right)+\tilde{x}_{3}\right)- \\
& m_{11}^{2} m_{22}^{-1} m_{33}^{-1}\left(\cos \left(\psi_{e}\right) x_{e}+\sin \left(\psi_{e}\right) y_{e}\right) \tilde{\psi}, \\
\Omega_{r}= & m_{11}^{2} m_{22}^{-1} m_{33}^{-1} \hat{u}\left(-\sin (\psi) \tilde{x}_{1}+\cos (\psi) \tilde{x}_{2}\right)- \\
& m_{22} m_{33}^{-1} \hat{v}\left(\cos (\psi) \tilde{x}_{1}+\sin (\psi) \tilde{x}_{2}\right)+  \tag{7.64}\\
& m_{22}^{-1} m_{33}^{-1}\left(0.5\left(m_{11}^{2}-m_{11} m_{22}\right) \sin (2 \psi) \Delta_{x d}+\right. \\
& \left.\left(m_{11} m_{22} \cos ^{2}(\psi)+m_{11}^{2} \sin ^{2}(\psi)\right) \Delta_{y d}\right) \tilde{x}+ \\
& m_{22}^{-1} m_{33}^{-1}\left(0.5\left(m_{11}^{2}-m_{11} m_{22}\right) \sin (2 \psi) \Delta_{y d}-\right. \\
& \left.\left(m_{11} m_{22} \sin ^{2}(\psi)+m_{11}^{2} \cos ^{2}(\psi)\right) \Delta_{x d}\right) \tilde{y}+ \\
& \left(m_{11}^{2} m_{33}^{-2}+1\right)\left(x_{e}^{2}+y_{e}^{2}+2 \sin \left(\psi_{d}\right)\left(x_{e} y_{d}-\right.\right. \\
& \left.\left.y_{e} x_{d}\right)+2 \cos \left(\psi_{d}\right)\left(x_{e} x_{d}+y_{e} y_{d}\right)\right) \tilde{\psi},
\end{align*}
$$

with $\Delta_{x d}=\cos \left(\psi_{d}\right) x_{e}-\sin \left(\psi_{d}\right) y_{e}$, and $\Delta_{y d}=\sin \left(\psi_{d}\right) x_{e}+\cos \left(\psi_{d}\right) y_{e}$.
The terms $h_{z}, h_{u}, h_{v}$, and $h_{r}$ containing the observer errors multiplied by bounded terms are defined as

$$
\begin{align*}
h_{z} & =k u_{d} \varpi_{1}^{-1}\left(\left(1+x_{e}^{2}\right) \varpi_{2}^{-2} h_{y}-\varpi_{2}^{-2} y_{e} x_{e} h_{x}\right)+h_{\psi}, \\
h_{u} & =\cos (\psi) \tilde{x}+\sin (\psi) \tilde{y}-m_{11} m_{33}^{-1}\left(\sin (\psi) x_{d}-\cos (\psi) y_{d}\right) \tilde{\psi}, \\
h_{v} & =-m_{11}^{2} m_{22}^{-1}(\sin (\psi) \tilde{x}-\cos (\psi) \tilde{y})-m_{11}^{2} m_{22}^{-1} m_{33}^{-1}\left(\cos (\psi) x_{d}+\sin (\psi) y_{d}\right) \tilde{\psi}, \\
h_{r} & =m_{22}^{-1} m_{33}^{-1}\left(0.5\left(m_{11}^{2}-m_{11} m_{22}\right) \sin (2 \psi)\left(x_{d} \tilde{x}-y_{d} \tilde{y}\right)+\left(m_{11} m_{22} \cos ^{2}(\psi)+\right.\right. \\
& \left.\left.m_{11}^{2} \sin ^{2}(\psi)\right) y_{d} \tilde{x}-\left(m_{11} m_{22} \sin ^{2}(\psi)+m_{11}^{2} \cos ^{2}(\psi)\right) x_{d} \tilde{y}\right)+m_{33}^{-2}\left(m_{11}^{2}+\right. \\
& \left.m_{33}^{2}\right)\left(x_{d}^{2}+y_{d}^{2}\right) \tilde{\psi} . \tag{7.65}
\end{align*}
$$

The problem of forcing the underactuated ship (7.1) to track the virtual ship (7.5) becomes one of stabilizing the system (7.62) at the origin. The effort, we have made so far is to obtain the stabilizing term $-k u_{d}^{2} y_{e} / \varpi_{2}$ in the $y_{e}$-dynamics.

### 7.3.3 Control Design

Before designing the control inputs, it is important to note that the terms $\Omega_{u}$ and $\Omega_{r}$ can be dominated by adding some nonlinear damping terms in the control inputs $\tau_{u}$ and $\tau_{r}$. However the term $\Omega_{v}$ cannot be dominated by any nonlinear damping terms in $\tau_{u}$ and $\tau_{r}$. Since $\Omega_{v}$ contains $u_{e} x_{e}$ and $u_{e} y_{e}$ multiplied by the observer errors, with $x$ and $y$ being substituted in from (7.58), if one designs a virtual control of $u_{e}$, which is linear in $x_{e}$ and $y_{e}$, the sway velocity dynamics might have a finite escape time due to the fact that separation principle does not hold for the nonlinear system in question. The coordinate change (7.58) allows us to design a virtual control of $u_{e}$ such that it is bounded for all $x_{e}$ and $y_{e}$ and stabilizes the $x_{e}$-dynamics at the origin. We design the controls $\tau_{u}$ and $\tau_{r}$ in two steps.

## Step 1

Define the virtual control errors as

$$
\begin{align*}
& \tilde{u}_{e}=u_{e}-u_{e}^{d} \\
& \tilde{r}_{e}=r_{e}-r_{e}^{d} \tag{7.66}
\end{align*}
$$

where $u_{e}^{d}$ and $r_{e}^{d}$ are the virtual velocity controls of $u_{e}$ and $r_{e}$, respectively. The virtual controls $u_{e}^{d}$ and $r_{e}^{d}$ are chosen as follows:

$$
\begin{align*}
& u_{e}^{d}=-\frac{k_{1} x_{e}}{\varpi_{2}} \\
& r_{e}^{d}=-k_{2} z_{e}-f_{z}-g_{z} u_{e}^{d} \tag{7.67}
\end{align*}
$$

where $k_{1}$ and $k_{2}$ are positive design constants to be specified later.

## Step 2

By differentiating (7.66) along the solutions of (7.62) and (7.67), the actual controls $\tau_{u}$ and $\tau_{r}$ with some nonlinear damping terms to overcome the effect of observer errors, and without canceling the useful damping terms, are chosen as

$$
\tau_{u}=m_{11}\left(-m_{22} m_{11}^{-1} \hat{v} \hat{r}+d_{11} m_{11}^{-1}\left(u_{e}^{d}+u_{d}\right)+\dot{u}_{d}+\frac{\partial u_{e}^{d}}{\partial x_{e}}\left(u_{e}+\left(u_{e}+u_{d}\right) \times\right.\right.
$$

$$
\begin{align*}
& \left.p_{x}+\left(v_{e}+v_{d}\right)\left(-p_{y}+k u_{d} \varpi_{2}^{-1} y_{e}\right)+r_{d} y_{e}\right) \frac{\partial u_{e}^{d}}{\partial y_{e}}\left(v_{e}+\left(v_{e}+v_{d}\right) p_{x}+\right. \\
& \left.\left(u_{e}+u_{d}\right) p_{y}-k u_{d} y_{e} \varpi_{2}^{-1}\left(u_{e}+u_{d}\right)-r_{d} x_{e}\right)-c_{1} \tilde{u}_{e}-g_{z} z_{e}+ \\
& \left.k_{3} k_{4}^{-1} m_{11} m_{22}^{-1} \hat{r} v_{e}-\delta_{1} \tilde{u}_{e} \tau_{u \mathrm{dam}}\right), \\
& \tau_{r}=m_{33}\left(-\left(m_{11}-m_{22}\right) m_{33}^{-1} \hat{u} \hat{v}+d_{33} m_{33}^{-1}\left(r_{e}^{d}+r_{d}\right)+\dot{r}_{d}+\right. \\
& \quad \frac{\partial r_{e}^{d}}{\partial x_{e}}\left(u_{e}+\left(u_{e}+u_{d}\right) p_{x}+\left(v_{e}+v_{d}\right)\left(-p_{y}+k u_{d} \varpi_{2}^{-1} y_{e}\right)+r_{d} y_{e}\right)+ \\
& \quad \frac{\partial r_{e}^{d}}{\partial y_{e}}\left(v_{e}+\left(v_{e}+v_{d}\right) p_{x}+\left(u_{e}+u_{d}\right) p_{y}-k u_{d} y_{e} \varpi_{2}^{-1}\left(u_{e}+u_{d}\right)-\right. \\
& \left.\quad r_{d} x_{e}\right)+\frac{\partial r_{e}^{d}}{\partial z_{e}}\left(r_{e}+f_{z}+g_{z} u_{e}\right)-\frac{\partial r_{e}^{d}}{\partial v_{e}}\left(d_{22} m_{22}^{-1} v_{e}+m_{11} m_{22}^{-1}\left(u_{e} r_{e}+\right.\right. \\
& \left.\left.\quad u_{d} r_{e}+u_{e} r_{d}\right)\right)+\frac{\partial r_{e}^{d}}{\partial u_{d}} \dot{u}_{d}+\frac{\partial r_{e}^{d}}{\partial \dot{u}_{d}} \ddot{u}_{d}+\frac{\partial r_{e}^{d}}{\partial v_{d}} \dot{v}_{d}+\frac{\partial r_{e}^{d}}{\partial r_{d}} \dot{r}_{d}-c_{2} \tilde{r}_{e}-z_{e}+ \\
& \left.k_{3} k_{4}^{-1} m_{11} m_{22}^{-1}\left(u_{e}^{d}+u_{d}\right) v_{e}-\delta_{2} \tilde{r}_{e} \tau_{r \mathrm{dam}}\right), \tag{7.68}
\end{align*}
$$

where $c_{1}, c_{2}, k_{3}$, and $k_{4}$ are positive constants to be specified later, $\delta_{1}$ and $\delta_{2}$ are arbitrarily positive constants. We introduced the ratio $k_{3} / k_{4}$ to enhance the feasibility of design constants. The nonlinear damping terms $\tau_{u \text { dam }}$ and $\tau_{r \text { dam }}$ are defined as:

$$
\begin{align*}
& \tau_{u \mathrm{dam}}=\left(\hat{v}^{2}+1\right)\left(x_{e}^{2}+y_{e}^{2}\right)+\hat{v}^{2}, \\
& \tau_{r \mathrm{dam}}=\hat{u}^{2}+\hat{v}^{2}+\left(x_{e}^{2}+y_{e}^{2}\right)^{2}+\left(x_{e}^{2}+y_{e}^{2}\right)\left(\hat{u}^{2}+1\right) . \tag{7.69}
\end{align*}
$$

We now state the second main result of this chapter, the proof of which is given in the next section.

Theorem 7.2. Under Assumption 7.2, assume that (a) there are no environmental disturbances; (b) the ship parameters are known; (c) the reference signals $\left(x_{d}\right.$, $y_{d}, \psi_{d}, v_{d}$ ) are generated by the virtual ship model (7.5), and Assumption 7.1 holds. If the output feedback control law (7.68) together with the observer (7.51) are applied to the ship system (7.1), then the tracking errors $\left(x(t)-x_{d}(t), y(t)-\right.$ $\left.y_{d}(t), \psi(t)-\psi_{d}(t), v(t)-v_{d}(t)\right)$ globally asymptotically and locally exponentially converge to zero with an appropriate choice of the design constants $k, c_{1}, c_{2}$, and $k_{i}, i=1, \ldots, 4$.

Remark 7.4. The main differences between the partial-state feedback and output feedback designs are the nonlinear coordinate transformations (7.6), (7.24), (7.43), (7.58), and (7.61). Furthermore, the partial-state feedback controller can allow the reference trajectory $\left(x_{d}, y_{d}\right)$ to exponentially grow but the output feedback controller cannot. This is because the observer errors of the partial-state feedback design do not depend on $x$ and $y$ while those of the output feedback design depend linearly on $x$ and $y$. Indeed, the output feedback control design can directly yield a
controller for the partial-state feedback case but not vice versa. We have, however, presented both control designs for the sake of completeness.

### 7.3.4 Stability Analysis

Substituting (7.68), (7.67), and (7.66) into (7.62) results in the following closed loop system:

$$
\begin{align*}
& \dot{x}_{e}=-k_{1} \varpi_{2}^{-1} x_{e}+\left(-k_{1} \varpi_{2}^{-1} x_{e}+u_{d}\right) p_{x}-\left(v_{e}+v_{d}\right) p_{y}+ \\
&\left(v_{e}+v_{d}\right) k u_{d} \varpi_{2}^{-1} y_{e}+r_{d} y_{e}+h_{x}+\tilde{u}_{e}\left(1+p_{x}\right), \\
& \dot{y}_{e}= v_{e}+\left(v_{e}+v_{d}\right) p_{x}+\left(-k_{1} \varpi_{2}^{-1} x_{e}+u_{d}\right) p_{y}-k u_{d}^{2} \varpi_{2}^{-1} y_{e}+ \\
& k_{1} k u_{d} \varpi_{2}^{-1} x_{e} y_{e}-r_{d} x_{e}+h_{y}+\tilde{u}_{e}\left(1-k u_{d} \varpi_{2}^{-1} y_{e}\right), \\
& \dot{z}_{e}=-k_{2} z_{e}+\Omega_{z}+h_{z}+g_{z} \tilde{u}_{e}+\tilde{r}_{e}, \\
& \dot{\tilde{u}}_{e}=-\left(c_{1}+d_{11} m_{11}^{-1}\right) \tilde{u}_{e}-g_{z} z_{e}+\frac{k_{3}}{k_{4}} \frac{m_{11}}{m_{22}} \hat{r} v_{e}-\delta_{1} \tilde{u}_{e} \tau_{u \mathrm{dam}}+ \\
& \quad \Omega_{u}+h_{u}-\frac{\partial u_{e}^{d}}{\partial x_{e}} h_{x}-\frac{\partial u_{e}^{d}}{\partial y_{e}} h_{y}, \\
& \dot{v}_{e}=-d_{22} m_{22}^{-1} v_{e}-m_{11} m_{22}^{-1}\left(u_{e}^{d} r_{e}^{d}+u_{d} r_{e}^{d}+u_{e}^{d} r_{d}\right)-m_{11} m_{22}^{-1} \times \\
& \quad\left(u_{e}^{d}+u_{d}\right) \tilde{r}_{e}-m_{11} m_{22}^{-1} \hat{r} \tilde{u}_{e}+\Omega_{v}+h_{v}, \\
& \dot{\tilde{r}}_{e}=-\left(c_{2}+d_{33} m_{33}^{-1}\right) \tilde{r}_{e}-z_{e}-\frac{\partial r_{e}^{d}}{\partial x_{e}} h_{x}-\frac{\partial r_{e}^{d}}{\partial y_{e}} h_{y}-\frac{\partial r_{e}^{d}}{\partial z_{e}}\left(\Omega_{z}+h_{z}\right)- \\
&  \tag{7.70}\\
& \frac{\partial r_{e}^{d}}{\partial v_{e}}\left(\Omega_{v}+h_{v}\right)+k_{3} k_{4}^{-1} m_{11} m_{22}^{-1}\left(u_{e}^{d}+u_{d}\right) v_{e}-\delta_{2} \tilde{r}_{e} \tau_{r \mathrm{dam}}+\Omega_{r}+h_{r},
\end{align*}
$$

where, for brevity, we did not substitute the expressions of $u_{e}^{d}$ and $r_{e}^{d}$ into the sway dynamics. To prove Theorem 7.2, we just need to show that the closed loop system (7.70) is globally asymptotically and locally exponentially stable at the origin. It is noted that $\Omega_{z}$ contains $x_{e}$ and $y_{e}$ multiplied by the observer errors, see (7.64) and (7.60). On the other hand, the $x_{e}$ and $y_{e}$-dynamics are stabilized by the terms $-k_{1} x_{e} / \varpi_{2}$ and $-k u_{d}^{2} y_{e} / \varpi_{2}$, respectively. This makes the stability analysis of (7.70) difficult, i.e., we cannot consider the $\left(x_{e}, y_{e}, v_{e}\right)$ and $\left(z_{e}, \tilde{u}_{e}, \tilde{r}_{e}\right)$-subsystems separately as is often done in applying stability results for cascade systems. To illustrate our idea of proving asymptotic stability of (7.70), we first give a simple example. For any initial conditions ( $\left.\xi_{1}\left(t_{0}\right), \xi_{2}\left(t_{0}\right)\right)$, the system

$$
\begin{align*}
& \dot{\xi}_{1}=-\frac{\xi_{1}}{\sqrt{1+\xi_{1}^{2}}}+\phi\left(\xi_{2}\right), \\
& \dot{\xi}_{2}=-\xi_{2}+\xi_{1} e^{-\sigma_{\xi_{2}}\left(t-t_{0}\right)} \tag{7.71}
\end{align*}
$$

is GAS at the origin for $t \geq t_{0} \geq 0$, any $\sigma_{\xi 2}>0$, and $\left|\phi\left(\xi_{2}(t)\right)\right| \leq A\left|\xi_{2}(t)\right|$ with $A$ being any positive constant. In the first step, we show that there exists a positive constant $\sigma_{\xi 1}<\sigma_{\xi 2}$ such that $\left|\xi_{1}(t)\right| \leq \lambda_{1}(\cdot) e^{\sigma_{\xi 1}\left(t-t_{0}\right)}$, with $\lambda_{1}(\cdot)$ being a nondecreasing function of $\left\|\left(\xi_{1}\left(t_{0}\right), \xi_{2}\left(t_{0}\right)\right)\right\|$. Take the Lyapunov function $W=0.5\left(\xi_{1}^{2}+K \xi_{2}^{2}\right)$, with $K$ being a positive constant, whose derivative along the solution of (7.71) satisfies

$$
\begin{equation*}
\dot{W} \leq \sigma_{\xi 1} \xi_{1}^{2}+\left(0.25 A^{2} / \sigma_{\xi 1}-K\right) \xi_{2}^{2}+K \xi_{1} \xi_{2} e^{-\sigma_{\xi 2}\left(t-t_{0}\right)} \tag{7.72}
\end{equation*}
$$

Pick $K$ such that $A^{2} /\left(4 \sigma_{\xi 1}\right)-K<0$, then (7.72) implies that

$$
\begin{equation*}
\dot{W} \leq 2 \sigma_{\xi 1} W+K \max (1, K) W e^{-\sigma_{\xi 2}\left(t-t_{0}\right)}, \tag{7.73}
\end{equation*}
$$

which in turn yields $\left|\xi_{1}(t)\right| \leq \lambda_{1}(\cdot) e^{\sigma_{\xi 1}\left(t-t_{0}\right)}$. Therefore the term $\xi_{1} e^{-\sigma_{\xi 2}\left(t-t_{0}\right)}$ in (7.71) globally exponentially converges to zero. The second step of proving asymptotic stability can be carried out easily by taking the Lyapunov function $W_{1}=\sqrt{1+\xi_{1}^{2}}-1+K_{1} \xi_{2}^{2}, K_{1}>0$. We now present the proof of asymptotic stability of the closed loop system (7.70) in two parts.

Part 1. In this part, we show that there exists a positive constant $\sigma_{1}<\sigma_{0}$ such that

$$
\begin{equation*}
\left\|\left(x_{e}(t), y_{e}(t)\right)\right\| \leq \gamma_{11}(\cdot) e^{\sigma_{1}\left(t-t_{0}\right)}+\gamma_{10}(\cdot) \tag{7.74}
\end{equation*}
$$

where $\gamma_{11}(\cdot)$ and $\gamma_{10}(\cdot)$ are nondecreasing functions of $\left\|\boldsymbol{\Xi}\left(t_{0}\right)\right\|$ with

$$
\begin{aligned}
& \boldsymbol{\Xi}\left(t_{0}\right):=\left[\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right), \boldsymbol{X}_{\boldsymbol{e}}\left(t_{0}\right)\right]^{T}, \\
& \boldsymbol{X}_{\boldsymbol{e}}:=\left[x_{e}, y_{e}, z_{e}, v_{e}, \tilde{u}_{e}, \tilde{r}_{e}\right]^{T} .
\end{aligned}
$$

Consider the following Lyapunov function:

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(x_{e}^{2}+y_{e}^{2}+k_{3} v_{e}^{2}+k_{4}\left(z_{e}^{2}+\tilde{u}_{e}^{2}+\tilde{r}_{e}^{2}\right)\right) \tag{7.75}
\end{equation*}
$$

with $k_{4}$ being a positive constant, whose time derivative along the solutions of (7.70), after some lengthy but simple calculation by completing squares satisfies

$$
\begin{align*}
\dot{V}_{1} \leq & -k_{1} \varpi_{2}^{-1} x_{e}^{2}-k u_{d}^{2} \varpi_{2}^{-1} y_{e}^{2}+7 \varepsilon_{1}\left(x_{e}^{2}+y_{e}^{2}\right)-a_{z} z_{e}^{2}-a_{v} v_{e}^{2}-a_{u} \tilde{u}_{e}^{2}- \\
& a_{r} \tilde{r}_{e}^{2}+\left(\chi_{11}(\cdot) V_{1}+\chi_{10}(\cdot)\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+a_{0}, \tag{7.76}
\end{align*}
$$

where

$$
\begin{align*}
& a_{z}=k_{2} k_{4}-0.25 \varepsilon_{1}^{-1}\left(k_{2} k_{3} m_{11} m_{22}^{-1}\left(k_{1}+\left|u_{d}\right|\right)\right)^{2}-\varepsilon_{1}^{-1}\left(k_{1}+\left|u_{d}\right|\right)^{2} \\
& a_{u}=k_{4}\left(c_{1}+d_{11} m_{11}^{-1}\right)-5 \varepsilon_{1}^{-1} \\
& a_{v}=k_{3} d_{22} m_{22}^{-1}-6.75 \varepsilon_{1}^{-1}-k_{3} m_{11} m_{22}^{-1}\left(k_{1}+\left|u_{d}\right|\right)\left(8+\left|k u_{d}\right|\right)\left|k u_{d}\right| \\
& a_{r}=k_{4}\left(c_{2}+d_{33} m_{33}^{-1}\right) \tag{7.77}
\end{align*}
$$

with $\varepsilon_{1}$ and $a_{0}$ being positive constants. We choose the design constants $k, c_{1}, c_{2}$, and $k_{i}, i=1, \ldots, 4$ such that

$$
\begin{equation*}
\varepsilon_{1}<\sigma_{0} / 7, a_{z}>0, a_{v}>0, a_{u}>0, a_{r}>0 \tag{7.78}
\end{equation*}
$$

It is not hard to show that there always exist $k, c_{1}, c_{2}$ and $k_{i}, i=1, \ldots, 4$ such that (7.78) holds for arbitrarily small $\varepsilon_{1}$. We will discuss more detail of (7.78) later. With the choice of (7.78), we can write (7.76) as

$$
\begin{align*}
\dot{V}_{1} \leq & 14 \varepsilon_{1}\left(V_{1}+\left(a_{0}+\chi_{10}(\cdot)\right) /\left(14 \varepsilon_{1}\right)\right)+ \\
& \chi_{11}(\cdot)\left(V_{1}+\left(a_{0}+\chi_{10}(\cdot)\right) /\left(14 \varepsilon_{1}\right)\right) e^{-\sigma_{0}\left(t-t_{0}\right)}, \tag{7.79}
\end{align*}
$$

which, together with (7.75), yields (7.74).
Part 2. We now prove that the closed loop system (7.70) is GAS at the origin by taking the Lyapunov function

$$
\begin{equation*}
V_{2}=\sqrt{1+x_{e}^{2}+y_{e}^{2}}-1+\frac{1}{2}\left(k_{3} v_{e}^{2}+k_{4}\left(z_{e}^{2}+\tilde{u}_{e}^{2}+\tilde{r}_{e}^{2}\right)\right) \tag{7.80}
\end{equation*}
$$

whose time derivative along the solutions of (7.70), after some lengthy but simple calculation by completing squares, and using (7.74), satisfies

$$
\begin{align*}
\dot{V}_{2} \leq & -\mu_{x} \varpi_{2}^{-2} x_{e}^{2}-\mu_{y}(t) \varpi_{2}^{-2} y_{e}^{2}-\mu_{z}(t) z_{e}^{2}-\mu_{v}(t) v_{e}^{2}-\mu_{u} \tilde{u}_{e}^{2}-\mu_{r} \tilde{r}_{e}^{2}+ \\
& \chi_{21}(\cdot) V_{2} e^{-\sigma_{2}\left(t-t_{0}\right)}+\chi_{20}(\cdot) e^{-\sigma_{2}\left(t-t_{0}\right)}, \tag{7.81}
\end{align*}
$$

where

$$
\begin{align*}
& \mu_{x}=k_{1}-6 \rho_{1}, \\
& \mu_{y}(t)=k u_{d}^{2}-\left(k u_{d}\right)^{2}\left(k_{1}+\left|u_{d}\right|+\left|v_{d}\right|+0.5\right)-\left|k u_{d}\right|\left(k_{1}+\rho_{1}\right)- \\
& \quad 0.25 \rho_{1}^{-1}\left(k u_{d} v_{d}\right)^{2}-10 \rho_{1}, \\
& \mu_{z}(t)=k_{2} k_{4}-\rho_{1}^{-1}\left(0.5\left(k_{1}+\left|u_{d}\right|\right)^{2}+5\left(1+v_{d}^{2}\right)+0.25\left(k_{3} m_{11} m_{22}^{-1} \times\right.\right. \\
& \left.\left.\quad\left(k_{1}+\left|u_{d}\right|\right)\right)^{2}\left(k_{2}^{2}+\left(k u_{d}\right)^{2}\left(k_{1}+\left|u_{d}\right|+v_{d}^{2}\right)+4+4 v_{d}^{2}\right)\right), \\
& \mu_{v}(t)=k_{3} d_{22} m_{22}^{-1}-7 \rho_{1}-0.25 \rho_{1}^{-1}\left(\left|k u_{d}\right|+1\right)-0.5\left(k u_{d}\right)^{2}- \\
& \quad k_{3} m_{11} m_{22}^{-1}\left(k_{1}+\left|u_{d}\right|\right)\left|k u_{d}\right|\left(8+\left|k u_{d}\right|\right)-0.25 \rho_{1}^{-1}\left(k_{3} m_{11} m_{22}^{-1} \times\right. \\
& \left.\quad\left(k_{1}+\left|u_{d}\right|\right)\right)^{2}\left(\left(k \dot{u}_{d}\right)^{2} /\left(1-\left(k u_{d}\right)^{2}\right)+\left(k u_{d}\right)^{4}\left(u_{d}^{2}+k_{1}^{2}+\right.\right. \\
& \left.\left.\quad\left(k u_{d} v_{d}\right)^{2}+r_{d}^{2}+2\right)+\left(k u_{d}\right)^{2}\right), \\
& \mu_{u}=k_{4}\left(c_{1}+d_{11} / m_{11}\right)-29 /\left(4 \rho_{1}\right), \\
& \mu_{r}=k_{4}\left(c_{2}+d_{33} / m_{33}\right), \tag{7.82}
\end{align*}
$$

where $\sigma_{2}=\sigma_{0}-\sigma_{1}, \rho_{1}$ is a positive constant, and $\chi_{21}(\cdot)$ and $\chi_{20}(\cdot)$ are nondecreasing functions of $\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right), \boldsymbol{X}_{\boldsymbol{e}}\left(t_{0}\right)\right)\right\|$. We now choose the constants $k, c_{1}, c_{2}$, and $k_{i}, i=1, \ldots, 4$, such that

$$
\begin{align*}
& \mu_{x} \geq \mu_{x}^{*}, \\
& \mu_{y}(t) \geq \mu_{y}^{*}, \\
& \mu_{z}(t) \geq \mu_{z}^{*}, \\
& \mu_{v}(t) \geq \mu_{v}^{*},  \tag{7.83}\\
& \mu_{u} \geq \mu_{u}^{*}, \\
& \mu_{r} \geq \mu_{r}^{*}
\end{align*}
$$

for all $t \geq 0$, where $\mu_{x}^{*}, \mu_{y}^{*}, \mu_{z}^{*}, \mu_{v}^{*}, \mu_{u}^{*}$, and $\mu_{r}^{*}$ are positive constants. Substituting (7.83) into (7.81) yields

$$
\begin{align*}
& \dot{V}_{2} \leq-\mu_{x}^{*} \varpi_{2}^{-2} x_{e}^{2}-\mu_{y}^{*} \varpi_{2}^{-2} y_{e}^{2}-\mu_{z}^{*} z_{e}^{2}-\mu_{v}^{*} v_{e}^{2}-\mu_{u}^{*} \tilde{u}_{e}^{2}-\mu_{r}^{*} \tilde{r}_{e}^{2}+ \\
& \quad \chi_{21}(\cdot) V_{2} e^{-\sigma_{2}\left(t-t_{0}\right)}+\chi_{20}(\cdot) e^{-\sigma_{2}\left(t-t_{0}\right)} . \tag{7.84}
\end{align*}
$$

From (7.84) we have $\dot{V}_{2} \leq \chi_{21}(\cdot) V_{2} e^{-\sigma_{2}\left(t-t_{0}\right)}+\chi_{20}(\cdot) e^{-\sigma_{2}\left(t-t_{0}\right)}$, which implies that $V_{2} \leq \chi_{22}(\cdot)$, with $\chi_{22}(\cdot)$ being a nondecreasing function of $\left\|\Xi\left(t_{0}\right)\right\|$. With $V_{2} \leq \chi_{22}(\cdot)$ in mind, one can show from (7.84) that there exists $\sigma_{3}>0$ depending on $\left\|\boldsymbol{\Xi}\left(t_{0}\right)\right\|$ such that $\left\|\boldsymbol{X}_{\boldsymbol{e}}(t)\right\| \leq \gamma_{2}(\cdot) e^{-\sigma_{3}\left(t-t_{0}\right)}$, where $\gamma_{2}(\cdot)$ is a nondecreasing function of $\left\|\boldsymbol{\Xi}\left(t_{0}\right)\right\|$, which in turn implies that the closed loop system (7.70) is asymptotically stable at the origin. However one can straightforwardly show that (7.70) is also locally exponentially stable at the origin. By carrying out a similar arguments in Section 7.2.5, one can show that there always exist the design constants $k, c_{1}$, and $c_{2}$ and $k_{i}, i=1,2,3,4$, such that $\left|k u_{d}(t)\right|<1$ and that the conditions (7.78) and (7.83) hold.

### 7.4 Robustness Discussion

In this section, we discuss robustness of the output feedback tracking controller. A discussion for the partial-state feedback can be carried out in a similar way. The control law (7.68) has been designed under the assumption that there are no environmental disturbances. Indeed, this assumption is unrealistic in practice. The aim of this section is to discuss the robustness property of our proposed controller in relation to environmental disturbances. Under additive environmental disturbances, it is not hard to show that the observer (7.51) guarantees that the observer errors ( $\tilde{\boldsymbol{\eta}}(t), \tilde{\boldsymbol{X}}(t))$ globally exponentially converge to a ball centered at the origin. Moreover, one can prove that the tracking error vector $\boldsymbol{X}_{\boldsymbol{e}}(t)$ also globally asymptotically converges to a ball centered at the origin. The radius of this ball can be adjusted by changing the control gains if the environmental disturbances are not too large. When the environmental disturbances are large enough, the observer (7.51) cannot provide a sufficiently accurate estimate of unmeasured velocities, and the control law (7.68) cannot compensate for considerable environmental disturbances acting on the sway axis. These will result in an unstable closed loop system, especially at a low forward speed. This phenomenon should not be surprising since the vessel in
question is not actuated in the sway axis and does not have velocities available for feedback. One can see this phenomenon by observing the simple example system (7.71) with some additive disturbance in the first equation. It is easy to show that when this additive disturbance has magnitude larger than 1, then the system (7.71) will be unstable. The robustness issue is still a challenging problem in control of underactuated ocean vehicles and underactuated systems in general.

### 7.5 Simulations

This section illustrates the soundness of the control laws (7.68) by simulating them on the same monohull ship in the previous two chapters. Details of the ship parameters are listed in Section 5.4.

The initial conditions of the reference trajectories are chosen as $\left(x_{d}(0), y_{d}(0)\right.$, $\left.\psi_{d}(0), v_{d}(0)\right)=\left(0 \mathrm{~m}, 0 \mathrm{~m}, 0 \mathrm{rad}, 0 \mathrm{~ms}^{-1}\right)$. The reference velocities are $u_{d}=4 \mathrm{~ms}^{-1}$ and $r_{d}=0 \mathrm{rads}^{-1}$ for the first 300 seconds, and $u_{d}=4 \mathrm{~ms}^{-1}$ and $r_{d}=0.02 \mathrm{rads}^{-1}$ for the rest of the simulation time. This choice means that the reference trajectory is a straight line for the first 300 seconds and then followed by a circle with a radius of 200 m . Indeed the above choice of reference velocities satisfies Assumption 7.1. All of the initial conditions of $\hat{\boldsymbol{\eta}}$ and $\hat{\boldsymbol{v}}$ are chosen to be zero. We first choose $k_{2}=5, c_{1}=1, c_{2}=2, \delta_{1}=\delta_{2}=0.1, \boldsymbol{K}_{\mathbf{0 1}}=10 \operatorname{diag}(1,1,1)$, $\boldsymbol{P}_{\mathbf{0 1}}=\boldsymbol{P}_{\mathbf{0 2}}=0.5 \operatorname{diag}(1,1,1)$, and $\boldsymbol{K}_{\mathbf{0 2}}=\left(\boldsymbol{J}(\boldsymbol{\eta}) \boldsymbol{P}^{\mathbf{- 1}}(\eta)\right)^{T}$. The other design constants are chosen as: $k=0.05, k_{1}=1, k_{3}=5$, and $k_{4}=100$. Simulation results are plotted in Figures 7.1 and 7.2 with the initial conditions of the ship: $x=-10 \mathrm{~m}, y=10 \mathrm{~m}, \psi=0.1 \mathrm{rad}, u=0 \mathrm{~ms}^{-1}, v=0 \mathrm{~ms}^{-1}$, and $r=0 \mathrm{rads}^{-1}$. It is seen from Figures 7.1a and 7.1b that all of the tracking and observer errors converge to zero. The control inputs, $\tau_{u}$ and $\tau_{r}$, are within their limits. As always, the magnitude of $\tau_{u}$ and $\tau_{r}$ can be reduced by adjusting the control gains such as $c_{1}$ and $c_{2}$. However, the trade-off is a longer transient response time. For clarity, we only plot the tracking errors for the first 180 seconds, and the observer errors for the first 10 seconds. To illustrate robustness of our proposed controller, we also simulate with the same control gains and initial conditions chosen as above, and the environmental disturbance vector $\boldsymbol{\tau}_{\boldsymbol{w}}(t)=0.5 \boldsymbol{M}\left([\sin (t), \cos (t), \sin (t)]^{T}+1.5\right)$, i.e., the last equation of (7.1) is in the form of $\boldsymbol{M} \dot{\boldsymbol{v}}=-\boldsymbol{C}(\boldsymbol{v}) \boldsymbol{v}-\boldsymbol{D} \boldsymbol{v}+\boldsymbol{\tau}+\boldsymbol{\tau}_{\boldsymbol{w}}(\boldsymbol{t})$. Simulation results are plotted in Figures 7.3 and 7.4. For clarity, we only plot the tracking errors for the first 300 seconds and the observer errors for the first 10 seconds. It can be seen that the environmental disturbances deteriorate the performance of the controlled loop system in the sense that the tracking errors do not converge to zero but to a ball centered at the origin. This shows an important property of robustness of the controlled system with respect to the environmental disturbances.

### 7.6 Conclusions

The key to the control developments is an introduction of the global nonlinear coordinate transformations (7.6), (7.24), (7.43), (7.58), and (7.61) to obtain an exponential observer and to transform the tracking error dynamics to a suitable nonlinear system, to which Lyapunov's direct method and the backstepping technique can be applied. The work presented in this chapter is based on [94, 114, 116, 117].


Figure 7.1 Simulation results without disturbances: a. Tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]:$ dashed line, $\left.\psi-\psi_{d} \quad[\mathrm{rad}]\right]:$ dotted line); b. Observer errors $\left(\sqrt{(x-\hat{x})^{2}+(y-\hat{y})^{2}+(\psi-\hat{\psi})^{2}}\right.$ : solid line, $\left.\sqrt{(u-\hat{u})^{2}+(v-\hat{v})^{2}+(r-\hat{r})^{2}}\right):$ dotted line; $\mathbf{c}$. Ship position and orientation in the $(x, y)$-plane


Figure 7.2 Simulation results without disturbances (control inputs): a. Surge force; b. Yaw moment


Figure 7.3 Simulation results with disturbances: a. Tracking errors $\left(x-x_{d}[\mathrm{~m}]\right.$ : solid line, $y-y_{d}[\mathrm{~m}]:$ dashed line, $\left.\psi-\psi_{d}[\mathrm{rad}]\right]:$ dotted line); b. Observer errors $\left(\sqrt{(x-\hat{x})^{2}+(y-\hat{y})^{2}+(\psi-\hat{\psi})^{2}}\right.$ : solid line, $\left.\sqrt{(u-\hat{u})^{2}+(v-\hat{v})^{2}+(r-\hat{r})^{2}}\right)$ : dotted line; $\mathbf{c}$. Ship position and orientation in the $(x, y)$-plane


Figure 7.4 Simulation results with disturbances (control inputs): a. Surge force; b. Yaw moment

# Chapter 8 <br> Path-tracking Control of Underactuated Ships 

This chapter deals with the problem of designing controllers to force an underactuated surface ship with standard dynamics to track a reference path. Both full-state feedback and output feedback cases are considered. In comparison with the preceding three chapters, the requirement that the reference trajectories be generated by virtual ships is relaxed, and the control design is simpler and more amenable for implementation in practice. The control development is based on a series of coordinate transformations, the backstepping technique, and utilizing the dynamic structure of the ship.

### 8.1 Full-State Feedback

### 8.1.1 Control Objective

In addition to the assumptions made in Section 3.4.1.1, we ignore the nonlinear damping terms. The case with the nonlinear damping terms requires more effort, and is treated in Section 8.2 where an output feedback control design is addressed. As such, the resulting mathematical model of the underactuated ship moving in surge, sway, and yaw is rewritten as

$$
\begin{align*}
& \dot{\eta}=\boldsymbol{J}(\psi) \boldsymbol{v}, \\
& \boldsymbol{M} \dot{\boldsymbol{v}}=-\boldsymbol{C}(\boldsymbol{v}) \boldsymbol{v}-\boldsymbol{D} \boldsymbol{v}+\boldsymbol{\tau} \tag{8.1}
\end{align*}
$$

with

$$
\eta=\left[\begin{array}{l}
x \\
y \\
\psi
\end{array}\right], \boldsymbol{v}=\left[\begin{array}{l}
u \\
v \\
r
\end{array}\right], \boldsymbol{\tau}=\left[\begin{array}{c}
\tau_{u} \\
0 \\
\tau_{r}
\end{array}\right],
$$

$$
\begin{gathered}
\boldsymbol{J}(\psi)=\left[\begin{array}{ccc}
\cos (\psi) & -\sin (\psi) & 0 \\
\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right], \\
\boldsymbol{M}=\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & m_{23} \\
0 & m_{32} & m_{33}
\end{array}\right], \boldsymbol{D}=\left[\begin{array}{ccc}
d_{11} & 0 & 0 \\
0 & d_{22} & d_{23} \\
0 & d_{32} & d_{33}
\end{array}\right], \\
\boldsymbol{C}(\boldsymbol{v})=\left[\begin{array}{ccc}
0 & 0 & -m_{22} v-\frac{1}{2}\left(m_{23}+m_{32}\right) r \\
0 & 0 & m_{11} u \\
m_{22} v+\frac{1}{2}\left(m_{23}+m_{32}\right) r-m_{11} u & 0
\end{array}\right],
\end{gathered}
$$

where a definition of all symbols is given in Section 3.3.
Moreover, in order to clearly illustrate our idea we did not include the disturbances in (8.1) but will address this issue in Section 8.1.5 by a simple modification of the control design. In this section, we consider the following control objective.

Control Objective. Design the controls $\tau_{u}$ and $\tau_{r}$ such that the position $(x, y)$ of the ship (8.1) globally tracks a reference path $\boldsymbol{\Omega}$ parameterized by ( $x_{d}(s), y_{d}(s)$ ) with $s$ being the path parameter, the ship total linear velocity is tangential to the reference trajectory $\boldsymbol{\Omega}$, and the desired surge velocity can be adjusted on-line.

Since the ship system (8.1) is underactuated, it is not expected to force the ship to track an arbitrary path $\boldsymbol{\Omega}$. We here impose the following sufficient conditions on the path $\boldsymbol{\Omega}$ :

Assumption 8.1. There exist strictly positive constants $\varepsilon_{i}, 1 \leq i \leq 4$, such that

$$
\begin{align*}
& x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s) \geq \varepsilon_{1}, \forall s \in \mathbb{R},  \tag{8.2}\\
& \frac{\sqrt{2}}{2} \alpha \bar{u}_{d}(t)-|\gamma| \geq \varepsilon_{2}, \quad \forall t \geq 0,  \tag{8.3}\\
& \beta \bar{u}_{d}(t)+\dot{\bar{u}}_{d}(t) \geq \varepsilon_{3}, \quad \forall t \geq 0, \tag{8.4}
\end{align*}
$$

and the solutions of the following differential equation

$$
\begin{align*}
& \left(\bar{u}_{d} \cos (\delta)(1-\alpha)+\gamma\right) \dot{\delta}=-\left(\beta \bar{u}_{d}+\dot{\bar{u}}_{d}\right) \sin (\delta)-\left(\alpha \bar{u}_{d} \cos (\delta)-\gamma\right) \bar{r}_{d} \\
& \delta\left(t_{0}\right)=0 \tag{8.5}
\end{align*}
$$

satisfy

$$
\begin{equation*}
|\delta(t)| \leq \frac{\pi}{4}-\varepsilon_{4}, \quad \forall t \geq t_{0} \geq 0 \tag{8.6}
\end{equation*}
$$

where $x_{d}^{\prime}(s), y_{d}^{\prime}(s), \bar{u}_{d}$ and $\bar{r}_{d}$ are defined as

$$
\begin{align*}
& x_{d}^{\prime}(s)=\frac{\partial x_{d}}{\partial s}, y_{d}^{\prime}(s)=\frac{\partial y_{d}}{\partial s} \\
& \bar{u}_{d}=\sqrt{x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)} \dot{s}  \tag{8.7}\\
& \bar{r}_{d}=\frac{\left(x_{d}^{\prime}(s) y_{d}^{\prime \prime}(s)-x_{d}^{\prime \prime}(s) y_{d}^{\prime}(s)\right)}{\left(x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)\right)} \dot{s},
\end{align*}
$$

and the constants $\alpha, \beta$, and $\gamma$ are

$$
\begin{equation*}
\alpha=\frac{m_{11}}{m_{22}}, \beta=\frac{d_{22}}{m_{22}}, \gamma=\frac{d_{22} m_{23}}{m_{22}^{2}-d_{23} / m_{22}} . \tag{8.8}
\end{equation*}
$$

## Remark 8.1.

1. Condition (8.2) implies that the path $\boldsymbol{\Omega}$ is regular with respect to the path parameter $s$. If the path is not regular, one can break it into different regular paths and consider each path separately.
2. Conditions (8.3) and (8.4) mean that the path $\boldsymbol{\Omega}$ does not contain or approach a set-point, i.e., we do not consider stabilization/regulation problems. These conditions also imply that $\bar{u}_{d}(t)$ should not vary too fast. For clarity we only consider the forward tracking case, i.e. the reference velocity $\bar{u}_{d}(t)$ is always larger than a positive constant. The backward tracking case (i.e. $\bar{u}_{d}(t)$ is always less than a negative constant) can be done similarly.
3. Conditions (8.5) and (8.6) imply that the desired surge velocity is always larger than the desired sway velocity. These conditions are reasonable since the ship in question is underactuated in the sway axis, one should not expect to force it to track a path with an arbitrarily large curvature. These conditions place certain restrictions on the path $\boldsymbol{\Omega}$ but allow many types of trajectories such as a straight line, an arc, and a sinusoidal path. For vessels with three planes of symmetry (like spherical underwater vehicles see Section 3.4.1.3, i.e., $\alpha=1, \gamma=0$ ), the differential equation (8.5) reduces to an algebraic equation, and condition (8.6) holds if

$$
\begin{equation*}
\frac{\bar{u}_{d}(t)\left|\bar{r}_{d}(t)\right|}{\beta \bar{u}_{d}(t)+\dot{\bar{u}}_{d}(t)} \leq 1-\varepsilon_{5}, \forall t \geq 0 \tag{8.9}
\end{equation*}
$$

where $\varepsilon_{5}$ is an arbitrarily small positive constant. For other cases, an application of Lemma 9.3 in [6] to (8.5) readily shows that sufficient conditions such that (8.6) holds can be expressed as

$$
\begin{gather*}
0.5 \sqrt{2} \bar{u}_{d}(t)(1-\alpha)-|\gamma|>\varepsilon_{6}, \forall t \geq 0,  \tag{8.10}\\
\frac{\alpha \bar{u}_{d}(t)+|\gamma|}{\beta \bar{u}_{d}(t)+\dot{\bar{u}}_{d}(t)}\left|\bar{r}_{d}(t)\right| \leq \frac{\sqrt{2}}{2}-\varepsilon_{7}, \forall t \geq 0, \tag{8.11}
\end{gather*}
$$

where $\varepsilon_{6}$ and $\varepsilon_{7}$ are arbitrarily small positive constants.
4. The variable $\dot{s}$ can be used to specify the desired forward speed.

The main ideas to solve the control objective are as follows:

1. Introduce a coordinate to change the ship position $(x, y)$ such that the ship model (8.1) can be transformed to a diagonal form to overcome difficulties caused by nonzero off-diagonal terms in the system matrices.
2. Interpret tracking errors in a frame attached to the path $\Omega$ such that the tracking error dynamics are of a triangular form to which the backstepping technique can be applied.
3. Use the orientation tracking error as a virtual control to stabilize the cross-track error.

### 8.1.2 Coordinate Transformations

### 8.1.2.1 Transform Ship Dynamics to a "Diagonal Form"

Introduce the following coordinate transformation (changing the ship position, see Figure 8.1):

$$
\begin{align*}
& \bar{x}=x+\varepsilon \cos (\psi), \\
& \bar{y}=y+\varepsilon \sin (\psi),  \tag{8.12}\\
& \bar{v}=v+\varepsilon r,
\end{align*}
$$

where $\varepsilon=m_{23} / m_{22}$. Using the above change of coordinates, the ship dynamics (8.1) can be written as

$$
\begin{align*}
& \dot{\bar{x}}=u \cos (\psi)-\bar{v} \sin (\psi), \\
& \dot{\bar{y}}=u \sin (\psi)+\bar{v} \cos (\psi), \\
& \dot{\psi}=r,  \tag{8.13}\\
& \dot{u}=\bar{\tau}_{u}, \\
& \dot{\bar{v}}=-\alpha u r-\beta \bar{v}+\gamma r, \\
& \dot{r}=\bar{\tau}_{r},
\end{align*}
$$

where we have chosen the primary controls $\tau_{u}$ and $\tau_{r}$ as:

$$
\begin{align*}
\tau_{u}= & m_{11} \bar{\tau}_{u}-m_{22} v r-\frac{m_{23}+m_{32}}{2} r^{2}+d_{11} u, \\
\tau_{r}= & \frac{m_{22} m_{33}-m_{23} m_{32}}{m_{22}} \bar{\tau}_{r}-\frac{1}{m_{22}}\left(\left(m_{11} m_{22}-m_{22}^{2}\right) u v+\left(m_{11} m_{32}-\right.\right. \\
& \left.m_{22} \frac{m_{23}+m_{32}}{2} u r+\left(m_{32} d_{22}-m_{22} d_{32}\right) v-\left(m_{22} d_{33}-m_{32} d_{23}\right) r\right), \tag{8.14}
\end{align*}
$$

with $\bar{\tau}_{u}$ and $\bar{\tau}_{r}$ being considered as new controls to be designed later.


Figure 8.1 Interpretation of tracking errors

Remark 8.2. After the ship system (8.1) is transformed to the diagonal form (8.13), the methods mentioned in the preceding three chapters can be used to obtain controllers for specific tasks such as stabilization and model reference trajectorytracking if $(\bar{x}, \bar{y})$ are considered as the ship position instead of $(x, y)$.

### 8.1.2.2 Transform Path-tracking Errors

We now interpret the path-tracking errors in a frame attached to the reference path $\boldsymbol{\Omega}$ as follows (see Figure 8.1):

$$
\left[\begin{array}{c}
x_{e}  \tag{8.15}\\
y_{e} \\
\bar{\psi}_{e}
\end{array}\right]=\boldsymbol{J}^{\boldsymbol{T}}(\psi)\left[\begin{array}{c}
\bar{x}-x_{d} \\
\bar{y}-y_{d} \\
\psi-\psi_{d}
\end{array}\right],
$$

where $\psi_{d}$ is the angle between the path and the $X$-axis defined by

$$
\begin{equation*}
\psi_{d}=\arctan \left(\frac{y_{d}^{\prime}(s)}{x_{d}^{\prime}(s)}\right), \tag{8.16}
\end{equation*}
$$

with $x_{d}^{\prime}(s)$ and $y_{d}^{\prime}(s)$ being defined in (8.7).
In Figure 8.1, $O_{E} X_{E} Y_{E}$ is the earth-fixed frame, $O_{d} X_{d} Y_{d}$ is a frame attached to the path $\Omega$ such that $O_{d} X_{d}$ and $O_{d} Y_{d}$ are parallel to the surge and sway axes of the ship, $\bar{u}_{d}$ is tangential to the path, $C G \equiv O_{b}$ is the center of gravity, and $O_{s}$ is referred to as the center of oscillation of the ship. Therefore $x_{e}, y_{e}$, and $\bar{\psi}_{e}$ can be
referred to as the tangential tracking error, cross-tracking error, and heading error, respectively. Differentiating (8.15) along the solutions of the first three equations of (8.13) results in the kinematic error dynamics:

$$
\begin{align*}
& \dot{x}_{e}=u-\bar{u}_{d} \cos \left(\bar{\psi}_{e}\right)+r y_{e} \\
& \dot{y}_{e}=\bar{v}+\bar{u}_{d} \sin \left(\bar{\psi}_{e}\right)-r x_{e}  \tag{8.17}\\
& \dot{\bar{\psi}}_{e}=r-\bar{r}_{d}
\end{align*}
$$

where $\bar{u}_{d}$ and $\bar{r}_{d}$ are given in (8.7). From (8.17), it can be seen that $\left(x_{e}, y_{e}, \bar{\psi}_{e}\right)=$ $(0,0,0)$ is the equilibrium point only if the transformed sway velocity is zero, i.e., $\bar{v}=0$, which means that the ship must move on a straight line at the steady state. Since we allow the reference path $\boldsymbol{\Omega}$ to be different from a straight line, but also include it, the transformed sway velocity is generally different from zero at the steady state, i.e., $\bar{v}$ can be different from zero. To resolve this difficulty we introduce an angle $\delta$ to the orientation error $\bar{\psi}_{e}$ by defining

$$
\begin{equation*}
\psi_{e}=\bar{\psi}_{e}+\delta \tag{8.18}
\end{equation*}
$$

Substituting (8.18) into (8.17) yields

$$
\begin{align*}
& \dot{x}_{e}=u-u_{d} \cos \left(\psi_{e}\right)-v_{d} \sin \left(\psi_{e}\right)+r y_{e} \\
& \dot{y}_{e}=u_{d} \sin \left(\psi_{e}\right)+v_{e}-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)-r x_{e}  \tag{8.19}\\
& \dot{\psi}_{e}=r-r_{d}
\end{align*}
$$

where

$$
\begin{align*}
& v_{e}=\bar{v}-v_{d} \\
& u_{d}=\bar{u}_{d} \cos (\delta), \\
& v_{d}=\bar{u}_{d} \sin (\delta),  \tag{8.20}\\
& r_{d}=\bar{r}_{d}-\dot{\delta}
\end{align*}
$$

We refer to $u_{d}, v_{d}$, and $r_{d}$ as the desired surge, sway, and yaw reference velocities. It can now be seen that $\left(x_{e}, y_{e}, \psi_{e}\right)=(0,0,0)$ is the equilibrium point of (8.19) if $v_{e}=0$. At the steady state, the surge and yaw velocities $(u, r)$ should converge to their desired values $\left(u_{d}, r_{d}\right)$ with the help of a control to be designed later. Therefore, to determine the angle $\delta$ that guarantees $v_{e}=0$ at the steady state, from the fifth equation of (8.13), the desired sway velocity should be chosen as

$$
\begin{equation*}
\dot{v}_{d}=-\alpha u_{d} r_{d}-\beta v_{d}+\gamma r_{d} \tag{8.21}
\end{equation*}
$$

Substituting (8.20) into (8.21) results in (8.5). We make the following important observations:

Remark 8.3. With $|\delta(t)|<\pi / 4, \forall t \geq t_{0} \geq 0$, we have $\left|v_{d}(t)\right|<\left|u_{d}(t)\right|$ since $u_{d}=$ $\bar{u}_{d} \cos (\delta)$ and $v_{d}=\bar{u}_{d} \sin (\delta)$. In addition, with the help of a control (to be designed
later), convergence of $\left(x_{e}, y_{e}, \psi_{e}\right)$ to zero implies that $\bar{\psi}_{e}$ converges to $-\delta(t)$. Since $|\delta(t)|<\pi / 4$, the ship will not turn around.

To prepare for the control design, we rewrite the path-tracking error dynamics as follows:

$$
\begin{align*}
& \dot{x}_{e}=u-u_{d} \cos \left(\psi_{e}\right)-v_{d} \sin \left(\psi_{e}\right)+r y_{e}, \\
& \dot{y}_{e}=u_{d} \sin \left(\psi_{e}\right)+v_{e}-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)-r x_{e}, \\
& \dot{\psi}_{e}=r-r_{d}, \\
& \dot{v}_{e}=-\alpha\left(u r-u_{d} r_{d}\right)-\beta v_{e}+\gamma\left(r-r_{d}\right),  \tag{8.22}\\
& \dot{u}=\bar{\tau}_{u}, \\
& \dot{r}=\bar{\tau}_{r} .
\end{align*}
$$

The control objective has been converted to one of stabilizing (8.22) at the origin. The tracking error dynamics (8.22) are similar to the ones in the preceding chapters. The techniques in these chapters can be applied to design the controls $\bar{\tau}_{u}$ and $\bar{\tau}_{r}$ but will result in an extremely complex procedure to determine the control gains. In this chapter we will propose a much simpler technique to design $\bar{\tau}_{u}$ and $\bar{\tau}_{r}$ to globally stabilize (8.22) at the origin.

### 8.1.3 Control Design

We divide the control design into two stages. At the first stage, we consider the first four equations of (8.22), with $u$ and $r$ being viewed as controls. At the second stage, the last two equations of (8.22) are considered to design the actual controls $\bar{\tau}_{u}$ and $\bar{\tau}_{r}$ using the backstepping technique.

### 8.1.3.1 Kinematic Control Design

Substep 1 (Stabilizing $x_{e}$-dynamics)
Define

$$
\begin{equation*}
u_{e}=u-\alpha_{u}, \tag{8.23}
\end{equation*}
$$

where $\alpha_{u}$ is a virtual control of $u$. The first equation of (8.22) suggests that we choose

$$
\begin{equation*}
\alpha_{u}=-k_{1} x_{e}+u_{d} \tag{8.24}
\end{equation*}
$$

where $k_{1}$ is a positive constant to be selected later. Substituting (8.23) and (8.24) into the first equation of (8.22) yields

$$
\begin{equation*}
\dot{x}_{e}=-k_{1} x_{e}-u_{d}\left(\cos \left(\psi_{e}\right)-1\right)-v_{d} \sin \left(\psi_{e}\right)+r y_{e}+u_{e} \tag{8.25}
\end{equation*}
$$

Note that we have chosen $\alpha_{u}$ not to cancel $-u_{d}\left(\cos \left(\psi_{e}\right)-1\right)-v_{d} \sin \left(\psi_{e}\right)+r y_{e}$ to simplify the control design.

Substep 2 (Stabilizing $\left(y_{e}, v_{e}\right)$-dynamics)
Introduce a coordinate change

$$
\begin{equation*}
z_{1}=y_{e}+\frac{v_{e}}{\beta} . \tag{8.26}
\end{equation*}
$$

Differentiating (8.26) along the solutions of the second and fourth equations of (8.22), and using (8.23) and (8.24) yields

$$
\begin{align*}
\dot{z}_{1}= & u_{d} \sin \left(\psi_{e}\right)-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)-r x_{e}- \\
& \frac{\alpha}{\beta}\left(\left(u_{e}-k_{1} x_{e}+u_{d}\right) r-u_{d} r_{d}\right)+\frac{\gamma}{\beta}\left(r-r_{d}\right) . \tag{8.27}
\end{align*}
$$

For simplicity, we now choose

$$
\begin{equation*}
k_{1}=\frac{\beta}{\alpha} \tag{8.28}
\end{equation*}
$$

to cancel the term $-r x_{e}$ on the right-hand side of (8.27) to have $z_{1}$-dynamics "independent" of $x_{e}$. Substituting (8.28) into (8.27) yields

$$
\begin{equation*}
\dot{z}_{1}=u_{d} \sin \left(\psi_{e}\right)-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)-\frac{1}{\beta}\left(\alpha u_{d}-\gamma\right)\left(r-r_{d}\right)-\frac{\alpha}{\beta} u_{e} r . \tag{8.29}
\end{equation*}
$$

It is now seen that the third equation of (8.22) and (8.29) are of a forward structure with $r$ being considered as a control. One can apply forwarding control design techniques to design a control law for $r$. However, it is difficult to find a proper Lyapunov function to handle cross terms. The reader is referred to [7] for more details on control design techniques for forward systems. Here we use the backstepping technique by using $\psi_{e}$ as a virtual control to stabilize $z_{1}$. Define

$$
\begin{equation*}
z_{2}=\psi_{e}-\alpha_{\psi_{e}} \tag{8.30}
\end{equation*}
$$

where $\alpha_{\psi_{e}}$ is a virtual control of $\psi_{e}$. By noticing (8.3), (8.7) and $|\delta(t)|<\pi / 4, \forall t \geq$ $t_{0} \geq 0$, we have $u_{d}(t)>0$. Therefore, a simple control law for $\alpha_{\psi_{e}}$ can be chosen as

$$
\begin{equation*}
\alpha_{\psi_{e}}=-\arcsin \left(\frac{k_{2} z_{1}}{\Delta}\right), \Delta=\sqrt{1+\left(k_{2} z_{1}\right)^{2}} \tag{8.31}
\end{equation*}
$$

where $k_{2}$ is a positive constant to be selected later. It is of interest to note that (8.31) is well defined for all $z_{1} \in \mathbb{R}$.

Substep 3 (Stabilizing $z_{2}$-dynamics)
Define

$$
\begin{equation*}
r_{e}=r-r_{d}-\alpha_{r} \tag{8.32}
\end{equation*}
$$

where $\alpha_{r}$ is a virtual control of $r$. Differentiating both sides of (8.30) along the solutions of the third equation of (8.22) and (8.29) yields

$$
\begin{align*}
\dot{z}_{2}= & \left(1-k_{2} \frac{\alpha u_{d}-\gamma}{\beta \Delta^{2}}\right)\left(r_{e}+\alpha_{r}\right)+ \\
& \frac{k_{2}}{\Delta^{2}}\left(u_{d} \sin \left(\psi_{e}\right)-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)\right)-\frac{k_{2} \alpha}{\beta \Delta^{2}} u_{e} r . \tag{8.33}
\end{align*}
$$

We choose the constant $k_{2}$ such that

$$
\begin{equation*}
1-k_{2} \frac{\alpha u_{d}-\gamma}{\beta \Delta^{2}}>0 \Rightarrow 0<k_{2} \leq \frac{\beta}{\alpha u_{d}-|\gamma|} . \tag{8.34}
\end{equation*}
$$

There always exists $k_{2}$ satisfying (8.34) under condition (8.3) since $|\delta(t)|<\pi / 4$ and $u_{d}=\bar{u}_{d} \cos (\delta), \forall t \geq t_{0} \geq 0$. We choose the virtual control $\alpha_{r}$ from (8.33) as

$$
\begin{equation*}
\alpha_{r}=\frac{1}{1-\frac{k_{2}}{\beta \Delta^{2}}\left(\alpha u_{d}-\gamma\right)}\left(-\frac{c_{1} z_{2}}{\sqrt{1+z_{2}^{2}}}-\frac{k_{2}}{\Delta^{2}}\left(u_{d} \sin \left(\psi_{e}\right)-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)\right)\right), \tag{8.35}
\end{equation*}
$$

where $c_{1}$ is a positive constant. It is noted that $\alpha_{r}$ is bounded by some constant for all $\left(z_{1}, z_{2}, \psi_{e}\right) \in \mathbb{R}^{3}$.

### 8.1.3.2 Kinetic Control Design

By differentiating both sides of (8.23) and (8.32) along the solutions of the first, third, fifth, and sixth equations of (8.22), the control $\bar{\tau}_{u}$ is chosen as

$$
\begin{equation*}
\bar{\tau}_{u}=-c_{2} u_{e}+\frac{\partial \alpha_{u}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial \alpha_{u}}{\partial u_{d}} \dot{u}_{d}+\frac{k_{2} \alpha}{\beta \Delta^{2}} z_{2} r, \tag{8.36}
\end{equation*}
$$

and the control $\bar{\tau}_{r}$ is chosen as

$$
\begin{align*}
\bar{\tau}_{r}= & \frac{\partial \alpha_{r}}{\partial z_{1}} \dot{z}_{1}+\frac{\partial \alpha_{r}}{\partial z_{2}} \dot{z}_{2}+\frac{\partial \alpha_{r}}{\partial \psi_{e}} \dot{\psi}_{e}+\frac{\partial \alpha_{r}}{\partial u_{d}} \dot{u}_{d}+\frac{\partial \alpha_{r}}{\partial v_{d}} \dot{v}_{d}- \\
& \left(1-\frac{k_{2}}{\beta \Delta^{2}}\left(\alpha u_{d}-\gamma\right)\right) z_{2}-c_{3} r_{e}+\dot{r}_{d}, \tag{8.37}
\end{align*}
$$

where $c_{2}$ and $c_{3}$ are positive constants. Substituting (8.30), (8.31), (8.35), (8.36), and (8.37) into (8.29) and (8.33), derivatives of $u_{e}$ and $r_{e}$ along the solutions of the last two equations of (8.22) gives the following closed loop system:

$$
\begin{align*}
\dot{x}_{e}= & -k_{1} x_{e}-u_{d}\left(\cos \left(\psi_{e}\right)-1\right)-v_{d} \sin \left(\psi_{e}\right)+\left(r_{d}+\alpha_{r}+r_{e}\right)\left(z_{1}-\right. \\
& \left.\beta^{-1} v_{e}\right)+u_{e} \\
\dot{v}_{e}= & -\alpha\left(\left(u_{e}+\alpha_{u}\right)\left(r_{d}+\alpha_{r}+r_{e}\right)-u_{d} r_{d}\right)-\beta v_{e}+\gamma\left(\alpha_{r}+r_{e}\right) \\
\dot{z}_{1}= & -k_{2} u_{d} z_{1} / \Delta-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)-\beta^{-1}\left(\alpha u_{d}-\gamma\right)\left(\alpha_{r}+r_{e}\right)+ \\
& u_{d}\left(\sin \left(z_{2}\right) \cos \left(\alpha_{\psi_{e}}\right)+\left(\cos \left(z_{2}\right)-1\right) \cos \left(\alpha_{\psi_{e}}\right)\right)-\alpha \beta^{-1} u_{e}\left(r_{d}+\right. \\
& \left.\alpha_{r}+r_{e}\right) \\
\dot{z}_{2}= & -c_{1} z_{2} / \sqrt{1+z_{2}^{2}}+\left(1-k_{2}\left(\alpha u_{d}-\gamma\right) \beta^{-1} \Delta^{-2}\right) r_{e}-k_{2} \alpha \beta^{-1} \Delta^{-2} u_{e} r \\
\dot{u}_{e}= & -c_{2} u_{e}+k_{2} \alpha \beta^{-1} \Delta^{-2} z_{2} r \\
\dot{r}_{e}= & -c_{3} r_{e}-\left(1-k_{2}\left(\alpha u_{d}-\gamma\right) \beta^{-1} \Delta^{-2}\right) z_{2} . \tag{8.38}
\end{align*}
$$

We now present the first main result of this chapter, the proof of which is given in the next section.

Theorem 8.1. Assume that the path $\boldsymbol{\Omega}$ satisfies all conditions specified in Assumption 8.1, then the controls $\tau_{u}$ and $\tau_{r}$ given by (8.14), (8.36), and (8.37) force the transformed tracking errors ( $x_{e}, y_{e}, \psi_{e}$ ) to converge to zero globally asymptotically and locally exponentially with an appropriate choice of $k_{2}$. As a result, the actual position tracking errors $\left(x-x_{d}, y-y_{d}\right)$ and orientation tracking error $\left(\psi-\psi_{d}\right)$ globally asymptotically and locally exponentially converge to balls with radii of $\left|m_{23} / m_{22}\right|$ and $|\delta(t)| \leq \arcsin \left(\left(\alpha \bar{u}_{d}+|\gamma|\right)\left|\bar{r}_{d}\right| /\left(\beta \bar{u}_{d}+\dot{\bar{u}}_{d}\right)\right)<\pi / 4$, respectively. Furthermore, the desired forward speed of the ship on the path can be adjusted by adjusting $\dot{s}$.

### 8.1.4 Stability Analysis

To prove Theorem 8.1, we first show that the closed loop system (8.38) is forward complete, i.e., there is no finite escape in the closed loop system. We then consider the $\left(z_{2}, u_{e}, r_{e}\right)$-, $z_{1-}$ and ( $x_{e}, v_{e}$ )-subsystems separately. We first consider the Lyapunov function

$$
\begin{equation*}
V_{1}=x_{e}^{2}+z_{1}^{2}+z_{2}^{2}+v_{e}^{2}+u_{e}^{2}+r_{e}^{2} \tag{8.39}
\end{equation*}
$$

A simple calculation shows that the derivative of $V_{1}$ along the solutions of (8.38) satisfies

$$
\begin{array}{ll} 
& \dot{V}_{1} \leq \rho_{11} V_{1}+\rho_{12} \\
\Rightarrow \quad & V_{1}(t) \leq\left(V_{1}\left(t_{0}\right)+\rho_{12} / \rho_{11}\right) e^{\rho_{11}\left(t-t_{0}\right)} \tag{8.40}
\end{array}
$$

where $\rho_{11}$ and $\rho_{12}$ are some positive constants, which implies that the closed loop system (8.38) is forward complete.

To investigate stability of the $\left(z_{2}, u_{e}, r_{e}\right)$-subsystem, we consider the Lyapunov function

$$
\begin{equation*}
V_{21}=\frac{1}{2}\left(z_{2}^{2}+u_{e}^{2}+r_{e}^{2}\right), \tag{8.41}
\end{equation*}
$$

whose derivative along the solutions of the last three equations of (8.38) satisfies

$$
\begin{equation*}
\dot{V}_{21}=-c_{1} \frac{z_{2}^{2}}{\sqrt{1+z_{2}^{2}}}-c_{2} u_{e}^{2}-c_{3} r_{e}^{2} \leq 0 \tag{8.42}
\end{equation*}
$$

Since $\dot{V}_{21} \leq 0$, we have $V_{21}(t) \leq \chi_{21}(\cdot)$ where $\chi_{21}(\cdot)$ is a class- $K_{\infty}$ function of $X_{1}\left(t_{0}\right)=\left\|\left(z_{2}\left(t_{0}\right), u_{e}\left(t_{0}\right), r_{e}\left(t_{0}\right)\right)\right\|$. With this upper bound of $V_{21}$, we have

$$
\begin{equation*}
\dot{V}_{21}=-c_{1} \frac{z_{2}^{2}}{\sqrt{1+2 \chi_{21}(\cdot)}}-c_{2} u_{e}^{2}-c_{3} r_{e}^{2} \tag{8.43}
\end{equation*}
$$

which implies that the $\left(z_{2}, u_{e}, r_{e}\right)$-subsystem is globally asymptotically and locally exponentially stable at the origin, i.e. there exist a class- $K$ function $\gamma_{21}(\cdot)$ and a constant $\sigma_{21}$ depending on $X_{1}\left(t_{0}\right)$ such that

$$
\begin{equation*}
\left\|\left(z_{2}(t), u_{e}(t), r_{e}(t)\right)\right\| \leq \gamma_{21}(\cdot) e^{-\sigma_{21}\left(t-t_{0}\right)} . \tag{8.44}
\end{equation*}
$$

For the $z_{1}$-subsystem, we take the Lyapunov function

$$
\begin{equation*}
V_{22}=\frac{1}{2} z_{1}^{2} . \tag{8.45}
\end{equation*}
$$

By substituting $\alpha_{r}$, see (8.35), into the derivative of $V_{22}$ along the solutions of the third equation of (8.38), we have

$$
\begin{align*}
\dot{V}_{22} \leq & -k_{2}\left(1-k_{2} \frac{\left|\alpha u_{d}-\gamma\right|}{\left(\beta-k_{2}\left|\alpha u_{d}-\gamma\right|\right.}\right) \frac{u_{d}-\left|v_{d}\right|}{\Delta} z_{1}^{2}+ \\
& \gamma_{21}(\cdot)\left(V_{22}+\gamma_{22}(\cdot)\right) e^{-\sigma_{21}\left(t-t_{0}\right)}, \tag{8.46}
\end{align*}
$$

where $\gamma_{21}(\cdot)$ and $\gamma_{22}(\cdot)$ are class- $K_{\infty}$ functions of $X_{1}\left(t_{0}\right)$. We now choose $k_{2}$ such that

$$
\begin{align*}
& 1-\frac{k_{2}\left(\alpha u_{d}-|\gamma|\right)}{\beta-k_{2}\left(\alpha u_{d}-|\gamma|\right)}>0 \\
\Rightarrow \quad & 0<k_{2}<0.5 \frac{\beta}{\alpha u_{d}-|\gamma|}, \tag{8.47}
\end{align*}
$$

then (8.46) together with conditions (8.3) and (8.10) in Assumption 8.1 imply that

$$
\begin{equation*}
\dot{V}_{22} \leq \gamma_{21}(\cdot)\left(V_{22}+\gamma_{22}(\cdot)\right) e^{-\sigma_{21}\left(t-t_{0}\right)} . \tag{8.48}
\end{equation*}
$$

Solving the above differential inequality results in $V_{22}(t) \leq \gamma_{23}(\cdot)$ with $\gamma_{23}(\cdot)$ being a class- $K_{\infty}$ function of $X_{2}\left(t_{0}\right)=\left\|\left(z_{1}\left(t_{0}\right), z_{2}\left(t_{0}\right), u_{e}\left(t_{0}\right), r_{e}\left(t_{0}\right)\right)\right\|$. Substituting this upper bound of $V_{22}$ into (8.46) yields

$$
\begin{align*}
\dot{V}_{22} & \leq-k_{2}\left(1-k_{2} \frac{\left|\alpha u_{d}-\gamma\right|}{\beta-k_{2}\left|\alpha u_{d}-\gamma\right|}\right) \frac{u_{d}-\left|v_{d}\right|}{\Delta} z_{1}^{2}+\gamma_{24}(\cdot) e^{-\sigma_{21}\left(t-t_{0}\right)} \\
& \leq \gamma_{24}(\cdot) e^{-\sigma_{21}\left(t-t_{0}\right)} \tag{8.49}
\end{align*}
$$

where $\gamma_{24}(\cdot)=\gamma_{21}(\cdot)\left(\gamma_{23}(\cdot)+\gamma_{22}(\cdot)\right)$. The second inequality of (8.49) implies that there exists a class- $K_{\infty}$ function $\gamma_{25}(\cdot)$ of $X_{2}\left(t_{0}\right)$ such that $\left|z_{1}(t)\right| \leq \gamma_{25}(\cdot)$. Substituting this upper bound into the expression of $\Delta$ and using the first inequality of (8.49) yields global asymptotic and local exponential stability of the $z_{1}$-subsystem, i.e., there exist a class- $K_{\infty}$ function of $X_{2}\left(t_{0}\right)$ and a constant $\sigma_{22}>0$ depending on $X_{2}\left(t_{0}\right)$ such that $\left|z_{1}(t)\right| \leq \gamma_{26}(\cdot) e^{-\sigma_{22}\left(t-t_{0}\right)}$. Note that condition (8.47) covers the condition (8.34).

For the $\left(x_{e}, v_{e}\right)$-subsystem, we consider the Lyapunov function

$$
\begin{equation*}
V_{3}=\frac{1}{2}\left(x_{e}^{2}+v_{e}^{2}\right) . \tag{8.50}
\end{equation*}
$$

By taking the derivative of this function along the solutions of the first two equations of (8.38), it can readily be shown that there exist class- $K_{\infty}$ functions $\gamma_{31}(\cdot)$ and $\gamma_{32}(\cdot)$ of $X_{2}\left(t_{0}\right)$ such that

$$
\begin{align*}
\dot{V}_{3} & \leq-\frac{\beta}{\alpha} x_{e}^{2}-\beta v_{e}^{2}+\gamma_{31}(\cdot)\left(V_{3}+\gamma_{32}(\cdot)\right) e^{-\sigma_{22}\left(t-t_{0}\right)} \\
& \leq \gamma_{31}(\cdot)\left(V_{3}+\gamma_{32}(\cdot)\right) e^{-\sigma_{22}\left(t-t_{0}\right)} . \tag{8.51}
\end{align*}
$$

The second inequality of (8.51) means that $V_{3}(t) \leq \gamma_{33}(\cdot)$ with $\gamma_{33}(\cdot)$ being a class$K_{\infty}$ function of $\left\|\left(x_{e}\left(t_{0}\right), v_{e}\left(t_{0}\right) z_{1}\left(t_{0}\right), z_{2}\left(t_{0}\right), u_{e}\left(t_{0}\right), r_{e}\left(t_{0}\right)\right)\right\|$. Substituting this upper bound into the first line of (8.51) readily yields global asymptotic and local exponential stability of the $\left(x_{e}, v_{e}\right)$-subsystem. We have so far proven that the transformed tracking errors ( $x_{e}, z_{1}, z_{2}, v_{e}, u_{e}, r_{e}$ ) globally asymptotically and locally exponentially converge to zero. By (8.26), convergence of $z_{1}$ and $v_{e}$ to zero implies that of $y_{e}$. Convergence of the actual position tracking error $\left(x-x_{d}, y-y_{d}\right)$ and orientation tracking error $\left(\psi-\psi_{d}\right)$ to balls with radii of $\left|m_{23} / m_{22}\right|$ and $|\delta(t)| \leq\left(\alpha \bar{u}_{d}+|\gamma|\right)\left|\bar{r}_{d}\right| /\left|\beta \bar{u}_{d}+\dot{\bar{u}}_{d}\right|<\pi / 4$ follows from (8.15) and (8.18). Finally, since $v_{d}=\bar{u}_{d} \sin (\delta), u_{d}=\bar{u}_{d} \cos (\delta)$ and $\bar{u}_{d}=\sqrt{x_{d}^{\prime 2}+y_{d}^{\prime 2}} \dot{s}$, we can see that the total linear velocity of the ship is tangential to the path and the desired forward speed $u_{d}$ can be adjusted by adjusting $\dot{s}$.

### 8.1.5 Dealing with Environmental Disturbances

To clearly illustrate our idea, we have not included constant (or at least slowly timevarying) environmental disturbances induced by waves, wind, and ocean currents in the control design. To use the proposed controller in practice, the environmental disturbances should essentially be taken into account in the control design. The aim
of this section is to discuss a simple way to modify the proposed control design in the previous section to handle these disturbances. When the constant disturbances are present, the last equation of (8.1) is of the form

$$
\begin{equation*}
M \dot{v}=-C(v) v-D v+\tau+\tau_{E} \tag{8.52}
\end{equation*}
$$

where $\boldsymbol{\tau}_{\boldsymbol{E}}=\left[\begin{array}{lll}\tau_{u E} & \tau_{v E} & \tau_{r E}\end{array}\right]^{T}$ with $\tau_{u E}, \tau_{v E}$ and $\tau_{r E}$ being the environmental disturbances acting on the surge, sway, and yaw axes. Processing the same coordinate transformations as in Section 8.1.2, the last three equations of (8.13) are written as

$$
\begin{align*}
& \dot{u}=\bar{\tau}_{u}+\bar{\tau}_{u E}, \\
& \dot{\bar{v}}=-\alpha u r-\beta \bar{v}+\gamma r+\bar{\tau}_{v E},  \tag{8.53}\\
& \dot{r}=\bar{\tau}_{r}+\bar{\tau}_{r E},
\end{align*}
$$

with

$$
\begin{align*}
\bar{\tau}_{u E} & =\frac{1}{m_{11}} \tau_{u E}, \\
\bar{\tau}_{v E} & =\frac{1}{m_{22}} \tau_{v E},  \tag{8.54}\\
\bar{\tau}_{r E} & =\frac{\left(-m_{23} \tau_{v E}+m_{22} \tau_{r E}\right)}{\left(m_{22} m_{33}-m_{23} m_{32}\right)} .
\end{align*}
$$

We first deal with $\bar{\tau}_{v E}$. The idea to handle $\bar{\tau}_{v E}$ is to introduce an angle to the yaw angle. This angle together with the ship forward speed will compensate for $\bar{\tau}_{v E}$. We design an observer to estimate $\bar{\tau}_{v E}$ as

$$
\begin{align*}
& \dot{\hat{\bar{v}}}=k_{01}(\bar{v}-\hat{\bar{v}})-\alpha u r-\beta \hat{\bar{v}}+\gamma r+\hat{\bar{\tau}}_{v E}  \tag{8.55}\\
& \hat{\bar{\tau}}_{v E}=-\gamma_{01} \operatorname{proj}\left(\bar{v}-\hat{\bar{v}}, \hat{\bar{\tau}}_{v E}\right)
\end{align*}
$$

where $k_{01}$ and $\gamma_{01}$ are positive constants, $\hat{\bar{\tau}}_{v E}$ is an estimate of $\bar{\tau}_{v E}$, the operator proj represents the Lipschitz projection algorithm [118] as

$$
\begin{aligned}
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \text { if } \Xi(\hat{\omega}) \leq 0, \\
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \text { if } \Xi(\hat{\omega}) \geq 0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi \leq 0, \\
& \operatorname{proj}(\varpi, \hat{\omega})=(1-\Xi(\hat{\omega})) \varpi \text { if } \Xi(\hat{\omega})>0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi>0,
\end{aligned}
$$

where $\Xi(\hat{\omega})=\left(\hat{\omega}^{2}-\omega_{M}^{2}\right) /\left(\xi^{2}+2 \xi \omega_{M}\right), \Xi_{\hat{\omega}}(\hat{\omega})=\partial \Xi(\hat{\omega}) / \partial \hat{\omega}, \xi$ is an arbitrarily small positive constant, $\hat{\omega}$ is an estimate of $\omega$, and $|\omega| \leq \omega_{M}$. The projection algorithm is such that if $\dot{\hat{\omega}}=\operatorname{proj}(\varpi, \hat{\omega})$ and $\hat{\omega}\left(t_{0}\right) \leq \omega_{M}$, then

1. $\hat{\omega}(t) \leq \omega_{M}+\xi, \forall 0 \leq t_{0} \leq t<\infty$,
2. $\operatorname{proj}(\varpi, \hat{\omega})$ is Lipschitz continuous,
3. $|\operatorname{proj}(\varpi, \hat{\omega})| \leq|\varpi|$,
4. $\tilde{\omega} \operatorname{proj}(\varpi, \hat{\omega}) \geq \tilde{\omega} \varpi$ with $\tilde{\omega}=\omega-\hat{\omega}$.

From (8.55), the second equation of (8.53), and property 4 of the proj operator, it is readily shown that $\lim _{t \rightarrow \infty} \tilde{\bar{\tau}}_{v E}(t)=0$ with $\tilde{\tilde{\tau}}_{v E}(t)=\bar{\tau}_{v E}-\hat{\bar{\tau}}_{v E}(t)$. Now, by comparing the transformed sway dynamics with disturbances (the second equation of (8.53)) and the one without disturbances (the fifth equation of (8.13)), the desired sway velocity $v_{d}$ dynamics should now be modified from (8.21) as

$$
\begin{equation*}
\dot{v}_{d}=-\alpha u_{d} r_{d}-\beta v_{d}+\gamma r_{d}+\hat{\bar{\tau}}_{v E} \tag{8.56}
\end{equation*}
$$

With this choice, we have

$$
\begin{equation*}
\dot{v}_{e}=-\alpha\left(u r-u_{d} r_{d}\right)-\beta v_{e}+\gamma\left(r-r_{d}\right)+\tilde{\bar{\tau}}_{v E} \tag{8.57}
\end{equation*}
$$

Notice that (8.56) is equivalent to

$$
\begin{gather*}
\left(\bar{u}_{d} \cos (\delta)(1-\alpha)+\gamma\right) \dot{\delta}=-\left(\beta \bar{u}_{d}+\dot{\bar{u}}_{d}\right) \sin (\delta)-  \tag{8.58}\\
\left(\alpha \bar{u}_{d} \cos (\delta)-\gamma\right) \bar{r}_{d}+\hat{\bar{\tau}}_{v E}, \delta\left(t_{0}\right)=0
\end{gather*}
$$

Therefore the tracking objective is solvable if condition (8.11) is replaced by

$$
\begin{equation*}
\frac{\left(\alpha \bar{u}_{d}+|\gamma|\right)\left|\bar{r}_{d}\right|+\max \left(\left|\bar{\tau}_{v E}\right|\right)+\xi_{01}}{\beta \bar{u}_{d}+\dot{\bar{u}}_{d}} \leq \frac{\sqrt{2}}{2}-\varepsilon_{7} \tag{8.59}
\end{equation*}
$$

where property 1 of the projection algorithm that guarantees $\left|\hat{\bar{\tau}}_{v E}\right| \leq \max \left(\left|\bar{\tau}_{v E}\right|\right)+$ $\xi_{01}$ has been utilized, with $\xi_{01}$ an arbitrarily small positive constant. Hence, if the disturbances are not too large and the path $\boldsymbol{\Omega}$ is feasible, i.e., Assumption 8.1 with condition (8.11) being replaced by (8.59) holds, then the control objective is solvable under constant disturbances. Now processing the same coordinate changes as in Section 8.1.2 results in the same tracking error dynamics as (8.22) but the last three equations of (8.22) are replaced by

$$
\begin{align*}
& \dot{v}_{e}=-\alpha\left(u r-u_{d} r_{d}\right)-\beta v_{e}+\gamma\left(r-r_{d}\right)+\tilde{\tilde{\tau}}_{v E} \\
& \dot{u}=\bar{\tau}_{u}+\bar{\tau}_{u E}  \tag{8.60}\\
& \dot{r}=\bar{\tau}_{r}+\bar{\tau}_{r E}
\end{align*}
$$

Since $\lim _{t \rightarrow \infty} \tilde{\bar{\tau}}_{v E}(t)=0$, and $\bar{\tau}_{u E}$ and $\bar{\tau}_{r E}$ satisfy a matching condition, the control design is much the same as in Section 8.1.3 and is briefly described in the following.

Kinematic Control Design. The same as in Section 8.1.3.
Kinetic Control Design. The controls $\bar{\tau}_{u}$ and $\bar{\tau}_{r}$ are almost the same as in Section 8.1.3 but we cannot directly use $\dot{z}_{1}$ and $\dot{z}_{2}$ since they contain an unmeasured term $\tilde{\tilde{\tau}}_{v E}(t)$ and we need to compensate $\bar{\tau}_{u E}$ and $\bar{\tau}_{r E}$. These controls are designed as follows:

$$
\begin{align*}
\bar{\tau}_{u}= & -c_{2} u_{e}+\frac{\partial \alpha_{u}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial \alpha_{u}}{\partial u_{d}} \dot{u}_{d}+\frac{k_{2} \alpha}{\beta \Delta^{2}} z_{2} r-\hat{\bar{\tau}}_{u E}, \\
& \bar{\tau}_{r}= \\
& -c_{3} r_{e}+\dot{r}_{d}+\frac{\partial \alpha_{r}}{\partial z_{1}}\left(u_{d} \sin \left(\psi_{e}\right)-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)-\right. \\
& \frac{1}{\beta}\left(\alpha u_{d}-\gamma\right)\left(r-r_{d}\right)-\frac{\alpha}{\beta} u_{e} r+\frac{\partial \alpha_{r}}{\partial z_{2}}\left(-\frac{c_{1} z_{2}}{\sqrt{1+z_{2}^{2}}}+\right. \\
& \left.\left(1-\frac{k_{2}\left(\alpha u_{d}-\gamma\right)}{\beta \Delta^{2}}\right) r_{e}-\frac{k_{2} \alpha}{\beta \Delta^{2}} u_{e} r\right)+\frac{\partial \alpha_{r}}{\partial \psi_{e}} \dot{\psi}_{e}+\frac{\partial \alpha_{r}}{\partial u_{d}} \dot{u}_{d}+  \tag{8.61}\\
& \frac{\partial \alpha_{r}}{\partial v_{d}} \dot{v}_{d}-\frac{\left(1-k_{2}\left(\alpha u_{d}-\gamma\right)\right.}{\beta \Delta^{2}} z_{2}-\hat{\bar{\tau}}_{r E},
\end{align*}
$$

where $\hat{\bar{\tau}}_{u E}$ and $\hat{\bar{\tau}}_{r E}$ are estimates of $\bar{\tau}_{u E}$ and $\bar{\tau}_{r E}$ updated by

$$
\begin{align*}
& \dot{\overline{\tilde{\tau}}}_{u E}=-\gamma_{02} \operatorname{proj}\left(u_{e}, \hat{\bar{\tau}}_{u E}\right), \\
& \dot{\bar{\tau}}_{r E}=-\gamma_{03} \operatorname{proj}\left(r_{e}, \hat{\bar{\tau}}_{r E}\right), \tag{8.62}
\end{align*}
$$

where $\gamma_{02}$ and $\gamma_{03}$ are positive constants. From the above control design, we have a closed loop system as

$$
\begin{align*}
\dot{x}_{e}= & -k_{1} x_{e}-u_{d}\left(\cos \left(\psi_{e}\right)-1\right)-v_{d} \sin \left(\psi_{e}\right)+\left(r_{d}+\alpha_{r}+r_{e}\right)\left(z_{1}-\right. \\
& \left.\beta^{-1} v_{e}\right)+u_{e}, \\
\dot{v}_{e}= & -\alpha\left(\left(u_{e}+\alpha_{u}\right)\left(r_{d}+\alpha_{r}+r_{e}\right)-u_{d} r_{d}\right)-\beta v_{e}+\gamma\left(\alpha_{r}+r_{e}\right)+\tilde{\bar{\tau}}_{v E}, \\
\dot{z}_{1}= & -k_{2} u_{d} z_{1} / \Delta-v_{d}\left(\cos \left(\psi_{e}\right)-1\right)-\beta^{-1}\left(\alpha u_{d}-\gamma\right)\left(\alpha_{r}+r_{e}\right)+ \\
& u_{d}\left(\sin \left(z_{2}\right) \cos \left(\alpha_{\psi_{e}}\right)+\left(\cos \left(z_{2}\right)-1\right) \cos \left(\alpha_{\psi_{e}}\right)\right)-\alpha \beta^{-1} u_{e}\left(r_{d}+\right. \\
& \left.\alpha_{r}+r_{e}\right)+\beta^{-1} \tilde{\bar{\tau}}_{v E}, \\
\dot{z}_{2}= & -c_{1} z_{2} / \sqrt{1+z_{2}^{2}}+\left(1-k_{2}\left(\alpha u_{d}-\gamma\right) \beta^{-1} \Delta^{-2}\right) r_{e}-k_{2} \alpha \beta^{-1} \Delta^{-2} u_{e} r+ \\
& k_{2} \beta^{-1} \Delta^{-2} \tilde{\bar{\tau}}_{v E}, \\
\dot{u}_{e}= & -c_{2} u_{e}+k_{2} \alpha \beta^{-1} \Delta^{-2} z_{2} r+\tilde{\tilde{\tau}}_{u E}, \\
\dot{r}_{e}= & -c_{3} r_{e}-\left(1-\frac{k_{2}}{\beta \Delta^{2}}\left(\alpha u_{d}-\gamma\right)\right) z_{2}-\frac{\tilde{\bar{\tau}}_{v E}}{\beta} \frac{\partial \alpha_{r}}{\partial z_{1}}-\frac{k_{2} \tilde{\tilde{\tau}}_{v E}}{\beta \Delta^{2}} \frac{\partial \alpha_{r}}{\partial z_{2}}+\tilde{\bar{\tau}}_{r E}, \tag{8.63}
\end{align*}
$$

where $\tilde{\bar{\tau}}_{u E}=\bar{\tau}_{u E}-\hat{\bar{\tau}}_{u E}$ and $\tilde{\bar{\tau}}_{r E}=\bar{\tau}_{r E}-\hat{\bar{\tau}}_{r E}$. We now present a result for the case with constant disturbances.

Theorem 8.2. Assume that the path $\boldsymbol{\Omega}$ satisfies all conditions specified in Assumption 8.1 except for (8.5) being replaced by (8.58) and that the disturbances are not too large, i.e., (8.59) holds, the controls $\tau_{u}$ and $\tau_{r}$ given by (8.14) and (8.61) force the transformed tracking errors $\left(x_{e}, y_{e}, \psi_{e}\right)$ to globally asymptotically converge to zero with an appropriate choice of $k_{2}$. As a result, the ac-
tual position tracking errors $\left(x-x_{d}, y-y_{d}\right)$ and the orientation tracking error $\left(\psi-\psi_{d}\right)$ globally asymptotically converge to balls with radii of $\left|m_{23} / m_{22}\right|$ and $\left.|\delta(t)| \leq \arcsin \left(\left(\alpha \bar{u}_{d}+|\gamma|\right)\left|\bar{r}_{d}\right|+\max \left(\left|\bar{\tau}_{v E}\right|\right)+\xi_{01}\right) /\left(\beta \bar{u}_{d}+\dot{\bar{u}}_{d}\right)\right)<\pi / 4$, respectively.

Note that convergence of the tracking errors in the case with disturbances is asymptotic (not locally exponential due to adaptation), and the magnitude of $\delta(t)$ is generally larger than the one in the case without disturbances. Since

$$
\lim _{t \rightarrow \infty}\left(\tilde{\bar{\tau}}_{u E}(t), \tilde{\bar{\tau}}_{u E}(t), \tilde{\bar{\tau}}_{r E}(t)\right)=0
$$

$\partial \alpha_{r} / \partial z_{1}$ and $\partial \alpha_{r} / \partial z_{2}$ are bounded by some constants, the structure of (8.63) is very similar to (8.38), and the proof of Theorem 8.2 follows the same lines as for Theorem 8.1 plus an application of Barbalat's lemma, we omit proof of Theorem 8.2.

### 8.1.6 Numerical Simulations

We perform some numerical simulations to illustrate the effectiveness of our proposed control laws given by (8.14), (8.36), and (8.37) on a supply vessel with the length of $L=76.2 \mathrm{~m}$ and a mass of $m=226 \times 10^{3} \mathrm{~kg}$. This vessel has a minimum turning circle with the radius of 250 m , a maximum surge force of $6.5 \times 10^{10} \mathrm{~N}$, and a maximum yaw moment of $3.2 \times 10^{12} \mathrm{Nm}$. The Bis-scale parameters of the vessel taken from [11], pp. 446, are as follows:

$$
\begin{aligned}
& m_{11}=1.1274, m_{22}=1.8902, m_{23}=-0.0744, m_{33}=0.1278 \\
& d_{11}=0.0358, d_{22}=0.1183, d_{23}=-0.0124, d_{32}=-0.0041, d_{33}=0.0308
\end{aligned}
$$

From these values, we have $\alpha=0.5964, \beta=0.0626$, and $\gamma=0.0041$. The reference path $\Omega$ is chosen as $x_{d}(s)=s, y_{d}(s)=100 \sin (0.01 s)$ and $\dot{s}=1$. Therefore, $\bar{u}_{d}=2 \sqrt{1+\cos (0.01 s)^{2}}$ and $\bar{r}_{d}=-0.02 \sin (0.01 s) /\left(1+\cos (0.01 s)^{2}\right)$. A simple calculation shows that all of conditions of Assumption 8.1 hold. In the simulations, the initial conditions are chosen as $x(0)=-20, y(0)=50, \psi(0)=0.5, u(0)=$ $0, v(0)=0.1$, and $r(0)=0$. The control gains are chosen as $k_{1}=\beta / \alpha=0.1049$, $k_{2}=0.025$, and $c_{1}=c_{2}=c_{3}=2$. A simple calculation shows that $k_{2}$ satisfies the condition specified in (8.47). Simulation results without disturbances are plotted in Figures 8.2 and 8.3. In these figures, all variables are converted back to their original values using the following table with $g$ being the gravitational acceleration:

The tracking errors are plotted in Figure 8.2a, and the angle $\delta$ is plotted in Figure 8.2 b . It is seen that the proposed controller nicely forces the ship to track the given reference path in the sense that the tracking errors asymptotically converge to zero. From (8.20) and the fact that the surge and sway velocities $(u, v)$ converge to their desired values $\left(u_{d}, v_{d}\right)$, the change of the angle $\delta$ reflects the change of the surge and sway velocities of the ship with the total ship velocity $\bar{u}=\sqrt{u^{2}+v^{2}}$ as time

Table 8.1 Normalization variables used for the Bis system

| Variable | Bis-system |
| :---: | :---: |
| Time | $\sqrt{L / g}$ |
| Position | $L$ |
| Angle | 1 |
| Linear velocity | $\sqrt{L g}$ |
| Angular velocity | $\sqrt{g / L}$ |
| Force | $m g$ |
| Moment | $m g L$ |

evolutes. The ship and reference trajectories in the $(x, y)$-plane, and the control inputs are plotted in Figure 8.3.

To illustrate how our controller in Section 8.1.5 can compensate for the disturbances, we also simulate with the disturbances as $\boldsymbol{\tau}_{\boldsymbol{E}}=[0.15,0.08,0.05]^{T}$ in the Bis system. This choice of $\boldsymbol{\tau}_{\boldsymbol{E}}$ implies that the disturbance forces on the surge and sway are 0.15 and 0.08 , and the disturbance moment on the yaw is 0.05 . The adaptation gains are chosen as $\gamma_{0 i}=1, k_{0 i}=1$, and $\xi_{0 i}=0.05, i=$ $1,2,3$. With these disturbances, one can verify that condition (8.59) holds, i.e., $\left.\left(\left(\alpha \bar{u}_{d}+|\gamma|\right)\left|\bar{r}_{d}\right|+\max \left(\left|\bar{\tau}_{v E}\right|\right)+\xi_{01}\right) /\left(\beta \bar{u}_{d}+\dot{\bar{u}}_{d}\right)\right) \leq 0.71$. Simulation results are plotted in Figures 8.4 and 8.5. It is seen that our controller forces the heading angle to a value $\delta(t)$ to compensate for the disturbances.


Figure 8.2 Simulation results without disturbances: a. Tracking errors $\left(\left(x-x_{d}\right),\left(y-y_{d}\right),(\psi-\right.$ $\left.\psi_{d}\right)$ ); b. Variation of angle $\delta$


Figure 8.3 Simulation results without disturbances: a. Ship position and orientation in the $(x, y)$ plane; $\mathbf{b}$. Control input $\tau_{u} ; \mathbf{c}$. Control input $\tau_{r}$


Figure 8.4 Simulation results with disturbances: a. Tracking errors $\left(\left(x-x_{d}\right),\left(y-y_{d}\right),(\psi-\right.$ $\left.\psi_{d}\right)$ ); b. Variation of angle $\delta$


Figure 8.5 Simulation results with disturbances: a. Ship position and orientation in the $(x, y)$ plane; $\mathbf{b}$. Control input $\tau_{u}$; $\mathbf{c}$. Control input $\tau_{r}$

### 8.2 Output Feedback

### 8.2.1 Control Objective

For the reader's convenience the mathematical model of an underactuated ship moving in a horizontal plane is rewritten here, see Section 3.4.1.1:

$$
\begin{align*}
& \dot{\eta}=\boldsymbol{J}(\psi) \boldsymbol{v} \\
& \boldsymbol{M} \dot{\boldsymbol{v}}=-\boldsymbol{C}(\boldsymbol{v}) \boldsymbol{v}-\left(\boldsymbol{D}+\boldsymbol{D}_{\boldsymbol{n}}(\boldsymbol{v})\right) \boldsymbol{v}+\boldsymbol{\tau} \tag{8.64}
\end{align*}
$$

where

$$
\begin{gathered}
\boldsymbol{\eta}=\left[\begin{array}{l}
x \\
y \\
\psi
\end{array}\right], \boldsymbol{v}=\left[\begin{array}{l}
u \\
v \\
r
\end{array}\right], \boldsymbol{\tau}=\left[\begin{array}{c}
\tau_{u} \\
0 \\
\tau_{r}
\end{array}\right], \\
\boldsymbol{J}(\psi)=\left[\begin{array}{ccc}
\cos (\psi)-\sin (\psi) & 0 \\
\sin (\psi) & \cos (\psi) & 0 \\
0 & 0 & 1
\end{array}\right], \boldsymbol{M}=\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & m_{23} \\
0 & m_{32} & m_{33}
\end{array}\right], \\
\boldsymbol{C}(\boldsymbol{v})=\left[\begin{array}{ccc}
0 & 0 & C_{13} \\
0 & 0 & C_{23} \\
C_{31} & C_{32} & 0
\end{array}\right], \boldsymbol{D}+\boldsymbol{D}_{\boldsymbol{n}}(\boldsymbol{v})=\left[\begin{array}{ccc}
D_{11} & 0 & 0 \\
0 & D_{22} & D_{23} \\
0 & D_{32} & D_{33}
\end{array}\right],
\end{gathered}
$$

with

$$
\begin{aligned}
m_{11} & =m-X_{\dot{u}}, m_{22}=m-Y_{\dot{v}}, m_{23}=m x_{g}-Y_{\dot{r}} \\
m_{32} & =m x_{g}-N_{\dot{v}}, m_{33}=I_{z}-N_{\dot{r}} \\
C_{13} & =-C_{31}=-\left(m-Y_{\dot{v}}\right) v-\left(m x_{g}-0.5\left(Y_{\dot{r}}+N_{\dot{v}}\right)\right) r \\
C_{23} & =-C_{32}=\left(m-X_{\dot{u}}\right) u \\
D_{11} & =-\left(X_{u}+X_{|u| u}|u|\right), D_{22}=-\left(Y_{v}+Y_{|v| v}|v|\right) \\
D_{23} & =-Y_{r}, D_{32}=-N_{v}, D_{33}=-\left(N_{r}+N_{|r| r}|r|\right)
\end{aligned}
$$

All the symbols used here are defined in Section 3.3. It is noted that we did not include environmental disturbances in the ship dynamics (8.64). In the control design, we will add integrators to the path-tracking error outputs to compensate for the disturbances. In this section, we consider the following control objective.

Control Objective. Assume that the ship velocities, $u, v$, and $r$ are not available and that Assumption 8.2 holds, design the controls $\tau_{u}$ and $\tau_{r}$ to force the ship (8.64) to follow a prescribed path $\boldsymbol{\Omega}$ parameterized by $\left(x_{d}(s), y_{d}(s)\right)$ with $s$ being the path parameter in the sense that the position of the ship (8.64) globally tracks the path $\boldsymbol{\Omega}$, the ship total linear velocity is tangential to the path $\boldsymbol{\Omega}$, and let the desired surge speed, $u_{d 0}(t)$, be adjusted on-line. The desired surge speed, $u_{d 0}(t)$, is assumed to be bounded and twice differentiable.

Assumption 8.2. There exist strictly positive constants $\varepsilon_{i}, 0 \leq i \leq 5$ such that

$$
\begin{gather*}
\varepsilon_{0} \leq x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s) \leq \varepsilon_{1}, \forall s \in \mathbb{R},  \tag{8.65}\\
u_{d 0}(t) \geq \varepsilon_{2}, \forall t \geq 0,  \tag{8.66}\\
N_{\dot{v}}=Y_{\dot{r}},  \tag{8.67}\\
1-\left(\alpha\left(u_{d 0}(t)+\varepsilon_{3}\right)+|\gamma|\right) /\left(u_{d 0}(t)-\varepsilon_{4}\right) \geq \varepsilon_{5}, \forall t \geq 0, \tag{8.68}
\end{gather*}
$$

where $x_{d}^{\prime}(s), y_{d}^{\prime}(s), \alpha$ and $\gamma$ are defined as

$$
\begin{align*}
& x_{d}^{\prime}(s)=\frac{\partial x_{d}}{\partial s}, y_{d}^{\prime}(s)=\frac{\partial y_{d}}{\partial s}, \\
& \alpha=\frac{m_{11}}{m_{22}}, \gamma=\frac{Y_{r}}{m_{22}} . \tag{8.69}
\end{align*}
$$

Remark 8.4. The following observations are made on Assumption 8.2.

1. Condition (8.65) implies that the path $\boldsymbol{\Omega}$ is regular with respect to the path parameter $s$. If the path is not regular, one can break it into different regular paths and consider each path separately.
2. Condition (8.66) means that the path $\boldsymbol{\Omega}$ does not contain or approach a set-point, i.e., we do not consider stabilization/regulation problems. For clarity we only consider forward tracking. Backward tracking can be done similarly.
3. Condition (8.67) implies that the added mass matrix is symmetric. This condition is needed for designing an observer. If this condition does not hold, we can use
an acceleration feedback in the yaw dynamics [12] to reshape the added mass matrix.
4. Condition (8.68) is needed to guarantee stability of the sway dynamics under the proposed controller (to be designed later). Since $\alpha$ is strictly smaller than 1 (added mass, $-X_{\dot{u}}$, in surge is smaller than that, $-Y_{\dot{v}}$, in sway) and $|\gamma|$ is small for most ships, Condition (8.68) is satisfied under (8.66) with $\varepsilon_{2}$ being strictly larger than $|\gamma|$.

The main ideas in solving the above control objective are as follows:

1. Choose an appropriate body-fixed frame origin to avoid the yaw moment control $\tau_{r}$ acting directly on the sway dynamics. This choice will overcome the difficulty caused by off-diagonal terms in the mass matrix.
2. Transform the ship dynamics to a system, which is linear and monotonic in unmeasured velocities, so that an observer can be easily designed.
3. Interpret path-tracking errors in a frame attached to the path $\boldsymbol{\Omega}$ such that the error dynamics are of a triangular form to which the backstepping technique can be applied.
4. Use the orientation error as a virtual control to stabilize the cross-track error where a filter of the sway velocity is designed and used in the virtual control. This filter and the derivative of the path parameter used as an additional control allow a global controller and desired path with arbitrary curvature.
5. Add projection integrators to the path-tracking error dynamics' output to compensate for constant bias of the disturbances.


Figure 8.6 Interpretation of path-tracking errors

### 8.2.2 Coordinate Transformations

### 8.2.2.1 Choosing the Body-fixed Frame Origin

To avoid the yaw moment control $\tau_{r}$ acting directly on the sway dynamics, we choose the body-fixed frame origin such that it is on the center line of the ship (see Figure 8.6), i.e., $y_{g}=0$ and that:

$$
\begin{equation*}
x_{g}=\frac{Y_{\dot{r}}}{m} \tag{8.70}
\end{equation*}
$$

With this choice, the ship model (8.64) can be written as:

$$
\begin{align*}
& \dot{\eta}=J(\psi) v  \tag{8.71}\\
& \bar{M} \dot{v}=-\bar{C}(v) v-\left(\bar{D}+\bar{D}_{n}(v)\right) v+\tau
\end{align*}
$$

where

$$
\begin{align*}
& \overline{\boldsymbol{M}}=\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & 0 \\
0 & 0 & m_{33}
\end{array}\right], \overline{\boldsymbol{D}}_{\boldsymbol{n}}(\boldsymbol{v})=\left[\begin{array}{ccc}
d_{n 1}|u| & 0 & 0 \\
0 & d_{n 2}|v| & 0 \\
0 & 0 & d_{n 3}|r|
\end{array}\right], \\
& \overline{\boldsymbol{D}}=\left[\begin{array}{ccc}
d_{11} & 0 & 0 \\
0 & d_{22} & d_{23} \\
0 & d_{32} & d_{33}
\end{array}\right], \overline{\boldsymbol{C}}(\boldsymbol{v})=\left[\begin{array}{ccc}
0 & 0 & -m_{22} v \\
0 & 0 & m_{11} u \\
m_{22} v & -m_{11} u & 0
\end{array}\right], \tag{8.72}
\end{align*}
$$

with

$$
\begin{aligned}
& d_{11}=-X_{u}, d_{22}=-Y_{v} \\
& d_{23}=-Y_{r}, d_{32}=-N_{v} \\
& d_{33}=-N_{r}, d_{n 1}=-X_{|u| u}, \\
& d_{n 2}=Y_{|v| v}, d_{n 3}=N_{|r| r}
\end{aligned}
$$

It is noted that the ship position $(x, y)$ is the coordinates of the body-fixed frame origin (not the ship's center of gravity) with respect to the earth-fixed frame.

### 8.2.2.2 Transformation of the Ship Dynamics to be Linear and Monotonic in Unmeasured State System

The transformation here is very similar to that in Section 7.3.1. For the sake of completeness, we give the full derivation of the transformation. As such, we first remove the term $\overline{\boldsymbol{C}}(\boldsymbol{v}) \boldsymbol{v}$ in the right-hand side of the second equation of (8.71), which causes difficulties in the observer design, by introducing the following coordinate:

$$
\begin{equation*}
X=P(\eta) v, \tag{8.73}
\end{equation*}
$$

where $\boldsymbol{P}(\boldsymbol{\eta}) \in \mathbb{R}^{3 \times 3}$ is a global invertible matrix to be designed. With (8.73), the second equation of (8.71) is written as

$$
\begin{align*}
\dot{X}= & {\left[\dot{P}(\eta) v-P(\eta) \bar{M}^{-1} \bar{C}(v) v\right]-} \\
& P(\eta) \bar{M}^{-1}\left(\bar{D}^{\prime} \bar{D}_{n}\left(P^{-1}(\eta) X\right)\right) P^{-1}(\eta) X+P(\eta) \bar{M}^{-1} \tau \tag{8.74}
\end{align*}
$$

It can be seen that in order to remove the term $\overline{\boldsymbol{C}}(\boldsymbol{v}) \boldsymbol{v}$, we need to find the matrix $\boldsymbol{P}(\boldsymbol{\eta})$ such that

$$
\begin{equation*}
\left[\dot{P}(\eta) v-P(\eta) \bar{M}^{-\mathbf{1}} \bar{C}(v) v\right]=0 \tag{8.75}
\end{equation*}
$$

for all $(\boldsymbol{\eta}, \boldsymbol{v}) \in \mathbb{R}^{6}$. Assuming that the elements of $\boldsymbol{P}(\boldsymbol{\eta})$ are $p_{i j}(\boldsymbol{\eta}), i=1,2,3, j=$ $1,2,3,(8.75)$ is expanded to:

$$
\begin{align*}
& \dot{p}_{i 1} u+\dot{p}_{i 2} v+\dot{p}_{i 3} r+\frac{m_{22}}{m_{11}} p_{i 1} v r-\frac{m_{11}}{m_{22}} p_{i 2} u r+\frac{m_{11}-m_{22}}{m_{33}} p_{i 3} u v=0, \\
& i=1,2,3, \forall(\boldsymbol{\eta}, u, v, r) \in \mathbb{R}^{6}, \tag{8.76}
\end{align*}
$$

where for brevity, we omit the argument $\eta$ of $p_{i j}(\eta)$. With the first equation of (8.71), we expand (8.76) as

$$
\begin{align*}
& \left(\frac{\partial p_{i 1}}{\partial x} \cos (\psi)+\frac{\partial p_{i 1}}{\partial y} \sin (\psi)\right) u^{2}+\left(-\frac{\partial p_{i 2}}{\partial x} \sin (\psi)+\frac{\partial p_{i 2}}{\partial y} \cos (\psi)\right) v^{2}+ \\
& \frac{\partial p_{i 3}}{\partial \psi} r^{2}+\left(-\frac{\partial p_{i 1}}{\partial x} \sin (\psi)+\frac{\partial p_{i 1}}{\partial y} \cos (\psi)+\frac{\partial p_{i 2}}{\partial x} \cos (\psi)+\frac{\partial p_{i 2}}{\partial y} \sin (\psi)+\right. \\
& \left.\frac{m_{11}-m_{22}}{m_{33}} p_{i 3}\right) u v+\left(\frac{\partial p_{i 1}}{\partial \psi}+\frac{\partial p_{i 3}}{\partial x} \cos (\psi)+\frac{\partial p_{i 3}}{\partial y} \sin (\psi)-\frac{m_{11}}{m_{22}} p_{i 2}\right) u r+ \\
& \left(\frac{\partial p_{i 2}}{\partial \psi}-\frac{\partial p_{i 3}}{\partial x} \sin (\psi)+\frac{\partial p_{i 3}}{\partial y} \cos (\psi)+\frac{m_{22}}{m_{11}} p_{i 1}\right) v r=0 . \tag{8.77}
\end{align*}
$$

Therefore (8.77) holds for all $(x, y, \psi, u, v, r) \in \mathbb{R}^{6}$ if

$$
\begin{align*}
& \frac{\partial p_{i 1}}{\partial x} \cos (\psi)+\frac{\partial p_{i 1}}{\partial y} \sin (\psi)=0 \\
& -\frac{\partial p_{i 2}}{\partial x} \sin (\psi)+\frac{\partial p_{i 2}}{\partial y} \cos (\psi)=0 \\
& \frac{\partial p_{i 3}}{\partial \psi}=0 \\
& \left(\frac{\partial p_{i 2}}{\partial y}-\frac{\partial p_{i 1}}{\partial x}\right) \sin (\psi)+\left(\frac{\partial p_{i 1}}{\partial y}+\frac{\partial p_{i 2}}{\partial x}\right) \cos (\psi)+\frac{m_{11}-m_{22}}{m_{33}} p_{i 3}=0, \\
& \frac{\partial p_{i 1}}{\partial \psi}+\frac{\partial p_{i 3}}{\partial x} \cos (\psi)+\frac{\partial p_{i 3}}{\partial y} \sin (\psi)-\frac{m_{11}}{m_{22}} p_{i 2}=0, \\
& \frac{\partial p_{i 2}}{\partial \psi}-\frac{\partial p_{i 3}}{\partial x} \sin (\psi)+\frac{\partial p_{i 3}}{\partial y} \cos (\psi)+\frac{m_{22}}{m_{11}} p_{i 1}=0 . \tag{8.78}
\end{align*}
$$

A family of solutions of the above set of six partial differential equations is

$$
\begin{align*}
& p_{i 1}=\frac{\left(\left(m_{11} C_{i 3} x+m_{33} C_{i 1}\right) \sin (\psi)-\left(m_{11} C_{i 3} y-m_{33} C_{i 2}\right) \cos (\psi)\right)}{m_{33}}, \\
& p_{i 2}=\frac{m_{22}\left(\left(m_{11} C_{i 3} x+m_{33} C_{i 1}\right) \cos (\psi)+\left(m_{11} C_{i 3} y-m_{33} C_{i 2}\right) \sin (\psi)\right)}{m_{11} m_{33}}, \\
& p_{i 3}=C_{i 3}, \tag{8.79}
\end{align*}
$$

where $C_{i 1}, C_{i 2}$, and $C_{i 3}$ are arbitrary constants.
We now choose the constants $C_{i 1}, C_{i 2}$, and $C_{i 3}$ such that the matrix $\boldsymbol{P}(\eta)$ is invertible for all $(x, y, \psi) \in \mathbb{R}^{3}$. A trivial choice of $C_{13}=C_{11}=0, C_{12}=1, C_{23}=$ $C_{22}=0, C_{21}=1, C_{31}=C_{32}=0$, and $C_{33}=1$ results in

$$
\boldsymbol{P}(\boldsymbol{\eta})=\left[\begin{array}{ccc}
\cos (\psi) & -\sin (\psi) \frac{m_{22}}{m_{11}} & 0  \tag{8.80}\\
\sin (\psi) & \cos (\psi) \frac{m_{22}}{m_{11}} & 0 \\
(\sin (\psi) x-\cos (\psi) y) \frac{m_{11}}{m_{33}} & (\cos (\psi) x+\sin (\psi) y) \frac{m_{22}}{m_{33}} & 1
\end{array}\right]
$$

Indeed this matrix is invertible for all $(x, y, \psi) \in \mathbb{R}^{3}$. Substituting (8.80) into (8.74) gives

$$
\begin{equation*}
\dot{X}=-P(\eta) \bar{M}^{-1}\left(\bar{D}+\bar{D}_{n}\left(P^{-1}(\eta) X\right)\right) P^{-1}(\eta) X+P(\eta) \bar{M}^{-1} \tau \tag{8.81}
\end{equation*}
$$

### 8.2.3 Observer Design

It can be seen that system (8.81) is linear and monotonic in the unmeasured state vector $\boldsymbol{X}$. We now design a simple observer to estimate $\boldsymbol{X}$ as:

$$
\begin{equation*}
\dot{\hat{X}}=-P(\eta) \bar{M}^{-1}\left(\bar{D}+\bar{D}_{n}\left(P^{-1}(\eta) \hat{X}\right)\right) P^{-1}(\eta) \hat{X}+P(\eta) \bar{M}^{-1} \tau \tag{8.82}
\end{equation*}
$$

where $\hat{\boldsymbol{X}}$ is an estimate of $\boldsymbol{X}$. From (8.81) and (8.82), we have

$$
\begin{align*}
\dot{\tilde{X}}= & -P(\eta) \bar{M}^{-\mathbf{1}} \bar{D} P^{-1}(\eta) \tilde{X}-P(\eta) \bar{M}^{-1} \times \\
& {\left.\left.\left[\bar{D}_{n}\left(P^{-1}(\eta) X\right)\right) P^{-1}(\eta) X-\bar{D}_{n}\left(P^{-1}(\eta) \hat{X}\right)\right) P^{-1}(\eta) \hat{X}\right], } \tag{8.83}
\end{align*}
$$

where $\tilde{\boldsymbol{X}}:=\left(\tilde{x}_{1}, \tilde{x}_{2}, \tilde{x}_{3}\right)^{T}=\boldsymbol{X}-\hat{\boldsymbol{X}}$. We will use (8.83) in stability analysis of the closed loop system later. Define $\hat{\boldsymbol{v}}=[\hat{u}, \hat{v}, \hat{r}]^{T}$ as being an estimate of the velocity vector $\boldsymbol{v}$ as:

$$
\begin{equation*}
\hat{v}=P^{-1}(\eta) \hat{X} . \tag{8.84}
\end{equation*}
$$

Then the velocity estimate error vector, $\tilde{\boldsymbol{v}}:=[\tilde{u}, \tilde{v}, \tilde{r}]^{T}=\boldsymbol{v}-\hat{\boldsymbol{v}}$, satisfies:

$$
\left[\begin{array}{c}
\tilde{u}  \tag{8.85}\\
\tilde{v} \\
\tilde{r}
\end{array}\right]=\left[\begin{array}{l}
\tilde{x}_{1} \cos (\psi)+\tilde{x}_{2} \sin (\psi) \\
\frac{m_{11}}{m_{22}}\left(-\tilde{x}_{1} \sin (\psi)+\tilde{x}_{2} \cos (\psi)\right) \\
\frac{m_{11}}{m_{33}}\left(\tilde{x}_{1} y-\tilde{x}_{2} x\right)+\tilde{x}_{3}
\end{array}\right]
$$

Using (8.84), we rewrite the first equation of (8.71) and (8.82) as:

$$
\begin{align*}
\dot{x} & =(\hat{u}+\tilde{u}) \cos (\psi)-(v+\tilde{v}) \sin (\psi), \\
\dot{y} & =(\hat{u}+\tilde{u}) \sin (\psi)+(v+\tilde{v}) \cos (\psi), \\
\dot{\psi} & =\hat{r}+\tilde{r}, \\
\dot{\hat{u}} & =\frac{m_{22}}{m_{11}} \hat{v} \hat{r}-\frac{d_{11}}{m_{11}} \hat{u}-\frac{d_{n 1}}{m_{11}}|\hat{u}| \hat{u}+\frac{1}{m_{11}} \tau_{u}+\frac{m_{22}}{m_{11}} \hat{v} \tilde{r},  \tag{8.86}\\
\dot{\hat{v}} & =-\frac{m_{11}}{m_{22}} \hat{u} \hat{r}-\frac{d_{22}}{m_{22}} \hat{v}-\frac{d_{23}}{m_{22}} \hat{r}-\frac{d_{n 2}}{m_{22}}|\hat{v}| \hat{v}-\frac{m_{11}}{m_{22}} \hat{u} \tilde{r}, \\
\dot{\hat{r}} & =\frac{m_{11}-m_{22}}{m_{33}} \hat{u} \hat{v}-\frac{d_{22}}{m_{33}} \hat{v}-\frac{d_{33}}{m_{33}} \hat{r}-\frac{d_{n 3}}{m_{33}}|\hat{r}| \hat{r}+ \\
& \frac{1}{m_{33}} \tau_{r}+\frac{m_{11}}{m_{33}} \hat{u} \tilde{v}-\frac{m_{22}}{m_{33}} \hat{v} \tilde{u} .
\end{align*}
$$

## Transformation of Path-tracking Errors

We now interpret the path-tracking errors in a frame attached to the reference path $\boldsymbol{\Omega}$ as follows (see Figure 8.6):

$$
\left[\begin{array}{l}
x_{e}  \tag{8.87}\\
y_{e} \\
\psi_{e}
\end{array}\right]=\boldsymbol{J}^{\boldsymbol{T}}(\psi)\left[\begin{array}{c}
x-x_{d} \\
y-y_{d} \\
\psi-\psi_{d}
\end{array}\right],
$$

where $\psi_{d}$ is the angle between the path and the X -axis defined by

$$
\begin{equation*}
\psi_{d}=\arctan \left(\frac{y_{d}^{\prime}(s)}{x_{d}^{\prime}(s)}\right) \tag{8.88}
\end{equation*}
$$

with $x_{d}^{\prime}(s)$ and $y_{d}^{\prime}(s)$ being defined in (8.69).
In Figure 8.6, $O_{E} X_{E} Y_{E}$ is the earth-fixed frame; $O_{p} X_{p} Y_{p}$ is a frame attached to the path $\Omega$ such that $O_{p} X_{p}$ and $O_{p} Y_{p}$ are parallel to the surge and sway axes of the ship, respectively, $u_{d}$ is tangential to the path, $C G$ is the center of gravity of the ship; and $O_{b} X_{b} Y_{b}$ is the body-fixed frame. Therefore $x_{e}, y_{e}$, and $\psi_{e}$ can be referred to as tangential, cross and heading errors, respectively. Differentiating both sides of (8.87) along the solutions of the first three equations of (8.86) results in the kinematic path-tracking errors:

$$
\begin{align*}
\dot{x}_{e} & =\hat{u}-u_{d} \cos \left(\psi_{e}\right)+(\hat{r}+\tilde{r}) y_{e}+\tilde{u}, \\
\dot{y}_{e} & =\hat{v}+u_{d} \sin \left(\psi_{e}\right)-(\hat{r}+\tilde{r}) x_{e}+\tilde{v}, \\
\dot{\psi}_{e} & =\hat{r}-r_{d}+\tilde{r}, \tag{8.89}
\end{align*}
$$

where $u_{d}$ and $r_{d}$ are given by:

$$
\begin{align*}
u_{d} & =\sqrt{x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)} \dot{s} \\
r_{d} & =\frac{x_{d}^{\prime}(s) y_{d}^{\prime \prime}(s)-x_{d}^{\prime \prime}(s) y_{d}^{\prime}(s)}{x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)} \dot{s} \tag{8.90}
\end{align*}
$$

### 8.2.4 Control Design with Integral Action

There are often two methods to let a control system compensate for constant (or at least slowly-varying) environmental disturbances: (1) deriving an adaptive control law to estimate a constant bias (as the one in Section 8.1.5), and (2) adding extra integrators to the output of the path-tracking error system. Here we use the second method for simplicity since the first one will result in a complicated control law due to the observer developed in the previous section, see Section 8.2.6.1 for a discussion. We divide the control design procedure into two separate stages. At the first stage, we consider (8.89) with $\hat{u}$ and $\hat{r}$ being viewed as controls. At the second stage, the last three equations of (8.86) are considered to design the actual controls $\tau_{u}$ and $\tau_{r}$ using the backstepping technique.

### 8.2.4.1 Kinematic Control Design

Substep 1 (Stabilizing $\left(x_{e}, y_{e}\right)$-dynamics)
First we add two projection integrators to the position path-tracking error dynamics as

$$
\begin{align*}
& \dot{\sigma}_{x e}=\operatorname{proj}\left(\gamma_{1} x_{e} / \Delta, \sigma_{x e}\right), \\
& \dot{\sigma}_{y e}=\operatorname{proj}\left(\gamma_{2} y_{e} / \Delta, \sigma_{y e}\right), \\
& \dot{x}_{e}=\hat{u}-u_{d} \cos \left(\psi_{e}\right)+(\hat{r}+\tilde{r}) y_{e}+\tilde{u}, \\
& \dot{y}_{e}=\hat{v}+u_{d} \sin \left(\psi_{e}\right)-(\hat{r}+\tilde{r}) x_{e}+\tilde{v}, \tag{8.91}
\end{align*}
$$

where $\Delta=\sqrt{1+x_{e}^{2}+y_{e}^{2}}, \gamma_{1}$ and $\gamma_{2}$ are positive constants and the operator proj represents the Lipschitz projection algorithm (repeated here for the reader's convenience)

$$
\begin{aligned}
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \text { if } \Xi(\hat{\omega}) \leq 0, \\
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \text { if } \Xi(\hat{\omega}) \geq 0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi \leq 0, \\
& \operatorname{proj}(\varpi, \hat{\omega})=(1-\Xi(\hat{\omega})) \varpi \text { if } \Xi(\hat{\omega})>0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi>0,
\end{aligned}
$$

where $\Xi(\hat{\omega})=\left(\hat{\omega}^{2}-\omega_{M}^{2}\right) /\left(\xi^{2}+2 \xi \omega_{M}\right), \Xi_{\hat{\omega}}(\hat{\omega})=\partial \Xi(\hat{\omega}) / \partial \hat{\omega}, \xi$ is an arbitrarily small positive constant and $\omega_{M}$ is the maximum value of the constant bias that $\hat{\omega}$ will compensate for. The projection algorithm (repeated here for convenience of the reader) is such that if $\dot{\hat{\omega}}=\operatorname{proj}(\varpi, \hat{\omega})$ and $\hat{\omega}\left(t_{0}\right) \leq \omega_{M}$ then:

1. $\hat{\omega}(t) \leq \omega_{M}+\xi, \forall 0 \leq t_{0} \leq t<\infty$,
2. $\operatorname{proj}(\varpi, \hat{\omega})$ is Lipschitz continuous,
3. $|\operatorname{proj}(\varpi, \hat{\omega})| \leq|\varpi|$,
4. $\tilde{\omega} \operatorname{proj}(\varpi, \hat{\omega}) \geq \tilde{\omega} \varpi$ with $\tilde{\omega}=\omega-\hat{\omega}$.

It is noted that we have added the integrators via the projection algorithm to guarantee that $\sigma_{x e}$ and $\sigma_{y e}$ are always bounded by some predefined constants. Define

$$
\begin{align*}
& u_{e}=\hat{u}-\alpha_{u} \\
& \bar{\psi}_{e}=\psi_{e}-\alpha_{\psi_{e}} \tag{8.92}
\end{align*}
$$

where $\alpha_{u}$ and $\alpha_{\psi_{e}}$ are virtual controls of $\hat{u}$ and $\psi_{e}$, respectively. Substituting (8.92) into the last two equations of (8.91) gives

$$
\begin{align*}
\dot{x}_{e}= & \alpha_{u}+u_{e}-u_{d} \cos \left(\alpha_{\psi_{e}}\right)-u_{d}\left(\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \cos \left(\alpha_{\psi_{e}}\right)-\right. \\
& \left.\sin \left(\bar{\psi}_{e}\right) \sin \left(\alpha_{\psi_{e}}\right)\right)+(\hat{r}+\tilde{r}) y_{e}+\tilde{u}, \\
\dot{y}_{e}= & \hat{v}+u_{d} \sin \left(\alpha_{\psi_{e}}\right)+u_{d}\left(\sin \left(\bar{\psi}_{e}\right) \cos \left(\alpha_{\psi_{e}}\right)+\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \times\right. \\
& \left.\sin \left(\alpha_{\psi_{e}}\right)\right)-(\hat{r}+\tilde{r}) x_{e}+\tilde{v} . \tag{8.93}
\end{align*}
$$

From the first equation of (8.91) and first equation of (8.93), we can easily design a control law for $\alpha_{u}$ to stabilize the $x_{e}$-dynamics as follows

$$
\begin{equation*}
\alpha_{u}=-\varphi_{1}-\sigma_{x e}+u_{d} \cos \left(\alpha_{\psi_{e}}\right) \tag{8.94}
\end{equation*}
$$

with

$$
\varphi_{1}=k_{1} \frac{x_{e}}{\Delta}
$$

where $k_{1}$ is a positive constant to be specified later. It is noted that we have chosen $\alpha_{u}$ not to cancel the known term, $-u_{d}\left(\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \cos \left(\alpha_{\psi_{e}}\right)-\sin \left(\bar{\psi}_{e}\right) \sin \left(\alpha_{\psi_{e}}\right)\right)+$ $\hat{r} y_{e}$, to simplify the controller expression.

However, stabilizing the second equation of (8.91) and second equation of (8.93) is more difficult. We cannot directly use $\alpha_{\psi_{e}}$ to cancel $\hat{v}$ since the term $u_{d} \sin \left(\alpha_{\psi_{e}}\right)$ cannot cancel $\hat{v}$ without imposing a restriction on the initial conditions, i.e., no global result can be obtained. To resolve this difficulty, the derivative of the path parameter is used as an additional control and is designed as:

$$
\begin{equation*}
\dot{s}=\frac{\sqrt{u_{d 0}^{2}+\sigma_{x e}^{2}+\left(\varphi_{2}+\sigma_{y e}+v_{d}\right)^{2}}}{\sqrt{x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)}} \tag{8.95}
\end{equation*}
$$

with $\varphi_{2}=k_{2} \frac{y_{e}}{\Delta}$, where from now we drop the argument $t$ of $u_{d 0}$ and its derivatives, $k_{2}$ is a positive constant to be specified later, and $v_{d}$ is a filtered value of $\hat{v}$ to be designed later. From (8.95) and (8.90), we have:

$$
\begin{equation*}
u_{d}=\sqrt{u_{d 0}^{2}+\sigma_{x e}^{2}+\left(\varphi_{2}+\sigma_{y e}+v_{d}\right)^{2}} \tag{8.96}
\end{equation*}
$$

From (8.96), the second equation of (8.93), and (8.90), we design a control law for $\alpha_{\psi_{e}}$ as:

$$
\begin{equation*}
\alpha_{\psi_{e}}=-\arctan \left(\frac{\varphi_{2}+\sigma_{y e}+v_{d}}{\Upsilon_{u_{d 0}}}\right) \tag{8.97}
\end{equation*}
$$

with

$$
\begin{equation*}
\Upsilon_{u_{d 0}}=\sqrt{u_{d 0}^{2}+\sigma_{x e}^{2}} \tag{8.98}
\end{equation*}
$$

Substituting (8.97), (8.95), and (8.94) into (8.93) results in

$$
\begin{align*}
\dot{x}_{e}= & -k_{1} x_{e} / \Delta-\sigma_{x e}+u_{e}-u_{d}\left(\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \cos \left(\alpha_{\psi_{e}}\right)-\right. \\
& \left.\sin \left(\bar{\psi}_{e}\right) \sin \left(\alpha_{\psi_{e}}\right)\right)+(\hat{r}+\tilde{r}) y_{e}+\tilde{u} \\
\dot{y}_{e}= & -k_{2} y_{e} / \Delta-\sigma_{y e}+v_{e}+u_{d}\left(\sin \left(\bar{\psi}_{e}\right) \cos \left(\alpha_{\psi_{e}}\right)+\right. \\
& \left.\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \sin \left(\alpha_{\psi_{e}}\right)\right)-(\hat{r}+\tilde{r}) x_{e}+\tilde{v}, \tag{8.99}
\end{align*}
$$

where $v_{e}=\hat{v}-v_{d}$.
Substep 2 (Stabilizing $\bar{\psi}_{e}$ dynamics)
Define

$$
\begin{equation*}
r_{e}=\hat{r}-\alpha_{r} \tag{8.100}
\end{equation*}
$$

where $\alpha_{r}$ is a virtual control of $\hat{r}$. Differentiating both sides of the second equation of (8.92) along the solutions of (8.97), and adding a projection integrator gives
$\dot{\sigma}_{\bar{\psi} e}=\operatorname{proj}\left(\gamma_{3} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}, \sigma_{\bar{\psi} e}\right)$,
$\dot{\bar{\psi}}_{e}=b_{1}\left(r_{e}+\alpha_{r}+\tilde{r}\right)-r_{d}+\frac{\Upsilon_{u_{d 0}}\left(\Pi+\dot{v}_{d}+\dot{\sigma}_{y e}\right)}{u_{d}^{2}}-$
$\frac{\left(\varphi_{2}+\sigma_{y e}+v_{d}\right)\left(\dot{u}_{d 0} u_{d 0}+\dot{\sigma}_{x e} \sigma_{x e}\right)}{u_{d}^{2} \Upsilon_{u_{d 0}}}-\frac{k_{2} \Upsilon_{u_{d 0}}\left(x_{e} y_{e}\left(u_{e}+\tilde{u}\right)+\left(1+x_{e}^{2}\right)\left(v_{e}+\tilde{v}\right)\right)}{u_{d}^{2} \Delta^{3}}$,
where $\gamma_{3}$ is a positive constant and

$$
\begin{align*}
b_{1} & =1-k_{2} \frac{\Upsilon_{u_{d 0}} x_{e}}{u_{d}^{2} \Delta^{3}} \\
\Pi & =k_{2}\left(\frac{v_{d}+u_{d} \sin \left(\psi_{e}\right)}{\Delta}-\frac{y_{e}\left(x_{e}\left(\alpha_{u}-u_{d} \cos \left(\psi_{e}\right)\right)+y_{e}\left(v_{d}+u_{d} \sin \left(\psi_{e}\right)\right)\right.}{\Delta^{3}}\right) \tag{8.102}
\end{align*}
$$

From the first equation of (8.102), we choose the design constant $k_{2}$ such that:

$$
\begin{equation*}
k_{2} \leq \min \left(u_{d 0}\right)-\varepsilon_{6}, \tag{8.103}
\end{equation*}
$$

where $\varepsilon_{6}$ is a strictly positive constant. Then $b_{1}$ is always larger than some strictly positive constant. Now the virtual control $\alpha_{r}$ can be designed from (8.101) as

$$
\begin{align*}
\alpha_{r}= & \frac{1}{b_{1}}\left(-\frac{k_{3} \bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}-\sigma_{\bar{\psi} e}+r_{d}-\frac{\Upsilon_{u_{d 0}}}{u_{d}^{2}}\left(\Pi+\dot{v}_{d}+\dot{\sigma}_{y e}\right)+\frac{\varphi_{2}+\sigma_{y e}+v_{d}}{u_{d}^{2} \Upsilon_{u_{d 0}}} \times\right. \\
& \left.\left(\dot{u}_{d 0} u_{d 0}+\dot{\sigma}_{x e} \sigma_{x e}\right)\right) \tag{8.104}
\end{align*}
$$

where $k_{3}$ is a positive constant. Substituting (8.104) into the second equation of (8.101) gives:

$$
\begin{align*}
\dot{\bar{\psi}}_{e}= & -k_{3} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}-\sigma_{\bar{\psi} e}-\frac{k_{2} \Upsilon_{u_{d 0}} x_{e} y_{e}}{u_{d}^{2} \Delta^{3}}\left(u_{e}+\tilde{u}\right)+\frac{k_{2} \Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right)}{u_{d}^{2} \Delta^{3}} \times \\
& \left(v_{e}+\tilde{v}\right)+b_{1}\left(r_{e}+\tilde{r}\right) . \tag{8.105}
\end{align*}
$$

To determine $v_{d}$, we differentiate $v_{e}=\hat{v}-v_{d}$ along the solutions of the fifth equation of (8.86) to obtain:

$$
\begin{align*}
\dot{v}_{e}=- & -\frac{m_{11}}{m_{22}}\left(\alpha_{u}+u_{e}\right)\left(\alpha_{r}+r_{e}\right)-\frac{d_{22}}{m_{22}}\left(v_{e}+v_{d}\right)-\frac{d_{23}}{m_{22}}\left(\alpha_{r}+r_{e}\right)- \\
& \frac{d_{n 2}}{m_{22}}|\hat{v}|\left(v_{e}+v_{d}\right)-\frac{m_{11}}{m_{22}} \hat{u} \tilde{r}-\dot{v}_{d} \tag{8.106}
\end{align*}
$$

which suggests that we choose:

$$
\begin{equation*}
\dot{v}_{d}=-\frac{m_{11} \alpha_{u}+d_{23}}{m_{22}} \alpha_{r}-\frac{d_{22}}{m_{22}} v_{d}-\frac{d_{n 2}}{m_{22}}\left|v_{d}\right| v_{d}+\frac{k_{2} \Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right)}{u_{d}^{2} \Delta^{3}} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}} \tag{8.107}
\end{equation*}
$$

where the last term on the right-hand side of (8.107) is added to take care of the second last term in the right-hand side of (8.105). Substituting (8.107) into (8.106) gives:

$$
\begin{align*}
\dot{v}_{e}= & -\frac{d_{22}}{m_{22}} v_{e}-\frac{d_{n 2}}{m_{22}}|\hat{v}| v_{e}-\frac{m_{11} \hat{u}+d_{23}}{m_{22}} r_{e}-\frac{m_{11}}{m_{22}} \hat{r} u_{e}-\frac{d_{23}}{m_{22}} r_{e}- \\
& \frac{m_{11}}{m_{22}}\left(u_{e}+\alpha_{u}\right) \tilde{r}-\frac{k_{2} \Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right)}{u_{d}^{2} \Delta^{3}} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}} . \tag{8.108}
\end{align*}
$$

Now notice that $\alpha_{r}$ depends on $\dot{v}_{d}$, see (8.104). Hence, we substitute (8.104) into (8.107) to obtain:

$$
\begin{align*}
& \dot{v}_{d}=\frac{1}{b_{2}}\left(-\frac{d_{22} v_{d}+d_{n 2}\left|v_{d}\right| v_{d}}{m_{22}}-\frac{m_{11} \alpha_{u}+d_{23}}{b_{1} m_{22}}\left(-\frac{k_{3} \bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}-\sigma_{\bar{\psi} e}+r_{d}-\right.\right. \\
& \left.\left.\frac{\Upsilon_{u_{d 0}}\left(\Pi+\dot{\sigma}_{y e}\right)}{u_{d}^{2}}+\frac{\varphi_{2}+\sigma_{y e}+v_{d}}{u_{d}^{2} \Upsilon_{u_{d 0}}}\left(\dot{u}_{d 0} u_{d 0}+\dot{\sigma}_{x e} \sigma_{x e}\right)\right)-k_{2} \frac{\Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right) \bar{\psi}_{e}}{u_{d}^{2} \Delta^{3} \sqrt{1+\bar{\psi}_{e}^{2}}}\right), \tag{8.109}
\end{align*}
$$

where

$$
\begin{equation*}
b_{2}=1-\frac{m_{11} \alpha_{u}+d_{23}}{m_{22}} \frac{\Upsilon_{u_{d 0}}}{u_{d}^{2} b_{1}} \Rightarrow b_{2} \geq 1-\frac{m_{11}\left(u_{d 0}+k_{1}\right)+\left|d_{23}\right|}{m_{22}\left(u_{d 0}-k_{2}\right)} \tag{8.110}
\end{equation*}
$$

From condition (8.68) in Assumption 8.2, if we pick $k_{1}$ and $k_{2}$ such that

$$
\begin{equation*}
k_{1} \leq \varepsilon_{3}, k_{2} \leq \varepsilon_{4} \tag{8.111}
\end{equation*}
$$

then we have $b_{2} \geq \varepsilon_{5}$, i.e., there is no singularity in (8.109), see Assumption 8.2 for the constants $\varepsilon_{3}, \varepsilon_{4}$ and $\varepsilon_{5}$.

### 8.2.4.2 Kinetic Control Design

Before designing the actual controls $\tau_{u}$ and $\tau_{r}$, we note that the virtual control $\alpha_{u}$ is a smooth function of $x_{e}, y_{e}, u_{d 0}$, and $\sigma_{x e}$, and the virtual control $\alpha_{r}$ is a smooth function of $x_{e}, y_{e}, \psi_{e}, \sigma_{x e}, \sigma_{y e}, \sigma_{\bar{\psi} e}, u_{d 0}, \dot{u}_{d 0}, v_{d}$, and $s$. By differentiating both sides of (8.100) and the first equation of (8.92), and noting (8.101) and (8.108), we choose the actual controls $\tau_{u}$ and $\tau_{r}$ as:

$$
\begin{aligned}
\tau_{u}= & -c_{1} m_{11} u_{e}-m_{22} \hat{v} \hat{r}+\left(d_{11}-d_{n 1}|\hat{u}|\right) \alpha_{u}+ \\
& m_{11}\left(\frac{\partial \alpha_{u}}{\partial x_{e}}\left(\hat{u}-u_{d} \cos \left(\psi_{e}\right)+\hat{r} y_{e}\right)+\frac{\partial \alpha_{u}}{\partial y_{e}}\left(\hat{v}+u_{d} \sin \left(\psi_{e}\right)-\hat{r} x_{e}\right)+\right. \\
& \left.\frac{\partial \alpha_{u}}{\partial \sigma_{x e}} \dot{\sigma}_{x e}+\frac{\partial \alpha_{u}}{\partial u_{d 0}} \dot{u}_{d 0}+\frac{k_{2} \Upsilon_{u_{d 0}} x_{e} y_{e}}{u_{d}^{2} \Delta^{3}} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}+\frac{m_{11}}{m_{22}} \hat{r} v_{e}\right),
\end{aligned}
$$

$$
\begin{align*}
\tau_{r}= & -c_{2} m_{33} r_{e}+\left(d_{33}+d_{n 3}|\hat{r}|\right) \alpha_{r}+d_{32} \hat{v}-\left(m_{11}-m_{22}\right) \hat{u} \hat{v}+ \\
& m_{33}\left(\frac{\partial \alpha_{r}}{\partial x_{e}}\left(\hat{u}-u_{d} \cos \left(\psi_{e}\right)+\hat{r} y_{e}\right)+\frac{\partial \alpha_{r}}{\partial y_{e}}\left(\hat{v}+u_{d} \sin \left(\psi_{e}\right)-\hat{r} x_{e}\right)+\right. \\
& \frac{\partial \alpha_{r}}{\partial s} \dot{s}+\frac{\partial \alpha_{r}}{\partial \psi_{e}}\left(\hat{r}+r_{d}\right)+\frac{\partial \alpha_{r}}{\partial \sigma_{x e}} \dot{\sigma}_{x e}+\frac{\partial \alpha_{r}}{\partial \sigma_{y e}} \dot{\sigma}_{y e}+\frac{\partial \alpha_{r}}{\partial \sigma_{\bar{\psi} e}} \dot{\sigma}_{\bar{\psi} e}+ \\
& \left.\frac{\partial \alpha_{r}}{\partial u_{d 0}} \dot{u}_{d 0}+\frac{\partial \alpha_{r}}{\partial \dot{u}_{d 0}} \ddot{u}_{d 0}+\frac{\partial \alpha_{r}}{\partial v_{d}} \dot{v}_{d}-\frac{b_{1} \bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}+\frac{\left(m_{11} \hat{u}+d_{23}\right) v_{e}}{m_{22}}\right), \tag{8.112}
\end{align*}
$$

where $c_{1}$ and $c_{2}$ are positive constants. The closed loop system is:

$$
\begin{aligned}
\dot{\sigma}_{x e}= & \operatorname{proj}\left(\frac{\gamma_{1} x_{e}}{\Delta}, \sigma_{x e}\right), \dot{\sigma}_{y e}=\operatorname{proj}\left(\frac{\gamma_{2} y_{e}}{\Delta}, \sigma_{y e}\right), \dot{\sigma}_{\bar{\psi} e}=\operatorname{proj}\left(\frac{\gamma_{3} \bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}, \sigma_{\bar{\psi} e}\right), \\
\dot{x}_{e}= & -k_{1} \frac{x_{e}}{\Delta}-\sigma_{x e}+u_{e}-u_{d}\left(\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \cos \left(\alpha_{\psi_{e}}\right)-\right. \\
& \left.\sin \left(\bar{\psi}_{e}\right) \sin \left(\alpha_{\psi_{e}}\right)\right)+(\hat{r}+\tilde{r}) y_{e}+\tilde{u}, \\
\dot{y}_{e}= & -k_{2} y_{e} / \Delta-\sigma_{y e}+v_{e}+u_{d}\left(\sin \left(\bar{\psi}_{e}\right) \cos \left(\alpha_{\psi_{e}}\right)+\right. \\
& \left.\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \sin \left(\alpha_{\psi_{e}}\right)\right)-(\hat{r}+\tilde{r}) x_{e}+\tilde{v}, \\
\dot{\bar{\psi}}_{e}= & -k_{3} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}-\sigma_{\bar{\psi} e}+b_{1}\left(r_{e}+\tilde{r}\right)-\frac{k_{2} \Upsilon_{u_{d 0}} x_{e} y_{e}}{u_{d}^{2} \Delta^{3}}\left(u_{e}+\tilde{u}\right)+ \\
& \frac{k_{2} \Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right)}{u_{d}^{2} \Delta^{3}}\left(v_{e}+\tilde{v}\right), \\
\dot{u}_{e}= & -\left(c_{1}+\frac{d_{11}+d_{n 1}|\hat{u}|}{m_{11}}\right) u_{e}+\frac{m_{22}}{m_{11}} \hat{v} \tilde{r}-\frac{\partial \alpha_{u}}{\partial x_{e}}\left(\tilde{u}+y_{e} \tilde{r}\right)- \\
& \frac{\partial \alpha_{u}}{\partial y_{e}}\left(\tilde{v}-x_{e} \tilde{r}\right)+\frac{k_{2} \Upsilon_{u_{d 0}} x_{e} y_{e}}{u_{d}^{2} \Delta^{3}} \frac{\overline{\psi_{e}}}{\sqrt{1+\bar{\psi}_{e}^{2}}+\frac{m_{11}}{m_{22}} \hat{r} v_{e},} \\
\dot{v}_{e}= & -\frac{d_{22}}{m_{22}} v_{e}-\frac{d_{n 2}}{m_{22}}|\hat{v}| v_{e}-\frac{m_{11} \hat{u}+d_{23}}{m_{22}} r_{e}-\frac{m_{11}}{m_{22}} \hat{r} u_{e}-\frac{d_{23}}{m_{22}} r_{e}- \\
& \frac{m_{11}}{m_{22}}\left(u_{e}+\alpha_{u}\right) \tilde{r}-\frac{k_{2} \Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right)}{u_{d}^{2} \Delta^{3}} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}, \\
\dot{r}_{e}= & -\left(c_{2}+\frac{d_{33}+d_{n 3}|\hat{r}|}{m_{33}}\right) r_{e}+\frac{m_{11}}{m_{33}} \hat{u} \tilde{v}-\frac{m_{22}}{m_{33}} \hat{v} \tilde{u}-\frac{\partial \alpha_{r}}{\partial x_{e}}\left(\tilde{u}+y_{e} \tilde{r}\right)- \\
& \frac{\partial \alpha_{r}}{\partial y_{e}}\left(\tilde{v}-x_{e} \tilde{r}\right)-\frac{\partial \alpha_{r}}{\partial \psi_{e}} \tilde{r}+\frac{\left(m_{11} \hat{u}+d_{23}\right){v_{e}}^{m_{22}}-\frac{b_{1} \overline{\psi_{e}}}{\sqrt{1+\bar{\psi}_{e}^{2}}},}{}
\end{aligned}
$$

$$
\begin{align*}
\dot{v}_{d}= & \frac{1}{b_{2}}\left(-\frac{d_{22} v_{d}+d_{n 2}\left|v_{d}\right| v_{d}}{m_{22}}-\frac{m_{11} \alpha_{u}+d_{23}}{b_{1} m_{22}}\left(-\frac{k_{3} \bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}-\sigma_{\bar{\psi} e}+\right.\right. \\
& \left.r_{d}-\frac{\Upsilon_{u_{d 0}}\left(\Pi+\dot{\sigma}_{y e}\right)}{u_{d}^{2}}+\frac{\varphi_{2}+\sigma_{y e}+v_{d}}{u_{d}^{2} \Upsilon_{u_{d 0}}}\left(\dot{u}_{d 0} u_{d 0}+\dot{\sigma}_{x e} \sigma_{x e}\right)\right)- \\
& \left.k_{2} \frac{\Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right) \bar{\psi}_{e}}{u_{d}^{2} \Delta^{3} \sqrt{1+\bar{\psi}_{e}^{2}}}\right), \\
\dot{x}_{d}= & \frac{\partial x_{d}}{\partial s} \frac{u_{d}}{\sqrt{x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)}}, \dot{y}_{d}=\frac{\partial y_{d}}{\partial s} \frac{u_{d}}{\sqrt{x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)}} . \tag{8.113}
\end{align*}
$$

We now present the second main result of this chapter, the proof of which is given in the next subsection.

Theorem 8.3. Assume that Assumption 8.2 holds, the controls $\tau_{u}$ and $\tau_{r}$ given by (8.112) solve the control objective with an appropriate choice of $k_{1}$ and $k_{2}$ such that (8.103) and (8.111) hold. Particularly, the transformed path-tracking errors $\left(x_{e}, y_{e}, \bar{\psi}_{e}\right)$ globally asymptotically converge to zero. As a result, the actual position path-tracking errors $\left(x-x_{d}, y-y_{d}\right)$ and orientation path-tracking error $\left(\psi-\psi_{d}\right)$ globally asymptotically converge to zero and to a ball with a radius of smaller than $0.5 \pi$, respectively. Furthermore, the desired forward speed of the ship on the path can be adjusted by adjusting $u_{d 0}(t)$ and the total linear velocity of the ship is tangential to the path.

### 8.2.5 Stability Analysis

To prove Theorem 8.3, we first show that the closed loop system (8.113) is forward complete (i.e. no finite escape in the closed loop system), and there exist an arbitrarily small positive constant $\phi_{0}$, and a class- $K$ function $\Xi_{0}$ of $\left\|\chi_{0}\left(t_{0}\right)\right\|$, with

$$
\begin{aligned}
\chi_{0}\left(t_{0}\right):= & {\left[x_{e}\left(t_{0}\right), y_{e}\left(t_{0}\right), \psi_{e}\left(t_{0}\right), x_{d}\left(t_{0}\right), y_{d}\left(t_{0}\right), \sigma_{x e}\left(t_{0}\right),\right.} \\
& \left.\sigma_{y e}\left(t_{0}\right), \sigma_{\bar{\psi} e}\left(t_{0}\right), u_{e}\left(t_{0}\right), v_{e}\left(t_{0}\right), r_{e}\left(t_{0}\right), v_{d}\left(t_{0}\right)\right]^{T}
\end{aligned}
$$

such that

$$
\left\|\left(x_{e}(t), y_{e}(t), x_{d}(t), x_{d}(t)\right)\right\| \leq \Xi_{0} e^{\phi_{0}\left(t-t_{0}\right)}, \forall t \geq t_{0}
$$

The reason for doing this is that the term, $\tilde{r}$, contains $(x, y)$, see (8.85). We then consider the $\left(\sigma_{\bar{\psi} e}, \bar{\psi}_{e}, u_{e}, v_{e}, r_{e}\right)$-subsystem and prove that $v_{d}$ is bounded by some constant. Finally we consider the ( $\sigma_{x e}, \sigma_{y e}, x_{e}, y_{e}$ )-subsystem.

To prove that the closed loop system (8.113) is forward complete, we consider the Lyapunov function

$$
\begin{equation*}
W=\frac{1}{2} \log \left(1+W_{0}\right)+\frac{1}{2} K_{2} \tilde{\boldsymbol{X}}^{\boldsymbol{T}} \tilde{\boldsymbol{X}} \tag{8.114}
\end{equation*}
$$

where

$$
W_{0}=x_{e}^{2}+y_{e}^{2}+x_{d}^{2}+y_{d}^{2}+\sigma_{x e}^{2}+\sigma_{y e}^{2}+\sigma_{\bar{\psi} e}^{2}+K_{1}\left(\bar{\psi}_{e}^{2}+u_{e}^{2}+v_{e}^{2}+r_{e}^{2}+v_{d}^{2}\right),
$$

$K_{1}$ and $K_{2}$ are positive constants to be picked later. Differentiating both sides of (8.114) along the solutions of (8.113) and (8.83), and using properties of the projection algorithm, after some calculations we obtain:

$$
\begin{equation*}
\dot{W} \leq \frac{1}{2} \phi_{0}+\frac{\phi_{1}}{1+W_{0}}+\phi_{2}\|\tilde{\boldsymbol{X}}\|^{2} \tag{8.115}
\end{equation*}
$$

with constants $\phi_{i}, i=0,1,2$ being defined as

$$
\begin{align*}
\phi_{0}= & 2\left(5 \varepsilon_{01}+4 \varepsilon_{02}+16 \varepsilon_{03}+\frac{1}{2 K_{1} \varepsilon_{01}}-\left(c_{12}+\frac{d_{11}}{m_{11}}+\frac{d_{22}}{m_{22}}+\frac{d_{33}}{m_{33}}\right)\right), \\
\phi_{1}= & \frac{\bar{\sigma}_{x e}^{2}+\bar{\sigma}_{y e}^{2}}{4 \varepsilon_{01}}+\frac{9}{\varepsilon_{01}}\left(\bar{u}_{d 0}^{2}+\bar{\sigma}_{x e}^{2}+3 \bar{\sigma}_{y e}^{2}+3 k_{2}^{2}\right)+\left(\bar{\sigma}_{x e}^{2}+\bar{\sigma}_{y e}^{2}+\bar{\sigma}_{\bar{\psi} e}^{2}\right) \times \\
& \max \left(\overline{\dot{\sigma}}_{x e}, \bar{\sigma}_{y e}, \overline{\dot{\sigma}}_{\bar{\psi} e}\right)+K_{1}\left(\frac{0.75}{\varepsilon_{03}}+\frac{\varepsilon_{04}}{\underline{b}_{2}}+\frac{1}{4 \varepsilon_{05}}\right)\left(\frac{\left|d_{23}\right|}{\underline{b}_{1} m_{22}}+\right. \\
& \left.\frac{m_{11}\left(k_{1}+\bar{u}_{d 0}\right)}{\underline{b}_{1} m_{22}}\left(k_{3}+\bar{\sigma}_{\bar{\psi} e}+\frac{\left(2+k_{2}+\overline{\dot{\sigma}}_{x e}+\overline{\dot{\sigma}}_{y e}+\bar{\sigma}_{x e}+\bar{\sigma}_{y e}\right)}{\frac{u}{d 0}}\right)\right) \\
\phi_{2}= & -K_{2} \lambda_{\min }\left(\overline{\boldsymbol{M}}^{-1} \overline{\boldsymbol{D}}\right)+\frac{1}{2 \varepsilon_{01}}+K_{1}\left(\frac{2.25}{\varepsilon_{02}}+\frac{1}{\varepsilon_{03}}\left(0.5 k_{1}^{2}+0.25\left(A_{1}^{2}+\right.\right.\right. \\
& \left.A_{2}^{2}\right)+8.25 \frac{m_{22}^{2}}{m_{33}^{2}}+\left(2 k_{1}^{2}+2 k_{2}^{2}+3\left(1+k_{1}+\bar{u}_{d 0}\right)^{2}+2 A_{1}^{2}+\right. \\
& \left.2\left(A_{2}^{2}+A_{3}^{2}\right)\right), \tag{8.116}
\end{align*}
$$

where $c_{12}=c_{1}+c_{2}, \varepsilon_{0 i}, i=1, \ldots, 5$ are positive constants. The constants $A_{i}, i=$ $1,2,3$ denote upper bounds of $\left|\partial \alpha_{r} / \partial x_{e}\right|,\left|\partial \alpha_{r} / \partial y_{e}\right|,\left|\partial \alpha_{r} / \partial \psi_{e}\right|$, respectively. These constants can be calculated by taking corresponding partial derivatives of $\alpha_{r}$, see (8.104). Inequality (8.115) implies that the solution of the closed loop system (8.113) exists.

Now we pick the constants $\varepsilon_{0 i}, i=1,2,3$ and $K_{1}$ such that $\phi_{0}$ is strictly less than $\lambda_{\text {min }}\left(\overline{\boldsymbol{M}}^{\mathbf{- 1}} \overline{\boldsymbol{D}}\right)$. Then picking the constant $K_{2}$ such that $\phi_{2} \leq 0$, we have

$$
\begin{equation*}
\dot{W} \leq \frac{1}{2} \phi_{0}+\frac{\phi_{1}}{1+W_{0}} . \tag{8.117}
\end{equation*}
$$

On the other hand, it is noted from (8.114) that

$$
\begin{equation*}
\dot{W} \geq \frac{\dot{W}_{0}}{2\left(1+W_{0}\right)}-\lambda_{\max }\left(\bar{M}^{-1} \bar{D}\right)\|\tilde{\boldsymbol{X}}\|^{2} \geq \frac{\dot{W}_{0}}{2\left(1+W_{0}\right)} \tag{8.118}
\end{equation*}
$$

where $\lambda_{\text {max }}\left(\overline{\boldsymbol{M}}^{\mathbf{- 1}} \overline{\boldsymbol{D}}\right)>0$ is the maximum eigenvalue of $\overline{\boldsymbol{M}}^{\mathbf{1}} \overline{\boldsymbol{D}}$. From (8.117) and (8.118), we have

$$
\begin{equation*}
\dot{W}_{0} \leq \phi_{0} W_{0}+2 \phi_{1}+\phi_{0} . \tag{8.119}
\end{equation*}
$$

From (8.119) and the expression for $W_{0}$, we have

$$
\left\|\left(x_{e}(t), y_{e}(t), x_{d}(t), y_{d}(t)\right)\right\| \leq \Xi_{0} e^{\phi_{0}\left(t-t_{0}\right)}, \forall t \geq t_{0}
$$

To show that $\tilde{\boldsymbol{X}}(t)$ globally exponentially converges to zero, we take the Lyapunov function

$$
V_{0}=\frac{1}{2} \tilde{X}^{\boldsymbol{T}} \tilde{X}
$$

whose derivative along the solutions of (8.83) satisfies

$$
\begin{align*}
& \dot{V}_{0} \leq-\lambda_{\min }\left(\overline{\boldsymbol{M}}^{-\mathbf{1}} \overline{\boldsymbol{D}}\right)\|\tilde{X}\|^{2} \\
\Rightarrow \quad & \|\tilde{\boldsymbol{X}}(t)\| \leq\left\|\tilde{\boldsymbol{X}}\left(t_{0}\right)\right\| e^{-\lambda_{\min }\left(\overline{\boldsymbol{M}}^{-1} \overline{\boldsymbol{D}}\right)\left(t-t_{0}\right)}, \quad \forall t \geq t_{0} \geq 0 \tag{8.120}
\end{align*}
$$

where Lemma 2.2 has been used. To investigate stability of the $\left(\sigma_{\bar{\psi} e}, \bar{\psi}_{e}, u_{e}, v_{e}, r_{e}\right)$ subsystem, we take the Lyapunov function

$$
\begin{equation*}
V_{1}=\sqrt{1+\bar{\psi}_{e}^{2}}-1+\frac{1}{2}\left(u_{e}^{2}+v_{e}^{2}+r_{e}^{2}\right)+\frac{1}{2 \gamma_{3}} \sigma_{\bar{\psi} e}^{2} \tag{8.121}
\end{equation*}
$$

whose derivative along the solutions of the third, sixth, seventh, eighth, and ninth equations of (8.113) satisfied

$$
\begin{align*}
\dot{V}_{1} \leq & -\left(k_{3}-\varepsilon_{06}\right) \frac{\bar{\psi}_{e}^{2}}{1+\bar{\psi}_{e}^{2}}-\left(c_{1}+\frac{d_{11}}{m_{11}}-\varepsilon_{06}\right) u_{e}^{2}- \\
& \left(\frac{d_{22}}{m_{22}}-\varepsilon_{06}\right) v_{e}^{2}-\left(c_{2}+\frac{d_{33}}{m_{33}}-\varepsilon_{06}\right) r_{e}^{2}+\Xi_{1} e^{-\delta_{1}\left(t-t_{0}\right)} \\
\leq & \Xi_{1} e^{-\delta_{1}\left(t-t_{0}\right)}, \tag{8.122}
\end{align*}
$$

with $\Xi_{1}=A_{4}\left\|\tilde{\boldsymbol{X}}\left(t_{0}\right)\right\|+A_{5}\left\|\tilde{\boldsymbol{X}}\left(t_{0}\right)\right\|^{2}$, where

$$
\begin{align*}
A_{4}= & \max \left(\underline{b}_{1}, k_{2} \underline{u}_{d 0}^{-1}\right), \\
A_{5}= & 0.25 \varepsilon_{06}^{-1} \max \left(k_{1}^{2}+\underline{b}_{1}^{-1} \underline{u}_{d 0}^{-1}\left(k_{1}+k_{2}\right)^{2}\left(\bar{u}_{d}+\bar{v}_{d}\right)^{2}+\right. \\
& \left(m_{11} m_{33}^{-1}\right)^{2}\left(1+\bar{v}_{d}\right), 2 k_{1}^{2}+\left(m_{11} m_{33}^{-1}\right)^{2}\left(1+\bar{\alpha}_{u}\right)+ \\
& \underline{b}_{1}^{-1} \underline{u}_{d 0}^{-1}\left(k_{1}+k_{2}\right)^{2}\left(\bar{u}_{d}+\bar{v}_{d}\right)^{2}, 2 k_{1}^{2}+\left(m_{11} m_{22}^{-1}\right)^{2} \times \\
& \left.\left(1+\bar{\alpha}_{u}\right)+\underline{b}_{1}^{-1} \underline{u}_{d 0}^{-1}\left(k_{1}+k_{2}\right)^{2}\left(\bar{u}_{d}+\bar{v}_{d}\right)^{2}+k_{3}^{2}\right), \tag{8.123}
\end{align*}
$$

and $\delta_{1}=\lambda_{\text {min }}\left(\overline{\boldsymbol{M}}^{-\mathbf{1}} \overline{\boldsymbol{D}}\right)-\phi_{0}, \varepsilon_{06}$ is a positive constant that is strictly smaller than $\min \left(k_{3}, c_{1}+d_{11} / m_{11}, d_{22} / m_{22}, c_{2}+d_{33} / m_{33}\right)$. Since we have already proved that $\phi_{0}$ is strictly less than $\lambda_{\text {min }}\left(\overline{\boldsymbol{M}}^{-\mathbf{1}} \overline{\boldsymbol{D}}\right), \delta_{1}$ is a strictly positive constant. The second inequality of (8.122) implies that $V_{1}(t)$, i.e., $\left(\sigma_{\bar{\psi} e}(t), \bar{\psi}_{e}(t), u_{e}(t), v_{e}(t), r_{e}(t)\right)$, is bounded. By integrating both sides of the first inequality of (8.122) and applying Barbalat's lemma, we have $\lim _{t \rightarrow \infty}\left(\bar{\psi}_{e}(t), u_{e}(t), v_{e}(t), r_{e}(t)\right)=0$.

To show that $v_{d}$ is bounded, we take the Lyapunov function

$$
V_{2}=\frac{1}{2} v_{d}^{2}
$$

whose derivative along the solutions of the last equation of (8.113) yields

$$
\begin{equation*}
\dot{V}_{2}=b_{2}^{-1}\left(-d_{n 2} m_{22}^{-1}\left|v_{d}\right| v_{d}^{2}-B_{1} v_{d}^{2}+B_{2}\left|v_{d}\right|+B_{3}\right), \tag{8.124}
\end{equation*}
$$

where $B_{i}, 0 \leq i \leq 3$ are given in (8.142). Since $d_{n 2} / m_{22}$ is a positive constant, an application of Theorem 4.18 in [6] shows that $v_{d}$ is bounded by some constant. It is further noted that if nonlinear damping terms are ignored, i.e., $d_{n i}=0, i=1,2,3$, then the path and desired surge velocity must satisfy an additional condition such that $B_{1}$ is strictly positive otherwise the filtered sway velocity dynamics will be unstable. This will result in an unstable closed loop system.

To investigate stability of the ( $\sigma_{x e}, \sigma_{y e}, x_{e}, y_{e}$ )-dynamics, we take the Lyapunov function

$$
V_{3}=\sqrt{1+x_{e}^{2}+y_{e}^{2}}-1+\frac{1}{2}\left(\frac{\sigma_{x e}^{2}}{\gamma_{1}}+\frac{\sigma_{y e}^{2}}{\gamma_{2}}\right)
$$

whose derivative along the solutions of the first four equations of (8.113) satisfies

$$
\begin{align*}
\dot{V}_{3} \leq & -\frac{\left(k_{1}-3 \varepsilon_{07}\right) x_{e}^{2}+\left(k_{2}-3 \varepsilon_{07}\right) y_{e}^{2}}{\Delta^{2}}+\frac{u_{e}^{2}+v_{e}^{2}+8 u_{d}^{2} \bar{\psi}_{e}^{2}}{4 \varepsilon_{07}}+ \\
& \frac{\Xi_{2}}{\varepsilon_{07}} e^{-\lambda_{\min }\left(\overline{\mathbf{M}}^{-1} \overline{\mathbf{D}}\right)\left(t-t_{0}\right)} \\
\leq & \frac{u_{e}^{2}+v_{e}^{2}+8 u_{d}^{2} \bar{\psi}_{e}^{2}}{4 \varepsilon_{07}}+\frac{\Xi_{2}}{\varepsilon_{07}} e^{-\lambda_{\min }\left(\overline{\mathbf{M}}^{-1} \overline{\mathbf{D}}\right)\left(t-t_{0}\right)}, \tag{8.125}
\end{align*}
$$

where $\varepsilon_{07}$ is a positive constant, $\Xi_{2}$ is a class- $K$ function of $\left\|\chi_{0}\left(t_{0}\right)\right\|$. We have picked this constant such that $\varepsilon_{07}<\min \left(k_{1}, k_{2}\right) / 3$. Since we have proved that $\left(\sigma_{\bar{\psi} e}(t), \bar{\psi}_{e}(t), u_{e}(t), v_{e}(t), v_{d}(t)\right)$ is bounded and $u_{d}$ is given by (8.96), the second inequality of (8.125) implies that $\left(\sigma_{x e}(t), \sigma_{y e}(t), x_{e}(t), y_{e}(t)\right)$ is bounded. Hence integrating both sides of the first inequality of (8.125) and applying Barbalat's lemma yield $\lim _{t \rightarrow \infty}\left(x_{e}(t), y_{e}(t)\right)=0$.

It now follows from (8.97) and the second equation of (8.92) that $\left|\psi_{e}(t)\right|$ converges to $\left|\arctan \left(\left(\varphi_{2}+\sigma_{y e}+v_{d}\right) / \Upsilon_{u_{d 0}}\right)\right|<0.5 \pi$, since we have already proven that $\lim _{t \rightarrow \infty}\left(x_{e}(t), y_{e}(t)\right)=0, v_{d}(t)$ is bounded and $u_{d 0}(t)$ is larger than some positive constant by assumption, see Assumption 8.2. Proof that the ship's total linear velocity is tangential to the path follows readily, since $\psi_{e}(t)$ converges to
$\arctan \left(\left(\varphi_{2}+\sigma_{y e}+v_{d}\right) / \Upsilon_{u_{d 0}}\right)$, and the proven fact that $\lim _{t \rightarrow \infty}\left(v(t)-v_{d}(t)\right)=0$. Finally, we note that all constants $\varepsilon_{0 i}, 1 \leq i \leq 7, A_{j}, j=1,2,3, K_{1}$, and $K_{2}$ are only used in the proof. They are not needed in controller implementation.

### 8.2.6 Discussion

### 8.2.6.1 Adding Integrator Approach Versus Adaptive Approach

We have used the projection integral actions instead of an adaptive approach to compensate for the disturbances. The trade-off is that the integral action approach does not give deep insight into the ship dynamics with disturbances. However, it results in a simple controller that is suitable for practical implementation. An adaptive approach would be extremely complicated (even when nonlinear damping terms are ignored) if the controller design follows the methodology proposed in this chapter. In an adaptive approach, the disturbances, $\boldsymbol{\tau}_{\boldsymbol{E}}$, can be considered directly in the ship dynamics as:

$$
\begin{equation*}
\bar{M} \dot{v}=-\bar{C}(v) v-\bar{D} v+\tau+\tau_{E} \tag{8.126}
\end{equation*}
$$

where the nonlinear damping terms are ignored. Using the same coordinate transformation (8.73) with $\boldsymbol{P}(\boldsymbol{\eta})$ given in (8.80) gives

$$
\begin{equation*}
\dot{X}=-P(\eta) \bar{M}^{-1} \bar{D} P^{-1}(\eta) X+P(\eta) \bar{M}^{-1}\left(\tau+\tau_{E}\right) \tag{8.127}
\end{equation*}
$$

Because of the unknown $\boldsymbol{\tau}_{\boldsymbol{E}}$, we would design a dynamic observer for $\boldsymbol{X}$ by defining

$$
\begin{equation*}
\tilde{X}=X-\hat{X}-\boldsymbol{W} \boldsymbol{\tau}_{\boldsymbol{E}}, \tag{8.128}
\end{equation*}
$$

where $\boldsymbol{W} \in \mathbb{R}^{3 \times 3}$ is a matrix to be determined and $\hat{\boldsymbol{X}}$ is an estimate of $\boldsymbol{X}$. Differentiating both sides of (8.128) and choosing

$$
\begin{align*}
& \dot{\hat{X}}=-P(\eta) \bar{M}^{-1} \bar{D} P^{-1}(\eta) \hat{X}+P(\eta) \bar{M}^{-1} \tau,  \tag{8.129}\\
& \dot{W}=-P(\eta) \bar{M}^{-1} \bar{D} P^{-1}(\eta) W+P(\eta) \bar{M}^{-1}
\end{align*}
$$

results in

$$
\begin{equation*}
\dot{\tilde{X}}=-Q(\eta) \bar{M}^{-1} \bar{D} Q^{-1}(\eta) \tilde{X} \tag{8.130}
\end{equation*}
$$

From (8.129) and the first equation of (8.71), a controller can be designed but this will be very complicated since some elements of $\boldsymbol{P}(\boldsymbol{\eta})$ depend on $(x, y)$ and $\boldsymbol{W}$ must be generated by the second equation of (8.129).

### 8.2.6.2 Dealing with Parameter Uncertainties

In the control design, it was assumed that all the system parameters are known. The ship's mass and added mass and linear damping coefficients are determined quite
accurately by using semi-empirical methods or hydrodynamic programs (such as MARINTEK). However, it is difficult to obtain nonlinear damping coefficients accurately. The purpose of this section is to discuss how the proposed observer and controller can be modified so that the inaccuracies of the nonlinear damping coefficients can be taken care of. Dealing with inaccuracies of the linear damping coefficients can be carried out similarly if these coefficients are not known accurately. The idea is to replace the real coefficients $d_{n i}, i=1,2,3$ in the nonlinear damping matrix $\overline{\boldsymbol{D}}_{\boldsymbol{n}}(\boldsymbol{v})$, see (8.72), by their estimates. These estimates will be used in the observer (8.82) instead of the real nonlinear damping coefficients. The estimates of the nonlinear damping coefficients are updated in such a way that the modified observer guarantees that the observer error vector $\tilde{\boldsymbol{X}}(t)$ globally asymptotically converges to zero. Then the proposed control design can be simply modified by replacing all the nonlinear damping coefficients $d_{n i}, i=1,2,3$ by their estimates to take care of inaccuracies in the nonlinear damping coefficients.

## Observer Design Modifications

Letting $\hat{d}_{n i}, i=1,2,3$ be estimates of $d_{n i}$ and defining the estimate errors as $\tilde{d}_{n i}=$ $d_{n i}-\hat{d}_{n i}$, we can write (8.81) as

$$
\begin{align*}
\dot{X}= & -P(\eta) \bar{M}^{-1}\left(\bar{D}+\hat{\bar{D}}_{n}\left(P^{-1}(\eta) X\right)\right) P^{-1}(\eta) X+P(\eta) \bar{M}^{-1} \tau- \\
& \left.P(\eta) \bar{M}^{-1} \tilde{\bar{D}}_{n}\left(P^{-1}(\eta) X\right)\right) P^{-1}(\eta) X \tag{8.131}
\end{align*}
$$

where

$$
\hat{\overline{\boldsymbol{D}}}_{\boldsymbol{n}}(\boldsymbol{v})=\left[\begin{array}{ccc}
\hat{d}_{n 1}|u| & 0 & 0 \\
0 & \hat{d}_{n 2}|v| & 0 \\
0 & 0 & \hat{d}_{n 3}|r|
\end{array}\right], \tilde{\overline{\boldsymbol{D}}}_{\boldsymbol{n}}(\boldsymbol{v})=\left[\begin{array}{ccc}
\tilde{d}_{n 1}|u| & 0 & 0 \\
0 & \tilde{d}_{n 2}|v| & 0 \\
0 & 0 & \tilde{d}_{n 3}|r|
\end{array}\right]
$$

with (8.73) having been used for a short notation. From (8.131), we design an observer to estimate $\boldsymbol{X}$ as

$$
\begin{equation*}
\dot{\hat{X}}=-P(\eta) \bar{M}^{-1}\left(\bar{D}+\hat{\bar{D}}_{n}\left(P^{-1}(\eta) \hat{X}\right)\right) P^{-1}(\eta) \hat{X}+P(\eta) \bar{M}^{-1} \tau \tag{8.132}
\end{equation*}
$$

From (8.131) and (8.132), we have the observer error dynamics

$$
\begin{align*}
\dot{\tilde{X}}= & -P(\eta) \bar{M}^{-1} \bar{D} P^{-1}(\eta) \tilde{X}-P(\eta) \bar{M}^{-1} \times \\
& {\left.\left[\hat{\bar{D}}_{n}\left(P^{-1}(\eta) X\right) P^{-1}(\eta) X-\hat{\bar{D}}_{n}\left(P^{-1}(\eta) \hat{X}\right)\right) P^{-1}(\eta) \hat{X}\right]- } \\
& \left.P(\eta) \bar{M}^{-1} \tilde{\bar{D}}_{n}\left(P^{-1}(\eta) X\right)\right) P^{-1}(\eta) X \\
= & -P(\eta) \bar{M}^{-1} \bar{D}^{-1}(\eta) \tilde{X}-Q(\eta) \bar{M}^{-1} \times \\
& {\left.\left[\bar{D}_{n}\left(P^{-1}(\eta) X\right) P^{-1}(\eta) X-\bar{D}_{n}\left(P^{-1}(\eta) \hat{X}\right)\right) P^{-1}(\eta) \hat{X}\right]-\Phi \tilde{\Theta}, } \tag{8.133}
\end{align*}
$$

where $\tilde{\boldsymbol{X}}=\boldsymbol{X}-\hat{\boldsymbol{X}}$ and we have defined

$$
\boldsymbol{\Phi}=\boldsymbol{P}(\boldsymbol{\eta}) \overline{\boldsymbol{M}}^{-\mathbf{1}}\left[\begin{array}{ccc}
|\hat{u}| \hat{u} & 0 & 0 \\
0 & |\hat{v}| \hat{v} & 0 \\
0 & 0 & |\hat{r}| \hat{r}
\end{array}\right], \tilde{\boldsymbol{\Theta}}=\left[\begin{array}{c}
\tilde{d}_{n 1} \\
\tilde{d}_{n 2} \\
\tilde{d}_{n 3}
\end{array}\right] .
$$

It is noted that we have used (8.84). To determine an update law for $\hat{d}_{n i}, i=1,2,3$, we take the following Lyapunov function

$$
\begin{equation*}
L_{0}=\frac{1}{2} \tilde{\boldsymbol{X}}^{\boldsymbol{T}} \tilde{\boldsymbol{X}}+\frac{1}{2} \tilde{\boldsymbol{\Theta}}^{\boldsymbol{T}} \boldsymbol{\Gamma}^{-\mathbf{1}} \tilde{\boldsymbol{\Theta}} \tag{8.134}
\end{equation*}
$$

where $\boldsymbol{\Gamma}$ is a positive definite matrix. Using Lemma 2.2, the derivative of $L_{0}$ along the solutions of (8.133) satisfies

$$
\begin{align*}
\dot{L}_{0}= & -\tilde{X}^{\boldsymbol{T}} \boldsymbol{P}(\eta) \bar{M}^{-1} \bar{D} P^{-1}(\eta) \tilde{X}-\tilde{X}^{\boldsymbol{T}} \boldsymbol{P}(\eta) \bar{M}^{-1} \times \\
& {\left.\left[\bar{D}_{n}\left(P^{-1}(\eta) X\right) P^{-1}(\eta) X-\bar{D}_{n}\left(P^{-1}(\eta) \hat{X}\right)\right) P^{-1}(\eta) \hat{X}\right]- } \\
& \tilde{X}^{\boldsymbol{T}} \boldsymbol{\Phi} \tilde{\Theta}, \\
\leq & -\lambda_{\min }\left(\bar{M}^{-1} \bar{D}\right)\|\tilde{X}\|^{2}-\tilde{X}^{\boldsymbol{T}} \boldsymbol{\Phi} \tilde{\Theta}-\dot{\hat{\Theta}}^{\boldsymbol{T}} \Gamma^{-1} \tilde{\Theta}, \tag{8.135}
\end{align*}
$$

where

$$
\hat{\boldsymbol{\Theta}}=\left[\begin{array}{lll}
\hat{d}_{n 1} & \hat{d}_{n 2} & \hat{d}_{n 3}
\end{array}\right]^{T},
$$

and $\lambda_{\text {min }}\left(\overline{\boldsymbol{M}}^{-\mathbf{1}} \overline{\boldsymbol{D}}\right)>0$ is the minimum eigenvalue of $\overline{\boldsymbol{M}}^{-\mathbf{1}} \overline{\boldsymbol{D}}$. From (8.135), choosing an update law for $\hat{\boldsymbol{\Theta}}$ as

$$
\begin{equation*}
\dot{\hat{\Theta}}=-\Gamma \Phi^{\boldsymbol{T}}(\cdot) \tilde{X} \tag{8.136}
\end{equation*}
$$

results in

$$
\begin{equation*}
\dot{L}_{0} \leq-\lambda_{\min }\left(\overline{\boldsymbol{M}}^{-1} \overline{\boldsymbol{D}}\right)\|\tilde{\boldsymbol{X}}\|^{2} \tag{8.137}
\end{equation*}
$$

which means that $\lim _{t \rightarrow \infty 0} \tilde{X}(t)=0$. Indeed, we can use the projection algorithm instead of (8.136). Now it is important to note that the update law (8.136) cannot be used since $\tilde{\boldsymbol{X}}$ contains the unknown vector $\boldsymbol{X}$. We need to get around this problem. Substituting $\tilde{\boldsymbol{X}}=\boldsymbol{X}-\hat{\boldsymbol{X}}$ into the right-hand side of (8.136) and integrating both sides gives

$$
\begin{equation*}
\hat{\boldsymbol{\Theta}}(t)=\hat{\boldsymbol{\Theta}}\left(t_{0}\right)-\boldsymbol{\Gamma} \int_{t_{0}}^{t} \boldsymbol{\Phi}^{\boldsymbol{T}}(\boldsymbol{X}(\tau)-\hat{\boldsymbol{X}}(\tau)) \mathrm{d} \tau \tag{8.138}
\end{equation*}
$$

On the other hand, from the first equation of (8.71) and (8.73), we have

$$
\begin{equation*}
\boldsymbol{X}=(\boldsymbol{J}(\psi) \boldsymbol{P}(\boldsymbol{\eta}))^{-1} \dot{\boldsymbol{\eta}} \tag{8.139}
\end{equation*}
$$

Substituting (8.139) into (8.138) gives

$$
\begin{equation*}
\hat{\boldsymbol{\Theta}}(t)=\hat{\boldsymbol{\Theta}}\left(t_{0}\right)+\boldsymbol{\Gamma} \int_{t_{0}}^{\boldsymbol{t}} \boldsymbol{\Phi}^{\boldsymbol{T}} \hat{\boldsymbol{X}}(\tau) \mathrm{d} \tau-\boldsymbol{\Gamma} \int_{\boldsymbol{\eta}\left(t_{0}\right)}^{\boldsymbol{\eta}(t)} \boldsymbol{\Phi}^{\boldsymbol{T}}(\boldsymbol{J}(\sigma) \boldsymbol{P}(\sigma))^{-1} \mathrm{~d} \sigma \tag{8.140}
\end{equation*}
$$

where we have abused a notation of $\boldsymbol{J}(\psi)$ as $\boldsymbol{J}(\boldsymbol{\eta})$. It is seen that (8.140) is a linear Volterra equation, which can be numerically solved by using a number of methods available in the literature, see [119]. Noticing that the right-hand side of (8.140) contains only known terms, we can use (8.140) to calculate $\hat{\boldsymbol{\Theta}}(t)$ instead of (8.136). We summarize this section as follows: The observer (8.82) is replaced by the adaptive observer consisting of (8.132) and (8.140) for the case where the nonlinear damping coefficients are not known accurately.

## Control Design Modifications

The modified observer, see (8.132) and (8.140), guarantees that $\lim _{t \rightarrow \infty 0} \tilde{X}(t)=0$. The control design is the same as in Section 8.2.4 but the nonlinear damping coefficients $d_{n i}, i=1,2,3$ are replaced by their estimates $\hat{d}_{n i}, i=1,2,3$. The result with the modified observer and controller as discussed above is almost the same as the one stated in Theorem 8.3. The only difference is that in the case where the nonlinear damping coefficients are known accurately, the observer error $\tilde{\boldsymbol{X}}(t)$ globally exponentially converges to zero but only globally asymptotically converges to zero in the case where the nonlinear damping coefficients are not known accurately due to the adaptation. Proof of the result under the modified observer and controller follows the same lines as that of Theorem 8.3, and is therefore excluded here.

### 8.2.6.3 Dealing with Actuator Saturation

In Section 8.2.4, the controller was designed without imposing any limitations on the magnitude and rate of the ship actuators. Actuator saturation restricts the reference path, deteriorates the overall path-tracking performance, and can even destabilize the closed loop system. These problems are difficult to deal with but it is important that they are taken care of. In many cases, global stabilization and/or tracking results cannot be achieved when actuators are subject to saturation. For example, consider the scalar system $\dot{x}=x^{2}+u$ with $x$ the state and $u$ the control. For this system, no bounded control $u$ can globally exponentially/asymptotically stabilize the system at the origin. In fact, any bounded control $u$ can result in a finite escape time. On the other hand, only restrictive tracking can be achieved for many systems such as a simple system consisting of an integrator chain, see [120]. In this section, we will discuss how to use the controller proposed in Section 8.2.4 when the actuator saturation presents. We will focus on the magnitude saturation. A discussion on the rate saturation can be carried out similarly. The idea is to set $\left|\tau_{u}\right| \leq \tau_{u}^{\max }$ and $\left|\tau_{r}\right| \leq$ $\tau_{r}^{\max }$ where $\tau_{u}^{\max }$ and $\tau_{r}^{\max }$ are the maximum values that the actuators can supply, then to find restrictions on the reference path, control gains, and initial conditions such that $\left|\tau_{u}\right| \leq \tau_{u}^{\max }$ and $\left|\tau_{r}\right| \leq \tau_{r}^{\max }$ hold. From the control design in Subsection 8.2.4, we calculate the upper bounds of $\alpha_{u}, \alpha_{\psi e}$, and $\alpha_{r}$ as

$$
\begin{align*}
\bar{\alpha}_{u}= & k_{1}+\bar{\sigma}_{x e}+\sqrt{\bar{u}_{d 0}^{2}+\bar{\sigma}_{x e}^{2}}, \bar{\alpha}_{\psi e}=\arctan \left(\frac{k_{2}+\bar{\sigma}_{y e}+\bar{v}_{d}}{\sqrt{\bar{u}_{d 0}^{2}+\bar{\sigma}_{x e}^{2}}}\right),  \tag{8.141}\\
\bar{\alpha}_{r}= & \underline{b}_{1}^{-1}\left(k_{3}+\bar{\sigma}_{\bar{\psi} e}+\bar{r}_{d}+\underline{u}_{d 0}^{-1}\left(\left(k_{2}+1\right)\left(\bar{v}_{d}+\bar{u}_{d}\right)+\right.\right. \\
& \left.\left.\bar{\alpha}_{u}+\bar{u}_{d}+\overline{\dot{v}}_{d}+\overline{\dot{\sigma}}_{y e}\right)+\underline{u}_{d 0}^{-2}\left(\overline{\dot{u}}_{d 0} \bar{u}_{d 0}+\overline{\dot{\sigma}}_{x e} \bar{\sigma}_{x e}\right)\right),
\end{align*}
$$

where

$$
\begin{gather*}
\bar{u}_{d}=\sqrt{\bar{u}_{d 0}^{2}+\bar{\sigma}_{x e}^{2}+\left(k_{2}+\bar{\sigma}_{y e}+\bar{v}_{d}\right)^{2}}, \\
\bar{v}_{d}=\max \left(\sqrt{\frac{B_{3}}{B_{1}}}, \sqrt{\frac{B_{2} m_{22}}{d_{n 2}}}\right) \\
\overline{\dot{v}}_{d}=\frac{\left(m_{11} \bar{\alpha}_{u}+\bar{d}_{23}\right) \bar{\alpha}_{r}+d_{22} \bar{v}_{d}+d_{n 2} \bar{v}_{d}^{2}+k_{2} \underline{u}_{d 0}^{-1}}{m_{22}}, \\
\bar{r}_{d}=\frac{\left|x_{d}^{\prime}(s) y_{d}^{\prime \prime}(s)-x_{d}^{\prime \prime}(s) y_{d}^{\prime}(s)\right|}{\sqrt{\left(x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)\right)^{3}}} \bar{u}_{d}, \\
B_{1}=\frac{d_{22}}{m_{22}}-2 B_{0}, \\
B_{2}=\frac{m_{11}\left(k_{1}+\bar{u}_{d 0}\right)+\left|d_{23}\right|}{\underline{b}_{1} m_{22}}\left(k_{3}+\bar{\sigma}_{\bar{\psi} e}+\frac{k_{2}\left(2+k_{1}+k_{2}\right)+\overline{\dot{\sigma}}_{y e}}{\underline{u}_{d 0}}+\right. \\
B_{3}=\frac{1}{2} B_{0}\left(\bar{u}_{d 0}^{2}+\bar{\sigma}_{x e}^{2}+3 k_{2}^{2}+3 \bar{\sigma}_{y e}^{2}\right), \\
\bar{u}_{d 0}+\bar{\sigma}_{x e} \bar{\sigma}_{x e} \\
B_{0}=\frac{u_{d 0}^{2}}{\underline{u}_{d 0}}, \frac{k_{2}}{\underline{u}_{d 0}\left(k_{1}+\bar{u}_{d 0}\right)+\bar{d}_{23}} \frac{\underline{b}_{1} m_{22}}{} \frac{x_{d}^{\prime}(s) y_{d}^{\prime \prime}(s)-x_{d}^{\prime \prime}(s) y_{d}^{\prime}(s) \mid}{\sqrt{\left(x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)\right)^{3}}} \tag{8.142}
\end{gather*}
$$

and the notations $\bar{\bullet}$ and $\bullet$ denote the upper and lower bounds of $|\bullet|$, respectively. By substituting $\hat{u}=u_{e}+\alpha_{u}, \hat{v}=v_{e}+v_{d}$, and $\hat{r}=r_{e}+\alpha_{r}$ into (8.112), we can calculate the upper bounds of $\tau_{u}$ and $\tau_{r}$ as

$$
\begin{align*}
& \left|\tau_{u}\right| \leq \vartheta_{1} u_{e}^{2}+\vartheta_{2} v_{e}^{2}+\vartheta_{3} r_{e}^{2}+\vartheta_{0} \\
& \left|\tau_{r}\right| \leq v_{1} u_{e}^{2}+v_{2} v_{e}^{2}+v_{3} r_{e}^{2}+v_{0} \tag{8.143}
\end{align*}
$$

where

$$
\begin{aligned}
& \vartheta_{1}=0.5\left(m_{11}\left(k_{1}+c_{1}\right)+d_{n 1}\right), \\
& \vartheta_{2}=m_{22}+0.5 k_{2}+m_{11} m_{22}^{-1},
\end{aligned}
$$

$$
\begin{align*}
\vartheta_{3}= & m_{22}+0.5\left(k_{1}+k_{2}+m_{11} m_{22}^{-1}\right) \\
\vartheta_{0}= & 0.5 c_{1} m_{11}+m_{22}\left(0.5\left(\bar{\alpha}_{r}^{2}+\bar{v}_{d}^{2}+\bar{v}_{d} \bar{\alpha}_{r}\right)+d_{11} \bar{\alpha}_{u}+1.5 d_{n 1} \bar{\alpha}_{u}^{2}+m_{11} \times\right. \\
& \left(\left(k_{1}+k_{2}\right)\left(1+\bar{u}_{d}+\bar{\alpha}_{r}\right)+\overline{2}_{x e}+\overline{\dot{u}}_{d 0}+k_{2} \underline{u}_{d 0}^{-1}+0.5 m_{11} m_{22}^{-1} \bar{\alpha}_{r}^{2}\right), \\
v_{1}= & \left|m_{11}-m_{22}\right|+\underline{b}_{1}^{-1} \underline{u}_{d 0}^{-1}\left(k_{1}+k_{2}\right)\left(\bar{u}_{d}+\bar{v}_{d}\right)+m_{11} m_{22}^{-1}, \\
v_{2}= & 0.5 \bar{d}_{32}+\left|m_{11}-m_{22}\right|+\underline{b}_{1}^{-1} \underline{u}_{d 0}^{-1} m_{33}\left(k_{1}+k_{2}\right)\left(\bar{u}_{d}+\bar{v}_{d}\right)+ \\
& m_{11} m_{33} m_{22}^{-1}+0.5 m_{33} m_{22}^{-1} \bar{d}_{23}, \\
v_{3}= & 0.5 c_{2} m_{33}+0.5 d_{n 3}+2_{1}^{-1} \underline{u}_{d 0}^{-1} m_{33}\left(k_{1}+k_{2}\right)\left(\bar{u}_{d}+\bar{v}_{d}\right)+0.5 \underline{b}_{1}^{-1} k_{3}, \\
v_{0}= & 0.5 c_{2} m_{33}+1.5 d_{n 3} \bar{\alpha}_{r}^{2}+0.5 \bar{d}_{32}+\bar{d}_{23} \bar{v}_{d}+\left|m_{11}-m_{22}\right|\left(0.5 \bar{v}_{d}^{2}+\right. \\
& \left.0.5 \bar{\alpha}_{u}^{2}+\bar{\alpha}_{u} \bar{v}_{d}\right)+m_{33}\left(\underline { b } _ { 1 } ^ { - 1 } \underline { u } _ { d 0 } ^ { - 1 } ( k _ { 1 } + k _ { 2 } ) ( \overline { u } _ { d } + \overline { v } _ { d } ) \left(2+\bar{\alpha}_{u}+2 \bar{u}_{d}+\right.\right. \\
& \left.2 \bar{\alpha}_{r}+\bar{v}_{d}\right)+\underline{b}_{1}^{-1} \max ^{\prime \prime}\left(\mid x_{d}^{\prime}(s) y_{d}^{\prime \prime}(s)-x_{d}^{\prime \prime \prime}(s) y_{d}^{\prime}(s)\right)\left(x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)\right)- \\
& \left.2\left(x_{d}^{\prime}(s) y_{d}^{\prime \prime}(s)-x_{d}^{\prime \prime}(s) y_{d}^{\prime}(s)\right)\left(x_{d}^{\prime}(s) x_{d}^{\prime \prime}(s)-y_{d}^{\prime}(s) y_{d}^{\prime \prime}(s)\right)\right) \mid \times \\
& \left.\left(x_{d}^{\prime 2}(s)+y_{d}^{\prime 2}(s)\right)^{-2}\right)+\bar{u}_{d}+{ }_{1}^{-1} k_{3}\left(\bar{r}_{d}+0.5+\bar{\alpha}_{r}\right)+\underline{b}_{1}^{-1}\left(\underline { u } _ { d 0 } ^ { - 1 } \left(\overline{\dot{\sigma}}_{x e}^{2}+\right.\right. \\
& \left.\left.\left.\dot{\dot{\sigma}}_{y e}\right)+\overline{\dot{\sigma}}_{\bar{\psi} e}+\overline{\dot{u}}_{d 0}^{2}+\overline{\dot{v}}_{d}\right)+0.5 m_{11} m_{22}^{-1} \bar{\alpha}_{u}^{2}+0.5 m_{22}^{-1} \bar{d}_{23}\right) . \tag{8.144}
\end{align*}
$$

On the other hand, integrating both sides of the second inequality of (8.122) results in

$$
\begin{equation*}
u_{e}^{2}(t)+v_{e}^{2}(t)+r_{e}^{2}(t) \leq \Phi\left(t_{0}\right), \forall t \geq t_{0} \geq 0 \tag{8.145}
\end{equation*}
$$

where

$$
\begin{align*}
\Phi\left(t_{0}\right)= & 2\left(\sqrt{1+\bar{\psi}_{e}^{2}\left(t_{0}\right)}-1+0.5\left(u_{e}^{2}\left(t_{0}\right)+v_{e}^{2}\left(t_{0}\right)+r_{e}^{2}\left(t_{0}\right)\right)+\right. \\
& \left.\frac{1}{2} \gamma_{3} \sigma_{\bar{\psi} e}^{2}\left(t_{0}\right)+\delta_{1}^{-1} A_{4}\left\|\tilde{\boldsymbol{X}}\left(t_{0}\right)\right\|+\delta_{1}^{-1} A_{5}\left\|\tilde{\boldsymbol{X}}\left(t_{0}\right)\right\|^{2}\right) \tag{8.146}
\end{align*}
$$

with $A_{4}, A_{5}$ and $\delta_{1}$ given in (8.123). Now using (8.143) and (8.145), and setting $\left|\tau_{u}\right| \leq \tau_{u}^{\mathrm{Max}}$ and $\left|\tau_{r}\right| \leq \tau_{r}^{\mathrm{Max}}$, we deduce that

$$
\begin{equation*}
\Phi\left(t_{0}\right) \leq \min \left(\frac{\tau_{u}^{\max }-\vartheta_{0}}{\max \left(\vartheta_{1}, \vartheta_{2}, \vartheta_{3}\right)}, \frac{\tau_{r}^{\max }-v_{0}}{\max \left(v_{1}, v_{2}, v_{3}\right)}\right) \tag{8.147}
\end{equation*}
$$

which implies that under the actuator magnitude saturation, restrictions must be placed on the type of reference paths (both desired velocity of the vessel on the path and the path curvature cannot be arbitrarily large), control gains, and initial conditions such that (8.147) holds, i.e., the proposed controller under the actuator magnitude saturation is not only local but also restricts the commanded paths.

### 8.2.7 Experimental Results

This section describes the experimental set-up and tested results performed on a model ship in the Swan River, Western Australia. The model ship has a length of


Figure 8.7 Equipment for experimental tests
1.2 m and a mass of 17.5 kg . The model ship is equipped with one differential global positioning system (DGPS) SF-2050G from NavCom, which can provide an accuracy of 25 cm , to obtain longitude, latitude and altitude, one compass TCM2-50 from PNI to measure yaw angle, two direct current motors Torpedo 850 with two remote speed controllers driving two propellers to provide surge force and yaw moment, three batteries to supply power for motors, DGPS and wireless communication equipment, and two sets of wireless transceivers to transmit and receive signals between the model ship and the host computer, see Figure 8.7. The parameters of the ship model are calculated by VERES:

$$
\begin{aligned}
& m_{11}=25.8, m_{22}=33.8, m_{33}=2.76, m_{23}=m_{32}=6.2, \\
& D_{11}=12+2.5|u|, D_{22}=17+4.5|v|, D_{33}=0.5+0.1|r| \\
& D_{23}=0.2, D_{32}=0.5
\end{aligned}
$$

The global coordinates (longitude, latitude, altitude) are transformed to the "river coordinates" based on a three-point algorithm as follows, see Figure 8.8:

1. The DGPS is first used to measure the global coordinates of two fixed points $M_{1}\left(h_{1}, l_{1}, \mu_{1}\right)$ and $M_{2}\left(h_{2}, l_{2}, \mu_{2}\right)$ with $h_{i}, l_{i}, \mu_{i}, i=1,2$, being the altitude, longitude and latitude of the point $M_{i}\left(h_{i}, l_{i}, \mu_{i}\right)$. These points are chosen along the river bank.
2. We calculate the distance, $\overline{M_{1} M_{2}}$ between these two fixed points by the following formula, see [11]:

$$
\begin{equation*}
\overline{M_{1} M_{2}}=\sqrt{\left(x_{e 1}-x_{e 2}\right)^{2}+\left(y_{e 1}-y_{e 2}\right)^{2}+\left(z_{e 1}-z_{e 2}\right)^{2}} \tag{8.148}
\end{equation*}
$$

where

$$
\begin{align*}
x_{e i} & =\left(N_{i}+h_{i}\right) \cos \left(\mu_{i}\right) \cos \left(l_{i}\right), \\
y_{e i} & =\left(N_{i}+h_{i}\right) \cos \left(\mu_{i}\right) \sin \left(l_{i}\right), \\
z_{e i} & =\left(\frac{r_{p}^{2}}{r_{e}^{2}} N_{i}+h_{i}\right) \sin \left(\mu_{i}\right), \\
N_{i} & =\frac{r_{e}^{2}}{\sqrt{r_{e}^{2} \cos ^{2}\left(\mu_{i}\right)+r_{p}^{2} \sin ^{2}\left(\mu_{i}\right)}} i=1,2, \\
r_{p} & =6356752, \\
r_{e} & =6378137 . \tag{8.149}
\end{align*}
$$

In the experimental set-up, we adjusted the fixed points $M_{1}\left(h_{1}, l_{1}, \mu_{1}\right)$ and $M_{2}\left(h_{2}, l_{2}, \mu_{2}\right)$ such that $\overline{M_{1} M_{2}}=100 \mathrm{~m}$.
3. The river coordinate system, $O_{E} X_{E} Y_{E}$, is formed such that the $O_{E} X_{E}$-axis coincides with the line connecting between $M_{1}\left(h_{1}, l_{1}, \mu_{1}\right)$ and $M_{2}\left(h_{2}, l_{2}, \mu_{2}\right)$, and that the origin O coincides with $M_{1}\left(h_{1}, l_{1}, \mu_{1}\right)$. The OY axis is perpendicular to the OX axis.
4. From this coordinate system, the position coordinates $(x, y)$ of the point, say $M_{j}\left(h_{j}, l_{j}, \underline{\mu_{j}}\right)$, of interest on the model ship are obtained by calculating the distances, $\overline{M_{j} M_{1}}$ and $\overline{M_{j} M_{2}}$, from this point to the points $M_{1}\left(h_{1}, l_{1}, \mu_{1}\right)$ and $M_{2}\left(h_{2}, l_{2}, \mu_{2}\right)$.


Figure 8.8 Experimental set-up

The position coordinates $(x, y)$ and the yaw angle $\psi$ are sent to the host computer by the transceivers. The control signal is calculated in the host computer and is sent
back to the remote speed controllers by a Futaba ${ }^{\circledR}$ transmitter connecting to the host computer via a data acquisition card. A program is written in LabWindows ${ }^{\mathrm{TM}} / \mathrm{CVI}$, a product of National Instruments ${ }^{\mathrm{TM}}$, to implement the control algorithm developed in this chapter. The reference path $\boldsymbol{\Omega}$ is chosen as follows: For the first 90 seconds, $x_{d}=s, y_{d}=60$, and $x_{d}=20 \sin (0.1 s)+60, y_{d}=20 \cos (0.1 s)+40$ for the remainder of testing time. This choice implies that the reference path is a straight line for the first 90 seconds and is a circle centered at ( $60 \mathrm{~m}, 40 \mathrm{~m}$ ) with a radius of 20 m . The desired forward speed is chosen as $u_{d 0}=0.5(\mathrm{~m} / \mathrm{s})$. Based on conditions specified in Theorem 8.3, the design constants and initial conditions of the observer and filters are chosen as $c_{1}=c_{2}=4, k_{1}=k_{2}=k_{3}=1.5, \gamma_{i}=0.1, \xi_{i}=0.1, i=1,2,3$, $v_{d}(0)=0, \sigma_{x e}(0)=0, \sigma_{y e}(0)=0, \sigma_{\bar{\psi} e}(0)=0, \hat{X}(0)=(0,0,0)^{T}, \bar{\sigma}_{x e}=\bar{\sigma}_{y e}=10$, and $\bar{\sigma}_{\bar{\psi} e}=6.5$.


Figure 8.9 A screen shot of the user interface with integral actions

The initial conditions of the model ship are $x=2, y=40, \psi=0.5$, and $u=v=r=0$. Experimental results are given in Figures 8.9 and 8.10. Figure 8.9 is a screen shot of the user interface of the developed program. The user interface displays general information of the ship model such as global coordinates, forward speed, and heading angle. The history of the ship position and the reference path are also displayed in the user interface panel. Figure 8.10 plots the position and orientation tracking errors. It can be seen that these errors are actually not zero as proved in Theorem 8.3. These nonzero errors are due to inaccuracy of the DGPS and compass. A close look at this figure shows that the position and orientation errors are in the range of the DGPS and compass accuracy. To illustrate the role


Figure 8.10 Position and orientation errors with integral actions
of integral actions, we also tested the control algorithm without integral actions by setting $\sigma_{x e}=\sigma_{y e}=\sigma_{\bar{\psi} e}=0$. The tested results without integral actions are given in Figures 8.11 and 8.12. It can be seen that the integral actions play an important role in experimental tests. Poor performance in the test without integral actions can be understood as follows: The disturbances result in velocities, for which the controller does not compensate since the integral actions are switched off. The velocities induced by disturbances result in poor performance and even destroy the stability of the path-tracking error dynamics when these disturbing velocities are large enough (larger than $\min \left(k_{i}\right), i=1,2,3$ in magnitude). Theoretically, we can see this phenomenon by looking at the fourth, fifth, and sixth equations of the closed loop system (8.113):

$$
\begin{align*}
\dot{x}_{e}= & -k_{1} \frac{x_{e}}{\Delta}-\sigma_{x e}+u_{e}-u_{d}\left(\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \cos \left(\alpha_{\psi_{e}}\right)-\right. \\
& \left.\sin \left(\bar{\psi}_{e}\right) \sin \left(\alpha_{\psi_{e}}\right)\right)+(\hat{r}+\tilde{r}) y_{e}+\tilde{u}+u_{\tau_{E}}, \\
\dot{y}_{e}= & -k_{2} \frac{y_{e}}{\Delta}-\sigma_{y e}+v_{e}+u_{d}\left(\sin \left(\bar{\psi}_{e}\right) \cos \left(\alpha_{\psi_{e}}\right)+\right. \\
& \left.\left(\cos \left(\bar{\psi}_{e}\right)-1\right) \sin \left(\alpha_{\psi_{e}}\right)\right)-(\hat{r}+\tilde{r}) x_{e}+\tilde{v}+v_{\tau_{E}}, \\
\dot{\bar{\psi}}_{e}= & -k_{3} \frac{\bar{\psi}_{e}}{\sqrt{1+\bar{\psi}_{e}^{2}}}-\sigma_{\bar{\psi} e}+b_{1}\left(r_{e}+\tilde{r}\right)-\frac{k_{2} \Upsilon_{u_{d 0}} x_{e} y_{e}}{u_{d}^{2} \Delta^{3}}\left(u_{e}+\tilde{u}\right)+ \\
& \frac{k_{2} \Upsilon_{u_{d 0}}\left(1+x_{e}^{2}\right)}{u_{d}^{2} \Delta^{3}}\left(v_{e}+\tilde{v}\right)+r_{\tau_{E}}, \tag{8.150}
\end{align*}
$$

where $u_{\tau_{E}}, v_{\tau_{E}}$, and $r_{\tau_{E}}$ are velocities induced by the disturbances $\boldsymbol{\tau}_{\boldsymbol{E}}$. From (8.150), we can see that when $u_{\tau_{E}}, v_{\tau_{E}}$ and $r_{\tau_{E}}$ are larger than $k_{1}, k_{2}$ and $k_{3}$, respectively, in magnitude, the $x_{e^{-}}, y_{e^{-}}$, and $\bar{\psi}_{e^{-} \text {-dynamics will be unstable, respec- }}$ tively. It is noted that $\sigma_{x e}, \sigma_{y e}$, and $\sigma_{\bar{\psi} e}$ do not compensate for $u_{\tau_{E}}, v_{\tau_{E}}$, and $r_{\tau_{E}}$ since they are switched off. The constants $k_{1}$ and $k_{2}$ cannot be arbitrarily large since they have to satisfy conditions (8.103) and (8.111).


Figure 8.11 A screen shot of the user interface without integral actions


Figure 8.12 Position and orientation errors without integral actions

### 8.3 Conclusions

This chapter provides global path-tracking controllers for underactuated ships. Both state feedback and output feedback cases have been considered. Of particular note is the use of a "nontraditional" adaptive observer and a simple modification of the proposed controller to deal with inaccuracies of the nonlinear damping coefficients. Moreover, the proposed controllers in this chapter are much simpler than those in the preceding chapters. The simplicity of the proposed controllers in this chapter makes them more suitable for use in practice than those proposed in the previous chapters. The work presented in this chapter is based on [121-125].

## Chapter 9

## Way-point Tracking Control of Underactuated Ships

This chapter presents state feedback and output feedback controllers that force underactuated ships to globally ultimately track a straight line under environmental disturbances induced by waves, wind, and ocean currents. When there are no environmental disturbances, the controllers are able to drive the heading angle and cross-tracking error to zero asymptotically. Based on the backstepping technique and several technical lemmas introduced for a nonlinear system with nonvanishing disturbances, a full state feedback controller is first designed. An output feedback controller is then developed by using a nonlinear observer, which globally exponentially estimates the unmeasured sway and yaw velocities from the measured sway displacement and the measured yaw angle.

### 9.1 Control Objective

In addition to the assumptions made in Section 3.4.1.1, we assume that the surge velocity is controlled by the main propulsion control system. As such, the resulting mathematical model of the underactuated ship moving in sway and yaw is rewritten as

$$
\begin{align*}
& \dot{y}=u \sin (\psi)+\cos (\psi) v \\
& \dot{\psi}=r \\
& \dot{v}=-\frac{m_{11} u}{m_{22}} r-\frac{d_{22}}{m_{22}} v-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t),  \tag{9.1}\\
& \dot{r}=\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} v-\frac{d_{33}}{m_{33}} r-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t),
\end{align*}
$$

where $y, v, \psi, r$, and $u$ are sway displacement, sway velocity, yaw angle, yaw velocity, and forward speed controlled by the main thruster control system, respectively. Without loss of generality, we assume that the forward speed $u$ is positive and if
time-varying, has a bounded derivative $\dot{u}(t)$, i.e., $0<u_{\text {min }} \leq u(t) \leq u_{\text {max }}<\infty$ and $|\dot{u}(t)| \leq M<\infty, \forall t \geq 0$. The positive constant terms $m_{j j}, 1 \leq j \leq 3$ denote the ship's inertia including added mass. The positive constant terms $d_{22}, d_{33}, d_{v i}$ and $d_{r i}, i \geq 2$ represent the hydrodynamic damping in sway and yaw. The bounded time-varying terms, $\tau_{w v}(t)$ and $\tau_{w r}(t)$, are the environmental disturbance moments induced by wave, wind, and ocean current with an assumption that $\left|\tau_{w v}(t)\right| \leq$ $\tau_{w v \max }<\infty$ and $\left|\tau_{w r}(t)\right| \leq \tau_{w r \max }<\infty$. In this chapter, we study two control objectives. The first is full state feedback. In this case, we assume that all states $y, v, \psi$, and $r$ are available for feedback. In the design of an output feedback controller, only sway and yaw displacements are measurable. For both full state and output feedback cases, we design a control law, $\tau_{r}$, that forces the ship to track a linear course with ultimate boundedness, i.e. the tracking errors are globally ultimately bounded. When there are no environmental disturbances, the sway displacement and velocity, $y$ and $v$, yaw angle and velocity, $\psi$ and $r$, asymptotically converge to zero.

### 9.2 Full-state Feedback

### 9.2.1 Control Design

We define the following coordinate transformation

$$
\begin{equation*}
z_{1}=\psi+\arcsin \left(\frac{k y}{\sqrt{1+(k y)^{2}}}\right) \tag{9.2}
\end{equation*}
$$

where $k$ is a positive constant to be selected later. Note that the convergence of $z_{1}$ and $y$ to zero implies that of $\psi$. Upon application of the coordinate transformation (9.2), the ship dynamics (9.1) are rewritten as

$$
\begin{align*}
\dot{y}= & -\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{v}{\sqrt{1+(k y)^{2}}}+\frac{u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)}{\sqrt{1+(k y)^{2}}}+ \\
& \frac{v\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right)}{\sqrt{1+(k y)^{2}}}, \\
\dot{v}= & -\frac{m_{11} u}{m_{22}} r-\frac{d_{22}}{m_{22}} v-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t), \\
\dot{z}_{1}= & r-\frac{k^{2} u y}{\left(1+(k y)^{2}\right)^{3 / 2}}+\frac{k v}{\left(1+(k y)^{2}\right)^{3 / 2}}+\frac{k u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}+ \\
& \frac{k v\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}, \\
\dot{r}= & \frac{\left(m_{11}-m_{22}\right) u}{m_{33}} v-\frac{d_{33}}{m_{33}} r-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r} r(t) . \tag{9.3}
\end{align*}
$$

Therefore the problem of stabilizing (9.1) at the origin becomes that of stabilizing (9.3) at the origin. The structure of the model (9.3) suggests that we design the control $\tau_{r}$ in two stages by applying the popular backstepping technique. At the first step, we design an intermediate control $r_{d}$ for $r$ and at the second step the actual control $\tau_{r}$ will be designed to eliminate the error between $r_{d}$ and $r$.

## Step 1

Define

$$
\begin{equation*}
z_{2}=r-r_{d} \tag{9.4}
\end{equation*}
$$

where $r_{d}$ is an intermediate control designed as

$$
\begin{align*}
r_{d}= & -k_{1} z_{1}+\frac{k^{2} u y}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k v}{\left(1+(k y)^{2}\right)^{3 / 2}}- \\
& \frac{k u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k v\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}, \tag{9.5}
\end{align*}
$$

where $k_{1}$ is a positive constant to be selected later.

## Step 2

With (9.5), the time derivative of (9.4) along the solutions of the last equation of (9.3) is

$$
\begin{align*}
\dot{z}_{2}= & \frac{\left(m_{11}-m_{22}\right) u}{m_{33}} v-\frac{d_{33}}{m_{33}} r-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t)- \\
& \frac{\partial r_{d}}{\partial u} \dot{u}-\frac{\partial r_{d}}{\partial z_{1}}\left(-k_{1} z_{1}+z_{2}\right)-\frac{\partial r_{d}}{\partial y}\left(-\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{v}{\sqrt{1+(k y)^{2}}}+\right. \\
& \left.\frac{u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)}{\sqrt{1+(k y)^{2}}}+\frac{v\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right)}{\sqrt{1+(k y)^{2}}}\right)- \\
& \frac{\partial r_{d}}{\partial v}\left(-\frac{m_{11} u}{m_{22}} r-\frac{d_{22}}{m_{22}} v-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t)\right), \tag{9.6}
\end{align*}
$$

where

$$
\begin{aligned}
& \frac{\partial r_{d}}{\partial u}=\frac{k^{2} y}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k}{\left(1+(k y)^{2}\right)^{3 / 2}}\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right) \\
& \frac{\partial r_{d}}{\partial z_{1}}=-k_{1}-\frac{k u\left(\cos \left(z_{1}\right)+k y \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k v\left(-\sin \left(z_{1}\right)+k y \cos \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}
\end{aligned}
$$

$$
\begin{align*}
\frac{\partial r_{d}}{\partial y}= & \frac{-3 k^{3} y\left(k u y-u\left(\sin \left(z_{1}\right)-y\left(\cos \left(z_{1}\right)-1\right)\right)\right)}{\left(1+(k y)^{2}\right)^{5 / 2}}- \\
& \frac{v\left(\left(\cos \left(z_{1}\right)-1\right)+y \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{5 / 2}}+\frac{k^{2}\left(u \cos \left(z_{1}\right)-v \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}, \\
\frac{\partial r_{d}}{\partial v}= & -\frac{k}{\left(1+(k y)^{2}\right)^{3 / 2}}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) . \tag{9.7}
\end{align*}
$$

We now choose the actual control without canceling the useful damping terms as

$$
\begin{align*}
\tau_{r}= & m_{33}\left[-z_{1}-k_{2} z_{2}-\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} v+\frac{d_{33}}{m_{33}} r_{d}+\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} r_{d}+\right. \\
& \frac{\partial r_{d}}{\partial u} \dot{u}+\frac{\partial r_{d}}{\partial z_{1}}\left(-k_{1} z_{1}+z_{2}\right)+\frac{\partial r_{d}}{\partial y}\left(-\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{v}{\sqrt{1+(k y)^{2}}}+\right. \\
& \left.\frac{u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)}{\sqrt{1+(k y)^{2}}}+\frac{v\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right)}{\sqrt{1+(k y)^{2}}}\right)+ \\
& \frac{\partial r_{d}}{\partial v}\left(-\frac{m_{11} u}{m_{22}} r-\frac{d_{22}}{m_{22}} v-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v\right)-  \tag{9.8}\\
& \left.\frac{1}{m_{33}} \tau_{w r \max } \tanh \left(\frac{z_{2}}{\rho_{1}}\right)-\frac{1}{m_{22}} \tau_{w v \max } \frac{\partial r_{d}}{\partial v} \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{z_{2}}{\rho_{2}}\right)\right],
\end{align*}
$$

where $k_{2}, \rho_{1}$, and $\rho_{2}$ are positive constants to be chosen later. Substituting (9.4), (9.5), and (9.8) into (9.3) results in the following closed loop system:

$$
\begin{align*}
\dot{y}= & -\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{v}{\sqrt{1+(k y)^{2}}}+\frac{u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)}{\sqrt{1+(k y)^{2}}}+ \\
& \frac{v}{\sqrt{1+(k y)^{2}}}\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right), \\
\dot{v}= & -\frac{d_{22}}{m_{22}} v-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v-\frac{m_{11} u}{m_{22}} \frac{k^{2} u y-k v}{\left(1+(k y)^{2}\right)^{3 / 2}}- \\
& \frac{m_{11} u}{m_{22}}\left(-k_{1} z_{1}+z_{2}-\frac{k u}{\left(1+(k y)^{2}\right)^{3 / 2}}\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)-\right. \\
& \left.\frac{k v}{\left(1+(k y)^{2}\right)^{3 / 2}}\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right)\right)+\frac{1}{m_{22}} \tau_{w v}(t), \\
\dot{z}_{1}= & -k_{1} z_{1}+z_{2}, \\
\dot{z}_{2}= & -z_{1}-k_{2} z_{2}-\frac{d_{33}}{m_{33}} z_{2}-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} z_{2}+\frac{1}{m_{33}}\left(\tau_{w r}(t)-\tau_{w r \max \times}\right. \\
& \left.\tanh \left(\frac{z_{2}}{\rho_{1}}\right)\right)+\frac{1}{m_{22}}\left(-\frac{\partial r_{d}}{\partial v} \tau_{w v}(t)-\frac{\partial r_{d}}{\partial v} \tau_{w v \max } \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{z_{2}}{\rho_{2}}\right)\right) . \tag{9.9}
\end{align*}
$$

### 9.2.2 Stability Analysis

The following two lemmas will be used extensively in stability analysis.
Lemma 9.1. Consider the following nonlinear system:

$$
\begin{equation*}
\dot{x}=f(t, x)+g(t, x, \xi(t)) \tag{9.10}
\end{equation*}
$$

where $x \in \mathbb{R}^{n}, \xi(t) \in \mathbb{R}^{m}, f(t, x)$ is piecewise continuous in $t$ and locally Lipschitz in $x$. If there exist positive constants $c_{i}, 1 \leq i \leq 4, \lambda_{j}, 1 \leq j \leq 2, \sigma_{0}, \varepsilon_{0}, \mu_{0}, c_{0}$, and a class- $K$ function $\alpha_{0}$ such that the following conditions are satisfied:

C1. There exists a proper function $V(t, x)$ satisfying:

$$
\begin{aligned}
& c_{1}\|x\|^{2} \leq V(t, x) \leq c_{2}\|x\|^{2}, \\
& \left\|\frac{\partial V}{\partial x}(t, x)\right\| \leq c_{3}\|x\|, \\
& \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x) \leq-c_{4}\|x\|^{2}+c_{0} .
\end{aligned}
$$

C2. The vector function $g(t, x, \xi(t))$ satisfies:

$$
\|g(t, x, \xi(t))\| \leq\left(\lambda_{1}+\lambda_{2}\|x\|\right)\|\xi(t)\|
$$

$C 3$. $\xi(t)$ globally exponentially converges to a ball centered at the origin:

$$
\|\xi(t)\| \leq \alpha_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+\varepsilon_{0}, \forall t \geq t_{0} \geq 0
$$

C4. The following gain condition is satisfied:

$$
c_{4}-\lambda_{2} c_{3} \varepsilon_{0}-\frac{\lambda_{1} c_{3} \varepsilon_{0}}{4 \mu_{0}}>0
$$

Then the solution $x(t)$ of (9.10) globally exponentially converges to a ball centered at the origin, i.e.,

$$
\begin{equation*}
\|x(t)\| \leq \alpha\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right) e^{-\sigma\left(t-t_{0}\right)}+\varepsilon, \forall t \geq t_{0} \geq 0 \tag{9.11}
\end{equation*}
$$

where $\varepsilon=\sqrt{a_{4} / c_{1} a_{1}}$ and

$$
\begin{aligned}
& \text { if } a_{1}=\sigma_{0} \text { then } \\
& \qquad \begin{aligned}
\alpha(s) & =\sqrt{\frac{e^{\frac{a_{2}(s)}{\sigma_{0}}}}{c_{1}}}\left(c_{2} s^{2}+\left(a_{3}(s)+a_{1}^{-1} a_{2}(s) a_{4}\right) \phi\right) \\
& \sigma=0.5\left(a_{1}-d\right)
\end{aligned}
\end{aligned}
$$

if $a_{1} \neq \sigma_{0}$ then

$$
\begin{aligned}
& \alpha(s)=\sqrt{\frac{e^{\frac{a_{2}(s)}{\sigma_{0}}}}{c_{1}}\left(c_{2} s^{2}+\frac{a_{1} a_{3}(s)+a_{2}(s) a_{4}}{a_{1}\left|a_{1}-\sigma_{0}\right|}\right)} \\
& \sigma=0.5 \min \left(a_{1},\left|a_{1}-\sigma_{0}\right|\right)
\end{aligned}
$$

with

$$
\begin{aligned}
& a_{1}=\frac{1}{c_{2}}\left(c_{4}-\lambda_{2} c_{3} \varepsilon_{0}-\frac{\lambda_{1} c_{3} \varepsilon_{0}}{4 \mu_{0}}\right), \\
& a_{2}(s)=\frac{c_{3}}{c_{1}}\left(\lambda_{1}+\lambda_{2}\right) \alpha_{0}(s), \\
& a_{3}(s)=\frac{\lambda_{1} c_{3}}{4} \alpha_{0}(s) \\
& a_{4}=c_{0}+\lambda_{1} c_{3} \varepsilon_{0} \mu_{0} \\
& 0<d<a_{1}, \quad \phi \geq\left(t-t_{0}\right) e^{-d\left(t-t_{0}\right)}, \quad \forall t \geq t_{0} \geq 0, s \geq 0 .
\end{aligned}
$$

When $c_{0}=0$ and $\varepsilon_{0}=0$, we have $\varepsilon=0$ and the system (9.10) is globally $K$ exponentially stable. Note that a finite value of the constant $\phi$ exists for an arbitrarily small positive d.

Proof. From conditions C1, C2, and C3, we have

$$
\begin{align*}
\dot{V}= & \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x)+\frac{\partial V}{\partial x} g(t, x, \xi(t)) \\
\leq & -\left(c_{4}-\lambda_{2} c_{3} \varepsilon_{0}-\lambda_{1} c_{3} \varepsilon_{0} / 4 \mu_{0}\right)\|x\|^{2}+ \\
& c_{3}\|x\|\left(\lambda_{1}+\lambda_{2}\|x\|\right) \alpha_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+ \\
& \lambda_{1} c_{3} \varepsilon_{0} \mu_{0}+c_{0} . \tag{9.12}
\end{align*}
$$

Upon application of the completing square, (9.12) can be rewritten as

$$
\begin{equation*}
\dot{V} \leq-\left(a_{1}-a_{2} e^{-\sigma_{0}\left(t-t_{0}\right)}\right) V+a_{3} e^{-\sigma_{0}\left(t-t_{0}\right)}+a_{4} \tag{9.13}
\end{equation*}
$$

Solving the above differential inequality results in

$$
\begin{equation*}
V(t) \leq V\left(t_{0}\right) e^{\frac{a_{2}}{\sigma_{0}}} e^{-a_{1}\left(t-t_{0}\right)}+\left(a_{3}+\frac{a_{2} a_{4}}{a_{1}}\right) e^{\frac{a_{2}}{\sigma_{0}}} e^{-a_{1} t+\sigma_{0} t_{0}} \int_{t_{0}}^{t} e^{\left(a_{1}-\sigma_{0}\right) \tau} \mathrm{d} \tau+\frac{a_{4}}{a_{1}} \tag{9.14}
\end{equation*}
$$

which yields (9.11) readily.
Lemma 9.2. Consider the following nonlinear system:

$$
\begin{equation*}
\dot{x}=f(t, x)+g(t, x, \xi(t)) \tag{9.15}
\end{equation*}
$$

where $x_{1} \in \mathbb{R}^{n_{1}}, x_{2} \in \mathbb{R}^{n_{2}}, x=\left[x_{1} x_{2}\right]^{T} \in \mathbb{R}^{n_{1}+n_{2}}, \xi(t) \in \mathbb{R}^{m}, f(t, x)$ is piecewise continuous in $t$ and locally Lipschitz in $x$. If there exist positive constants $c_{0}, c_{1}, c_{2}$, $c_{31}, c_{32}, \lambda_{i}, 0 \leq i \leq 2, \sigma_{0}, \varepsilon_{0}, c_{0}$, and a class- $K$ function $\alpha_{0}$ such that the following conditions are satisfied.

C1. There exists a proper function $V(t, x)$ such that:

$$
\begin{aligned}
& c_{1}\|x\|^{2} \leq V(t, x) \leq c_{2}\|x\|^{2}, \\
& \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x) \leq-c_{31}\left\|x_{1}\right\|^{2}-\frac{c_{32}\left\|x_{2}\right\|^{2}}{\sqrt{1+c_{4}\left\|x_{2}\right\|^{2}}}+c_{0}, \\
& \left\|\frac{\partial V}{\partial x} g(t, x, \xi(t))\right\| \leq\left(\lambda_{0}+\lambda_{1}\left\|x_{1}\right\|^{2}+\frac{\lambda_{2}\left\|x_{2}\right\|^{2}}{\sqrt{1+c_{4}\left\|x_{2}\right\|^{2}}}\right)\|\xi(t)\| .
\end{aligned}
$$

$C 2 . \xi(t)$ globally exponentially converges to a ball centered at the origin:

$$
\|\xi(t)\| \leq \alpha_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+\varepsilon_{0}, \forall t \geq t_{0} \geq 0
$$

C3. The following gain conditions are satisfied:

$$
c_{31}-\lambda_{1} \varepsilon_{0}>0 \quad \text { and } \quad c_{32}-\lambda_{2} \varepsilon_{0}>0
$$

C4. $x_{2}(t)$ is bounded:

$$
\left\|x_{2}(t)\right\| \leq \varpi
$$

where $\varpi$ is a nondecreasing function of $\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|$,
then the solution $x(t)$ of (9.15) globally asymptotically converges to a ball centered at the origin, i.e.,

$$
\begin{equation*}
\|x(t)\| \leq \alpha\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right) e^{-\sigma\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right)\left(t-t_{0}\right)}+\varepsilon(s), \quad \forall t \geq t_{0} \geq 0 \tag{9.16}
\end{equation*}
$$

where $\varepsilon(s)=\sqrt{\frac{a_{4}}{c_{1} a_{1}(s)}}$ and

$$
\begin{aligned}
& \text { if } a_{1}(s)=\sigma_{0} \text { then } \\
& \alpha(s)=\sqrt{c_{1}^{-1} e^{a_{2}(s) / \sigma_{0}}\left(c_{2} s^{2}+\left(a_{3}(s)+a_{1}^{-1}(s) a_{2}(s) a_{4}\right) \phi\right)} \\
& \sigma(s)=0.5\left(a_{1}(s)-d\right) ; \\
& \text { if } a_{1}(s) \neq \sigma_{0} \text { then } \\
& \alpha(s)=\sqrt{c_{1}^{-1} e^{a_{2}(s) / \sigma_{0}}\left(c_{2} s^{2}+\frac{a_{1}(s) a_{3}(s)+a_{2}(s) a_{4}}{a_{1}(s)\left|a_{1}(s)-\sigma_{0}\right|}\right)} \\
& \sigma(s)=0.5 \min \left(a_{1}(s),\left|a_{1}(s)-\sigma_{0}\right|\right) ;
\end{aligned}
$$

with

$$
\begin{aligned}
& a_{1}(s)=\frac{1}{c_{2}} \min \left(c_{31}-\lambda_{1} \varepsilon_{0}, \frac{c_{32}-\lambda_{2} \varepsilon_{0}}{\sqrt{1+c_{4} \varpi^{2}(s)}}\right), \\
& a_{2}(s)=\frac{1}{c_{1}} \max \left(\lambda_{1}, \lambda_{2}\right) \alpha_{0}(s), \\
& a_{3}(s)=\lambda_{0} \alpha_{0}(s), \\
& a_{4}=c_{0}+\lambda_{0} \varepsilon_{0}, \\
& 0<d<a_{1}(s), \phi \geq\left(t-t_{0}\right) e^{-d\left(t-t_{0}\right)}, \quad \forall t \geq t_{0} \geq 0, s \geq 0 .
\end{aligned}
$$

When $c_{0}=0$ and $\varepsilon_{0}=0$, we have $\varepsilon=0$ and the system (9.15) is GAS. Note that a finite value of the constant $\phi$ exists for an arbitrarily small positive $d$.

Proof. The proof of this lemma is similar to that of Lemma 9.1.
Remark 9.1. It is important to note that the rate $\sigma>0$ in (9.16) and $a_{1}$ depend on the initial conditions. In addition, around the origin, both $\sigma$ and $a_{1}$ are bounded below from zero.

We first need to show that the closed loop system (9.9) is forward complete. It is straightforward to show that the derivative of the function $V_{0}=z_{1}^{2}+z_{2}^{2}+v^{2}+y^{2}$ along the solutions of the closed loop system (9.9) satisfies $\dot{V}_{0} \leq a_{0} V_{0}+b_{0}$ where $a_{0}$ and $b_{0}$ are nonnegative constants. The inequality $\dot{V}_{0} \leq a_{0} V_{0}+b_{0}$ implies that the closed loop system (9.9) is forward complete. We now apply Lemmas 9.1 and 9.2 to analyze the closed loop system (9.9). We view $\left(z_{1}, z_{2}\right)$ as $\xi(t), v$ as $x$ in Lemma 9.1, and $(v, y)$ as $x$ in Lemma 9.2. Hence it is necessary to verify all the conditions of Lemmas 9.1 and 9.2.

## $\left(z_{1}, z_{2}\right)$-subsystem

We take the following quadratic function:

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(z_{1}^{2}+z_{2}^{2}\right) \tag{9.17}
\end{equation*}
$$

whose time derivative along the solutions of the last two equations of (9.9) satisfies

$$
\begin{aligned}
& \dot{V}_{1}=-k_{1} z_{1}^{2}-k_{2} z_{2}^{2}-\frac{d_{33}}{m_{33}} z_{2}^{2}-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} z_{2}^{2}+\frac{z_{2}}{m_{33}}\left(\tau_{w r}(t)-\tau_{w r \max } \times\right. \\
& \left.\quad \tanh \left(\frac{z_{2}}{\rho_{1}}\right)\right)+\frac{z_{2}}{m_{22}}\left(-\frac{\partial r_{d}}{\partial v} \tau_{w v}(t)-\frac{\partial r_{d}}{\partial v} \tau_{w v \max } \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{z_{2}}{\rho_{2}}\right)\right)
\end{aligned}
$$

$$
\begin{align*}
\leq & -k_{1} z_{1}^{2}-k_{2} z_{2}^{2}+\frac{\tau_{w r \max }}{m_{33}} \times \\
& \left(\left|z_{2}\right|-z_{2} \tanh \left(\frac{z_{2}}{\rho_{1}}\right)\right)+\frac{\tau_{w v \max }}{m_{22}}\left(\left|\frac{\partial r_{d}}{\partial v} z_{2}\right|-\frac{\partial r_{d}}{\partial v} z_{2} \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{z_{2}}{\rho_{2}}\right)\right) \\
\leq & -k_{1} z_{1}^{2}-k_{2} z_{2}^{2}+0.2785\left(\frac{1}{m_{33}} \tau_{w r \max } \rho_{1}+\frac{1}{m_{22}} \tau_{w v \max } \rho_{2}\right) \tag{9.18}
\end{align*}
$$

where we have used $|x|-x \tanh (x / \lambda) \leq 0.2785 \lambda, \forall x \in \mathbb{R}$ and $\lambda>0$. From (9.17) and (9.18), it can be shown that

$$
\begin{equation*}
\|z(t)\| \leq\left\|z\left(t_{0}\right)\right\| e^{-\sigma_{0}\left(t-t_{0}\right)}+\varepsilon_{0} \forall t \geq t_{0} \geq 0 \tag{9.19}
\end{equation*}
$$

where $z=\left[\begin{array}{ll}z_{1} & z_{2}\end{array}\right]^{T}$ and

$$
\begin{align*}
\sigma_{0} & =\min \left(k_{1}, k_{2}\right) \\
\varepsilon_{0} & =\sqrt{\frac{0.2785\left(\tau_{w r \max } \rho_{1} / m_{33}+\tau_{w v \max } \rho_{2} / m_{22}\right)}{\sigma_{0}}} \tag{9.20}
\end{align*}
$$

Therefore the $\left(z_{1}, z_{2}\right)$-subsystem is globally ultimately stable at the origin. Furthermore, (9.19) implies that $\xi(t):=\left(z_{1}, z_{2}\right)^{T}$ globally exponentially converges to a ball centered at the origin. The radius of this ball can be made arbitrarily small by increasing $k_{1}$ and $k_{2}$ and/or reducing $\rho_{1}$ and $\rho_{2}$.

## Boundedness of $v$

To prove that $v$ is bounded, we consider the second equation of (9.9). In order to apply Lemma 9.1, define $x=v, \xi(t)=\left[z_{1} z_{2}\right]^{T}$ and consider $y$ as a function of time $t$,

$$
\begin{align*}
f(\bullet) & =-\frac{d_{22}}{m_{22}} v-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v- \\
& \frac{m_{11} u}{m_{22}}\left(\frac{k^{2} u y}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k v}{\left(1+(k y)^{2}\right)^{3 / 2}}\right)+\frac{1}{m_{22}} \tau_{w v}(t) \\
g(\bullet) & =-\frac{m_{11} u}{m_{22}}\left(-k_{1} z_{1}+z_{2}-\right.  \tag{9.21}\\
& \left.\frac{k\left(u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)+v\left(\left(\cos \left(z_{1}\right)-1\right)+k y \sin \left(z_{1}\right)\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}\right) .
\end{align*}
$$

This abuse of notation is introduced for simplicity and is possible because:

$$
0 \leq \frac{1}{\left(1+(k y)^{2}\right)^{3 / 2}} \leq 1,0 \leq\left|\frac{k y}{\left(1+(k y)^{2}\right)^{3 / 2}}\right|<1, \forall y \in \mathbb{R}
$$

and we have shown that the closed loop system is forward complete. We now verify all of the conditions of Lemma 9.1.

Verifying Condition C1. We take the function $V_{2}=0.5 v^{2}$ whose time derivative along the solutions of the differential equation $\dot{v}=f(t, v)$, see (9.21), satisfies

$$
\begin{align*}
& \dot{V}_{2}=-\frac{d_{22}}{m_{22}} v^{2}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v^{2}-\frac{m_{11} u v}{m_{22}} \frac{k^{2} u y-k v}{\left(1+(k y)^{2}\right)^{3 / 2}}+\frac{v}{m_{22}} \tau_{w v}(t) \\
& \leq-\left(\frac{d_{22}}{m_{22}}-\frac{m_{11} k u_{\max }}{m_{22}}-\frac{m_{11} k u_{\max }^{2} \mu_{1}}{m_{22}}-\frac{\mu_{1}}{m_{22}}\right) v^{2}+\frac{m_{11} k u_{\max }^{2}+\tau_{w v \max }^{2}}{4 \mu_{1} m_{22}} . \tag{9.22}
\end{align*}
$$

Hence, the condition C 1 is satisfied with

$$
\begin{align*}
& c_{0}=\frac{1}{4 \mu_{1} m_{22}}\left(m_{11} k u_{\max }^{2}+\tau_{w v \max }^{2}\right), c_{1}=c_{2}=0.5, c_{3}=1, \\
& c_{4}=\frac{d_{22}}{m_{22}}-\frac{m_{11} k u_{\max }^{2}}{m_{22}} \mu_{1}-\frac{m_{11} k u_{\max }}{m_{22}}-\frac{\mu_{1}}{m_{22}} \tag{9.23}
\end{align*}
$$

where $\mu_{1}>0$ and $k>0$ are chosen such that $c_{4}>0$.
Verifying Condition C2. It is directly shown from (9.21) that

$$
\begin{equation*}
|g(t, v, z(t))| \leq\left(\lambda_{1}+\lambda_{2}|v|\right)\|z(t)\|, \tag{9.24}
\end{equation*}
$$

where

$$
\begin{equation*}
\lambda_{1}=\frac{m_{11} u_{\max }}{m_{22}}\left(1+k_{1}+2 k u_{\max }\right), \lambda_{2}=\frac{2 k m_{11} u_{\max }}{m_{22}} \tag{9.25}
\end{equation*}
$$

Verifying Condition C3. This condition follows directly from (9.19).
Verifying Condition C4. It can be shown from (9.20), (9.23), and (9.25) that we can find positive constant $k$ such that the condition C 4 is satisfied, i.e.,

$$
\begin{equation*}
c_{4}-\lambda_{2} c_{3} \varepsilon_{0}-\frac{\lambda_{1} c_{3} \varepsilon_{0}}{4 \mu_{0}}>0 \tag{9.26}
\end{equation*}
$$

All of the conditions of Lemma 9.1 have been verified, hence the sway velocity is bounded and satisfies

$$
\begin{equation*}
|v(t)| \leq \alpha_{1}\left(\left\|\left(v\left(t_{0}\right), z\left(t_{0}\right)\right)\right\|\right) e^{\sigma_{1}\left(t-t_{0}\right)}+\varepsilon_{1} \tag{9.27}
\end{equation*}
$$

where $\varepsilon_{1}, \sigma_{1}$, and $\alpha_{1}$ are calculated as in Lemma 9.1, and the constants $c_{i}, 1 \leq i \leq 4$, $\lambda_{j}, 1 \leq j \leq 2, \sigma_{0}, \varepsilon_{0}, \mu_{0}$, and $c_{0}$ are given in (9.20), (9.23), and (9.26).

## ( $v, y$ )-subsystem

In this section, we will apply Lemma 9.2 to prove global ultimate boundedness of the $(v, y)$-subsystem. It can be seen that the first two equations of (9.9) are in the form of the system in Lemma 9.2 with $x_{1}=v, x_{2}=y, \xi(t)=z(t)$, and

$$
\begin{align*}
& f(\bullet)=\left[\begin{array}{c}
-\frac{d_{22}}{m_{22}} v-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v-\frac{m_{11} u}{m_{22}} \frac{k^{2} u y-k v}{\left(1+(k y)^{2}\right)^{3 / 2}}+\frac{1}{m_{22}} \tau_{w v}(t) \\
-\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{v}{\sqrt{1+(k y)^{2}}}
\end{array}\right], \\
& g(\bullet)=\left[\begin{array}{c}
\frac{m_{11} u}{m_{22}}\left(k \frac{\sin \left(z_{1}\right)(u+k v y)+\left(\cos \left(z_{1}\right)-1\right)(v-k u y)}{\left(1+(k y)^{2}\right)^{3 / 2}}+k_{1} z_{1}-z_{2}\right) \\
\frac{u\left(\sin \left(z_{1}\right)-\left(\cos \left(z_{1}\right)-1\right) k y\right)}{\sqrt{1+(k y)^{2}}}+\frac{v\left(\left(\cos \left(z_{1}\right)-1\right)+\sin \left(z_{1}\right) k y\right)}{\sqrt{1+(k y)^{2}}}
\end{array}\right] . \tag{9.28}
\end{align*}
$$

We now need to verify all of the conditions of Lemma 9.2.
Verifying Condition C1. To verify this condition, we take the function $V_{3}=$ $0.5\left(v^{2}+y^{2}\right)$. It can be directly shown that this function satisfies condition C 1 with $|v(t)| \leq \alpha_{1}\left(\left\|\left(v\left(t_{0}\right), z\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{1}\left(t-t_{0}\right)}+\varepsilon_{1}$ and

$$
\begin{align*}
& c_{0}=\frac{\tau_{w v \max }^{2}}{4 \mu_{3} m_{22}}, c_{1}=c_{2}=0.5, \\
& c_{31}=\frac{d_{22}}{m_{22}}-\frac{m_{11} k u_{\max }}{m_{22}}-\mu_{2}-\frac{m_{11} k^{2} u_{\max }^{2}}{m_{22}} \mu_{2}-\frac{\mu_{3}}{m_{22}}, \\
& c_{32}=k u_{\min }-\frac{1}{4 \mu_{2}}\left(1+\frac{m_{11} k^{2} u_{\max }^{2}}{m_{22}}\right), \\
& \lambda_{0}=\frac{m_{11} u_{\max }}{4 \mu_{4} m_{22}}\left(1+k_{1}+2 k u_{\max }\right)+\frac{u_{\max }}{4 \mu_{4}},  \tag{9.29}\\
& \lambda_{1}=\frac{m_{11} u_{\max }}{m_{22}}\left(2 k u_{\max }+k_{1}+1\right) \mu_{4}+\frac{2 k m_{11} u_{\max }}{m_{22}}+\mu_{4}, \\
& \lambda_{2}=\frac{1}{4 \mu_{4}}+u_{\max } \mu_{4}+k u_{\max }+k\left(\alpha_{1}+\varepsilon_{1}\right),
\end{align*}
$$

where $k>0$ and $\mu_{2}>0$ are chosen such that $c_{31}>0$ and $c_{32}>0$.
Verifying Condition C2. This condition follows directly from (9.19).
Verifying Condition C3. It can be shown that there exists a positive constant $k$ such that the condition C3 satisfies, i.e.,

$$
\begin{equation*}
c_{31}-\lambda_{1} \varepsilon_{0}>0, c_{32}-\lambda_{2} \varepsilon_{0}>0 \tag{9.30}
\end{equation*}
$$

Verifying Condition C4. From the boundedness of the sway velocity $v(t)$ proven in the previous subsection and noting that $\varepsilon_{0}$ in (9.19) can be made arbitrarily small, it is shown that there exists a nondecreasing function $\varpi$ of $\left\|\left(\left(v\left(t_{0}\right), y\left(t_{0}\right)\right), z\left(t_{0}\right)\right)\right\|$ such that $|y(t)| \leq \varpi$ by applying Lemma 9.1 to the first equation of (9.9) with the Lyapunov function $V_{y}=0.5 y^{2}$.

All of the conditions of Lemma 9.2 have been verified. Therefore we have

$$
\begin{equation*}
\|(v(t), y(t))\| \leq \alpha_{2}\left(\left\|\left(\left(v\left(t_{0}\right), y\left(t_{0}\right)\right), z\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{2}\left(t-t_{0}\right)}+\varepsilon_{2}, \forall t \geq t_{0} \geq 0 \tag{9.31}
\end{equation*}
$$

where $\varepsilon_{2}, \sigma_{2}$, and $\alpha_{2}$ are calculated as in Lemma 9.2, and all other constants given in (9.29). It can be seen that when there are no environmental disturbances, since $\varepsilon_{2}=0,(v(t), y(t))$ globally asymptotically converges to zero. We have thus proven the first main result of this chapter.

Theorem 9.1. The full-state feedback control problem stated in Section 9.1 is solved by the control law (9.8) as long as the design constants $k, k_{1}$, and $k_{2}$ are chosen such that (9.26) and (9.30) hold.

### 9.3 Output Feedback

This section is devoted to the development of an output feedback controller to fulfill the output feedback control objective. A nonlinear observer is first designed so that it globally exponentially drives the observer error dynamics to a ball centered at the origin. When there are no environmental disturbances, the observer error dynamics are GES at the origin. A controller is then designed based on the approach in the preceding section and the proposed observer. Before designing an observer and output feedback controller, we impose the following assumption, see [12].

Assumption 9.1. For the ship model (9.1), the matrix

$$
\boldsymbol{K}_{\mathbf{2}}=\left[\begin{array}{ll}
-\frac{d_{22}}{m_{22}} & -\frac{m_{11} u}{m_{22}}  \tag{9.32}\\
\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} & -\frac{d_{33}}{m_{33}}
\end{array}\right]
$$

is Hurwitz.
The above assumption implies that the ship (when the nonlinear damping terms $\sum_{i \geq 2} \frac{d_{v i}}{m_{33}}|v|^{i-1} v$ and $\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} r$ are ignored) is dynamic stable in straight-line motion. Straight-line stability physically implies that a new path of the ship will be a straight line after an action in yaw. The direction of the new path will usually be different from that of the initial path, as mentioned in [12]. On the other hand,
unstable ships will go into a starboard or port turn without any rudder deflection. We impose Assumption 9.1 to make our observer design possible. Note that this assumption does not hold for several types of surface ships such as large tankers and high-speed crafts with sufficiently small ratios $d_{22} / m_{22}$ and $d_{33} / m_{33}$, and the added mass in the sway axis sufficiently larger than the added mass in the surge axis. Consequently, for these ships the real part of at least one of the eigenvalues of the matrix $\boldsymbol{K}_{\mathbf{2}}$ is positive.

### 9.3.1 Observer Design

The ship dynamics (9.1) represent some difficulties for output feedback control design. These difficulties are mainly due to the nonlinear terms $\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|v|^{i-1} v$ and $\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|r|^{i-1} r$, the nonlinear kinematic term $\cos (\psi)$, and the underactuated situation. However we first observe that the nonlinear terms are monotonic, i.e., they satisfy

$$
\begin{align*}
& \left(v_{1}-v_{2}\right)\left(\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}\left|v_{1}\right|^{i-1} v_{1}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}\left|v_{2}\right|^{i-1} v_{2}\right) \geq 0, \forall v_{1} \in \mathbb{R}, v_{2} \in \mathbb{R}, \\
& \left(r_{1}-r_{2}\right)\left(\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}\left|r_{1}\right|^{i-1} r_{1}-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}\left|r_{2}\right|^{i-1} r_{2}\right) \geq 0, \forall r_{1} \in \mathbb{R}, r_{2} \in \mathbb{R} . \tag{9.33}
\end{align*}
$$

Based on the structure of the underatuated ship dynamics (9.1) and property (9.33), we propose the following nonlinear observer:

$$
\begin{align*}
\dot{\hat{y}}= & u \sin (\psi)+\cos (\psi) \hat{v}+k_{11}(y-\hat{y})+k_{12}(\psi-\hat{\psi}), \\
\dot{\hat{\psi}}= & \hat{r}+k_{21}(y-\hat{y})+k_{22}(\psi-\hat{\psi}), \\
\dot{\hat{v}}= & -\frac{m_{11} u}{m_{22}} \hat{r}-\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}+k_{31}(y-\hat{y})+  \tag{9.34}\\
& \left(k_{13}+\cos (\psi)\right)(y-\hat{y}), \\
\dot{\hat{r}}= & \frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \hat{v}-\frac{d_{33}}{m_{33}} \hat{r}-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} \hat{r}+\frac{1}{m_{33}} \tau_{r}+ \\
& k_{42}(\psi-\hat{\psi})+\left(k_{24}+1\right)(\psi-\hat{\psi}),
\end{align*}
$$

where $\hat{y}, \hat{\psi}, \hat{v}$, and $\hat{r}$ are the estimate of $y, \psi, v$ and $r$ respectively. All the constants $k_{11}, k_{12}, k_{21}, k_{22}, k_{13}, k_{31}$, and $k_{42}$ will be chosen later.

By defining the observer errors as $\tilde{y}=y-\hat{y}, \tilde{\psi}=\psi-\hat{\psi}, \tilde{v}=v-\hat{v}$, and $\tilde{r}=$ $r-\hat{r}$, the observer error dynamics can be rewritten as

$$
\begin{align*}
\dot{\tilde{y}}= & -k_{11} \tilde{y}-k_{12} \tilde{\psi}-k_{13} \tilde{v}+\left(k_{13}+\cos (\psi)\right) \tilde{v}, \\
\dot{\tilde{\psi}}= & -k_{21} \tilde{y}-k_{22} \tilde{\psi}-k_{24} \tilde{r}+\left(k_{24}+1\right) \tilde{r}, \\
\dot{\tilde{v}}= & -k_{31} \tilde{y}-\frac{d_{22}}{m_{22}} \tilde{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}\left(|v|^{i-1} v-|\hat{v}|^{i-1} \hat{v}\right)-\frac{m_{11} u}{m_{22}} \tilde{r}- \\
& \left(k_{13}+\cos (\psi)\right) \tilde{y}+\frac{1}{m_{22}} \tau_{w v}(t), \\
\dot{\tilde{r}}= & -k_{42} \tilde{\psi}-\left(k_{24}+1\right) \tilde{\psi}+\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \tilde{v}-\frac{d_{33}}{m_{33}} \tilde{r}- \\
& \sum_{i \geq 2} \frac{d_{r i}}{m_{33}}\left(|r|^{i-1} r-|\hat{r}|^{i-1} \hat{r}\right)+\frac{1}{m_{33}} \tau_{w r}(t) . \tag{9.35}
\end{align*}
$$

We now show that there exist suitable observer gains $k_{11}, k_{12}, k_{13}, k_{21}, k_{22}$, $k_{24}, k_{31}$, and $k_{42}$ such that the observer error dynamics (9.35) is globally ultimately stable. Consider the Lyapunov function

$$
\begin{equation*}
V_{\mathrm{obs}}=\frac{1}{2} \tilde{x}^{T} \tilde{x} \tag{9.36}
\end{equation*}
$$

where $\tilde{x}=\left[\begin{array}{llll}\tilde{y} & \tilde{\psi} & \tilde{v} & \tilde{r}\end{array}\right]^{T}$. The time derivative of (9.36) along the solutions of (9.35) and property (9.33) results in

$$
\begin{equation*}
\dot{V}_{\mathrm{obs}} \leq-p_{0}\|\tilde{x}\|^{2}+q_{0} \tag{9.37}
\end{equation*}
$$

where

$$
\left.\begin{array}{rl}
p_{0} & =-\lambda_{\max }(A)-\max \left(\frac{\tau_{w v \max }}{m_{22}}, \frac{\tau_{w r \max }}{m_{33}}\right) \frac{1}{4 \mu_{0}}, \\
q_{0} & =\max \left(\frac{\tau_{w v \max }}{m_{22}}, \frac{\tau_{w r \max }}{m_{33}}\right) \mu_{0}, \mu_{0}>0, \\
A & =\left[\begin{array}{ccc}
-k_{11}-k_{12} & -k_{13} & 0 \\
-k_{21}-k_{22} & 0 & -k_{24} \\
-k_{31} & 0 & -\frac{d_{22}}{m_{22}} \\
0 & -k_{42} & -\frac{\left(m_{11} u\right.}{m_{22}} \\
m_{33}
\end{array}\right]
\end{array}\right] .
$$

The above matrix $A$ is made negative definite by choosing

$$
\begin{align*}
& k_{13}=k_{24}=k_{31}=k_{42}, \\
& K_{1}:=\left[\begin{array}{l}
-k_{11}-k_{12} \\
-k_{21}-k_{22}
\end{array}\right]<0,  \tag{9.38}\\
& K_{2}=K_{12} K_{1}^{-1} K_{12}<0,
\end{align*}
$$

where

$$
K_{12}=\left[\begin{array}{ll}
-k_{13} & 0  \tag{9.39}\\
0 & -k_{24}
\end{array}\right]
$$

and $K_{2}$ is defined in (9.32). Here are details of choosing the observer gains such that (9.38) holds. The condition (9.38) is expanded as

$$
\begin{aligned}
& {\left[\begin{array}{l}
-k_{11}-k_{12} \\
-k_{21}-k_{22}
\end{array}\right]<0,} \\
& {\left[\begin{array}{cc}
-\frac{d_{22}}{m_{22}}+\frac{k_{13}^{2} k_{22}}{k_{11} k_{22}-k_{12} k_{21}} & -\frac{m_{11} u}{m_{22}}-\frac{k_{13} k_{12} k_{24}}{k_{11} k_{22}-k_{12} k_{21}} \\
\frac{\left(m_{11}-m_{22}\right) u}{m_{33}}-\frac{k_{24} k_{21} k_{13}}{k_{11} k_{22}-k_{12} k_{21}} & -\frac{d_{33}}{m_{33}}+\frac{k_{24}^{2} k_{11}}{k_{11} k_{22}-k_{12} k_{21}}
\end{array}\right]<0 .}
\end{aligned}
$$

From (9.40), it suffices that

$$
\begin{align*}
& \frac{d_{22}}{m_{22}}-\frac{k_{13}^{2} k_{22}}{k_{11} k_{22}-k_{12} k_{21}}>0 \\
& \frac{d_{33}}{m_{33}}-\frac{k_{24}^{2} k_{11}}{k_{11} k_{22}-k_{12} k_{21}}>0, \\
& \frac{m_{11} u}{m_{22}}=-\frac{k_{13} k_{12} k_{24}}{k_{11} k_{22}-k_{12} k_{21}},  \tag{9.41}\\
& \frac{\left(m_{11}-m_{22}\right) u}{m_{33}}=\frac{k_{13} k_{21} k_{24}}{k_{11} k_{22}-k_{12} k_{21}}, \\
& k_{11}>0, k_{22}>0 \\
& k_{11} k_{22}-k_{12} k_{21}>0 .
\end{align*}
$$

For simplicity, we choose

$$
\begin{align*}
& k_{13}=k_{24}=\rho \sqrt{u}, \\
& k_{11} k_{22}-k_{12} k_{21}=\rho, \tag{9.42}
\end{align*}
$$

where $\rho>0$ is to be selected later. Substituting (9.42) into (9.41) yields

$$
0<k_{22}<\frac{d_{22}}{\rho u_{\max } m_{22}}, 0<k_{11}<\frac{d_{33}}{\rho u_{\max } m_{33}}
$$

$$
\begin{align*}
& k_{12}=-\frac{m_{11}}{\rho m_{22}}, k_{21}=\frac{\left(m_{11}-m_{22}\right)}{\rho m_{33}} \\
& k_{11} k_{22}>-\frac{m_{11}}{\rho^{2} m_{22}} \frac{\left(m_{11}-m_{22}\right)}{m_{33}} \tag{9.43}
\end{align*}
$$

Hence, under Assumption 9.1, we can always pick a suitable constant $\rho>0$ such that (9.43) holds. In summary, the observer gains $k_{11}, k_{12}, k_{13}, k_{21}, k_{22}, k_{24}, k_{31}$, and $k_{42}$ are chosen such that (9.42) and (9.43) hold.

We choose $A$ and $\mu_{0}$ such that $p_{0}>0$. Hence (9.36) and (9.37) yield

$$
\begin{equation*}
\|\tilde{x}(t)\| \leq\left\|\tilde{x}\left(t_{0}\right)\right\| e^{-\eta\left(t-t_{0}\right)}+\eta_{0}, \forall t \geq t_{0} \geq 0 \tag{9.44}
\end{equation*}
$$

with $\eta_{0}=\sqrt{q_{0} / p_{0}}$ and $\eta=p_{0}$. When there are no environmental disturbances, we have $\eta_{0}=0$. The observer error dynamics (9.35) is thus GES at the origin.

### 9.3.2 Control Design

We use the coordinate transformation (9.2) to rewrite the ship dynamics (9.3) in conjunction with (9.34) as follows

$$
\begin{align*}
\dot{y}= & -\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{\hat{v}}{\sqrt{1+(k y)^{2}}}+\frac{u\left(\sin \left(z_{1}\right)-k y\left(\cos \left(z_{1}\right)-1\right)\right)}{\sqrt{1+(k y)^{2}}}+ \\
& \frac{\hat{v}\left(\left(\cos \left(z_{1}\right)-1\right)+\sin \left(z_{1}\right) k y\right)}{\sqrt{1+(k y)^{2}}}+\frac{\tilde{v}}{\sqrt{1+(k y)^{2}}}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right), \\
\dot{z}_{1}= & \hat{r}+\tilde{r}-\frac{k^{2} u y-k \hat{v}}{\left(1+(k y)^{2}\right)^{3 / 2}}+\frac{k u\left(\sin \left(z_{1}\right)-k y\left(\cos \left(z_{1}\right)-1\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}+ \\
& \frac{k \hat{v}\left(\left(\cos \left(z_{1}\right)-1\right)+\sin \left(z_{1}\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}+\frac{k \tilde{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}, \\
\dot{\hat{v}}= & -\frac{m_{11} u}{m_{22}} \hat{r}-\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}+\left(k_{31}+k_{13}+\cos (\psi)\right) \tilde{y}, \\
\dot{\hat{r}}= & \frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \hat{v}-\frac{d_{33}}{m_{33}} \hat{r}-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} \hat{r}+\frac{1}{m_{33}} \tau_{r}+\left(k_{42}+k_{24}+1\right) \tilde{\psi} . \tag{9.45}
\end{align*}
$$

Similarly to the full state feedback case, we design the control law $\tau_{r}$ in two steps.

## Step 1

Define

$$
\begin{equation*}
z_{2}=\hat{r}-\hat{r}_{d} \tag{9.46}
\end{equation*}
$$

where $\hat{r}_{d}$ is an intermediate control designed as

$$
\begin{equation*}
\hat{r}_{d}=-k_{1} z_{1}-\frac{k u\left(\sin \left(z_{1}\right)-k y \cos \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k \hat{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}} \tag{9.47}
\end{equation*}
$$

with $k_{1}$ being a positive constant to be selected later.

## Step 2

The first time derivative of (9.46) along the solutions of the last equation of (9.45) together with (9.47) is

$$
\begin{align*}
\dot{z}_{2}= & \frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \hat{v}-\frac{d_{33}}{m_{33}} \hat{r}-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} \hat{r}+\frac{1}{m_{33}} \tau_{r}+ \\
& \left(k_{42}+k_{24}+1\right) \tilde{\psi}-\frac{\partial \hat{r}_{d}}{\partial u} \dot{u}-\frac{\partial \hat{r}_{d}}{\partial z_{1}}\left(-k_{1} z_{1}+z_{2}\right)- \\
& \frac{\partial \hat{r}_{d}}{\partial y}\left(\frac{u\left(\sin \left(z_{1}\right)-k y \cos \left(z_{1}\right)\right)}{\sqrt{1+(k y)^{2}}}+\frac{\hat{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\sqrt{1+(k y)^{2}}}\right)- \\
& \frac{\partial \hat{r}_{d}}{\partial \hat{v}}\left(-\frac{m_{11} u}{m_{22}} \hat{r}-\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}\right)-  \tag{9.48}\\
& \frac{\partial \hat{r}_{d}}{\partial z_{1}}\left(\frac{k}{\left(1+(k y)^{2}\right)^{3 / 2}}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}+\tilde{r}\right)- \\
& \frac{\partial \hat{r}_{d}}{\partial y} \frac{1}{\sqrt{1+(k y)^{2}}}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}-\frac{\partial \hat{r}_{d}}{\partial \hat{v}}\left(k_{31}+k_{13}+\cos (\psi)\right) \tilde{y},
\end{align*}
$$

where

$$
\begin{align*}
\frac{\partial \hat{r}_{d}}{\partial u} & =-\frac{k\left(-\cos \left(z_{1}\right) k y+\sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}, \frac{\partial \hat{r}_{d}}{\partial \hat{v}}=-\frac{k\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}} \\
\frac{\partial \hat{r}_{d}}{\partial z_{1}} & =-k_{1}-\frac{k u\left(\cos \left(z_{1}\right)+k y \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k \hat{v}\left(-\sin \left(z_{1}\right)+k y \cos \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}} \\
\frac{\partial \hat{r}_{d}}{\partial y} & =-\frac{3 k^{3} y\left(k u y-u\left(\sin \left(z_{1}\right)-(y+\hat{v})\left(\cos \left(z_{1}\right)-1\right)\right)-\hat{v} y \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{5 / 2}}+ \\
& \frac{k^{2}\left(u \cos \left(z_{1}\right)-\hat{v} \sin \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}} . \tag{9.49}
\end{align*}
$$

We now choose the actual control without canceling the useful damping terms as

$$
\begin{align*}
\tau_{r}= & m_{33}\left(-z_{1}-k_{2} z_{2}-\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \hat{v}+\frac{d_{33}}{m_{33}} \hat{r}_{d}+\right. \\
& \sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} \hat{r}_{d}+\frac{\partial \hat{r}_{d}}{\partial u} \dot{u}+\frac{\partial \hat{r}_{d}}{\partial z_{1}}\left(-k_{1} z_{1}+z_{2}\right)+ \\
& \frac{\partial \hat{r}_{d}}{\partial y}\left(\frac{u\left(\sin \left(z_{1}\right)-k y \cos \left(z_{1}\right)\right)}{\sqrt{1+(k y)^{2}}}+\frac{\hat{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\sqrt{1+(k y)^{2}}}\right)+ \\
& \left.\frac{\partial \hat{r}_{d}}{\partial \hat{v}}\left(-\frac{m_{11} u}{m_{22}} \hat{r}-\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}\right)\right), \tag{9.50}
\end{align*}
$$

where $k_{2}$ is a positive constant to be chosen later. Substituting (9.46), (9.47), and (9.50) into (9.45) results in the following closed loop system:

$$
\begin{align*}
\dot{y}= & -\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{\hat{v}}{\sqrt{1+(k y)^{2}}}+\frac{u\left(\sin \left(z_{1}\right)-k y\left(\cos \left(z_{1}\right)-1\right)\right)}{\sqrt{1+(k y)^{2}}}+ \\
& \frac{\hat{v}\left(\left(\cos \left(z_{1}\right)-1\right)+\sin \left(z_{1}\right) k y\right)}{\sqrt{1+(k y)^{2}}}+\frac{\tilde{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\sqrt{1+(k y)^{2}}}, \\
\dot{\hat{v}}= & -\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}-\frac{m_{11} u}{m_{22}}\left(-k_{1} z_{1}+z_{2}-\right. \\
& \left.\frac{k u\left(\sin \left(z_{1}\right)-k y \cos \left(z_{1}\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}-\frac{k \hat{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}\right)+ \\
& \left(k_{31}+k_{13}+\cos (\psi)\right) \tilde{y}, \\
\dot{z}_{1}= & -k_{1} z_{1}+z_{2}+\frac{k\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}}{\left(1+(k y)^{2}\right)^{3 / 2}}+\tilde{r}, \\
\dot{z}_{2}= & -z_{1}-k_{2} z_{2}-\frac{d_{33}}{m_{33}} z_{2}-\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} z_{2}- \\
& \frac{\partial \hat{r}_{d}}{\partial z_{1}}\left(\frac{k\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}}{\left(1+(k y)^{2}\right)^{3 / 2}}+\tilde{r}\right)+\left(k_{42}+k_{24}+1\right) \tilde{\psi}- \\
& \frac{\partial \hat{r}_{d}}{\partial \hat{v}}\left(k_{31}+k_{13}+\cos (\psi)\right) \tilde{y}-\frac{\partial \hat{r}_{d}}{\partial y} \frac{1\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}}{\sqrt{1+(k y)^{2}}} . \tag{9.51}
\end{align*}
$$

### 9.3.3 Stability Analysis

It is not difficult to show that the closed loop system (9.51) is forward complete. We now use Lemmas 9.1 and 9.2 to prove that the closed loop (9.51) is globally ultimately stable. From (9.49), it can be seen that the closed loop (9.51) is different from (9.9) since $\hat{v}$ enters the $\left(z_{1}, z_{2}\right)$-subsystem. To remove this obstacle, we first
prove that $\hat{v}$ is bounded. We then prove the convergence of $\left(z_{1}, z_{2}\right)$ and finally $\hat{v}$ and $y$.

## Boundedness of $\hat{v}$

To prove that $\hat{v}$ is bounded, we view the last three equations of (9.51) as the system studied in Lemma 9.1 with $x_{1}=\left[\begin{array}{lll}\hat{v} & z_{1} & z_{2}\end{array}\right]^{T}$ as $x, \tilde{x}$ as $\xi(t)$ and

$$
\begin{equation*}
\dot{x}_{1}=f_{1}\left(t, x_{1}\right)+g_{1}\left(t, x_{1}, \tilde{x}\right) \tag{9.52}
\end{equation*}
$$

where

$$
\begin{align*}
& f_{1}\left(t, x_{1}\right)=\left[\begin{array}{c}
\Omega_{1} \\
-k_{1} z_{1}+z_{2} \\
-z_{1}-\left(k_{2}+\frac{d_{33}}{m_{33}}+\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1}\right) z_{2}
\end{array}\right], \\
& g_{1}\left(t, x_{1}, \tilde{x}\right)=\left[\begin{array}{c}
\left(k_{31}+k_{13}+\cos (\psi)\right) \tilde{y} \\
\frac{k \tilde{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}+\tilde{r} \\
\Omega_{2}
\end{array}\right] \tag{9.53}
\end{align*}
$$

with

$$
\begin{aligned}
\Omega_{1}= & -\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}-\frac{m_{11} u}{m_{22}} \times \\
& \left(-k_{1} z_{1}+z_{2}-\frac{k\left(u\left(\sin \left(z_{1}\right)-k y \cos \left(z_{1}\right)\right)+\hat{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}\right), \\
\Omega_{2}= & \left(k_{42}+k_{24}+1\right) \tilde{\psi}-\frac{\partial \hat{r}_{d}}{\partial z_{1}}\left(\frac{k\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}}{\left(1+(k y)^{2}\right)^{3 / 2}}+\tilde{r}\right)- \\
& \frac{\partial \hat{r}_{d}}{\partial \hat{v}}\left(k_{31}+k_{13}+\cos (\psi)\right) \tilde{y}-\frac{\partial \hat{r}_{d}}{\partial y} \frac{\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}}{\sqrt{1+(k y)^{2}}}
\end{aligned}
$$

where again with abuse of notation, $\hat{r}$ is considered as a function of time. We now need to verify all of the conditions of Lemma 9.1.

Verifying Condition C1. We take the following Lyapunov function:

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(\hat{v}^{2}+\delta_{1}\left(z_{1}^{2}+z_{2}^{2}\right)\right) \tag{9.54}
\end{equation*}
$$

where $\delta_{1}$ is a positive constant. The first time derivative of (9.54) along the solutions of the differential equation $\dot{x}_{1}=f_{1}\left(t, x_{1}\right)$, see (9.52) and (9.53), satisfies

$$
\begin{align*}
\dot{V}_{1}= & -\frac{d_{22}}{m_{22}} \hat{v}^{2}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}^{2}-\frac{m_{11} u \hat{v}}{m_{22}}\left(-k_{1} z_{1}+z_{2}-\right. \\
& \left.\frac{k\left(u\left(\sin \left(z_{1}\right)-k y \cos \left(z_{1}\right)\right)+\hat{v}\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}\right)- \\
& \delta_{1} k_{1} z_{1}^{2}-\delta_{1} k_{2} z_{2}^{2}-\delta_{1} \frac{d_{33}}{m_{33}} z_{2}^{2}-\delta_{1} \sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} z_{2}^{2} \\
\leq & -\left(\frac{d_{22}}{m_{22}}-\frac{m_{11} u_{\max }}{m_{22}}\left(\mu_{1}+k_{1} \mu_{1}+4 k+2 k \mu_{1} u_{m x}\right)\right) \hat{v}^{2}- \\
& \left(\delta_{1} k_{1}-\frac{m_{11} k_{1} u_{\max }}{4 \mu_{1} m_{22}}\right) z_{1}^{2}-\left(\delta_{1} k_{2}-\frac{m_{11} u_{\max }}{4 \mu_{1} m_{22}}\right) z_{2}^{2}+\frac{5 m_{11} k u_{\max }^{2}}{2 m_{22}} . \tag{9.55}
\end{align*}
$$

Hence the condition C1 of Lemma 9.1 is verified with

$$
\begin{align*}
c_{0}= & \frac{5 m_{11} k u_{\max }^{2}}{2 m_{22}}, c_{1}=\frac{1}{2} \min \left(1, \delta_{1}\right), \\
c_{2}= & \frac{1}{2} \max \left(1, \delta_{1}\right), c_{3}=\max \left(1, \delta_{1}\right), \\
c_{4}= & \min \left(\left(\frac{d_{22}}{m_{22}}-\frac{m_{11} u_{\max }}{m_{22}}\left(\mu_{1}+k_{1} \mu_{1}+4 k+2 k \mu_{1} u_{m x}\right)\right),\right. \\
& \left.\left(\delta_{1} k_{1}-\frac{m_{11} k_{1} u_{\max }}{4 \mu_{1} m_{22}}\right),\left(\delta_{1} k_{2}-\frac{m_{11} u_{\max }}{4 \mu_{1} m_{22}}\right)\right), \tag{9.56}
\end{align*}
$$

where $\mu_{1}>0$ and $k>0$ are chosen such that $c_{4}>0$.
Verifying Condition C2. To verify this condition of Lemma 9.1, we note from (9.49) that

$$
\begin{align*}
& \left|\frac{\partial \hat{r}_{d}}{\partial z_{1}}\right| \leq k_{1}+2 k\left(u_{\max }+|\hat{v}|\right) \\
& \left|\frac{\partial \hat{r}_{d}}{\partial \hat{v}}\right| \leq 2 k,\left|\frac{\partial \hat{r}_{d}}{\partial y}\right| \leq 3 k^{2}\left(u_{\max }+3 k u_{\max }+3 k|\hat{v}|\right) \tag{9.57}
\end{align*}
$$

From (9.56) and (9.57), a simple calculation shows that the condition C2 of Lemma 9.1 is satisfied with

$$
\begin{align*}
\lambda_{1}= & (2 k+1)\left(k_{31}+k_{13}+1\right)+k_{42}+k_{24}+1+ \\
& (2 k+1)\left(k_{1}+2 k u_{\max }+1\right)+6 k^{2} u_{\max }(3 k+1) \\
\lambda_{2}= & 2 k(2 k+1)+18 k^{3} . \tag{9.58}
\end{align*}
$$

Verifying Condition C3. This condition follows directly from (9.44).
Verifying Condition C4. It can be shown that there exists a positive constant $k$ such that the condition C4 of Lemma 9.1 satisfies

$$
\begin{equation*}
c_{4}-\lambda_{2} c_{3} \varepsilon_{0}-\frac{\lambda_{1} c_{3} \varepsilon_{0}}{4 \mu_{0}}>0 \tag{9.59}
\end{equation*}
$$

where $c_{4}, c_{3}, \varepsilon_{0}, \lambda_{1}$, and $\lambda_{2}$ given in (9.44), (9.56) and (9.58).
All of the conditions of Lemma 9.1 have been verified, therefore we have

$$
\begin{equation*}
|\hat{v}| \leq\left\|x_{1}(t)\right\| \leq \alpha_{1}\left(\left\|\left(x_{1}\left(t_{0}\right), \tilde{x}\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{1}\left(t-t_{0}\right)}+\varepsilon_{1}, \forall t \geq t_{0} \geq 0 \tag{9.60}
\end{equation*}
$$

where $\alpha_{1}, \sigma_{1}$, and $\varepsilon_{1}$ are in the form of $\alpha, \sigma$, and $\varepsilon$ in Lemma 9.1 with all constants given in (9.44), (9.56), and (9.58).

## ( $z_{1}, z_{2}$ )-subsystem

Having proven that $\hat{v}$ is bounded in the previous section, we now apply Lemma 9.1 to the $\left(z_{1}, z_{2}\right)$-subsystem. It is clear that the last two equations of (9.51) are in the form of the system studied in Lemma 9.1 with $z=\left[\begin{array}{ll}z_{1} & z_{2}\end{array}\right]^{T}$ as $x, \tilde{x}$ as $\xi(t)$ and

$$
\begin{equation*}
\dot{z}=f_{z}(t, z)+g_{z}(t, y, z, \tilde{x}), \tag{9.61}
\end{equation*}
$$

where

$$
\begin{gather*}
f_{z}(t, z)=\left[\begin{array}{c}
-k_{1} z_{1}+z_{2} \\
-z_{1}-\left(k_{2}+\frac{d_{33}}{m_{33}}+\sum_{i \geq 2} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1}\right) z_{2}
\end{array}\right], \\
g_{z}(t, y, z, \tilde{x})=\left[\begin{array}{c}
\frac{k\left(\cos \left(z_{1}\right)+\sin \left(z_{1}\right) k y\right) \tilde{v}}{\left(1+(k y)^{2}\right)^{3 / 2}}+\tilde{r} \\
\Omega_{2}
\end{array}\right] \tag{9.62}
\end{gather*}
$$

Proceeding with the same steps as in the previous section, it is shown that all the conditions of Lemma 9.1 hold with the Lyapunov function $V_{2}=0.5\left(z_{1}^{2}+z_{2}^{2}\right)$ and

$$
\begin{aligned}
& c_{0}=0, c_{1}=c_{2}=0.5, c_{3}=1, c_{4}=\min \left(k_{1}, k_{2}\right), \\
& \lambda_{1}=2 k\left(k_{13}+k_{31}+1\right)+(2 k+1)\left(2+2 k u_{\max }\right)+k_{42}+
\end{aligned}
$$

$$
\begin{align*}
& k_{24}+1+6 k^{2} u_{\max }(1+3 k)+\left(18 k^{3}+2 k(2 k+1)\right)\left(\alpha_{1}+\varepsilon_{1}\right), \\
\lambda_{2}= & 0 \tag{9.63}
\end{align*}
$$

where $\alpha_{1}$ and $\varepsilon_{1}$ are given in (9.60). The condition C 4 of Lemma 9.1 becomes

$$
\begin{equation*}
c_{4}-\frac{\lambda_{1} c_{3} \varepsilon_{0}}{4 \mu_{0}}>0 \tag{9.64}
\end{equation*}
$$

where $c_{4}, c_{3}, \varepsilon_{0}$, and $\lambda_{2}$ are calculated from in (9.44) and (9.63).
All of the conditions of Lemma 9.1 have been verified, therefore we have

$$
\begin{equation*}
\|z(t)\| \leq \alpha_{2}\left(\left\|\left(z\left(t_{0}\right), \tilde{x}\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{2}\left(t-t_{0}\right)}+\varepsilon_{2}, \forall t \geq t_{0} \geq 0 \tag{9.65}
\end{equation*}
$$

where $\alpha_{2}, \sigma_{2}$, and $\varepsilon_{2}$ are in the form of $\alpha, \sigma$, and $\varepsilon$ in Lemma 9.1 with all constants given in (9.63).

## $(y, \hat{v})$-subsystem

It can be seen that the first two equations of (9.51) are in the form of the system studied in Lemma 9.2, i.e., $x_{3}=\left[\begin{array}{ll}\hat{v} & y\end{array}\right]^{T}, \tilde{x}_{3}=\left[\begin{array}{lll}z_{1} & z_{2} & \tilde{y} \\ \tilde{v}\end{array}\right]^{T}$, and

$$
\begin{equation*}
\dot{x}_{3}=f_{3}\left(t, x_{3}\right)+g_{3}\left(t, x_{3}, \tilde{x}_{3}\right) \tag{9.66}
\end{equation*}
$$

where

$$
\begin{align*}
& f_{3}\left(t, x_{3}\right)=\left[\begin{array}{c}
-\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i \geq 2} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}-\frac{m_{11} u}{m_{22}} \frac{k^{2} u y-k \hat{v}}{\left(1+(k y)^{2}\right)^{3 / 2}} \\
-\frac{k u y}{\sqrt{1+(k y)^{2}}}+\frac{\hat{v}}{\sqrt{1+(k y)^{2}}}
\end{array}\right], \\
& g_{3}\left(t, x_{3}, \tilde{x}_{3}\right)=\left[\begin{array}{c}
\Omega_{31} \\
\Omega_{32}
\end{array}\right], \tag{9.67}
\end{align*}
$$

with

$$
\begin{aligned}
\Omega_{31}= & -\frac{m_{11} u}{m_{22}}\left(-k_{1} z_{1}+z_{2}-\frac{k u\left(\sin \left(z_{1}\right)-k y\left(\cos \left(z_{1}\right)-1\right)\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}-\right. \\
& \left.\frac{k \hat{v}\left(\left(\cos \left(z_{1}\right)-1\right)+\sin \left(z_{1}\right) k y\right)}{\left(1+(k y)^{2}\right)^{3 / 2}}\right)+\left(k_{31}+k_{13}+\cos (\psi)\right) \tilde{y}, \\
\Omega_{32}= & \frac{u\left(\sin \left(z_{1}\right)-k y\left(\cos \left(z_{1}\right)-1\right)\right)}{\sqrt{1+(k y)^{2}}}+\frac{\hat{v}\left(\left(\cos \left(z_{1}\right)-1\right)+\sin \left(z_{1}\right) k y\right)}{\sqrt{1+(k y)^{2}}}+ \\
& \frac{\left(\cos \left(z_{1}\right)-\sin \left(z_{1}\right) k y\right) \tilde{v}}{\sqrt{1+(k y)^{2}}} .
\end{aligned}
$$

We now need to verify all conditions of Lemma 9.2 for the system (9.66).
Verifying Condition C1. To verify this condition of Lemma 9.2, we take the following proper function

$$
\begin{equation*}
V_{3}=\frac{1}{2}\left(\hat{v}^{2}+y^{2}\right), \tag{9.68}
\end{equation*}
$$

whose time derivative along (9.67) satisfies

$$
\begin{equation*}
\dot{V}_{3} \leq-c_{31} \hat{v}^{2}-c_{32} \frac{y^{2}}{\sqrt{1+c_{4} y^{2}}} \tag{9.69}
\end{equation*}
$$

where

$$
\begin{align*}
& c_{31}=\frac{d_{22}}{m_{22}}-\frac{m_{11} k u_{\max }}{m_{22}}-\mu_{2}-\frac{m_{11} k^{2} u_{\max }^{2}}{m_{22}} \mu_{3} \\
& c_{32}=k u_{\min }-\frac{1}{4 \mu_{2}}-\frac{1}{4 \mu_{3}} \frac{m_{11} k^{2} u_{\max }^{2}}{m_{22}} \tag{9.70}
\end{align*}
$$

with $\mu_{2}>0$ and $\mu_{3}>0$ chosen such that $c_{31}>0$ and $c_{32}>0$.
From (9.67) and (9.68), it is easy to show that

$$
\begin{equation*}
\left|\frac{\partial V_{3}}{\partial x_{3}} g_{3}\left(t, x_{3}, \tilde{x}_{3}\right)\right| \leq\left(\lambda_{0}+\lambda_{1} \hat{v}^{2}+\lambda_{2} \frac{y^{2}}{\sqrt{1+c_{4} y^{2}}}\right)\left\|\tilde{x}_{3}\right\| \tag{9.71}
\end{equation*}
$$

with

$$
\begin{align*}
\lambda_{0}= & \frac{1}{4 \mu_{4}}\left(\frac{m_{11} u_{\max }}{m_{22}}\left(k_{1}+1+2 k u_{\max }\right)+\left(k_{31}+k_{13}+1\right)+\right. \\
& \left.u_{\max }+1+\left(\alpha_{1}+\varepsilon_{1}\right)^{2}\right) \\
\lambda_{1}= & \frac{m_{11} u_{\max } \mu_{4}}{m_{22}}\left(k_{1}+1+2 k u_{\max }\right)+\frac{2 m_{11} k u_{\max }}{m_{22}}+\mu_{4}\left(k_{31}+k_{13}+1\right), \\
\lambda_{2}= & \left(k+\mu_{4}\right) u_{\max }+2 \mu_{4}+k\left(\alpha_{1}+\varepsilon_{1}+\alpha_{2}+\varepsilon_{2}\right) \tag{9.72}
\end{align*}
$$

where $\mu_{4}>0, \alpha_{1}$, and $\varepsilon_{1}$ are given in (9.60), and $\alpha_{2}$ and $\varepsilon_{2}$ are given in (9.65).
Verifying Condition C2. To verify this condition, we note that

$$
\begin{equation*}
\left\|\tilde{x}_{3}(t)\right\| \leq\left\|\binom{z_{1}(t)}{z_{2}(t)}\right\|+\left\|\binom{\tilde{v}(t)}{\tilde{y}(t)}\right\| . \tag{9.73}
\end{equation*}
$$

Therefore we can write (9.73) from (9.44) and (9.65) as

$$
\begin{equation*}
\left\|\tilde{x}_{3}(t)\right\| \leq \alpha_{3}\left(\left\|\left(z\left(t_{0}\right), \tilde{x}\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{3}\left(t-t_{0}\right)}+\varepsilon_{3}, \tag{9.74}
\end{equation*}
$$

where

$$
\begin{align*}
& \alpha_{3}\left(\left\|\left(z\left(t_{0}\right), \tilde{x}\left(t_{0}\right)\right)\right\|\right)=\left\|\left(z\left(t_{0}\right), \tilde{x}\left(t_{0}\right)\right)\right\|+\alpha_{2}\left(\left\|\left(z\left(t_{0}\right), \tilde{x}\left(t_{0}\right)\right)\right\|\right),  \tag{9.75}\\
& \sigma_{3}=\min \left(\eta, \sigma_{2}\right), \varepsilon_{3}=\eta_{0}+\varepsilon_{2},
\end{align*}
$$

with $\alpha_{2}$ and $\varepsilon_{2}$ given in (9.65), $\eta_{0}$ and $\eta$ given in (9.44).
Verifying Condition C3. This condition is satisfied if

$$
\begin{equation*}
c_{31}-\lambda_{1} \varepsilon_{3}>0 \text { and } c_{32}-\lambda_{2} \varepsilon_{3}>0 \tag{9.76}
\end{equation*}
$$

where $c_{31}, c_{32}, \lambda_{1}, \lambda_{2}$, and $\varepsilon_{3}$ are given in (9.72) and (9.75). After some lengthy but simple calculation, it can be shown that, under the assumption of small enough environmental disturbances, the condition (9.76) holds for a suitable choice of the observer gains $k_{11}, k_{12}, k_{13}, k_{21}, k_{22}, k_{24}, k_{31}$, and $k_{42}$, and the control gains $k$, $k_{1}$, and $k_{2}$.

Verifying Condition C4. From the boundedness of the sway velocity estimate, $\hat{v}(t)$, proven in the previous section and noting that $\varepsilon_{2}$ in (9.65) can be made arbitrarily small, it is directly shown that there exists a nondecreasing function $\varpi$ of $\left\|\left(\left(v\left(t_{0}\right), y\left(t_{0}\right)\right), z\left(t_{0}\right)\right)\right\|$ such that $|y(t)| \leq \varpi$ by applying Lemma 9.1 to the first equation of (9.51) with the Lyapunov function $V_{y}=0.5 y^{2}$.

All of the conditions of Lemma 9.2 have been verified, the closed loop (9.51) is globally ultimately stable, i.e.,

$$
\begin{equation*}
\left\|x_{3}(t)\right\| \leq \alpha_{4}\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{4}\left(t-t_{0}\right)}+\varepsilon_{4}, \forall t \geq t_{0} \geq 0 \tag{9.77}
\end{equation*}
$$

where $\alpha_{4}, \sigma_{4}$, and $\varepsilon_{4}$ are calculated as in Lemma 9.2.
It is noted that when there are no environmental disturbances, $\varepsilon=0$. Therefore the closed loop (9.51) is GAS. We note that the convergence of $z_{1}$ and $z_{2}$ implies the convergence of $\hat{r}$ and $\psi$. The convergence of $v$ and $r$ results from that of $\hat{v}$ and $\hat{r}$ due to the global exponential property of the observer. We have thus proven the second main result of this chapter.

Theorem 9.2. Under Assumption 9.1, the output feedback control problem stated in Section 9.1 is solved by the observer (9.34) and the control law (9.50) as long as the observer gains $k_{11}, k_{12}, k_{13}, k_{21}, k_{22}, k_{24}, k_{13}, k_{31}$, and $k_{42}$, and the control gains $k, k_{1}$, and $k_{2}$ are chosen such that (9.64), (9.70), (9.76), and (9.43) hold.

Remark 9.2. Due to underactuaction and nonzero-mean environmental disturbances in the sway dynamic, our controller is only able to force the sway and its velocity to converge to a ball centered at the origin. The radius of this ball cannot be made arbitrarily small. This phenomenon should not be surprising since there is no control force in the sway direction. In addition, the yaw angle cannot be made arbitrarily small due to the effect of the sway. In fact, to guarantee the sway displacement bounded under nonzero-mean environmental disturbances acting on the sway dynamics, our controller forces the heading angle to a small value. This value together with the forward speed will prevent the sway from growing unbounded.

Remark 9.3. The choice of $k$ depends on the ship parameters and forward speed, which coincides with the steering practice of a helmsman. The helmsman uses the ship's course angle to steer the ship toward the straight line rather than use the sway velocity, which will cause the ship to glide sideways. Furthermore, the design constant $k$ is reduced when the ship forward speed is large, see (9.2), (9.23), (9.30), (9.56), and (9.59), otherwise the ship will miss the point on the straight line and slide in the sway direction.

Remark 9.4. By setting the value of $k$ equal to zero, our proposed controller reduces to a course-keeping controller. In this case, the heading angle can be made arbitrarily small. However the sway will grow linearly unbounded under nonzero-mean environmental disturbances, see Figures 9.3 and 9.6.

### 9.4 Simulations

This section validates the control laws (9.8) and (9.50) for both cases of state and output feedback on a monohull ship with the parameters given in Section 5.4. The ship surge velocity is chosen as $u=10+0.5 \sin (3 t) \mathrm{ms}^{-1}$. The environmental disturbances $\tau_{w v}(t)$ and $\tau_{w r}(t)$ are taken as $\tau_{w v}(t)=10^{5} \times 0.5 \times(1+\operatorname{rand}(\cdot))$ and $\tau_{w r}(t)=1.5 \times 10^{7} \times \operatorname{rand}(\cdot)$, with $\operatorname{rand}(\cdot)$ being zero mean random noise with the uniform distribution on the interval [ -0.50 .5 ]. We run simulations for both state feedback and output feedback cases.

### 9.4.1 State Feedback Simulation Results

The control design parameters are chosen as $k=0.05, k_{1}=0.2, k_{2}=0.5$, and $\rho_{1}=\rho_{2}=0.05$. It can be directly verified that this choice satisfies all the conditions stated in Theorem 9.1. The initial values are

$$
[y(0), v(0), \psi(0), r(0)]=\left[15 \mathrm{~m}, 0.2 \mathrm{~ms}^{-1},-0.5 \mathrm{rad}, 0.1 \mathrm{rads}^{-1}\right]
$$

Simulation results are plotted in Figure 9.1 for the case without disturbances. In this case, it can be seen that all sway displacement, sway velocity, and yaw angle converge to zero as desired. The large control effort is due to the fact that we simulate our controllers on a real surface ship but it is within the limit of the maximum yaw moment. For the case with disturbances, simulation results are plotted in Figure 9.2. In this case, all the states converge to a ball centered at the origin as proven in Theorem 9.1. To illustrate Remark 9.4, we simulate our controller with the design constant $k=0$. The simulation results for this case are given in Figure 9.3. The sway displacement $y$ grows linearly unbounded due to nonvanishing environmental disturbances. It should be noted that all of the course-keeping controllers, see for


Figure 9.1 State feedback control results without disturbances: a. Sway displacement $y$; b. Heading angle $\psi ; \mathbf{c}$. Sway velocity $v ; \mathbf{d}$. Yaw velocity $r ; \mathbf{e}$. Yaw moment $\tau_{r}$
example [12], which do not take the sway displacement into account, will result in similar unboundedness of the sway that was pointed out in Remark 9.4.

### 9.4.2 Output Feedback Simulation Results

The control design parameters are chosen as $k=0.05, k_{1}=0.2, k_{2}=0.5$, and $\rho_{1}=\rho_{2}=0.05$. The observer gains are selected as $k_{11}=k_{22}=2, k_{12}=-\frac{m_{11}}{\rho m_{22}}$, $k_{21}=\frac{m_{11}-m_{22}}{\rho m_{33}}, k_{31}=k_{13}=k_{24}=k_{42}=\rho \sqrt{u}$, and $\rho=0.015$. A calculation shows that this choice satisfies all the conditions stated in Theorem 9.2. The initial values are

$$
\begin{aligned}
& {[y(0), v(0), \psi(0), r(0)]=\left[15 \mathrm{~m}, 0.2 \mathrm{~ms}^{-1},-0.5 \mathrm{rad}, 0.1 \mathrm{rads}^{-1}\right],} \\
& {[\hat{y}(0), \hat{v}(0), \hat{\psi}(0), \hat{r}(0)]=\left[10 \mathrm{~m}, 0 \mathrm{~ms}^{-1},-0.2 \mathrm{rad}, 0.2 \mathrm{rads}^{-1}\right] .}
\end{aligned}
$$

Simulation results are plotted in Figure 9.4 for the case without disturbances and in Figure 9.5 for the case with disturbances. From Figure 9.4, it is seen that all sway displacement, sway velocity, and yaw angle converge to zero asymptotically. It is also observed that the observer states (dotted lines) exponentially converge to


Figure 9.2 State feedback control results with disturbances: a. Sway displacement $\boldsymbol{y}$; b. Heading angle $\psi ; \mathbf{c}$. Sway velocity $v ; \mathbf{d}$. Yaw velocity $r$; e. Yaw moment $\tau_{r}$
their unknown estimated ones (solid lines). For the case with disturbances, all the states converge to a ball centered at the origin as proven in Theorem 9.2. The simulation results with the design constant $k=0$ are plotted in Figure 9.6. Again, the sway displacement $y$ grows linearly unbounded due to nonvanishing environmental disturbances as mentioned in Remark 9.4.

### 9.5 Conclusions

The control design was based on the idea of an interaction between the ship behavior and the action of a helmsman on a linear course. Although our proposed state feedback controller has been designed by using precise knowledge of the ship parameters, we can easily change them to an adaptive version to take inaccurate knowledge of the system parameters into account, see (9.6). However, for the case of output feedback, an adaptive observer will be required, see (9.34) and (9.45). This chapter is based on [116, 126, 127].


Figure 9.3 State feedback control results with disturbances and $k=0$ : a. Sway displacement $y$; b. Heading angle $\psi$; c. Sway velocity $v$; d. Yaw velocity $r$; e. Yaw moment $\tau_{r}$


Figure 9.4 Output feedback control results without disturbances: a. Sway displacement $y$; b. Heading angle $\psi$; c. Sway velocity $v$ (solid line) and its estimate $\hat{v}$ (dotted line); d. Yaw velocity $r$ (solid line) and its estimate $\hat{r}$ (dotted line); e. Yaw moment $\tau_{r}$


Figure 9.5 Output feedback control results with disturbances: a. Sway displacement $\boldsymbol{y}$; $\mathbf{b}$. Heading angle $\psi$; c. Sway velocity $v$ (solid) and its estimate $\hat{v}$ (dot); d. Yaw velocity $r$ (solid) and its estimate $\hat{r}$ (dot); $\mathbf{e}$. Yaw moment $\tau_{r}$


Figure 9.6 Output feedback control results with disturbances and $k=0$ : a. Sway displacement $y$; b. Heading angle $\psi$; c. Sway velocity $v$ (solid line) and its estimate $\hat{v}$ (dotted line); d. Yaw velocity $\boldsymbol{r}$ (solid line) and its estimate $\hat{r}$ (dotted line); e. Yaw moment $\tau_{r}$

# Chapter 10 <br> Path-following of Underactuated Ships Using Serret-Frenet Coordinates 

This chapter presents state feedback and output feedback controllers that force an underactuated surface ship to follow a predefined path under the presence of environmental disturbances induced by waves, wind, and ocean currents. The control solutions originated an observation that it is reasonable in practice to steer a vessel such that it is on the reference path and its total velocity is tangent to the path, and that the vessel's forward speed is independently controlled by the main propulsion control system. The proposed controllers are designed using Lyapunov's direct method, the popular backstepping technique, and the Serret-Frenet frame. The unmeasured sway and yaw velocities are estimated by introducing a novel nonlinear passive observer.

### 10.1 Control Objective

In addition to the assumptions made in Section 3.4.1.1, we assume that the surge velocity is a positive constant and is independently controlled by the main propulsion control system. The resulting mathematical model of the underactuated ship moving in surge, sway, and yaw is rewritten as

$$
\begin{align*}
\dot{x} & =u \cos (\psi)-v \sin (\psi), \\
\dot{y} & =u \sin (\psi)+v \cos (\psi), \\
\dot{\psi} & =r, \\
\dot{v} & =-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t),  \tag{10.1}\\
\dot{r} & =\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t),
\end{align*}
$$

where all symbols in (10.1) are defined in Section 3.4.1.2. The positive constant terms $d_{v}, d_{r}, d_{v i}$, and $d_{r i}, i=2,3$ represent the hydrodynamic damping in surge, sway and yaw. The bounded time-varying terms, $\tau_{w v}(t)$ and $\tau_{w r}(t)$, are the environmental disturbances induced by wave, wind, and ocean current with $\left|\tau_{w v}(t)\right| \leq$ $\tau_{w v \text { max }}<\infty$ and $\left|\tau_{w r}(t)\right| \leq \tau_{w r \text { max }}<\infty$. The available control is the yaw moment $\tau_{r}$. Since the sway control force is not available in the sway dynamics, the ship model (10.1) is again underactuated.

The control objective of this chapter is to design the yaw moment $\tau_{r}$ to force the underactuated ship (10.1) to follow a specified path $\boldsymbol{\Omega}$, see Figure 10.1 , where $M$ is the orthogonal projection of the ship point $P$ on $\Omega, x_{n}$ and $x_{t}$ are the normal and the tangent unit vectors to the path at $M$, respectively. We assume that this point is uniquely defined. This assumption holds if the interior of any circle tangenting $\Omega$ at two or more points does not contain any point of the path and the distance between the ship and $\Omega$ is not too large. Furthermore, it is assumed that the radius of any osculating circle of the path is larger than or equal to $R_{\min }$ which is feasible for the ship to follow. Let the signed distance between $M$ and $P$ be $z_{e}$. Also $s$ is the signed distance along the path between some arbitrary fixed point on the path and $M, \psi_{d}$ is the angle between $x_{t}$ and $X_{b}, c(s)$ is the curvature of the path at the point $M$. We assume that $c(s)$ is uniformly bounded and differentiable. Let the total velocity of the ship be $u_{t}$ and $\psi_{e}=\psi-\psi_{d}$. The variables $s, z_{e}$, and $\psi_{e}$ form a new set of state coordinates for the ship. It can be seen that when the path $\boldsymbol{\Omega}$ coincides with the $X_{E}$-axis, the above variables coincide with the ship variables $x, y$, and $\psi$.

By using the above parameterization, it is straightforward, see [128, 129], to transform the kinematics of (10.1) to

$$
\begin{align*}
& \dot{z}_{e}=u \sin \left(\psi_{e}\right)+v \cos \left(\psi_{e}\right) \\
& \dot{\psi}_{e}=r-\frac{c(s)}{1-c(s) z_{e}}\left(u \cos \left(\psi_{e}\right)-v \sin \left(\psi_{e}\right)\right)  \tag{10.2}\\
& \dot{s}=\frac{1}{1-c(s) z_{e}}\left(u \cos \left(\psi_{e}\right)-v \sin \left(\psi_{e}\right)\right)
\end{align*}
$$

Note that the above transformation is singular when $z_{e} c(s)=1$. We first assume that

$$
\begin{equation*}
1-z_{e} c(s) \geq \delta^{*}>0 \tag{10.3}
\end{equation*}
$$

We then find the initial conditions after the controllers are designed such that this hypothesis holds. Ideally, we want the ship to be on the path and tangent to it, i.e., $z_{e}=\psi_{e}=0$.

However when the curvature of the path is different from zero, the sway velocity $v$ is also different from zero. Hence $z_{e}=\psi_{e}=0$ cannot be the equilibrium point of (10.2) with $r$ as the control input. This feature distinguishes the ship from mobile robots, see [15, 130-133] for some work on controlling mobile robots. If one designs a controller as was proposed in [129], to achieve $\psi_{e}=0$, then the error $z_{e}$ might be very large (depending on $v$ ). This phenomenon can be seen by substituting $\psi_{e}=0$ into the first equation of (10.2). Therefore, in this chapter, we formulate the


Figure 10.1 General framework of ship path-following
control problem such that $\left(z_{e}, \psi_{e}\right)=(0,-\phi)$ is the equilibrium point of (10.2), with $\phi$ being the angle between the surge velocity and the total velocity. This is desirable in practice since it guarantees that the ship is on the path and the ship's total velocity is tangent to the path. Define

$$
\begin{equation*}
\psi_{e}^{*}=\psi_{e}+\phi, \tag{10.4}
\end{equation*}
$$

where $\psi_{e}^{*}$ is referred to as the modified heading error. With (10.4), the last four equations of the ship model (10.1) with (10.2) are transformed to

$$
\begin{align*}
\dot{z}_{e}= & u_{t} \sin \left(\psi_{e}^{*}\right), \\
\dot{\psi}_{e}^{*}= & r\left(1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{u_{t}^{2}}\right)-\frac{c(s) u_{t}}{1-c(s) z_{e}} \cos \left(\psi_{e}^{*}\right)- \\
& \frac{u}{u_{t}^{2}}\left(\frac{d_{22}}{m_{22}} v+\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v-\frac{1}{m_{22}} \tau_{w v}(t)\right), \\
\dot{v}= & -\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t),  \tag{10.5}\\
\dot{r}= & \frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t),
\end{align*}
$$

where

$$
\begin{equation*}
u_{t}=\sqrt{u^{2}+v^{2}} . \tag{10.6}
\end{equation*}
$$

As a result, the path-following objective has been converted to a problem of stabilizing $\left(z_{e}, \psi_{e}^{*}\right)$ in (10.5) at the origin. We study this stabilization problem under either Assumption 10.1 or Assumption 10.2 as follows.

Assumption 10.1. The ship parameters satisfy the condition $m_{11}<m_{22}$. All of the ship states (position, orientation, and velocities) are measurable.

Assumption 10.2. The ship parameters satisfy the condition $m_{11}<m_{22}$. The sway and yaw velocities are not available for feedback.

Assumption 10.1 means a robust state feedback path-following problem, while Assumption 10.2 gives a robust output feedback path-following problem. It is noted that the condition $m_{11}<m_{22}$ always holds for surface ships since the added mass in sway is larger than that in surge. The triangular structure of (10.5) suggests that we design the actual control $\tau_{r}$ in two stages. First, we design the virtual velocity control $r$ to stabilize $z_{e}$ and $\psi_{e}^{*}$ at the origin. Based on the backstepping technique, the control $\tau_{r}$ will then be designed to make the error between the virtual velocity control and its actual values exponentially tend to a small ball centered at the origin. A new nonlinear observer is introduced to estimate the sway and yaw velocities, and an output feedback controller is then designed.

### 10.2 State Feedback

### 10.2.1 Control Design

Observing from the first equation of (10.5) that the $z_{e}$-dynamics have to be stabilized by using the angle $\psi_{e}^{*}$, we introduce the following coordinate transformation

$$
\begin{equation*}
w_{e}^{*}=\psi_{e}^{*}+\arcsin \left(\frac{k z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}\right) \tag{10.7}
\end{equation*}
$$

where $k$ is a positive constant to be selected later.
Remark 10.1. The above coordinate transformation is well defined and convergence of $w_{e}^{*}$ and $z_{e}$ implies that of $\psi_{e}^{*}$. The function arcsin is not unique. It can be replaced by some other smooth bounded functions such as tanh, arctan or an identity function. We here use the arcsin function due to its simplicity. Using (10.7) instead of a linear coordinate change such as $w_{e}^{*}=\psi_{e}^{*}+k z_{e}$, we avoid the ship whirling around when $z_{e}$ is large. In addition the nonlinear change of coordinate (10.7) will result in a "global design".

Upon an application of the coordinate transformation (10.7), the system (10.5) is rewritten as

$$
\begin{align*}
\dot{z}_{e}= & -\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}+\frac{u_{t} \sin \left(w_{e}^{*}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}-\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}\left(\cos \left(w_{e}^{*}\right)-1\right), \\
\dot{w}_{e}^{*}= & r\left(1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{u_{t}^{2}}\right)+\frac{k u_{t} \sin \left(\psi_{e}^{*}\right)}{1+\left(k z_{e}\right)^{2}}-\frac{c(s) u_{t} \cos \left(\psi_{e}^{*}\right)}{1-c(s) z_{e}}- \\
& \frac{u}{u_{t}^{2}}\left(\frac{d_{22}}{m_{22}} v+\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v-\frac{1}{m_{22}} \tau_{w v}(t)\right), \\
\dot{v}= & -\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t),  \tag{10.8}\\
\dot{r}= & \frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t),
\end{align*}
$$

where we leave the variable $\psi_{e}^{*}$ in the second equation of (10.8) for simplicity of presentation. As discussed in Remark 10.1, we will design the control $\tau_{r}$ to stabilize (10.8) at the origin.

## Step 1

Define

$$
\begin{equation*}
b=1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{u_{t}^{2}} \tag{10.9}
\end{equation*}
$$

The first condition of Assumptions 10.1 implies that

$$
\begin{equation*}
b \geq 1-\frac{m_{11}}{m_{22}}:=b^{*}>0 \tag{10.10}
\end{equation*}
$$

Introduce the yaw velocity control error $\tilde{r}$ as

$$
\begin{equation*}
\tilde{r}=r-r_{d}, \tag{10.11}
\end{equation*}
$$

where $r_{d}$ is a virtual control of the yaw velocity $r$. From the second equation of (10.8), we choose the virtual control $r_{d}$ as

$$
\begin{align*}
r_{d}= & -\frac{k_{1}}{b} w_{e}^{*}-\frac{1}{b} \frac{k u_{t} \sin \left(\psi_{e}^{*}\right)}{1+\left(k z_{e}\right)^{2}}+\frac{1}{b} \frac{c(s) u_{t} \cos \left(\psi_{e}^{*}\right)}{1-c(s) z_{e}}+ \\
& \frac{1}{b} \frac{u}{u_{t}^{2}}\left(\frac{d_{22}}{m_{22}} v+\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v\right)-  \tag{10.12}\\
& \frac{1}{b} \frac{u}{u_{t}^{2}} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{w_{e}^{*}}{\varepsilon_{1}} \frac{\tau_{w v \max }}{m_{22}}\right),
\end{align*}
$$

where $k_{1}$ and $\varepsilon_{1}$ are positive design constants to be selected later. It is of interest to note that $r_{d}$ is a smooth function of $s, z_{e}, v$, and $\psi_{e}$. Substituting (10.12) and
(10.11) into the second equation of (10.8) yields

$$
\begin{equation*}
\dot{w}_{e}^{*}=-k_{1} w_{e}^{*}+b \tilde{r}+\frac{u}{u_{t}^{2}} \frac{\tau_{w v}(t)}{m_{22}}-\frac{u}{u_{t}^{2}} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{w_{e}^{*}}{\varepsilon_{1}} \frac{\tau_{w v \max }}{m_{22}}\right) \tag{10.13}
\end{equation*}
$$

## Step 2

Differentiating both sides of (10.11) along the solutions of the last equation of (10.8) and (10.12) yields

$$
\begin{align*}
\dot{\tilde{r}}= & \frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t)- \\
& \frac{\partial r_{d}}{\partial s} \frac{1}{1-c(s) z_{e}}\left(u \cos \left(\psi_{e}\right)-v \sin \left(\psi_{e}\right)\right)-\frac{\partial r_{d}}{\partial z_{e}}\left(u \sin \left(\psi_{e}\right)+v \cos \left(\psi_{e}\right)\right)- \\
& \frac{\partial r_{d}}{\partial v}\left(-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t)\right)- \\
& \frac{\partial r_{d}}{\partial \psi_{e}}\left(r-\frac{c(s)}{1-c(s) z_{e}}\left(u \cos \left(\psi_{e}\right)-v \sin \left(\psi_{e}\right)\right)\right) . \tag{10.14}
\end{align*}
$$

From (10.14), we choose the actual control $\tau_{r}$ without canceling the useful damping terms as

$$
\begin{align*}
\tau_{r} & =m_{33}\left(-k_{2} \tilde{r}-b w_{e}^{*}-\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v+\frac{d_{33}}{m_{33}} r_{d}+\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1} r_{d}+\right. \\
& \frac{\partial r_{d}}{\partial s} \frac{1}{1-c(s) z_{e}}\left(u \cos \left(\psi_{e}\right)-v \sin \left(\psi_{e}\right)\right)+\frac{\partial r_{d}}{\partial z_{e}}\left(u \sin \left(\psi_{e}\right)+v \cos \left(\psi_{e}\right)\right)+ \\
& \frac{\partial r_{d}}{\partial v}\left(-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v\right)+ \\
& \frac{\partial r_{d}}{\partial \psi_{e}}\left(r-\frac{c(s)}{1-c(s) z_{e}}\left(u \cos \left(\psi_{e}\right)-v \sin \left(\psi_{e}\right)\right)\right)- \\
& \left.\frac{\tau_{w r \max }}{m_{33}} \tanh \left(\frac{\tilde{r}}{\varepsilon_{2}} \frac{\tau_{w r \max }}{m_{33}}\right)-\frac{\partial r_{d}}{\partial v} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{\tilde{r}}{\varepsilon_{3}} \frac{\tau_{w v \max }}{m_{22}}\right)\right), \tag{10.15}
\end{align*}
$$

where $k_{2}, \varepsilon_{2}$, and $\varepsilon_{3}$ are positive constants to be selected later. Substituting (10.15) into (10.14) yields

$$
\dot{\tilde{r}}=-\left(k_{2}+\frac{d_{33}}{m_{33}}\right) \tilde{r}-b w_{e}^{*}-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1} \tilde{r}+\frac{1}{m_{33}} \tau_{w r}(t)-
$$

$$
\begin{align*}
& \frac{\tau_{w r \max }}{m_{33}} \tanh \left(\frac{\tilde{r}}{\varepsilon_{2}} \frac{\tau_{w r \max }}{m_{33}}\right)-\frac{\partial r_{d}}{\partial v} \frac{1}{m_{22}} \tau_{w v}(t)- \\
& \frac{\partial r_{d}}{\partial v} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{\tilde{r}}{\varepsilon_{3}} \frac{\tau_{w v \max }}{m_{22}}\right) . \tag{10.16}
\end{align*}
$$

We now present the first main result of this chapter, the proof of which is given in the next section.

Theorem 10.1. Under Assumption 10.1, if the state feedback control law (10.15) is applied to the ship system (10.1) then there exist feasible initial conditions such that the regulation errors $\left(z_{e}(t), \psi_{e}^{*}(t)\right)$ converge to a small ball centered at the origin with an appropriate choice of the design constants $k, k_{1}, k_{2}$ and $\varepsilon_{i}, 1 \leq i \leq 3$. Furthermore if there are no environmental disturbances, the regulation errors converge to zero asymptotically. In addition, the sway velocity $v(t)$ is always bounded.

### 10.2.2 Stability Analysis

For the reader's convenience we write the closed loop consisting of the first and third equations of (10.8), (10.13), and (10.16) as follows:

$$
\begin{align*}
\dot{z}_{e}= & -\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}+\frac{u_{t} \sin \left(w_{e}^{*}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}-\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}\left(\cos \left(w_{e}^{*}\right)-1\right), \\
\dot{w}_{e}^{*}= & -k_{1} w_{e}^{*}+b \tilde{r}+\frac{u}{u_{t}^{2}} \frac{\tau_{w v}(t)}{m_{22}}-\frac{u}{u_{t}^{2}} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{w_{e}^{*}}{\varepsilon_{1}} \frac{\tau_{w v \max }}{m_{22}}\right), \\
\dot{v}= & -\frac{m_{11}}{m_{22}} u\left(\tilde{r}+r_{d}\right)-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t),  \tag{10.17}\\
\dot{\tilde{r}}= & -\left(k_{2}+\frac{d_{33}}{m_{33}}\right) \tilde{r}-b w_{e}^{*}-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1} \tilde{r}+\frac{1}{m_{33}} \tau_{w r}(t)- \\
& \frac{\tau_{w r \max }}{m_{33}} \tanh \left(\frac{\tilde{r}}{\varepsilon_{2}} \frac{\tau_{w r \max }}{m_{33}}\right)-\frac{\partial r_{d}}{\partial v} \frac{1}{m_{22}} \tau_{w v}(t)- \\
& \frac{\partial r_{d}}{\partial v} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{\tilde{r}}{\varepsilon_{3}} \frac{\tau_{w v \max }}{m_{22}}\right),
\end{align*}
$$

where $r_{d}$ is given in (10.12). A direct calculation of the Lyapunov function candidate $V_{0}=z_{e}^{2}+w_{e}^{* 2}+v^{2}+\tilde{r}^{2}$ along the solutions of the closed loop system (10.17) shows that there exist nonnegative constants $a_{0}$ and $b_{0}$ such that $\dot{V}_{0} \leq a_{0} V_{0}+b_{0}$. This means that the closed loop system (10.17) is forward complete. Therefore, to prove Theorem 10.1, we can first consider the ( $w_{e}^{*}, \tilde{r}$ )-dynamics, then move to the $\left(z_{e}, v\right)$-dynamics.
$\left(w_{e}^{*}, \tilde{r}\right)$-dynamics
Consider the following Lyapunov function:

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(w_{e}^{* 2}+\tilde{r}^{2}\right) . \tag{10.18}
\end{equation*}
$$

It is not hard to show that the time derivative of (10.18) along the solutions of (10.13) and (10.16) satisfies

$$
\begin{equation*}
\dot{V}_{1} \leq-k_{1} w_{e}^{* 2}-\left(k_{2}+\frac{d_{33}}{m_{33}}\right) \tilde{r}^{2}+0.2785 \sum_{i=1}^{3} \varepsilon_{i} \tag{10.19}
\end{equation*}
$$

which in turn implies that

$$
\begin{equation*}
\left\|\left(w_{e}^{*}(t), \tilde{r}(t)\right)\right\| \leq \alpha_{w r}(\cdot) e^{-\sigma_{w r}\left(t-t_{0}\right)}+\rho_{w r}, \tag{10.20}
\end{equation*}
$$

where

$$
\begin{align*}
& \alpha_{w r}(\cdot)=\left\|\left(w_{e}^{*}\left(t_{0}\right), \tilde{r}\left(t_{0}\right)\right)\right\| \\
& \sigma_{w r}=\min \left(k_{1},\left(k_{2}+\frac{d_{33}}{m_{33}}\right)\right),  \tag{10.21}\\
& \rho_{w r}=\sqrt{\frac{0.2785}{\sigma_{1}} \sum_{i=1}^{3} \varepsilon_{i}}
\end{align*}
$$

From (10.21), we can see that $\rho_{w r}$ can be made arbitrarily small by adjusting the design constants, $k_{1}, k_{2}$ and $\varepsilon_{i}, 1 \leq i \leq 3$. This observation is important in stability analysis of the $\left(z_{e}, v\right)$-dynamics.

## $\left(z_{e}, v\right)$-dynamics

For the reader's convenience, we rewrite the $\left(z_{e}, v\right)$-dynamics from the first and third equations of the closed loop system (10.17) with the virtual control $r_{d}$ substituted in the third equation of the closed loop system (10.17) from (10.12) as

$$
\begin{aligned}
\dot{z}_{e}= & -\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}+\frac{u_{t} \sin \left(w_{e}^{*}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}-\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}\left(\cos \left(w_{e}^{*}\right)-1\right), \\
\dot{v}= & -\frac{d_{22}}{m_{22}}\left(1+\frac{m_{11} u^{2}}{b m_{22} u_{t}^{2}}\right) v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}\left(1+\frac{m_{11} u^{2}}{b m_{22} u_{t}^{2}}\right)|v|^{i-1} v+ \\
& \frac{1}{m_{22}} \tau_{w v}(t)-\frac{m_{11} u \tilde{r}}{m_{22}}-\frac{m_{11} u}{m_{22} b}\left(-k_{1} w_{e}^{*}-\frac{k u_{t} \sin \left(\psi_{e}^{*}\right)}{1+\left(k z_{e}\right)^{2}}+\right.
\end{aligned}
$$

$$
\begin{equation*}
\left.\frac{c(s) u_{t} \cos \left(\psi_{e}^{*}\right)}{1-c(s) z_{e}}-\frac{u}{u_{t}^{2}} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{w_{e}^{*}}{\varepsilon_{1}} \frac{\tau_{w v \max }}{m_{22}}\right)\right) . \tag{10.22}
\end{equation*}
$$

If we view that (10.22) and the $\left(w_{e}^{*}, \tilde{r}\right)$-subsystem are in a cascade form, we might think that stability results developed for cascade systems, see Section 2.1.3 can be applied. However, stability results in that section are developed for the cascade systems without nonvanishing disturbances. In fact, nonvanishing disturbances may destroy stability of a cascade system that satisfies all conditions stated in Section 2.1.3. To illustrate this fact, we give the following simple example.

$$
\begin{align*}
& \dot{x}_{1}=-\frac{x_{1}}{\sqrt{1+x_{1}^{2}}}+\frac{x_{1} x_{2}}{\sqrt{1+x_{1}^{2}}} \\
& \dot{x}_{2}=-x_{2}+d(t) \tag{10.23}
\end{align*}
$$

When there is no disturbance $d(t)$, by applying stability results in Section 2.1.3, the cascade system (10.23) is GAS at the origin. However, whenever the magnitude of bounded disturbance $d(t)$ is larger than 1 , we have the fact that $x_{1}$ grows unbounded although $x_{2}$ is bounded. This fact results from the first equation being not globally ISS with respect to $x_{2}$ as input, see Section 2.2. Therefore, to analyze stability of (10.22) we first present the following lemma.

Lemma 10.1. Consider the following nonlinear system:

$$
\begin{equation*}
\dot{x}=f(t, x)+g(t, x, \xi(t)) \tag{10.24}
\end{equation*}
$$

where $x \in \mathbb{R}^{n}, \xi(t) \in \mathbb{R}^{m}, f(t, x)$ is piecewise continuous in $t$ and locally Lipschitz in $x$. If there exist $\sigma_{0}>0$, positive constants $c_{i}, 1 \leq i \leq 4, \lambda_{j}, 1 \leq j \leq 2, c, \varepsilon_{0}, m u_{0}$, and $c_{0}$, and a class-K function $\alpha_{0}$ such that the following conditions hold:
C1. There exists a proper function $V(t, x)$ such that:

$$
\begin{aligned}
& c_{1}\|x\|^{2} \leq V(t, x) \leq c_{2}\|x\|^{2} \\
& \left\|\frac{\partial V}{\partial x}(t, x)\right\| \leq c_{3}\|x\| \\
& \frac{\partial V}{\partial t}+\frac{\partial V}{\partial x} f(t, x) \leq-\frac{c_{4}\|x\|^{2}}{\sqrt{1+c\|x\|^{2}}}+\frac{c_{0}}{\sqrt{1+c\|x\|^{2}}}
\end{aligned}
$$

C2. $g(t, x, \xi(t))$ satisfies:

$$
\|g(t, x, \xi(t))\| \leq \frac{1}{\sqrt{1+c\|x\|^{2}}}\left(\lambda_{1}+\lambda_{2}\|x\|\right)\|\xi(t)\|
$$

C3. $\xi(t)$ globally asymptotically converges to a ball centered at the origin:

$$
\|\xi(t)\| \leq \alpha_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+\varepsilon_{0}, \forall t \geq t_{0} \geq 0
$$

C4. The following gain condition is satisfied:

$$
c_{4}-c_{3} \varepsilon_{0}\left(\lambda_{2}+\frac{\lambda_{1}}{4 \mu_{0}}\right)>0
$$

Then the solution $x(t)$ of (10.24) globally asymptotically converges to a ball centered at the origin, i.e.,

$$
\begin{equation*}
\|x(t)\| \leq \alpha\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right) e^{-\sigma\left(t-t_{0}\right)}+\varepsilon\left(\left\|\left(x\left(t_{0}\right), \xi\left(t_{0}\right)\right)\right\|\right), \forall t \geq t_{0} \geq 0 \tag{10.25}
\end{equation*}
$$

where $\varepsilon(s)=\sqrt{\frac{a_{4}}{c_{1} a_{1}} \sqrt{1+c x_{m}^{2}(s)}}, \sigma$ and $\alpha$ are given as follows:

$$
\begin{aligned}
& \text { if } \frac{a_{1}}{\sqrt{1+c x_{m}^{2}}}=\sigma_{0} \text { then } \\
& \begin{array}{l}
\alpha(s)=\sqrt{\frac{e^{\frac{a_{2}}{\sigma_{0}}}}{c_{1}}\left(c_{2} s^{2}+\left(a_{3}+\frac{a_{2} a_{4}}{a_{1}} \sqrt{1+c x_{m}^{2}}\right) \phi\right),} \\
\quad \sigma=\frac{1}{2}\left(a_{1}-d\right) ; \\
\text { if } \frac{a_{1}}{\sqrt{1+c x_{m}^{2}}} \neq \sigma_{0} \text { then }
\end{array},
\end{aligned}
$$

$$
\alpha(s)=\sqrt{\frac{e^{\frac{a_{2}}{\sigma_{0}}}}{c_{1}}\left(c_{2} s^{2}+\frac{\left(a_{1} a_{3}+a_{2} a_{4} \sqrt{1+c x_{m}^{2}}\right)}{a_{1}\left|a_{1}-\sigma_{0} \sqrt{1+c x_{m}^{2}}\right|} \sqrt{1+c x_{m}^{2}}\right)},
$$

$$
\sigma(s)=\frac{1}{2} \min \left(\frac{a_{1}}{\sqrt{1+c x_{m}^{2}}},\left|\frac{a_{1}}{\sqrt{1+c x_{m}^{2}}}-\sigma_{0}\right|\right)
$$

with

$$
\begin{aligned}
& a_{1}=\frac{1}{c_{2}}\left(c_{4}-\lambda_{2} c_{3} \varepsilon_{0}-\frac{\lambda_{1} c_{3} \varepsilon_{0}}{4 \mu_{0}}\right), a_{2}=\frac{c_{3}}{c_{1}}\left(\lambda_{1}+\lambda_{2}\right) \alpha_{0}(s), \\
& a_{3}=\frac{\lambda_{1} c_{3}}{4} \alpha_{0}(s), a_{4}=c_{0}+\lambda_{1} c_{3} \varepsilon_{0} \mu_{0}, \\
& 0<d<a_{1}, \phi \geq\left(t-t_{0}\right) e^{-d\left(t-t_{0}\right)}, \\
& x_{m}(s)=\sqrt{\frac{1}{c_{1}}\left(\frac{a_{4}}{a_{1}}+\sqrt{W_{m}(s)}\right)}, \\
& W_{m}(s)=\left(V(s)-\frac{a_{4}}{a_{1}}\right) e^{\left(\frac{2\left(a_{2}+a_{3}+a_{2} a_{4} / a_{1}\right)}{\sigma_{0}}\right)}+ \\
& \quad \frac{a_{3}+a_{2} a_{4} / a_{1}}{8\left(a_{2}(s)+a_{3}(s)+a_{2}(s) a_{4} / a_{1}\right)}\left(e^{\left(\frac{2\left(a_{2}+a_{3}+a_{2} a_{4} / a_{1}\right)}{\sigma_{0}}\right)}-1\right) .
\end{aligned}
$$

When $c_{0}=0$ and $\varepsilon_{0}=0$, we have $\varepsilon=0$ and the system (10.24) is GAS. Note that a finite value of $\phi$ exists for an arbitrarily small positive $d$ and the convergence rate $\sigma$ depends on the initial conditions.

Proof. We first prove that $x$ is bounded. From conditions C1, C2, and C3, we have

$$
\begin{align*}
\dot{V} \leq & -\left(c_{4}-c_{3} \varepsilon_{0}\left(\lambda_{2}+\frac{\lambda_{1}}{4 \mu_{0}}\right)\right) \frac{\|x\|^{2}}{\sqrt{1+c\|x\|^{2}}}+\frac{c_{3}\|x\|}{\sqrt{1+c\|x\|^{2}}} \times \\
& \left(\lambda_{1}+\lambda_{2}\|x\|\right) \alpha_{0}\left(\left\|\xi\left(t_{0}\right)\right\|\right) e^{-\sigma_{0}\left(t-t_{0}\right)}+\frac{\lambda_{1} c_{3} \varepsilon_{0} \mu_{0}+c_{0}}{\sqrt{1+c\|x\|^{2}}} \tag{10.26}
\end{align*}
$$

Upon application of completing squares, (10.26) can be rewritten as

$$
\begin{align*}
\dot{V} & \leq-\left(a_{1}-a_{2} e^{-\sigma_{0}\left(t-t_{0}\right)}\right) \frac{V}{\sqrt{1+c\|x\|^{2}}}+\frac{a_{3} e^{-\sigma_{0}\left(t-t_{0}\right)}}{\sqrt{1+c\|x\|^{2}}}+\frac{a_{4}}{\sqrt{1+c\|x\|^{2}}} \\
& \leq-\left(a_{1}-a_{2} e^{-\sigma_{0}\left(t-t_{0}\right)}\right) \frac{\left(V-a_{4} / a_{1}\right)}{\sqrt{1+c\|x\|^{2}}}+\frac{\left(a_{3}+a_{2} a_{4} / a_{1}\right) e^{-\sigma_{0}\left(t-t_{0}\right)}}{\sqrt{1+c\|x\|^{2}}} \tag{10.27}
\end{align*}
$$

Now consider the differential equation

$$
\begin{equation*}
\dot{\kappa}=-\left(a_{1}-a_{2} e^{-\sigma_{0}\left(t-t_{0}\right)}\right) \frac{\kappa}{\sqrt{1+c\|x\|^{2}}}+\frac{\left(a_{3}+a_{2} a_{4} / a_{1}\right) e^{-\sigma_{0}\left(t-t_{0}\right)}}{\sqrt{1+c\|x\|^{2}}} \tag{10.28}
\end{equation*}
$$

Take the following Lyapunov function:

$$
\begin{equation*}
W=\frac{1}{2} \kappa^{2} \tag{10.29}
\end{equation*}
$$

whose time derivative along the solutions of (10.28) with the use of condition C 4 satisfies

$$
\begin{align*}
\dot{W} & \leq a_{2} e^{-\sigma_{0}\left(t-t_{0}\right)} \frac{\kappa^{2}}{\sqrt{1+c\|x\|^{2}}}+\frac{\left(a_{3}+a_{2} a_{4} / a_{1}\right) \kappa e^{-\sigma_{0}\left(t-t_{0}\right)}}{\sqrt{1+c\|x\|^{2}}} \\
& \leq 2\left(a_{2}+a_{3}+\frac{a_{2} a_{4}}{a_{1}}\right) W e^{-\sigma_{0}\left(t-t_{0}\right)}+\frac{\left(a_{3}+\frac{a_{2} a_{4}}{a_{1}}\right)}{4} e^{-\sigma_{0}\left(t-t_{0}\right)} . \tag{10.30}
\end{align*}
$$

Hence $W(t) \leq W_{m}$ which implies from (10.30) and the comparison principle found in [6] that $V(t) \leq a_{4} / a_{1}+\sqrt{2 W_{m}}$ or $\|x(t)\| \leq x_{m}$. Substituting this inequality into (10.27) yields

$$
\begin{equation*}
\dot{V} \leq-\left(\frac{a_{1}}{\sqrt{1+c x_{m}^{2}}}-a_{2} e^{-\sigma_{0}\left(t-t_{0}\right)}\right) V+a_{3} e^{-\sigma_{0}\left(t-t_{0}\right)}+a_{4} . \tag{10.31}
\end{equation*}
$$

Solving the inequality (10.31) readily yields (10.25).
To investigate stability of the $\left(z_{e}, v\right)$-dynamics, we apply Lemma 10.1.

## $z_{e}$-dynamics

We will apply Lemma 10.1 to estimate convergence of the path-following error, $z_{e}$, dynamics. Indeed, it is seen that $z_{e}$ - dynamics, the first equation of (10.22), is in the form of the system studied in Lemma 10.1, i.e.,

$$
\begin{equation*}
\dot{z}_{e}=f_{z e}(\cdot)+g_{z e}(\cdot), \tag{10.32}
\end{equation*}
$$

where

$$
\begin{align*}
& f_{z e}(\cdot)=-\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}} \\
& g_{z e}(\cdot)=\frac{u_{t} \sin \left(w_{e}^{*}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}-\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}\left(\cos \left(w_{e}^{*}\right)-1\right) \tag{10.33}
\end{align*}
$$

We now verify all conditions of Lemma 10.1.
Verifying Condition C1. By taking the proper function $V_{2}=0.5 z_{e}^{2}$, it is not hard to show that this condition holds with

$$
\begin{equation*}
c=k^{2}, c_{1}=c_{2}=0.5, c_{3}=1, c_{4}=k u_{t}, c_{0}=0 \tag{10.34}
\end{equation*}
$$

Verifying Condition C2. By noting that

$$
\begin{gather*}
\left|\frac{u_{t} \sin \left(w_{e}^{*}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}-\frac{k u_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}\left(\cos \left(w_{e}^{*}\right)-1\right)\right| \leq \\
\frac{1}{\sqrt{1+\left(k z_{e}\right)^{2}}} u_{t}\left(1+k\left|z_{e}\right|\right)\left|w_{e}^{*}\right| \tag{10.35}
\end{gather*}
$$

condition C2 holds with

$$
\begin{align*}
& \lambda_{1}=u_{t}, \\
& \lambda_{2}=k u_{t} . \tag{10.36}
\end{align*}
$$

Verifying Condition C3. This condition holds directly from (10.20).
Verifying Condition C4. From (10.34), (10.36), and (10.20), this condition becomes

$$
\begin{equation*}
k u_{t}-\rho_{w r}\left(u_{t}+\frac{k u_{t}}{4 \mu_{0}}\right)>0 \Leftrightarrow k-\rho_{w r}\left(1+\frac{k}{4 \mu_{0}}\right)>0 . \tag{10.37}
\end{equation*}
$$

Since $\mu_{0}$ is an arbitrarily positive constant, see Lemma 10.1, and $\rho_{w r}$ can be made arbitrarily small, we can always pick a positive constant $k$ such that (10.37) holds. Therefore, we have

$$
\begin{equation*}
\left|z_{e}(t)\right| \leq \alpha_{z}(\cdot) e^{-\sigma_{z}\left(t-t_{0}\right)}+\rho_{z} \tag{10.38}
\end{equation*}
$$

where $\alpha_{z}(\cdot), \sigma_{z}$ and $\rho_{z}$ are calculated as in Lemma 10.1. From (10.38), the requirement of initial conditions such that $1-c(s) z_{e} \geq \delta^{*}>0$ becomes

$$
\begin{equation*}
\alpha_{z}(\cdot) \leq \frac{1-\delta^{*}}{|c(s)|}-\rho_{z} . \tag{10.39}
\end{equation*}
$$

## Boundedness of $v$

To show that $v$ is bounded, we take the following Lyapunov function

$$
\begin{equation*}
V_{3}=\frac{1}{2} v^{2}, \tag{10.40}
\end{equation*}
$$

whose time derivative along the solutions of the second equation of (10.22), after some simple calculation, satisfies

$$
\begin{align*}
\dot{V}_{3} \leq & -\frac{d_{22}}{m_{22}}\left(1+\frac{m_{11} u^{2}}{b m_{22} u_{t}^{2}}\right) v^{2}+v\left(\frac{1}{m_{22}} \tau_{w v}(t)-\frac{m_{11} u \tilde{r}}{m_{22}}-\right. \\
& \frac{m_{11} u}{m_{22} b}\left(-k_{1} w_{e}^{*}-\frac{k u_{t} \sin \left(\psi_{e}^{*}\right)}{1+\left(k z_{e}\right)^{2}}+\frac{c(s) u_{t} \cos \left(\psi_{e}^{*}\right)}{1-c(s) z_{e}}-\right. \\
& \left.\left.\frac{u}{u_{t}^{2}} \frac{\tau_{w v \max }}{m_{22}} \tanh \left(\frac{w_{e}^{*}}{\varepsilon_{1}} \frac{\tau_{w v \max }}{m_{22}}\right)\right)\right) . \tag{10.41}
\end{align*}
$$

From (10.41), (10.20), (10.38), and (10.7), it can be shown that

$$
\begin{equation*}
|v(t)| \leq \alpha_{v}(\cdot) e^{-\sigma_{v}\left(t-t_{0}\right)}+\rho_{v}, \tag{10.42}
\end{equation*}
$$

where $\alpha_{v}(\cdot)$ is a nondecreasing function of $\left\|\left(w_{e}^{*}\left(t_{0}\right), \tilde{r}\left(t_{0}\right), z_{e}\left(t_{0}\right)\right)\right\|$, and $\sigma_{v}$ and $\rho_{v}$ are positive constants. The constant $\rho_{v}$ depends on the environmental disturbances and the path curvature. Indeed, this constant is equal to zero when there are no environmental disturbances and $c(s)=0$.

Remark 10.2. From (10.39), it is seen that the initial conditions are mainly limited by the path curvature $c(s)$. If the initial conditions do not satisfy (10.39), one can generate a smooth curve segment, $\boldsymbol{\Omega}_{\mathbf{0}}$, such that it is tangent to the path $\boldsymbol{\Omega}$ and that (10.39) holds, see Figure 10.2.


Figure 10.2 Overcoming "bad" initial conditions

### 10.3 Output Feedback

### 10.3.1 Observer Design

To design an observer that estimates the sway and yaw velocities from the ship position and orientation measurements, we rewrite (10.1) as follows

$$
\begin{align*}
& \dot{x}=u \cos (\psi)-v \sin (\psi) \\
& \dot{\boldsymbol{x}}_{\mathbf{1}}=\boldsymbol{F}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right)+\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right) \boldsymbol{x}_{\mathbf{2}}  \tag{10.43}\\
& \dot{\boldsymbol{x}}_{\mathbf{2}}=\boldsymbol{K}_{\mathbf{2}} \boldsymbol{x}_{\mathbf{2}}-\boldsymbol{D}_{\mathbf{2}}\left(\boldsymbol{x}_{\mathbf{2}}\right) \boldsymbol{x}_{\mathbf{2}}+\boldsymbol{\tau}_{\mathbf{2}}+\boldsymbol{\tau}_{\boldsymbol{w} \mathbf{2}}(t)
\end{align*}
$$

where

$$
\begin{align*}
\boldsymbol{x}_{\mathbf{1}} & =\left[\begin{array}{c}
y \\
\psi
\end{array}\right], \boldsymbol{x}_{\mathbf{2}}=\left[\begin{array}{c}
y \\
\psi
\end{array}\right], \boldsymbol{F}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right)=\left[\begin{array}{c}
u \sin (\psi) \\
0
\end{array}\right], \\
\boldsymbol{\tau}_{\mathbf{2}} & =\left[\begin{array}{c}
0 \\
\frac{1}{m_{33}} \tau_{r}
\end{array}\right], \boldsymbol{\tau}_{\boldsymbol{w} \mathbf{2}}(t)=\left[\begin{array}{c}
\frac{1}{m_{22}} \tau_{w v}(t) \\
\frac{1}{m_{33}} \tau_{w r}(t)
\end{array}\right], \boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right)=\left[\begin{array}{cc}
\cos (\psi) & 0 \\
0 & 1
\end{array}\right], \\
\boldsymbol{D}_{\mathbf{2}}\left(\boldsymbol{x}_{\mathbf{2}}\right) & =\left[\begin{array}{cc}
\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} & 0 \\
0 & \sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|r|^{i-1}
\end{array}\right], \boldsymbol{K}_{\mathbf{2}}=\left[\begin{array}{cc}
-\frac{d_{22}}{m_{22}} & -\frac{m_{11} u}{m_{22}} \\
\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} & -\frac{d_{33}}{m_{33}}
\end{array}\right] . \tag{10.44}
\end{align*}
$$

If we use the observer designed in $[116,126]$ and Chapter 9 , then the matrix $\boldsymbol{K}_{\mathbf{2}}$ has to be Hurwitz, i.e., the ship has to possess straight-line stability. This limitation motivates us to seek a new observer in this section. The idea is to find a coordinate transformation to transform the system (10.1) to a form that does not require $\boldsymbol{K}_{\mathbf{2}}$ to be a Hurwitz matrix to design an observer. To this end, we define

$$
\begin{equation*}
z_{2}=x_{2}-\chi_{2}\left(x, x_{1}\right), \tag{10.45}
\end{equation*}
$$

where $\chi_{2}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)$ is a locally Lipschitz vector function, which will be determined later. Using (10.45), we write the last two equations of (10.43) as follows:

$$
\begin{align*}
\dot{x}_{1}= & F_{1}\left(x_{1}\right)+J_{1}\left(x_{1}\right) \chi_{2}\left(x, x_{1}\right)+J_{1}\left(x_{1}\right) z_{2}, \\
\dot{z}_{2}= & \left(K_{2}+\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x} g_{2}\left(x_{1}\right)-\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x_{1}} J_{1}\left(x_{1}\right)\right) z_{2}-D_{2}\left(z_{2}+\right. \\
& \left.\chi_{2}\left(x, x_{1}\right)\right)\left(z_{2}+\chi_{2}\left(x, x_{1}\right)\right)+\Phi\left(x, x_{1}\right)+\tau_{2}+\tau_{w 2}(t), \tag{10.46}
\end{align*}
$$

where

$$
\begin{align*}
& \boldsymbol{\Phi}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)= \boldsymbol{K}_{\mathbf{2}} \boldsymbol{\chi}_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)-\frac{\partial \chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)}{\partial x}\left(u \cos (\psi)-\boldsymbol{g}_{\mathbf{2}}\left(\boldsymbol{x}_{\mathbf{1}}\right) \chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)- \\
& \frac{\partial \chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)}{\partial \boldsymbol{x}_{\mathbf{1}}}\left(\boldsymbol{F}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right)+\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right) \chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right), \\
& \boldsymbol{g}_{\mathbf{2}}\left(\boldsymbol{x}_{\mathbf{1}}\right)= {[\sin (\psi) 0]^{T}, }  \tag{10.47}\\
& \boldsymbol{D}_{\mathbf{2}}\left(\boldsymbol{z}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)= \\
& {\left[\begin{array}{ll}
\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}\left|z_{21}+\chi_{\mathbf{2 1}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right|^{i-1} 0 \\
0 & \left.\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}\left|z_{22}+\chi_{\mathbf{2 1}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right|^{i-1}\right],
\end{array}\right] }
\end{align*}
$$

where $z_{2 i}$ and $\chi_{\mathbf{2 i}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right), i=1,2$, are the first and second elements of $\boldsymbol{z}_{\mathbf{2}}$ and $\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)$, respectively. From (10.46), one can design a reduced order observer. However it is often noise-sensitive. We here propose the following observer:

$$
\begin{align*}
\dot{\hat{x}}_{\mathbf{1}}= & \boldsymbol{F}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right)+\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right) \chi_{2}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)+\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right) \hat{z}_{\mathbf{2}}+\boldsymbol{K}_{01}\left(\boldsymbol{x}_{\mathbf{1}}-\hat{\boldsymbol{x}}_{\mathbf{1}}\right), \\
\dot{\hat{z}}_{\mathbf{2}}= & \left(\boldsymbol{K}_{\mathbf{2}}+\frac{\partial \chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)}{\partial x} \boldsymbol{g}_{2}\left(\boldsymbol{x}_{\mathbf{1}}\right)-\frac{\partial \chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)}{\partial \boldsymbol{x}_{\mathbf{1}}} \boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{x}_{\mathbf{1}}\right)\right) \hat{z}_{\mathbf{2}}-\boldsymbol{D}_{\mathbf{2}}\left(\hat{z}_{\mathbf{2}}+\right. \\
& \left.\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)\left(\hat{z}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)+\boldsymbol{\Phi}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)+\boldsymbol{\tau}_{\mathbf{2}}+\boldsymbol{K}_{\mathbf{0}}\left(\boldsymbol{x}_{\mathbf{1}}-\hat{\boldsymbol{x}}_{\mathbf{1}}\right) \tag{10.48}
\end{align*}
$$

where $\hat{\boldsymbol{x}}_{\mathbf{1}}:=[\hat{y} \hat{\psi}]^{T}$ and $\hat{\boldsymbol{z}}_{\mathbf{2}}$ are estimates of $\boldsymbol{x}_{\mathbf{1}}$ and $\boldsymbol{z}_{\mathbf{2}}$, respectively, and $\boldsymbol{K}_{\mathbf{0 1}}$ and $\boldsymbol{K}_{02}$ are the observer gains to be selected later. From (10.46) and (10.48), we have the observer error dynamics:

$$
\begin{align*}
\dot{\tilde{x}}_{\mathbf{1}}= & -K_{01} \tilde{x}_{1}+J_{1}\left(\boldsymbol{x}_{1}\right) \tilde{z}_{2}, \\
\dot{\tilde{z}}_{2}= & -K_{02} \tilde{x}_{1}+\left(\boldsymbol{K}_{2}+\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x} g_{2}\left(\boldsymbol{x}_{\mathbf{1}}\right)-\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x_{1}} J_{1}\left(x_{1}\right)\right) \tilde{z}_{2}- \\
& \left(D_{2}\left(z_{2}+\chi_{2}\left(x, x_{1}\right)\right)\left(z_{2}+\chi_{2}\left(x, x_{1}\right)\right)-D_{2}\left(\hat{z}_{2}+\right.\right. \\
& \left.\left.\chi_{2}\left(x, x_{1}\right)\right)\left(\hat{z}_{2}+\chi_{2}\left(x, x_{1}\right)\right)\right)+\tau_{w 2}(t), \tag{10.49}
\end{align*}
$$

where $\tilde{\boldsymbol{x}}_{\mathbf{1}}:=[\tilde{y} \tilde{\psi}]^{T}=\boldsymbol{x}_{\mathbf{1}}-\hat{\boldsymbol{x}}_{\mathbf{1}}$, and $\tilde{\boldsymbol{z}}_{\mathbf{2}}=\boldsymbol{z}_{\mathbf{2}}-\hat{\boldsymbol{z}}_{\mathbf{2}}$. Now, considering the Lyapunov function

$$
\begin{equation*}
V_{0}=\tilde{\boldsymbol{x}}_{\mathbf{1}}^{\boldsymbol{T}} \boldsymbol{P}_{\mathbf{0} 1} \tilde{\boldsymbol{x}}_{\mathbf{1}}+\tilde{z}_{\mathbf{2}}^{\boldsymbol{T}} \boldsymbol{P}_{\mathbf{0} 2} \tilde{z}_{\mathbf{2}} \tag{10.50}
\end{equation*}
$$

where $\boldsymbol{P}_{\mathbf{0 1}}$ and $\boldsymbol{P}_{\mathbf{0 2}}$ are positive definite matrices, whose time derivative along the solutions of (10.49) satisfies:

$$
\begin{equation*}
\dot{V}_{0} \leq-\tilde{x}_{1}^{T} Q_{01} \tilde{x}_{1}-\tilde{z}_{2}^{T} Q_{02} \tilde{z}_{2}+\tau_{w 2}^{T}(t) P_{02} \tilde{z}_{2}+\tilde{z}_{2}^{T} P_{02} \tau_{w 2}(t), \tag{10.51}
\end{equation*}
$$

where

$$
\begin{align*}
& Q_{01}=K_{01}^{T} P_{01}+P_{01} K_{01}, \\
& Q_{02}=-\left(K_{2}+\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x} g_{2}\left(x_{1}\right)-\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x_{1}} J_{1}\left(x_{1}\right)\right)^{T} P_{02}- \\
& \quad P_{02}\left(K_{2}+\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x} g_{2}\left(x_{1}\right)-\frac{\partial \chi_{2}\left(x, x_{1}\right)}{\partial x_{1}} J_{1}\left(x_{1}\right)\right), \tag{10.52}
\end{align*}
$$

and we have chosen

$$
\begin{equation*}
P_{02} K_{02}-J_{1}^{T}\left(x_{1}\right) P_{01}=0 \tag{10.53}
\end{equation*}
$$

and used the following inequality:

$$
\begin{align*}
& \left(\boldsymbol{D}_{\mathbf{2}}\left(\boldsymbol{z}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)\left(\boldsymbol{z}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)-\boldsymbol{D}_{\mathbf{2}}\left(\hat{z}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)\right. \\
& \left.\left(\hat{z}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)\right) \boldsymbol{P}_{\mathbf{0 2}} \tilde{z}_{\mathbf{2}}+\tilde{\boldsymbol{z}}_{\mathbf{2}}^{\boldsymbol{T}} \boldsymbol{P}_{\mathbf{0 2}}\left(\boldsymbol { D } _ { \mathbf { 2 } } ( \boldsymbol { z } _ { \mathbf { 2 } } + \chi _ { \mathbf { 2 } } ( x , \boldsymbol { x } _ { \mathbf { 1 } } ) ) \left(\boldsymbol{z}_{\mathbf{2}}+\right.\right. \\
& \left.\left.\boldsymbol{\chi}_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)-\boldsymbol{D}_{\mathbf{2}}\left(\hat{z}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)\left(\hat{z}_{\mathbf{2}}+\chi_{2}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)\right)\right) \geq 0 . \tag{10.54}
\end{align*}
$$

It is not hard to show that there always exist $\boldsymbol{K}_{\mathbf{0 1}}, \boldsymbol{K}_{\mathbf{0 2}}$ and $\boldsymbol{\chi}_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)$ such that $\boldsymbol{Q}_{\mathbf{0 1}}$ and $\boldsymbol{Q}_{\mathbf{0 2}}$ are positive definite and (10.53) holds. For example, one can take

$$
\begin{align*}
& \boldsymbol{K}_{\mathbf{0 1}}=\boldsymbol{K}_{\mathbf{0 1}}^{\boldsymbol{T}}>0, \\
& \boldsymbol{K}_{\mathbf{0 2}}=\boldsymbol{P}_{\mathbf{0 2}}^{-1} \boldsymbol{J}_{\mathbf{1}}^{\boldsymbol{T}}\left(\boldsymbol{x}_{\mathbf{1}}\right) \boldsymbol{P}_{\mathbf{0 1}}, \\
& \chi_{2}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)=\left[-\frac{m_{11} u}{m_{22}} \psi 0\right]^{T} . \tag{10.55}
\end{align*}
$$

Therefore (10.51) implies that

$$
\begin{equation*}
\left\|\tilde{\mathbf{x}}_{\mathbf{0}}(t)\right\| \leq \varphi_{0}\left\|\tilde{\mathbf{x}}_{\mathbf{0}}\left(t_{0}\right)\right\| e^{-\eta_{0}\left(t-t_{0}\right)}+\xi_{0} \tag{10.56}
\end{equation*}
$$

where $\varphi_{0}, \eta_{0}$, and $\xi_{0}$ are some positive constants, and $\tilde{\boldsymbol{x}}_{\boldsymbol{0}}=\left[\tilde{\boldsymbol{x}}_{\mathbf{1}} \tilde{\boldsymbol{z}}_{\mathbf{2}}\right]^{\boldsymbol{T}}$. It is noted that $\xi_{0}$ cannot be made arbitrarily small since the ratio $\frac{\lambda_{\min }\left(\boldsymbol{Q}_{\mathbf{0 2}}\right)}{\lambda_{\max }\left(\boldsymbol{P}_{\mathbf{0 2}}\right)}$ is maximized with the choice of $\boldsymbol{Q}_{\mathbf{0 2}}=\operatorname{diag}(1,1)$. By defining an estimate of the velocity vector as

$$
\begin{equation*}
\hat{\boldsymbol{x}}_{\mathbf{2}}=\hat{\boldsymbol{z}}_{\mathbf{2}}+\chi_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right), \tag{10.57}
\end{equation*}
$$

where $\hat{\mathbf{x}}_{\mathbf{2}}=\left[\begin{array}{ll}\hat{v} & \hat{r}\end{array}\right]^{T}$, then the velocity observer errors $\tilde{\mathbf{x}}_{\mathbf{2}}=\left[\begin{array}{ll}\tilde{v} & \tilde{r}\end{array}\right]^{T}$ satisfy:

$$
\begin{equation*}
\|(\tilde{v}(t), \tilde{r}(t))\| \leq \varphi_{0}\left\|\tilde{\boldsymbol{x}}_{\mathbf{0}}\left(t_{0}\right)\right\| e^{-\eta_{0}\left(t-t_{0}\right)}+\xi_{0} . \tag{10.58}
\end{equation*}
$$

It is seen that when there are no environmental disturbances, the observer errors globally exponentially tend to zero.

### 10.3.2 Control Design

The control design in this section is very similar to the one of the state feedback controller. We therefore present it briefly. Define

$$
\begin{align*}
& \hat{\phi}=\arctan \left(\frac{\hat{v}}{u}\right), \\
& \hat{u}_{t}=\sqrt{u^{2}+\hat{v}^{2}}  \tag{10.59}\\
& \hat{\psi}_{e}^{*}=\psi_{e}+\hat{\phi}
\end{align*}
$$

By denoting $\tilde{\phi}=\phi-\hat{\phi}$, using the Taylor series expansion, it can be shown that

$$
\begin{equation*}
|\tilde{\phi}(t)|=\left|\frac{\arctan (v / u)-\arctan (\hat{v} / u)}{(v-\hat{v}) / u} \frac{v-\hat{v}}{u}\right| \leq \frac{1}{u}|\tilde{v}(t)| . \tag{10.60}
\end{equation*}
$$

Similar to the state feedback control design, we apply the following coordinate transformation

$$
\begin{equation*}
\hat{w}_{e}^{*}=\hat{\psi}_{e}^{*}+\arcsin \left(\frac{k z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}\right) \tag{10.61}
\end{equation*}
$$

to (10.48) in conjunction with (10.59). After some simple calculation, we have

$$
\begin{aligned}
& \dot{z}_{e}=-\frac{k \hat{u}_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}+\frac{\hat{u}_{t} \sin \left(\hat{w}_{e}^{*}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}-\frac{\left(\cos \left(\hat{w}_{e}^{*}\right)-1\right) k \hat{u}_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}+ \\
& \frac{\tilde{v} \cos \left(\hat{w}_{e}^{*}-\hat{\phi}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}+\frac{\tilde{v} \sin \left(\hat{w}_{e}^{*}-\hat{\phi}\right) k z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}, \\
& \dot{\hat{w}}_{e}^{*}= \hat{r}\left(1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{\hat{u}_{t}^{2}}\right)-\frac{c(s) \hat{u}_{t} \cos \left(\hat{\psi}_{e}^{*}\right)}{1-c(s) z_{e}}-
\end{aligned}
$$

$$
\begin{gather*}
\quad \frac{u}{\hat{u}_{t}^{2}}\left(\frac{d_{22}}{m_{22}} \hat{v}+\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}\right)+\frac{k \hat{u}_{t} \sin \left(\hat{\psi}_{e}^{*}\right)}{1+\left(k z_{e}\right)^{2}}+ \\
\frac{u}{\hat{u}_{t}^{2}} \cos (\psi) \tilde{y}+\tilde{r}\left(1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{\hat{u}_{t}^{2}}\right)+\frac{k \tilde{v} \cos \left(\psi_{e}\right)}{1+\left(k z_{e}\right)^{2}}+\frac{c(s) \tilde{v} \sin \left(\psi_{e}\right)}{1-c(s) z_{e}}, \\
\dot{\hat{v}}=-\frac{m_{11} u}{m_{22}} \hat{r}-\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}+\cos (\psi) \tilde{y}-\frac{m_{11}}{m_{22}} u \tilde{r}, \\
\dot{\hat{r}}=\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \hat{v}-\frac{d_{33}}{m_{33}} \hat{r}-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} \hat{r}+\frac{1}{m_{33}} \tau_{r}+\tilde{\psi} \tag{10.62}
\end{gather*}
$$

where for simplicity, we have taken

$$
\begin{align*}
& \boldsymbol{K}_{\mathbf{0 2}}=\boldsymbol{J}_{\mathbf{1}}^{\boldsymbol{T}}\left(\boldsymbol{x}_{\mathbf{1}}\right), \\
& \boldsymbol{\chi}_{\mathbf{2}}\left(x, \boldsymbol{x}_{\mathbf{1}}\right)=\left[-\frac{m_{11} u}{m_{22}} \psi 0\right]^{T} . \tag{10.63}
\end{align*}
$$

The control design consists of two steps as follows:

## Step 1

Introduce the virtual yaw velocity error as

$$
\begin{equation*}
r_{e}=\hat{r}-\hat{r}_{d}, \tag{10.64}
\end{equation*}
$$

where $\hat{r}_{d}$ is the virtual yaw velocity control which is chosen as

$$
\begin{align*}
\hat{r}_{d}= & \frac{1}{1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{\hat{u}_{t}^{2}}}\left(-k_{1} \hat{w}_{e}^{*}+\frac{c(s) \hat{u}_{t}}{1-c(s) z_{e}} \cos \left(\hat{\psi}_{e}^{*}\right)+\right. \\
& \left.\frac{u}{\hat{u}_{t}^{2}}\left(\frac{d_{22}}{m_{22}} \hat{v}+\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}\right)-\frac{k \hat{u}_{t} \sin \left(\hat{\psi}_{e}^{*}\right)}{1+\left(k z_{e}\right)^{2}}-\frac{d_{1} \hat{w}_{e}^{*}}{\left(1-c(s) z_{e}\right)^{2}}\right), \tag{10.65}
\end{align*}
$$

where $k_{1}$ and $d_{1}$ are positive constants. The term multiplied by $d_{1}$ is the nonlinear damping term to overcome the effect of the observer errors. It is seen that $\hat{r}_{d}$ is a smooth function of $s, z_{e}, \psi_{e}$, and $\hat{v}$. Substituting (10.64) and (10.65) into the second equation of (10.62) results in

$$
\begin{equation*}
\dot{\hat{w}}_{e}^{*}=-k_{1} \hat{w}_{e}^{*}+r_{e}\left(1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{\hat{u}_{t}^{2}}\right)+\hat{\boldsymbol{f}}_{\mathbf{1}}^{\boldsymbol{T}}(\cdot) \tilde{\boldsymbol{x}}-\frac{d_{1} \hat{w}_{e}^{*}}{\left(1-c(s) z_{e}\right)^{2}}, \tag{10.66}
\end{equation*}
$$

where

$$
\hat{\boldsymbol{f}}_{\mathbf{1}}(\cdot)=\left[\begin{array}{c}
\frac{u \cos (\psi)}{\hat{u}_{t}^{2}}  \tag{10.67}\\
0 \\
\frac{c(s) \sin \left(\psi_{e}\right)}{1-c(s) z_{e}}+\frac{k \cos \left(\psi_{e}\right)}{1+\left(k z_{e}\right)^{2}} \\
1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{\hat{u}_{t}^{2}}
\end{array}\right], \tilde{\boldsymbol{x}}=\left[\begin{array}{c}
\tilde{y} \\
\tilde{\psi} \\
\tilde{v} \\
\tilde{r}
\end{array}\right] .
$$

## Step 2

Differentiating both sides of (10.64) yields

$$
\begin{align*}
\dot{r}_{e}= & \frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \hat{v}-\frac{d_{33}}{m_{33}} \hat{r}-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} \hat{r}+\frac{1}{m_{33}} \tau_{r}- \\
& \frac{\partial \hat{r}_{d}}{\partial s} \frac{\cos \left(\hat{\psi}_{e}^{*}\right)}{1-c(s) z_{e}}-\frac{\partial \hat{r}_{d}}{\partial z_{e}} \sin \left(\hat{\psi}_{e}^{*}\right)-\frac{\partial \hat{r}_{d}}{\partial \psi_{e}}\left(\hat{r}-\frac{c(s) \cos \left(\hat{\psi}_{e}^{*}\right)}{1-c(s) z_{e}}\right)- \\
& \frac{\partial \hat{r}_{d}}{\partial \hat{v}}\left(-\frac{m_{11} u}{m_{22}} \hat{r}-\frac{d_{22}}{m_{22}} \hat{v}-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}\right)+\hat{\boldsymbol{f}}_{2}^{\boldsymbol{T}}(\cdot) \tilde{\boldsymbol{x}} \tag{10.68}
\end{align*}
$$

where

$$
\hat{f}_{2}(\cdot)=\left[\begin{array}{c}
-\frac{\partial \hat{r}_{d}}{\partial \hat{v}} \cos (\psi)  \tag{10.69}\\
1 \\
\frac{\partial \hat{r}_{d}}{\partial s} \frac{\sin \left(\psi_{e}\right)}{1-c(s) z_{e}}-\frac{\partial \hat{r}_{d}}{\partial z_{e}} \cos \left(\psi_{e}\right)-\frac{\partial \hat{r}_{d}}{\partial \psi_{e}} \frac{c(s) \sin \left(\psi_{e}\right)}{1-c(s) z_{e}} \\
-\frac{\partial \hat{r}_{d}}{\partial \psi_{e}}-\frac{\partial \hat{r}_{d}}{\partial \hat{v}} \frac{m_{11}}{m_{22}} u
\end{array}\right]
$$

From (10.68), the control $\tau_{r}$ is designed without canceling the useful nonlinear damping terms as follows:

$$
\begin{align*}
& \tau_{r}=m_{33}\left(-k_{2} r_{e}-\frac{\left(m_{11}-m_{22}\right) u}{m_{33}} \hat{v}+\frac{d_{33}}{m_{33}} \hat{r}_{d}+\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} \hat{r}_{d}+\right. \\
& \quad \frac{\partial \hat{r}_{d}}{\partial s} \frac{\cos \left(\hat{\psi}_{e}^{*}\right)}{1-c(s) z_{e}}+\frac{\partial \hat{r}_{d}}{\partial z_{e}} \hat{u}_{t} \sin \left(\hat{\psi}_{e}^{*}\right)+\frac{\partial \hat{r}_{d}}{\partial \psi_{e}}\left(\hat{r}-\frac{c(s) \cos \left(\hat{\psi}_{e}^{*}\right)}{1-c(s) z_{e}}\right)-\frac{\partial \hat{r}_{d}}{\partial \hat{v}} \times \\
& \left.\quad\left(\frac{m_{11} u}{m_{22}} \hat{r}+\frac{d_{22}}{m_{22}} \hat{v}+\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|\hat{v}|^{i-1} \hat{v}\right)-\sum_{i=1}^{4} d_{2} \hat{f}_{2 i}^{2}(\cdot) r_{e}-\hat{w}_{e}^{*}\left(1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{\hat{u}_{t}^{2}}\right)\right), \tag{10.70}
\end{align*}
$$

where $k_{2}$ and $d_{2}$ are positive design constants, $\hat{f_{2 i}}(\cdot)$ is the $i$ th component of $\hat{\boldsymbol{f}_{2}}(\cdot)$. Substituting (10.70) into (10.68) yields

$$
\begin{align*}
\dot{r}_{e}= & -k_{2} r_{e}-\frac{d_{33}}{m_{33}} r_{e}-\sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} r_{e}-\left(1-\frac{m_{11}}{m_{22}} \frac{u^{2}}{\hat{u}_{t}^{2}}\right) \hat{w}_{e}^{*}+ \\
& \hat{\boldsymbol{f}}_{\mathbf{2}}^{\boldsymbol{T}}(\cdot) \tilde{\boldsymbol{x}}-\sum_{i=1}^{4} d_{2} \hat{f}_{2 i}^{2}(\cdot) r_{e} . \tag{10.71}
\end{align*}
$$

We now present the second main result of this chapter, the proof of which is given in the next section.

Theorem 10.2. Under Assumption 10.2, if the output feedback control law (10.70) and the observer (10.48) are applied to the ship system (10.1) then there exist feasible initial conditions such that the regulation errors $\left(z_{e}(t), \psi_{e}^{*}(t)\right)$ converge to a ball centered at the origin with an appropriate choice of the design constants $k, k_{i}$ and $d_{i}$ with $i=1,2$. Furthermore when there are no environmental disturbances, the regulation errors converge to zero asymptotically. In addition, the velocity $v(t)$ is always bounded.

### 10.3.3 Stability Analysis

To prove Theorem 10.2, we first consider the $\left(\hat{w}_{e}^{*}, r_{e}\right)$-dynamics then move to the $\left(\hat{v}, z_{e}\right)$ dynamics.

## $\left(\hat{w}_{e}^{*}, r_{e}\right)$-dynamics

Consider the following Lyapunov function

$$
\begin{equation*}
V_{1}=\frac{1}{2} \hat{w}_{e}^{* 2}+\frac{1}{2} r_{e}^{2} \tag{10.72}
\end{equation*}
$$

whose time derivative along the solutions of (10.66) and (10.71) satisfies

$$
\begin{align*}
\dot{V}_{1}= & -k_{1} \hat{w}_{e}^{* 2}+\hat{w}_{e}^{*} \hat{\boldsymbol{f}}_{\mathbf{1}}^{\boldsymbol{T}}(\cdot) \tilde{\boldsymbol{x}}-\frac{d_{1} \hat{w}_{e}^{* 2}}{\left(1-c(s) z_{e}\right)^{2}}-k_{2} r_{e}^{2}-\frac{d_{33}}{m_{33}} r_{e}^{2}- \\
& \sum_{i=2}^{3} \frac{d_{r i}}{m_{33}}|\hat{r}|^{i-1} r_{e}^{2}+r_{e} \hat{\boldsymbol{f}}_{\mathbf{T}}^{\boldsymbol{T}}(\cdot) \tilde{\boldsymbol{x}}-d_{2} r_{e}^{2} \sum_{i=1}^{4} \hat{f}_{2 i}^{2}(\cdot) \\
\leq & -\left(k_{1}-\mu_{11}\right) \hat{w}_{e}^{* 2}-k_{2} r_{e}^{2}+\mu_{12}\|\tilde{\boldsymbol{x}}\|^{2}, \tag{10.73}
\end{align*}
$$

where $\mu_{11}$ and $\mu_{12}$ are some positive constants. From (10.73) and (10.56), it is direct to show that

$$
\begin{equation*}
\left\|\left(\hat{w}_{e}^{*}(t), r_{e}(t)\right)\right\| \leq\left\|\left(\hat{w}_{e}^{*}\left(t_{0}\right), r_{e}\left(t_{0}\right)\right)\right\| e^{-\sigma_{1}\left(t-t_{0}\right)}+\rho_{1}, \tag{10.74}
\end{equation*}
$$

for some positive constants $\rho_{1}$ and $\sigma_{1}$. It is noted that $\rho_{1}$ can be made arbitrarily small by adjusting the design constants $k_{1}, k_{2}, d_{1}$, and $d_{2}$.

## $\left(\hat{v}, z_{e}\right)$-dynamics

To obtain convergence of $z_{e}$, we apply Lemma 10.1 to the first equation of (10.62) by verifying all conditions of this lemma. Indeed, the first equation of (10.62) is in the form of the system studied in Lemma 10.1 by writing this equation as

$$
\begin{equation*}
\dot{z}_{e}=f_{z e}(\cdot)+g_{z e}(\cdot), \tag{10.75}
\end{equation*}
$$

where

$$
\begin{align*}
f_{z e}(\cdot)= & -\frac{k \hat{u}_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}, \\
g_{z e}(\cdot)= & \frac{\hat{u}_{t} \sin \left(\hat{w}_{e}^{*}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}-\left(\cos \left(\hat{w}_{e}^{*}\right)-1\right) \frac{k \hat{u}_{t} z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}}+ \\
& \frac{\tilde{v} \cos \left(\hat{w}_{e}^{*}-\hat{\phi}\right)}{\sqrt{1+\left(k z_{e}\right)^{2}}}+\frac{\tilde{v} \sin \left(\hat{w}_{e}^{*}-\hat{\phi}\right) k z_{e}}{\sqrt{1+\left(k z_{e}\right)^{2}}} . \tag{10.76}
\end{align*}
$$

Verifying Condition C1. By taking the Lyapunov function $V_{2}=0.5 z_{e}^{2}$, it can be shown that this condition holds with

$$
\begin{equation*}
c_{0}=0, c_{1}=c_{2}=0.5, c_{3}=1, c=k^{2}, c_{4}=k \hat{u}_{t} \tag{10.77}
\end{equation*}
$$

Verifying Condition C2. From (10.76), we have

$$
\begin{equation*}
\left|g_{z e}(\cdot)\right| \leq \frac{1}{\sqrt{1+\left(k z_{e}\right)^{2}}}\left(\hat{u}_{t}+1+\left(\hat{u}_{t}+1\right) k\left|z_{e}\right|\right)\left\|\left(\tilde{v}, \hat{w}_{e}^{*}\right)\right\|, \tag{10.78}
\end{equation*}
$$

which implies that condition C 2 holds with

$$
\begin{equation*}
\lambda_{1}=\hat{u}_{t}+1, \lambda_{2}=k\left(\hat{u}_{t}+1\right) . \tag{10.79}
\end{equation*}
$$

Verifying Condition C3. This condition directly holds from (10.74) and (10.56).
Verifying Condition C4. From (10.77), (10.74), and (10.56), this condition becomes

$$
\begin{equation*}
k \hat{u}_{t}-\left(\rho_{1}+\xi_{0}\right)\left(\frac{\hat{u}_{t}+1}{4 \mu_{0}}+k\left(\hat{u}_{t}+1\right)\right)>0 . \tag{10.80}
\end{equation*}
$$

It is noted that $\rho_{1}$ can be made arbitrarily small but $\xi_{0}$ cannot, see (10.56). Therefore Condition C 4 is satisfied if the environmental disturbances are not too large and
the ship forward speed $u$ is high enough such that (10.80) holds for some positive constant $\mu_{0}$. This can be explained as follows: When there are large environmental disturbances, the observer cannot estimate sufficiently accurate velocities. In addition, the low ship speed cannot compensate for the environmental disturbances. Consequently, the ship under large environmental disturbances in question will diverge from the path. All the conditions of Lemma 10.1 have been verified, we have

$$
\begin{equation*}
\left|z_{e}(t)\right| \leq \alpha_{z}(\cdot) e^{-\sigma_{z}\left(t-t_{0}\right)}+\rho_{z} \tag{10.81}
\end{equation*}
$$

where $\alpha_{z}, \sigma_{z}$ and $\rho_{z}$ are calculated as in Lemma 10.1. The feasibility of initial conditions such that $1-c(s) z_{e} \geq \delta^{*}>0$ is similar to the case of the state feedback design, see (10.39).

Similar to the state feedback control design, it is not hard to show that

$$
\begin{equation*}
|v(t)| \leq \alpha_{v}(\cdot) e^{-\sigma_{v}\left(t-t_{0}\right)}+\rho_{v} \tag{10.82}
\end{equation*}
$$

where $\alpha_{v}, \sigma_{v}$ and $\rho_{v}$ are calculated similarly.

### 10.4 Simulations

This section validates the control laws (10.15) and (10.70) by simulating them on a monohull ship with the parameters given in Section 5.4. In the simulation, we assume that the disturbances are $\tau_{w v}=26 \times 10^{4}(1+\operatorname{rand}(\cdot))$, and $\tau_{w r}=$ $950 \times 10^{5}(1+\operatorname{rand}(\cdot))$, where $\operatorname{rand}(\cdot)$ is random zero-mean noise with the uniform distribution on the interval [ -0.50 .5 ]. This choice results in nonzero-mean disturbances. It should be noted that only boundaries of the environmental disturbances are needed in our proposed controllers. We simulate both state feedback and output feedback cases. For both cases, the ship forward speed is $4 \mathrm{~m} / \mathrm{s}$.

### 10.4.1 State Feedback Simulation Results

The initial conditions are chosen as

$$
[x(0), y(0), \psi(0), v(0), r(0)]=[-250,50,-0.5,0,0] .
$$

The reference path is a circle centered at the origin with a radius of 200 m . The control parameters are taken as $k=0.5, k_{1}=k_{2}=10$, and $\varepsilon_{i}=0.2$. Figure 10.3 plots the ship positions and orientation in the $(x, y)$-plane. The path-following errors are plotted in Figure 10.4. It can be seen from this figure that the orientation error $\psi_{e}^{*}$ converges close to zero but not $\psi_{e}$ due to nonzero sway velocity, and that the position path-following error, $z_{e}$, also converges closely to zero. The sway and yaw velocities and control torque $\tau_{r}$ are drawn in Figure 10.5. The high control effort
is due to the fact that we simulate on data of a real ship and large environmental disturbances. However, the control magnitude is within the limit of the maximum yaw moment. It is clearly seen from Figure 10.3 that the controller forces the vessel to move from its initial position to the path in the direction perpendicular to the circle. When the vessel closes to the path, the controller is able to drive it along the path in the desired direction. These observations coincide with our control objective. All the simulation results plotted in Figures 10.3, 10.4, and 10.5 illustrate that our control goal is achieved as stated in Theorem 10.1. It is noted that for clarity, we only plot the path-following errors, sway, and yaw velocities, and control input for the first 20 seconds.

### 10.4.2 Output Feedback Simulation Results

The initial conditions are

$$
\begin{aligned}
& {[x(0), y(0), \psi(0), v(0), r(0)]=[-250,50,-0.5,-0.2,-0.2]} \\
& {[\hat{x}(0), \hat{y}(0), \hat{\psi}(0), \hat{v}(0), \hat{r}(0)]=[-220,40,-0.2,-0.2,0.1,0.1] .}
\end{aligned}
$$

The observer gain matrix is matrix chosen as $\boldsymbol{K}_{\mathbf{0 1}}=\operatorname{diag}(1.5,1.5)$. The control gains are $k=0.5, k_{1}=k_{2}=15, d_{1}=d_{2}=0.5$, and $\varepsilon_{i}=0.2$. The environmental disturbances are the same as full state feedback case. Simulation results are plotted in Figures 10.6-10.8. All the comments on the results are similar to the state feedback simulation results. However, due to the effect of the observer, the performance of the output feedback case is slightly worse than the performance of the state feedback case.

### 10.5 Conclusions

Although the sway velocity $v$ is controlled in some conventional course-keeping control systems, see for example [12, 134], using Nomoto's second-order model, the sway displacement $y$ is not controlled and can be unbounded. This phenomenon can be seen from the second equation of (10.1), i.e., $\dot{y}=u \sin (\psi)+v \cos (\psi)$, that boundedness of the sway velocity $v$ does not imply that of the sway displacement $y$. The controllers proposed in this chapter did control the sway displacement $y$ and kept the sway velocity $v$ bounded, and cover the ones in [126] where the linear course stabilization of the underactuated surface ships is addressed. This can be seen from (10.2) by setting $z_{e}=y$ and the curvature of the path $c(s)$ equal zero. The work in this chapter is based on [135]. A combination of the results in this chapter with the results in [127] is straightforward to cover the case where the roll and pitch modes are not ignored. By setting the value of $k$ equal zero, the proposed controllers in this chapter reduce to the conventional course-keeping controllers,


Figure 10.3 State feedback, ship positions and orientation in the $(x, y)$-plane: Real path (solid line) and reference path (dash-dotted line)
see for example [11, 12, 134]. However, the sway displacement will grow linearly unbounded under nonzero-mean environmental disturbances.

The main limitation of both state feedback and output feedback controllers designed in this chapter is that the ship must not be too far away from the reference path at the initial time, i.e., the condition (10.3) holds at the initial time $t_{0}$. This limitation will be removed in the next chapter where a different approach for solving the path-following problem is proposed.


Figure 10.4 State feedback, path-following errors: a. Position error $z_{e}$; b. Modified heading error $\psi_{e}^{*} ; \mathbf{c}$. True heading error $\psi_{e}$


Figure 10.5 State feedback: a. Sway velocity $v ; \mathbf{b}$. Yaw velocity $r$; c. Control torque $\tau_{r}$


Figure 10.6 Output feedback, ship positions and orientation in the $(x, y)$-plane: Real path (solid line) and reference path (dash-dotted line)


Figure 10.7 Output feedback, path-following errors: a. Position error $z_{e}$; b. Modified heading error $\psi_{e}^{*}$; c. True heading error $\psi_{e}$


Figure 10.8 Output feedback results: a. Sway velocity $v$ (solid line) and its estimate $\hat{v}$ (dashdotted line); b. Yaw velocity $\boldsymbol{r}$ (solid line) and its estimate $\hat{r}$ (dash-dotted line); control torque $\tau_{r}$

## Chapter 11

## Path-following of Underactuated Ships Using Polar Coordinates

This chapter is devoted to a different path-following approach from the preceding chapter. The approach is motivated by the practical experience of steering a ship in the sense that when traveling in open sea the helmsman first looks at the weather map, then generates way-points to avoid the vessel moving into bad weather areas. A continuous reference path curve is then generated so that it goes via (almost) all of the way-points. It is then practical to steer a vessel so that it is in a tube of nonzero adjustable diameter centered on the reference path, and moves along the path with the desired speed. In comparison with the approach in the preceding chapter, the control system developed in this chapter allows the ship to be far away from the desired path at the initial time, and covers situations of practical importance such as parking and point-to-point navigation.

### 11.1 Control Objective

For the reader's convenience, the mathematical model of the underactuated ship moving in surge, sway, and yaw, see Sections 3.4.1.1 and 3.4.1.2, is rewritten as

$$
\begin{align*}
& \dot{x}=u \cos (\psi)-v \sin (\psi), \\
& \dot{y}=u \sin (\psi)+v \cos (\psi), \\
& \dot{\psi}=r, \\
& \dot{u}=\frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u-\sum_{i=2}^{3} \frac{d_{u i}}{m_{11}}|u|^{i-1} u+\frac{1}{m_{11}} \tau_{u}+\frac{1}{m_{11}} \tau_{w u}(t),  \tag{11.1}\\
& \dot{v}=-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t), \\
& \dot{r}=\frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r-\sum_{i=2}^{3} \frac{d_{r i}}{m_{11}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t),
\end{align*}
$$

where all of symbols have been defined in Section 3.3. The bounded time-varying terms, $\tau_{w u}(t), \tau_{w v}(t)$, and $\tau_{w r}(t)$, are the environmental disturbances induced by waves, wind, and ocean currents with $\left|\tau_{w u}(t)\right| \leq \tau_{w u \max }<\infty,\left|\tau_{w v}(t)\right| \leq \tau_{w v \max }<$ $\infty$, and $\left|\tau_{w r}(t)\right| \leq \tau_{w r \max }<\infty$.

In this chapter, we consider a control objective of designing the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$ to force the underactuated ship (11.1) to follow a specified path $\boldsymbol{\Omega}$, see Figure 11.1. In this figure, $P$ is the ship's center of mass and $P_{d}$ is a point attached to the virtual ship, which moves along the path with speed of $u_{0}$. If we are able to steer the ship to closely follow a virtual ship that moves along the path with a desired speed $u_{0}$, then the control objective is fulfilled, i.e., the ship is in a tube of nonzero diameter centered on the reference path and moves along the specified path at the speed $u_{0}$. Roughly speaking, the approach is to steer the ship so that it heads toward the virtual ship and eliminates the distance between itself and the virtual ship.


Figure 11.1 General framework of ship path-following

Define

$$
\begin{align*}
& x_{e}=x_{d}-x \\
& y_{e}=y_{d}-x \\
& \psi_{e}=\psi-\psi_{d}  \tag{11.2}\\
& z_{e}=\sqrt{x_{e}^{2}+y_{e}^{2}}
\end{align*}
$$

where

$$
\begin{equation*}
\psi_{d}=\arcsin \left(\frac{y_{e}}{z_{e}}\right) \tag{11.3}
\end{equation*}
$$

With (11.2), the above control objective is mathematically stated as follows.
Control Objective. Under Assumption 11.1, design the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$ to make the underactuated ship (11.1) follow the path $\Omega$ given by

$$
\begin{equation*}
x_{d}=x_{d}(s), y_{d}=y_{d}(s) \tag{11.4}
\end{equation*}
$$

where $s$ is the path parameter variable, such that

$$
\begin{align*}
& \lim _{t \rightarrow \infty} z_{e}(t) \leq \overline{z_{e}} \\
& \lim _{t \rightarrow \infty}\left|\psi_{e}(t)\right| \leq \overline{\psi_{e}} \tag{11.5}
\end{align*}
$$

with $\overline{z_{e}}$ and $\overline{\psi_{e}}$ being arbitrarily small positive constants.

## Assumption 11.1.

1. The reference path is regular, i.e., $0<R_{\min } \leq\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2} \leq R_{\max }<\infty$.
2. The minimum radius of the osculating circle of the path is larger than or equal to the minimum possible turning radius of the ship.

## Remark 11.1.

1. Assumption 11.1 ensures that the path is feasible for the ship to follow.
2. The angle $\psi_{d}$ is not defined at $z_{e}=0$ but $\lim _{z_{e} \rightarrow 0} \psi_{d}=\theta_{s}$ with $\theta_{s}$ being the orientation angle of the virtual ship on the path, see Figure 11.1. Therefore, the fulfillment of the control objective guarantees that the ship closely follows the path in terms of both position and orientation. If the reference path is not regular, then we can often split it into regular pieces and consider each of them separately. This is a case of point-to-point navigation.
3. The path parameter, $s$, is not the arclength of the path in general. For example, a circle with radius of $R$ centered at the origin can be described as $x_{d}=R \cos (s)$ and $y_{d}=R \sin (s)$, see [136] for more details.

If one differentiates both sides of $\psi_{e}=\psi-\psi_{d}$ to get the $\dot{\psi}_{e}$-dynamics, there will be discontinuity in the $\dot{\psi}_{e}$-dynamics on the $y_{e}$-axis. This discontinuity will cause difficulties in applying the backstepping technique. To get around this problem, we compute the $\dot{\psi}_{e}$-dynamics based on

$$
\begin{align*}
& \sin \left(\psi_{e}\right)=\frac{x_{e} \sin (\psi)-y_{e} \cos (\psi)}{z_{e}} \\
& \cos \left(\psi_{e}\right)=\frac{x_{e} \cos (\psi)+y_{e} \sin (\psi)}{z_{e}} \tag{11.6}
\end{align*}
$$

We now use (11.2) and (11.6) to transform (11.1) to

$$
\dot{z}_{e}=-\cos \left(\psi_{e}\right) u+\sin \left(\psi_{e}\right) v+\left(\frac{x_{e}}{z_{e}} \frac{\partial x_{d}}{\partial s}+\frac{y_{e}}{z_{e}} \frac{\partial y_{d}}{\partial s}\right) \dot{s},
$$

$$
\begin{align*}
\dot{\psi}_{e}= & r+\frac{\dot{s}}{\cos \left(\psi_{e}\right)}\left(\frac{\partial x_{d}}{\partial s}\left(\frac{\sin (\psi)}{z_{e}}-\frac{x_{e} \sin \left(\psi_{e}\right)}{z_{e}^{2}}\right)-\right. \\
& \left.\frac{\partial y_{d}}{\partial s}\left(\frac{\cos (\psi)}{z_{e}}+\frac{y_{e} \sin \left(\psi_{e}\right)}{z_{e}^{2}}\right)\right)+\frac{\sin \left(\psi_{e}\right) u+\cos \left(\psi_{e}\right) v}{z_{e}} \\
\dot{u}= & \frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u-\sum_{i=2}^{3} \frac{d_{u i}}{m_{11}}|u|^{i-1} u+\frac{1}{m_{11}} \tau_{u}+\frac{1}{m_{11}} \tau_{w u}(t), \\
\dot{v}= & -\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\frac{1}{m_{22}} \tau_{w v}(t) \\
\dot{r}= & \frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r-\sum_{i=2}^{3} \frac{d_{r i}}{m_{11}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+\frac{1}{m_{33}} \tau_{w r}(t) . \tag{11.7}
\end{align*}
$$

It is noted that the $\dot{\psi}_{e}$-dynamic is not defined at $\psi_{e}= \pm 0.5 \pi$. However, our controller will guarantee $\left|\psi_{e}(t)\right|<0.5 \pi, \forall 0 \leq t<\infty$ for feasible initial conditions. Therefore, we will design the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$ for (11.7) to yield the control objective. A procedure to design a stabilizer for the path-following error system (11.7) will be presented in detail. The triangular structure of (11.7) suggests that we design the actual controls $\tau_{u}$ and $\tau_{r}$ in two stages. First, we design the virtual velocity controls for $u, r$ and choose $\dot{s}$ to ultimately stabilize $z_{e}$ and $\psi_{e}$ at the origin. Based on the backstepping technique, controls $\tau_{u}$ and $\tau_{r}$ will then be designed to make the errors between the virtual velocity controls, and their actual values exponentially tend to a small ball centered at the origin. Since the ship parameters are unknown, an adaptation scheme is also introduced in this step to estimate their values used in the control laws. The nonzero lower bound of $z_{e}$ to guarantee the existence of $\psi_{d}$ and the boundedness of the sway velocity, $v$, are analyzed.

### 11.2 Control Design

### 11.2.1 Step 1

The $z_{e}$ and $\psi_{e}$ dynamics have three inputs that can be chosen to stabilize $z_{e}$ and $\psi_{e}$, namely $\dot{s}, u$, and $r$. The input $r$ should be designed to stabilize the $\psi_{e}$-dynamics at the origin. Therefore, two inputs, $\dot{s}, u$, can be used to ultimately stabilize $z_{e}$ at the origin. We can either choose the input $u$ or $\dot{s}$ and then design the remaining input. If we fix $\dot{s}$, then the virtual ship is allowed to move at a desired speed. The real ship will follow the virtual one on the path by the controller, and vice versa. In this chapter we choose to fix $\dot{s}$. This allows us to adjust the initial conditions in most cases without moving the ships, see Section 11.3. Since the transformed system (11.7) is not defined at $z_{e}=0$ and $\psi_{e}= \pm 0.5 \pi$, we first assume in the following
that $z_{e}(t) \geq z_{e}^{*}>0$ and $\left|\psi_{e}(t)\right|<0.5 \pi, \forall 0 \leq t<\infty$. In Section 11.3, we will then show that there exist initial conditions such that this hypothesis holds. Define

$$
\begin{align*}
& \tilde{u}=u-u_{d},  \tag{11.8}\\
& \tilde{r}=r-r_{d},
\end{align*}
$$

where $u_{d}$ and $r_{d}$ are the virtual controls of $u$ and $r$, respectively. As discussed above, we choose the virtual controls $u_{d}$ and $r_{d}$, and $\dot{s}$ as follows:

$$
\begin{equation*}
u_{d}=k_{1}\left(z_{e}-\delta_{e}\right)+\left(\frac{x_{e}}{z_{e}} \frac{\partial x_{d}}{\partial s}+\frac{y_{e}}{z_{e}} \frac{\partial y_{d}}{\partial s}\right) \frac{u_{0}\left(t, z_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}}+\tan \left(\psi_{e}\right) v \tag{11.9}
\end{equation*}
$$

where

$$
\begin{align*}
r_{d}= & -\left(\frac{\partial x_{d}}{\partial s}\left(\frac{\sin (\psi)}{z_{e}}-\frac{x_{e} \sin \left(\psi_{e}\right)}{z_{e}^{2}}\right)-\frac{\partial y_{d}}{\partial s}\left(\frac{\cos (\psi)}{z_{e}}+\frac{y_{e} \sin \left(\psi_{e}\right)}{z_{e}^{2}}\right)\right) \times \\
& \frac{u_{0}\left(t, z_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}}-\frac{\sin \left(\psi_{e}\right) u_{d}+\cos \left(\psi_{e}\right) v}{z_{e}}-k_{2} \psi_{e} \tag{11.10}
\end{align*}
$$

and

$$
\begin{equation*}
\dot{s}=\frac{u_{0}\left(t, z_{e}\right) \cos \left(\psi_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}} \tag{11.11}
\end{equation*}
$$

where $k_{1}, k_{2}$, and $\delta_{e}$ are positive constants to be selected later, $u_{0}\left(t, z_{e}\right) \neq 0, \forall t \geq$ $t_{0} \geq 0$, and $z_{e} \in \mathbb{R}$, is the speed of the virtual ship on the path. Indeed, one can choose this speed to be a constant. However, the time-varying speed and position pathfollowing dependence of the virtual ship on the path is more desirable, especially when the ship starts to follow the path. For example, one might choose

$$
\begin{equation*}
u_{0}\left(t, z_{e}\right)=u_{0}^{*}\left(1-\chi_{1} e^{-\chi_{2}\left(t-t_{0}\right)}\right) e^{-\chi_{3} z_{e}} \tag{11.12}
\end{equation*}
$$

where $u_{0}^{*} \neq 0, \chi_{i}>0, i=1,2,3, \chi_{1}<1$. The choice of $u_{0}\left(t, z_{e}\right)$ in (11.12) has the following desired feature: When the path-following error, $z_{e}$, is large, the virtual ship will wait for the real one; when $z_{e}$ is small, the virtual ship will move along the path at the speed closed to $u_{0}^{*}$ and the real one follows it within the specified look-ahead distance. This feature is suitable in practice because it avoids using a high gain control for large signal $z_{e}$.

Remark 11.2. If the sway velocity is assumed to be bounded by the surge velocity as in [129], the term $\sin \left(\psi_{e}\right) v$ is not required to be canceled. This controller can be designed similarly to the one in this chapter. Lemma 10.1 can be directly applied to the stability analysis. It is noted that the sway velocity is not needed to be bounded by the surge velocity with a relatively small constant.

Substituting (11.8)-(11.11) into the first two equations of (11.7) results in

$$
\begin{align*}
& \dot{z}_{e}=-k_{1} \cos \left(\psi_{e}\right)\left(z_{e}-\delta_{e}\right)-\cos \left(\psi_{e}\right) \tilde{u} \\
& \dot{\psi}_{e}=-k_{2} \psi_{e}+\frac{\sin \left(\psi_{e}\right)}{z_{e}} \tilde{u}+\tilde{r} \tag{11.13}
\end{align*}
$$

### 11.2.2 Step 2

By noting that with $z_{e}(t) \geq z_{e}^{*}>0$ and $\left|\psi_{e}(t)\right|<0.5 \pi, \forall 0 \leq t<\infty$, the virtual controls $u_{d}$ and $r_{d}$ are smooth functions of $x_{e}, y_{e}, s, u_{0}, \psi$ and $v$, differentiating both sides of (11.8) with (11.9) and (11.10) yields

$$
\begin{align*}
\dot{\tilde{u}}= & \frac{m_{22}}{m_{11}} v r-\frac{d_{11}}{m_{11}} u-\sum_{i=2}^{3} \frac{d_{u i}}{m_{11}}|u|^{i-1} u+\frac{1}{m_{11}} \tau_{u}+\frac{1}{m_{11}} \tau_{w u}(t)- \\
& \frac{\partial u_{d}}{\partial x_{e}} \dot{x}_{e}-\frac{\partial u_{d}}{\partial y_{e}} \dot{y}_{e}-\frac{\partial u_{d}}{\partial s} \dot{s}-\frac{\partial u_{d}}{\partial \psi} \dot{\psi}-\frac{\partial u_{d}}{\partial u_{0}} \dot{u}_{0}-\frac{\partial u_{d}}{\partial v} \times \\
& \left(-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v\right)-\frac{\partial u_{d}}{\partial v} \frac{1}{m_{22}} \tau_{w v}(t), \\
\dot{\tilde{r}}= & \frac{\left(m_{11}-m_{22}\right)}{m_{33}} u v-\frac{d_{33}}{m_{33}} r-\sum_{i=2}^{3} \frac{d_{r i}}{m_{11}}|r|^{i-1} r+\frac{1}{m_{33}} \tau_{r}+  \tag{11.14}\\
& \frac{1}{m_{33}} \tau_{w r}(t)-\frac{\partial r_{d}}{\partial x_{e}} \dot{x}_{e}-\frac{\partial r_{d}}{\partial y_{e}} \dot{y}_{e}-\frac{\partial r_{d}}{\partial s} \dot{s}-\frac{\partial r_{d}}{\partial \psi} \dot{\psi}-\frac{\partial r_{d}}{\partial u_{0}} \dot{u}_{0}- \\
& \frac{\partial r_{d}}{\partial v}\left(-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v\right)-\frac{\partial r_{d}}{\partial v} \frac{1}{m_{22}} \tau_{w v}(t),
\end{align*}
$$

where for convenience of choosing $u_{0}$, the terms $\frac{\partial u_{d}}{\partial x_{e}}, \frac{\partial u_{d}}{\partial y_{e}}, \frac{\partial r_{d}}{\partial x_{e}}$ and $\frac{\partial r_{d}}{\partial y_{e}}$ do not include $\frac{\partial u_{0}}{\partial x_{e}}$ and $\frac{\partial u_{0}}{\partial y_{e}}$, which are lumped into $\dot{u}_{0}$. From (11.14), we choose the actual controls $\tau_{u}$ and $\tau_{r}$ without canceling useful nonlinear damping terms as

$$
\begin{align*}
\tau_{u}= & -k_{3} \tilde{u}-\hat{\theta}_{1}^{T} f_{1}(\cdot)-\frac{\sin \left(\psi_{e}\right)}{z_{e}} \psi_{e}-\hat{\theta}_{31} \tanh \left(\frac{\tilde{u} \hat{\theta}_{31}}{\varepsilon_{1}}\right)- \\
& \hat{\theta}_{32} \frac{\partial u_{d}}{\partial v} \tanh \left(\frac{\partial u_{d}}{\partial v} \frac{\tilde{u} \hat{\theta}_{32}}{\varepsilon_{2}}\right),  \tag{11.15}\\
\tau_{r}= & -k_{4} \tilde{r}-\hat{\theta}_{2}^{T} f_{2}(\cdot)-\psi_{e}-\hat{\theta}_{33} \tanh \left(\frac{\tilde{r} \hat{\theta}_{33}}{\varepsilon_{3}}\right)-\hat{\theta}_{34} \frac{\partial r_{d}}{\partial v} \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{\tilde{r} \hat{\theta}_{34}}{\varepsilon_{4}}\right),
\end{align*}
$$

and the update laws as

$$
\begin{align*}
& \dot{\hat{\theta}}_{1 j}=\gamma_{1 j} \operatorname{proj}\left(\tilde{u} f_{1 j}(\cdot), \hat{\theta}_{1 j}\right), 1 \leq j \leq 9, \\
& \dot{\hat{\theta}}_{2 j}=\gamma_{2 j} \operatorname{proj}\left(\tilde{r} f_{2 j}(\cdot), \hat{\theta}_{2 j}\right), 1 \leq j \leq 9, \\
& \dot{\hat{\theta}}_{31}=\gamma_{31} \operatorname{proj}\left(|\tilde{u}|, \hat{\theta}_{31}\right), \\
& \dot{\hat{\theta}}_{32}=\gamma_{32} \operatorname{proj}\left(\left|\tilde{u} \frac{\partial u_{d}}{\partial v}\right|, \hat{\theta}_{32}\right),  \tag{11.16}\\
& \dot{\hat{\theta}}_{33}=\gamma_{33} \operatorname{proj}\left(|\tilde{r}|, \hat{\theta}_{33}\right), \\
& \dot{\hat{\theta}}_{34}=\gamma_{34} \operatorname{proj}\left(\left|\tilde{r} \frac{\partial r_{d}}{\partial v}\right|, \hat{\theta}_{34}\right)
\end{align*}
$$

where $k_{3}, k_{4}, \varepsilon_{i}, \gamma_{1 j}, \gamma_{2 j}, \gamma_{3 i}, 1 \leq i \leq 4$, and $1 \leq j \leq 9$ are positive constants to be selected later, and $f_{1 j}(\cdot)$ and $f_{2 j}(\cdot)$ are the $j$ th elements of $f_{1}(\cdot)$ and $f_{2}(\cdot)$, respectively, with

$$
\begin{align*}
f_{1}(\cdot)= & {\left[v r-u_{d}-|u| u_{d}-u^{2} u_{d}-\left(\frac{\partial u_{d}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial u_{d}}{\partial y_{e}} \dot{y}_{e}+\frac{\partial u_{d}}{\partial s} \dot{s}+\right.\right.} \\
& \left.\left.\frac{\partial u_{d}}{\partial \psi} \dot{\psi}+\frac{\partial u_{d}}{\partial u_{0}} \dot{u}_{0}\right) \frac{\partial u_{d}}{\partial v} u r \frac{\partial u_{d}}{\partial v} v \frac{\partial u_{d}}{\partial v}|v| v \frac{\partial u_{d}}{\partial v}|v|^{2} v\right]^{T}, \\
f_{2}(\cdot)= & {\left[u v-r_{d}-|r| r_{d}-r^{2} r_{d}-\left(\frac{\partial r_{d}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial r_{d}}{\partial y_{e}} \dot{y}_{e}+\frac{\partial r_{d}}{\partial s} \dot{s}+\right.\right.} \\
& \left.\left.\frac{\partial r_{d}}{\partial \psi} \dot{\psi}+\frac{\partial r_{d}}{\partial u_{0}} \dot{u}_{0}\right) \frac{\partial r_{d}}{\partial v} u r \frac{\partial r_{d}}{\partial v} v \frac{\partial r_{d}}{\partial v}|v| v \frac{\partial r_{d}}{\partial v}|v|^{2} v\right]^{T}, \tag{11.17}
\end{align*}
$$

$\hat{\theta}_{i j}, 1 \leq i \leq 3$ is the $j$ th element of $\hat{\theta}_{i}$, which is an estimate of $\theta_{i}$ with

$$
\begin{align*}
& \theta_{1}=\left[\begin{array}{llllll}
m_{22} & d_{11} & d_{u 2} & d_{u 3} m_{11} & \frac{m_{11}^{2}}{m_{22}} & \frac{d_{22} m_{11}}{m_{22}} \\
\frac{d_{v 2} m_{11}}{m_{22}} & \frac{d_{v 3} m_{11}}{m_{22}}
\end{array}\right]^{T}, \\
& \theta_{2}=\left[\begin{array}{lllll}
\left(m_{11}-m_{22}\right) d_{33} & d_{r 2} & d_{r 3} m_{33} & \frac{m_{11} m_{33}}{m_{22}} & \frac{d_{22} m_{33}}{m_{22}} \\
\frac{d_{v 2} m_{33}}{m_{22}} & \frac{d_{v 3} m_{33}}{m_{22}}
\end{array}\right]^{T}, \\
& \theta_{3}
\end{align*}=\left[\begin{array}{lll}
\tau_{w u \max } & \frac{m_{11}}{m_{22}} \tau_{w v \max } & \tau_{w r \max } \frac{m_{33}}{m_{22}} \tau_{w v \max } \tag{11.18}
\end{array}\right]^{T} . \quad .
$$

The operator, proj, is the Lipschitz continuous projection algorithm (repeated here for the reader's convenience) as follows:

$$
\begin{align*}
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \quad \text { if } \quad \Xi(\hat{\omega}) \leq 0, \\
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \quad \text { if } \quad \Xi(\hat{\omega}) \geq 0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi \leq 0,  \tag{11.19}\\
& \operatorname{proj}(\varpi, \hat{\omega})=(1-\Xi(\hat{\omega})) \varpi \quad \text { if } \Xi(\hat{\omega})>0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi>0,
\end{align*}
$$

where $\Xi(\hat{\omega})=\frac{\hat{\omega}^{2}-\omega_{M}^{2}}{\mu^{2}+2 \mu \omega_{M}}, \Xi_{\hat{\omega}}(\hat{\omega})=\frac{\partial \Xi(\hat{\omega})}{\partial \hat{\omega}}, \mu$ is an arbitrarily small positive constant, $\hat{\omega}$ is an estimate of $\omega$ and $|\omega| \leq \omega_{M}$.

The projection algorithm is such that if $\dot{\hat{\omega}}=\operatorname{proj}(\varpi, \hat{\omega})$ and $\hat{\omega}\left(t_{0}\right) \leq \omega_{M}$ then

1. $\hat{\omega}(t) \leq \omega_{M}+\xi, \forall 0 \leq t_{0} \leq t<\infty$,
2. $\operatorname{proj}(\varpi, \hat{\omega})$ is Lipschitz continuous,
3. $|\operatorname{proj}(\varpi, \hat{\omega})| \leq|\varpi|$,
4. $\tilde{\omega} \operatorname{proj}(\varpi, \hat{\omega}) \geq \tilde{\omega} \varpi$ with $\tilde{\omega}=\omega-\hat{\omega}$.

Substituting (11.15) into (11.14) yields

$$
\begin{align*}
\dot{\tilde{u}}= & -\left(\frac{k_{3}}{m_{11}}+\frac{d_{11}}{m_{11}}+\sum_{i=2}^{3} \frac{d_{u i}}{m_{11}}|u|^{i-1}\right) \tilde{u}-\frac{1}{m_{11}} \frac{\sin \left(\psi_{e}\right)}{z_{e}} \psi_{e}- \\
& \frac{1}{m_{11}} \hat{\theta}_{1}^{T} f_{1}(\cdot)+\frac{1}{m_{11}} \theta_{1}^{T} f_{1}(\cdot)-\frac{1}{m_{11}} \hat{\theta}_{31} \tanh \left(\frac{\tilde{u} \hat{\theta}_{31}}{\varepsilon_{1}}\right)+ \\
& \frac{1}{m_{11}} \tau_{w u}(t)-\hat{\theta}_{32} \frac{\partial u_{d}}{\partial v} \tanh \left(\frac{\partial u_{d}}{\partial v} \frac{\tilde{u} \hat{\theta}_{32}}{\varepsilon_{2}}\right)-\frac{\partial u_{d}}{\partial v} \frac{1}{m_{22}} \tau_{w v}(t), \\
\dot{\tilde{r}}= & -\left(\frac{k_{4}}{m_{33}}+\frac{d_{33}}{m_{33}}-\sum_{i=2}^{3} \frac{d_{r i}}{m_{11}}|r|^{i-1}\right) \tilde{r}-\frac{1}{m_{33}} \psi_{e}-\frac{1}{m_{33}} \hat{\theta}_{2}^{T} f_{2}(\cdot)+ \\
& \frac{1}{m_{33}} \theta_{2}^{T} f_{2}(\cdot)-\frac{1}{m_{33}} \hat{\theta}_{33} \tanh \left(\frac{\tilde{r} \hat{\theta}_{33}}{\varepsilon_{3}}\right)+\frac{1}{m_{33}} \tau_{w r}(t)- \\
& \frac{1}{m_{33}} \hat{\theta}_{34} \frac{\partial r_{d}}{\partial v} \tanh \left(\frac{\partial r_{d}}{\partial v} \frac{\tilde{r} \hat{\theta}_{34}}{\varepsilon_{4}}\right)-\frac{\partial r_{d}}{\partial v} \frac{1}{m_{22}} \tau_{w v}(t) . \tag{11.20}
\end{align*}
$$

We now present the following result, the proof of which is given in the next section.
Theorem 11.1. Assume that (a) the ship inertia, added mass and damping matrices are diagonal and (b) Assumption 11.1 is satisfied. If the adaptive state feedback control law (11.15) and adaptation law (11.16) are applied to the ship system (11.1) then there exist feasible initial conditions such that the path-following errors $\left(z_{e}(t), \psi_{e}(t)\right)$ converge to a small ball centered at the origin with an appropriate choice of the design constants $k_{i}, \varepsilon_{i}, \gamma_{1 j}, \gamma_{2 j}$, and $\gamma_{3 i}$ for all $1 \leq i \leq 4,1 \leq j \leq 9$.

### 11.3 Stability Analysis

To prove Theorem 11.1, we first consider the ( $\left.\psi_{e}, \tilde{u}, \tilde{r}\right)$-dynamics, then move to the $z_{e}$ - and $v$-dynamics.

## $\left(\psi_{e}, \tilde{u}, \tilde{r}\right)$-dynamics

Consider the following Lyapunov function

$$
\begin{equation*}
V_{1}=\frac{1}{2} \psi_{e}^{2}+\frac{m_{11}}{2} \tilde{u}^{2}+\frac{m_{33}}{2} \tilde{r}^{2}+\frac{1}{2} \sum_{i=1}^{3} \tilde{\theta}_{i}^{T} \Gamma_{i}^{-1} \tilde{\theta}_{i} \tag{11.21}
\end{equation*}
$$

where $\tilde{\theta}_{i}=\theta_{i}-\hat{\theta}_{i}$ and $\Gamma_{i}=\operatorname{diag}\left(\gamma_{i j}\right)$. Differentiating both sides of (11.21) along (11.14), (11.15), and (11.16) yields

$$
\begin{equation*}
\dot{V}_{1} \leq-k_{2} \psi_{e}^{2}-\left(k_{3}+d_{11}\right) \tilde{u}^{2}-\left(k_{4}+d_{33}\right) \tilde{r}^{2}+0.2785 \sum_{i=1}^{4} \varepsilon_{i} . \tag{11.22}
\end{equation*}
$$

By subtracting and adding $\frac{1}{2} \sum_{i=1}^{3} \tilde{\theta}_{i}^{T} \Gamma_{i}^{-1} \tilde{\theta}_{i}$ to the right-hand side of (11.22), we arrive at

$$
\begin{equation*}
\dot{V}_{1} \leq-\delta V_{1}+\rho, \tag{11.23}
\end{equation*}
$$

where

$$
\begin{align*}
& \delta=\min \left(1,2 k_{2}, \frac{2\left(k_{3}+d_{11}\right)}{m_{11}}, \frac{2\left(k_{4}+d_{33}\right)}{m_{33}}\right), \\
& \rho=\frac{1}{2} \sum_{i=1}^{3} \tilde{\theta}_{i}^{T} \Gamma_{i}^{-1} \tilde{\theta}_{i}+0.2785 \sum_{i=1}^{4} \varepsilon_{i} . \tag{11.24}
\end{align*}
$$

From (11.23), it is direct to show that

$$
\begin{equation*}
V_{1}(t) \leq V_{1}\left(t_{0}\right) e^{-\delta\left(t-t_{0}\right)}+\frac{\rho}{\delta}, \forall t \geq t_{0} \geq 0 \tag{11.25}
\end{equation*}
$$

which further yields

$$
\begin{align*}
& \min \left(\frac{1}{2}, \frac{m_{11}}{2}, \frac{m_{33}}{2}\right)\left\|X_{1}(t)\right\|^{2} \leq \\
& \quad\left\|X_{1}^{*}\left(t_{0}\right)\right\|^{2} \max \left(\frac{1}{2}, \frac{m_{11}}{2}, \frac{m_{33}}{2}, \frac{1}{2 \gamma_{1 j}}, \frac{1}{2 \gamma_{2 j}}, \frac{1}{2 \gamma_{3 j}}\right) e^{-\delta\left(t-t_{0}\right)}+\frac{\rho}{\delta}, \tag{11.26}
\end{align*}
$$

where

$$
\begin{aligned}
& X_{1}(t)=\left[\psi_{e}(t) \tilde{u}(t) \tilde{r}(t)\right]^{T} \\
& X_{1}^{*}\left(t_{0}\right)=\left[\psi_{e}\left(t_{0}\right) \tilde{u}\left(t_{0}\right) \tilde{r}\left(t_{0}\right) \tilde{\theta}_{1}^{T}\left(t_{0}\right) \tilde{\theta}_{2}^{T}\left(t_{0}\right) \tilde{\theta}_{3}^{T}\left(t_{0}\right)\right]^{T}
\end{aligned}
$$

Therefore we have

$$
\begin{equation*}
\left\|X_{1}(t)\right\| \leq \alpha_{1} e^{-\sigma_{1}\left(t-t_{0}\right)}+\rho_{1}, \forall t \geq t_{0} \geq 0 \tag{11.27}
\end{equation*}
$$

where

$$
\begin{align*}
p_{1} & =\sqrt{\frac{\max \left(\frac{1}{2}, \frac{m_{11}}{2}, \frac{m_{33}}{2}, \frac{1}{2 \gamma_{1 j}}, \frac{1}{2 \gamma_{2 j}}, \frac{1}{2 \gamma_{3 j}}\right)}{0.5 \min \left(1, m_{11}, m_{33}\right)}}, \\
\alpha_{1} & =p_{1}\left\|X_{1}^{*}\left(t_{0}\right)\right\| \\
\sigma_{1} & =\frac{\delta}{2} \\
\rho_{1} & =\sqrt{\frac{\rho}{0.5 \delta \min \left(1, m_{11}, m_{33}\right)}} . \tag{11.28}
\end{align*}
$$

Remark 11.3. It is important to note that, due to the use of the projection algorithm, by adjusting $k_{3}, k_{4}, \varepsilon_{i}, \gamma_{1 j}, \gamma_{2 j}, \gamma_{3 i}, 1 \leq i \leq 4,1 \leq j \leq 9$, we can make $\rho_{1}$ arbitrarily small. This observation plays a crucial role in the stability analysis of the $z_{e}$-dynamics.

## $z_{e}$-dynamics

Lower-bound of $z_{e}$. It can be seen from (11.13) and (11.27) that

$$
\begin{equation*}
\dot{\tilde{z}}_{e} \geq-k_{1} \tilde{z}_{e}-\left(\alpha_{1} e^{-\sigma_{1}\left(t-t_{0}\right)}+\rho_{1}\right) \tag{11.29}
\end{equation*}
$$

where $\tilde{z}_{e}=z_{e}-\delta_{e}$, which, with $\sigma_{1} \neq k_{1}$, further yields

$$
\begin{align*}
\tilde{z}_{e}(t) & \geq \tilde{z}_{e}\left(t_{0}\right) e^{-k_{1}\left(t-t_{0}\right)}+\frac{p_{1}\left\|X_{1}^{*}\left(t_{0}\right)\right\|}{\sigma_{1}-k_{1}}\left(e^{-\sigma_{1}\left(t-t_{0}\right)}-e^{-k_{1}\left(t-t_{0}\right)}\right)- \\
& \frac{\rho_{1}}{k_{1}}\left(1-e^{-k_{1}\left(t-t_{0}\right)}\right) \tag{11.30}
\end{align*}
$$

Hence

$$
\begin{align*}
& z_{e}(t) \geq\left(z_{e}\left(t_{0}\right)+\delta_{e}\right) e^{-k_{1}\left(t-t_{0}\right)}+\frac{p_{1}\left\|X_{1}^{*}\left(t_{0}\right)\right\|}{\sigma_{1}-k_{1}}\left(e^{-\sigma_{1}\left(t-t_{0}\right)}-e^{-k_{1}\left(t-t_{0}\right)}\right)- \\
& \quad \frac{\rho_{1}}{k_{1}}\left(1-e^{-k_{1}\left(t-t_{0}\right)}\right)+\delta_{e} \tag{11.31}
\end{align*}
$$

Therefore, the condition $z_{e}(t) \geq z_{e}^{*}>0$ holds when

$$
\begin{align*}
& \sigma_{1}>k_{1}, \delta_{e} \geq z_{e}^{*}+\frac{\rho_{1}}{k_{1}}, \\
& z_{e}\left(t_{0}\right) \geq \frac{p_{1}\left\|X_{1}^{*}\left(t_{0}\right)\right\|}{\sigma_{1}-k_{1}}-\delta_{e} \tag{11.32}
\end{align*}
$$

Upper Bound of $z_{e}$. We rewrite the first equation of (11.13) as

$$
\begin{equation*}
\dot{\tilde{z}}_{e}=-k_{1} \tilde{z}_{e}-\left(\cos \left(\psi_{e}\right)-1\right) k_{1} \tilde{z}_{e}-\cos \left(\psi_{e}\right) \tilde{u} \tag{11.33}
\end{equation*}
$$

It can be seen that (11.33) is of the form of the system studied in Lemma 10.1. Therefore, we will apply Lemma 10.1 to investigate stability of (11.33). We need to verify all conditions C1-C4 of Lemma 10.1 to (11.33).

Verifying Condition C1. Take the following Lyapunov function:

$$
\begin{equation*}
V_{2}=\frac{1}{2} z_{e}^{2} \tag{11.34}
\end{equation*}
$$

It is direct to show that C 1 holds with

$$
\begin{equation*}
c_{0}=0, c_{1}=c_{2}=0.5, c_{3}=1, c_{4}=k_{1} \tag{11.35}
\end{equation*}
$$

Verifying Condition C2. By noting that

$$
\left|\left(\cos \left(\psi_{e}\right)-1\right) k_{1} \tilde{z}_{e}-\cos \left(\psi_{e}\right) \tilde{u}\right| \leq\left(k_{1}\left|\tilde{z}_{e}\right|+1\right)\left\|X_{1}(t)\right\|,
$$

we have

$$
\begin{equation*}
\lambda_{1}=1, \lambda_{2}=k_{1} \tag{11.36}
\end{equation*}
$$

Verifying Condition C3. This condition directly holds from (11.27).
Verifying Condition C4. From (11.35), (11.36), and (11.27), condition C4 becomes

$$
\begin{equation*}
k_{1}-k_{1} \rho_{1}-0.25 \rho_{1} / \mu_{0}>0 \tag{11.37}
\end{equation*}
$$

From Remark 11.3 and noting that $\mu_{0}$ is an arbitrarily positive constant, we can see that there always exists $k_{1}$ such that (11.37) holds. All the conditions of Lemma 10.1 have been verified, and we therefore have

$$
\begin{equation*}
\left|z_{e}(t)\right| \leq \alpha_{2}\left(\left\|\left(X_{1}^{*}\left(t_{0}\right), z_{e}\left(t_{0}\right)\right)\right\|\right) e^{-\sigma_{2}\left(t-t_{0}\right)}+\rho_{2} \tag{11.38}
\end{equation*}
$$

where $\alpha_{2}, \sigma_{2}$, and $\rho_{2}$ are calculated as in Lemma 10.1.

## $v$-dynamics

From (11.9) and (11.10) the sway velocity dynamics can be rewritten as

$$
\begin{equation*}
\dot{v}=\phi_{1}(\cdot) v+\phi_{2}(\cdot) v^{2}-\frac{d_{v 2}}{m_{22}}|v| v-\frac{d_{v 3}}{m_{22}} v^{3}+\phi_{3}(\cdot), \tag{11.39}
\end{equation*}
$$

where

$$
\begin{align*}
\phi_{1}(\cdot)= & -\frac{d_{22}}{m_{22}}-\frac{m_{11}}{m_{22}}\left(\tan \left(\psi_{e}\right) \tilde{r}-\frac{1}{z_{e} \cos \left(\psi_{e}\right)} \tilde{u}+\right. \\
& \left.\left(a_{2}-\frac{\sin \left(\psi_{e}\right)}{z_{e}} a_{1}\right) \tan \left(\psi_{e}\right)-\frac{a_{1}}{z_{e} \cos \left(\psi_{e}\right)}\right), \\
\phi_{2}(\cdot)= & -\frac{m_{11}}{m_{22}} \frac{\tan \left(\psi_{e}\right)}{z_{e} \cos \left(\psi_{e}\right)}, \\
\phi_{3}(\cdot)= & -\frac{m_{11}}{m_{22}}\left(\tilde{u} \tilde{r}+a_{1} \tilde{r}+\left(a_{2}+\frac{\sin \left(\psi_{e}\right)}{z_{e}} a_{1}\right) \tilde{u}+\right. \\
& \left.\left(a_{2}+\frac{\sin \left(\psi_{e}\right)}{z_{e}} a_{1}\right) a_{1}\right)+\frac{1}{m_{22}} \tau_{w v}(t), \tag{11.40}
\end{align*}
$$

with

$$
\begin{aligned}
a_{1}= & k_{1}\left(z_{e}-\delta_{e}\right)+\left(\frac{x_{e}}{z_{e}} \frac{\partial x_{d}}{\partial s}+\frac{y_{e}}{z_{e}} \frac{\partial y_{d}}{\partial s}\right) \frac{u_{0}\left(t, z_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}} \\
a_{2}= & -\left(\frac{\partial x_{d}}{\partial s}\left(\frac{\sin (\psi)}{z_{e}}-\frac{x_{e} \sin \left(\psi_{e}\right)}{z_{e}^{2}}\right)-\frac{\partial y_{d}}{\partial s}\left(\frac{\cos (\psi)}{z_{e}}+\frac{y_{e} \sin \left(\psi_{e}\right)}{z_{e}^{2}}\right)\right) \times \\
& \frac{u_{0}\left(t, z_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}}-k_{2} \psi_{e}
\end{aligned}
$$

To show that $v$ is bounded, we take the following Lyapunov function:

$$
\begin{equation*}
V_{3}=\frac{1}{2} v^{2} \tag{11.41}
\end{equation*}
$$

whose derivative along the solutions of (11.39), for any $\eta>0$, satisfies

$$
\begin{equation*}
\dot{V}_{3} \leq-\left(\frac{d_{v 3}}{m_{22}}-\phi_{2}^{M} \eta\right) v^{4}+\left(\phi_{1}^{M}+\frac{\phi_{2}^{M}+\phi_{3}^{M}}{4 \eta}\right) v^{2}+\phi_{3}^{M} \eta \tag{11.42}
\end{equation*}
$$

where $\phi_{i}^{M}, 1 \leq i \leq 3$, are the upper bounds of $\phi_{i}(\cdot)$. It can be seen from (11.40) that $\phi_{i}^{M}$ exist and are finite since their arguments are bounded as shown above. Therefore, we can pick $\eta>0$ such that $d_{v 3} / m_{22}-\phi_{2}^{M} \eta>\eta^{*}>0$; then, see [6], the inequality (11.42) guarantees a finite upper-bound of the sway velocity $v$.

### 11.4 Discussion of the Initial Condition

We now discuss how to obtain the initial conditions such that $|\psi(t)|<0.5 \pi$ and the last inequality of (11.32) holds. Since $0.5 \psi^{2}(t) \leq V_{1}(t)$, from (11.25), the condition $|\psi(t)|<0.5 \pi$ can be rewritten as

$$
\begin{equation*}
\sqrt{\psi_{e}^{2}\left(t_{0}\right)+m_{11} \tilde{u}^{2}\left(t_{0}\right)+m_{33} \tilde{r}^{2}\left(t_{0}\right)+\sum_{i=1}^{3} \tilde{\theta}_{i}^{T}\left(t_{0}\right) \Gamma_{i}^{-1} \tilde{\theta}_{i}\left(t_{0}\right)+\rho / \delta}<\frac{1}{2} \pi \tag{11.43}
\end{equation*}
$$

On the other hand, the term $\sum_{i=1}^{3} \tilde{\theta}_{i}^{T}\left(t_{0}\right) \Gamma_{i}^{-1} \tilde{\theta}_{i}\left(t_{0}\right)+\frac{\rho}{\delta}$ can be made arbitrarily small. Noticing that $x_{e}\left(t_{0}\right)=x\left(t_{0}\right)-x_{d}\left(s\left(t_{0}\right)\right)$ and $y_{e}\left(t_{0}\right)=y\left(t_{0}\right)-y_{d}\left(s\left(t_{0}\right)\right)$, it can be seen that the initial value, $s\left(t_{0}\right)$, can be adjusted such that (11.43) holds if:

1. the ship heads toward "almost" of the half-plane containing the initial part of the path to be followed,
2. the path satisfies Assumption 11.1, see Figure 11.2,
3. the initial velocities $u\left(t_{0}\right), v\left(t_{0}\right)$ and $r\left(t_{0}\right)$ are not too large, and
4. the design constants are chosen such that $k_{1}$ is small and $k_{2}$ is large.

The angle $\delta_{0}$ (see Figure 11.2) should be increased if the initial velocities $u\left(t_{0}\right), v\left(t_{0}\right)$, and $r\left(t_{0}\right)$ are large. Otherwise the ship might cross the edge line of the plane in question, which might result in $\psi_{e}= \pm 0.5 \pi$. Similarly the last inequality of (11.32) is rewritten as

$$
\begin{equation*}
\sqrt{\left(x\left(t_{0}\right)-x_{d}\left(s\left(t_{0}\right)\right)\right)^{2}+\left(y\left(t_{0}\right)-y_{d}\left(s\left(t_{0}\right)\right)\right)^{2}} \geq \frac{p_{1}\left\|X_{1}^{*}\left(t_{0}\right)\right\|}{\sigma_{1}-k_{1}}-\delta_{e} . \tag{11.44}
\end{equation*}
$$

Again $s\left(t_{0}\right)$ can be adjusted such that (11.44) holds. It is noted that if the abovementioned conditions are not satisfied, one might generate an additional path segment such that it makes the above conditions hold and "smoothly" connects to the path to be followed.


Figure 11.2 Feasible initial conditions

### 11.5 Parking and Point-to-point Navigation

### 11.5.1 Parking

Parking Objective. Design the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$ to park the underactuated ship (11.1) from the initial position and orientation, $\left(x\left(t_{0}\right), y\left(t_{0}\right)\right.$, $\left.\psi\left(t_{0}\right)\right)$, to the desired parking position and orientation of $\left(x_{p}, y_{p}, \psi_{p}\right)$ under the following conditions:

1. There exists a large enough positive constant $\varpi_{p}$ such that

$$
\sqrt{\left(x\left(t_{0}\right)-x_{p}\right)^{2}+\left(y\left(t_{0}\right)-y_{p}\right)^{2}} \geq \varpi_{p} .
$$

2. The ship heads toward "almost" of the half-plane of the desired parking orientation.
3. At the desired parking position and orientation, the environmental disturbances are negligible.

The above conditions normally hold for parking practice. However, if the first two conditions do not hold, one can apply the strategy in the preceding section to move the ship until they hold. Having formulated the parking problem as above, one might claim that the path-following controller proposed in Section 11.2 can be applied by setting $u_{0}$ equal to zero. However, this will result in a yaw angle that may be very different from the desired parking one, at the desired parking position, since our proposed path-following controller is designed to drive $z_{\boldsymbol{e}}$ to a small ball, not to zero for reasons of robustness. To resolve this problem, we first generate a regular curve, $\boldsymbol{\Omega}_{\boldsymbol{p}}\left(x_{d}, y_{d}\right)$, which goes via the parking position and its tangent angle at the parking position is equal to the desired parking yaw angle, see Figure 11.3. For simplicity of calculation, the curve can be taken as a straight line in almost all cases of the vessel's initial conditions. Then the proposed path-following controller can be used to make the vessel follow $\boldsymbol{\Omega}_{\boldsymbol{p}}\left(x_{d}, y_{d}\right)$. In this case, the velocity $u_{0}$ should be chosen such that it is equal to zero or tends to zero when the virtual ship tends to the desired parking position, i.e., $\lim _{z_{e p} \rightarrow 0} u_{0}=0$ with $z_{e p}=\sqrt{\left(x_{d}-x_{p}\right)^{2}+\left(y_{d}-y_{p}\right)^{2}}$.

A simple choice can be taken as

$$
\begin{equation*}
u_{0}=u_{0}^{*}\left(1-e^{-\chi_{1} z_{e p}}\right) e^{-\chi_{2} z_{e}} \tag{11.45}
\end{equation*}
$$

where $\chi_{i}>0, i=1,2$. Special care should be taken when choosing the initial values of $\left(x_{d}\left(t_{0}\right), y_{d}\left(t_{0}\right)\right)$ and the sign of $u_{0}$ such that they result in a short parking time.

Remark 11.4. At the desired parking position and orientation, if there are large environmental disturbances, there will be an oscillatory behavior in the yaw dynamics and the ship might diverge from its desired position. This phenomenon is well known in dynamic positioning systems, [11]. However, for parking practice, we here assume that the environmental disturbances are negligible since the parking place is usually in a harbor.


Figure 11.3 Ship parking problem

### 11.5.2 Point-to-point Navigation

As seen in Section 11.1, the requirement of the reference path to be a regular curve might be too cumbersome in practice, since this curve has to go via desired points generated by the helmsman, and its derivatives are needed in the path-following controller. These restrictions motivate us to consider the point-to-point navigation problem as follows:

Point-to-point Navigation Objective. Design the surge force $\tau_{u}$ and the yaw moment $\tau_{r}$ to make the underactuated ship (11.1) go from the initial position and orientation, $\left(x\left(t_{0}\right), y\left(t_{0}\right), \psi\left(t_{0}\right)\right)$, via desired points generated by a path planner.

To solve this control objective, we first assume that the path planner generates desired points, which are feasible for the ship to be navigated through. We then apply the proposed path-following controller in Section 11.2 to each regular curve segment connecting desired points in sequence, see Figure 11.4. The regular curve segments can be straight line, arc, or known regular curves. A fundamental difference between point-to-point navigation and the proposed smooth path-following is that there are a finite number of "spikes", equal to the number of points, in the errors $z_{e}$ and $\psi_{e}$. This phenomenon is due to the path being nonsmooth in the orientation at the points.

### 11.6 Numerical Simulations

The same 32 m long monohull ship used in Section 5.4 is also used in this section. The ship parameters are again listed below for the convenience of the reader:


Figure 11.4 Point-to-point navigation

$$
\begin{aligned}
& m_{11}=120 \times 10^{3} \mathrm{~kg}, m_{22}=177.9 \times 10^{3} \mathrm{~kg}, m_{33}=636 \times 10^{5} \mathrm{kgm}^{2}, \\
& d_{u}=215 \times 10^{2} \mathrm{kgs}^{-1}, d_{u 2}=43 \times 10^{2} \mathrm{kgm}^{-1}, d_{u 3}=21.5 \times 10^{2} \mathrm{kgsm}^{-2}, \\
& d_{v}=117 \times 10^{3} \mathrm{kgs}^{-1}, d_{v 2}=23.4 \times 10^{3} \mathrm{kgm}^{-1}, d_{v 3}=11.7 \times 10^{3} \mathrm{kgsm}^{-2} \\
& d_{r}=802 \times 10^{4} \mathrm{kgm}^{2} s^{-1}, d_{r 2}=160.4 \times 10^{4} \mathrm{kgm}^{2}, d_{r 3}=80.2 \times 10^{4} \mathrm{kgm}^{2} s, \\
& d_{u i}=0, d_{v i}=0, d_{r i}=0, \forall i>3
\end{aligned}
$$

This ship has a minimum turning circle with a radius of 150 m , a maximum surge force of $5.2 \times 10^{9} \mathrm{~N}$, and the maximum yaw moment of $8.5 \times 10^{8} \mathrm{Nm}$. The above ship parameters are assumed to be those of real ships and are estimated on-line by adaptation laws (11.16). We assume that these parameters fluctuate around the above values $\pm 15 \%$ to calculate the maximum and minimum values used in the choice of the design constants in (11.37). We assume that the environmental disturbances are $\tau_{w u}=11 \times 10^{4}(1+\operatorname{rand}(\cdot)), \tau_{w v}=26 \times 10^{4}(1+\operatorname{rand}(\cdot))$ and $\tau_{w r}=950 \times 10^{5}(1+\operatorname{rand}(\cdot))$, where $\operatorname{rand}(\cdot)$ is random noise with the uniform distribution on the interval $[-0.50 .5]$. This choice results in nonzero-mean disturbances. In practice, the environmental disturbances may be different. We take the above disturbances for simplicity of generation. It should be noted that only the boundaries of the environmental disturbances are needed in our proposed controller. In simulations, the control parameters are taken as $k_{1}=1.5, k_{2}=7.5, k_{3}=6 \times 10^{5}$, $k_{4}=3 \times 10^{8}, \gamma_{1 j}=\gamma_{2 j}=\gamma_{3 j}=2 \times 10^{7}, \varepsilon_{i}=0.2$ and $\delta_{e}=0.95$. The initial conditions are $[x(0), y(0), \psi(0), u(0), v(0), r(0), s(0)]=[-200,-20,-0.4,0,0,0,0]$, and all initial values of parameter estimates are taken to be $70 \%$ of their assumed true ones.


Figure 11.5 Path-following: a. Ship position and orientation in the $(x, y)$-plane; $\mathbf{b}$. Ship position error $z_{e}$; c. Ship orientation error $\psi_{e}$

### 11.6.1 Path-following Simulation Results

The reference is $\left(x_{d}=s, y_{d}=200 \tanh (0.02 s)\right)$ for the first 120 seconds, then followed by a circle with a radius of 200 m . The virtual ship velocity on the path is taken as $u_{0}\left(t, z_{e}\right)=5\left(1-0.8 e^{-2 t}\right) e^{-0.5 z_{e}}$. The simulation results of the ship's position and orientation, path-following error $z_{e}$, and path-following orientation error $\psi_{e}$ are plotted in Figure 11.5. The control inputs $\tau_{u}$ and $\tau_{r}$ are plotted in Figure 11.6. For simplicity of presentation, we plot some samples of parameter estimates in Figure 11.7. It can be seen that the projection algorithm clearly prevents instability in adaptation due to non-vanishing disturbances.

### 11.6.2 Point-to-point Simulation Results

For the case of point-to-point navigation, we want the ship to go via points: $((0 \mathrm{~m}, 0 \mathrm{~m}),(400 \mathrm{~m}, 200 \mathrm{~m}),(1000 \mathrm{~m}, 200 \mathrm{~m}),(1400 \mathrm{~m}, 0 \mathrm{~m}))$, then follow a horizontal straight line. For simplicity, we use straight-line segments to connect the above desired points. The virtual ship velocity is taken to be the same as for the case of path-following. The simulation results of the ship position and orientation, path-


Figure 11.6 Path-following: a. Surge force $\tau_{u}[\mathrm{~N}]$; b. Yaw moment $\tau_{r}[\mathrm{Nm}]$
following error $z_{e}$, and path-following orientation error $\psi_{e}$ are plotted in Figures 11.8 and 11.9. The "spikes" in $z_{e}, \psi_{e}, \tau_{u}$ and $\tau_{r}$ are due to the nonsmooth reference path. However, these spikes are moderate in magnitude.

### 11.6.3 Parking Simulation Results

We assume that the parking position is at the origin and the parking orientation is $\psi_{p}=\pi / 4$ without environmental disturbances at the desired parking place. We choose $u_{0}=5\left(1-e^{-0.2 z_{e p}}\right) e^{-0.2 z_{e}}$ and a straight-line segment starting at $(-20,-20)$ with a slope of 1 . The simulation results are plotted in Figures 11.10 and 11.11.

In summary, it can be seen from Figures 11.5-11.11 that our proposed controller is able to force the underactuated ship in question to follow the predefined paths. The path-following position error $z_{e}$ converges to a small nonzero value specified by $\delta_{e}$ and does not cross zero as expected in the control design. It can be seen from Figures 11.5 and 11.8 that, under the nonvanishing environmental disturbances, the proposed controller forces the yaw angle to a small value. This value together with the ship forward speed prevents the ship from drifting away from the path. The high magnitude of controls results from the fact that we simulate the proposed con-


Figure 11.7 Path-following: a. Estimate $\hat{\theta}_{11}$; b. Estimate $\hat{\theta}_{21}$; c. Estimate $\hat{\theta}_{31}$
troller on a real ship and environmental disturbances with large magnitude. Indeed, the magnitude of the control inputs, can be reduced by adjusting the control and adaptation gains. However, this will result in a slow transient response. Finally, the transient response of $\psi_{e}$ is much shorter than that of $z_{e}$ due to the control gains chosen such that $k_{2} \gg k_{1}$ to guarantee a nonzero lower bound of $z_{e}$.

### 11.7 Conclusions

The proposed results in this chapter can be readily combined with the observer designs presented in Chapter 7 to design an output feedback path-following system for an underactuated ship. It is seen from Section 11.2 that the ship should not be too close to the desired path at the initial time. However, the strategy in Section 11.4 can be used to overcome the mentioned limitation of the control system proposed in this chapter. The work presented in this chapter is based on [137, 138]. The approach in this chapter will be extended to the case of the underactuated underwater vehicle in six degrees of freedom in Chapter 13.


Figure 11.8 Point-to-point navigation: a. Ship position and orientation in the $(x, y)$ plane; $\mathbf{b}$. Ship position error $z_{e} ; \mathbf{c}$. Ship orientation error $\psi_{e}$


Figure 11.9 Point-to-point navigation: a. Surge force $\tau_{u}[\mathrm{~N}]$; b.Yaw moment $\tau_{r}[\mathrm{Nm}]$


Figure 11.10 Parking: a. Ship position and orientation in the $(x, y)$ plane; $\mathbf{b}$. Ship position error; c. Ship orientation error in the $(x, y)$ plane


Figure 11.11 Parking: a. Surge force $\tau_{u}[\mathrm{~N}]$; b. Yaw moment $\tau_{r}[\mathrm{Nm}]$

## Part IV <br> Control of Underactuated Underwater Vehicles

# Chapter 12 <br> Trajectory-tracking Control of Underactuated Underwater Vehicles 

This chapter addresses the problem of trajectory-tracking control of underactuated underwater vehicles. These vehicles do not have independent actuators in the sway and heave axes. Based on the techniques developed for underactuated surface ships in Chapters 5, 6, and 7, we present a method to design a controller for an underactuated underwater vehicle to globally asymptotically track a reference trajectory generated by a suitable virtual underwater vehicle. The yaw and pitch reference velocities do not have to satisfy a persistently exciting condition as was often required in previous literature. Due to the complex dynamics of the underwater vehicles in comparison with that of the ships, the control design and stability analysis require more complicated coordinate transformations and techniques than those developed for underactuated ships in Chapters 5, 6, and 7.

### 12.1 Control Objective

In this chapter, we consider the following mathematical model of an underactuated underwater vehicle when nonlinear hydrodynamic damping terms and roll motion are ignored, see Section 3.4.2.2:

1. Kinematics

$$
\begin{align*}
\dot{x} & =\cos (\psi) \cos (\theta) u-\sin (\psi) v+\sin (\theta) \cos (\psi) w, \\
\dot{y} & =\sin (\psi) \cos (\theta) u+\cos (\psi) v+\sin (\theta) \sin (\psi) w, \\
\dot{z} & =-\sin (\theta) u+\cos (\theta) w, \\
\dot{\theta} & =q, \\
\dot{\psi} & =\frac{r}{\cos (\theta)} . \tag{12.1}
\end{align*}
$$

2. Kinetics

$$
\begin{align*}
& \dot{u}=\frac{m_{22}}{m_{11}} v r-\frac{m_{33}}{m_{11}} w q-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}, \\
& \dot{v}=-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v, \\
& \dot{w}=\frac{m_{11}}{m_{33}} u q-\frac{d_{33}}{m_{33}} w, \\
& \dot{q}=\frac{m_{33}-m_{11}}{m_{55}} u w-\frac{d_{55}}{m_{55}} q-\frac{\rho g \nabla \overline{G M_{L}} \sin (\theta)}{m_{55}}+\frac{1}{m_{55}} \tau_{q}, \\
& \dot{r}=\frac{m_{11}-m_{22}}{m_{66}} u v-\frac{d_{66}}{m_{66}} r+\frac{1}{m_{66}} \tau_{r} . \tag{12.2}
\end{align*}
$$

In (12.1) and (12.2) the symbols (see Figure 12.1) $\theta, \psi, q$, and $r$ denote the roll, pitch, and yaw angles and velocities while $x, y, z, u, v$, and $w$ are the surge, sway, and heave displacements and velocities, respectively. The available control inputs are $\tau_{u}, \tau_{q}$, and $\tau_{r}$. Since the sway and heave control forces are not available in the sway and heave dynamics, the underwater vehicle in question is underactuated. Notice that (12.1) is not defined when the pitch angle is equal to $\pm 90^{\circ}$. However, during practical operations with underwater vehicles, this problem is unlikely to happen due to the metacentric restoring forces. One way to avoid the singularities is to use a four-parameter description known as the quaternion. Here, we use the above Euler parameters because of their physical representation and computational efficiency. We first assume that $\left|\theta\left(t_{0}\right)\right|<0.5 \pi$. Then we find feasible initial conditions such


Figure 12.1 Motion variables of an underwater vehicle
that our proposed controller guarantees $|\theta(t)|<0.5 \pi, \forall t \geq t_{0} \geq 0$.
We consider a control objective of designing the control inputs $\tau_{u}, \tau_{q}$, and $\tau_{r}$ to force the underactuated vehicle given in (12.1) and (12.2) to asymptotically track a
reference trajectory generated by the following virtual vehicle

$$
\begin{align*}
\dot{x}_{d} & =\cos \left(\psi_{d}\right) \cos \left(\theta_{d}\right) u_{d}-\sin \left(\psi_{d}\right) v_{d}+\sin \left(\theta_{d}\right) \cos \left(\psi_{d}\right) w_{d}, \\
\dot{y}_{d} & =\sin \left(\psi_{d}\right) \cos \left(\theta_{d}\right) u_{d}+\cos \left(\psi_{d}\right) v_{d}+\sin \left(\theta_{d}\right) \sin \left(\psi_{d}\right) w_{d}, \\
\dot{z}_{d} & =-\sin \left(\theta_{d}\right) u_{d}+\cos \left(\theta_{d}\right) w_{d}, \\
\dot{\theta}_{d} & =q_{d}, \\
\dot{\psi}_{d} & =\frac{r_{d}}{\cos \left(\theta_{d}\right)}, \\
\dot{u}_{d} & =\frac{m_{22}}{m_{11}} v_{d} r_{d}-\frac{m_{33}}{m_{11}} w_{d} q_{d}-\frac{d_{11}}{m_{11}} u_{d}+\frac{1}{m_{11}} \tau_{u d},  \tag{12.3}\\
\dot{v}_{d} & =-\frac{m_{11}}{m_{22}} u_{d} r_{d}-\frac{d_{22}}{m_{22}} v_{d}, \\
\dot{w}_{d} & =\frac{m_{11}}{m_{33}} u_{d} q_{d}-\frac{d_{33}}{m_{33}} w_{d}, \\
\dot{q}_{d} & =\frac{m_{33}-m_{11}}{m_{55}} u_{d} w_{d}-\frac{d_{55}}{m_{55}} q_{d}-\frac{\rho g \nabla \overline{G M_{L}} \sin \left(\theta_{d}\right)}{m_{55}}+\frac{1}{m_{55}} \tau_{q d}, \\
\dot{r}_{d} & =\frac{m_{11}-m_{22}}{m_{66}} u_{d} v_{d}-\frac{d_{66}}{m_{66}} r_{d}+\frac{1}{m_{66}} \tau_{r d},
\end{align*}
$$

where all of the symbols in (12.3) have the same meaning as the real vehicle. In this section, we impose the following assumption on the reference model (12.3):

## Assumption 12.1.

1. The reference signals $u_{d}, q_{d}, r_{d}, \dot{u}_{d}, \dot{r}_{d}$, and $\dot{q}_{d}$ are bounded. There exists a strictly positive constant $u_{d \text { min }}$, such that $\left|u_{d}(t)\right| \geq u_{d \text { min }}, \forall t \geq 0$. The reference sway and heave velocities satisfy: $\left|v_{d}(t)\right|<\left|u_{d}(t)\right|,\left|w_{d}(t)\right|<\left|u_{d}(t)\right|, \forall t \geq 0$.
2. The reference pitch angle satisfies $\left|\theta_{d}(t)\right|<\frac{1}{2} \pi, \forall t \geq 0$.

Remark 12.1. The condition $\left|u_{d}(t)\right| \geq u_{d \min }, \forall t \geq 0$, covers both forward and backward tracking since the reference surge velocity is always nonzero but can be either positive or negative. Indeed, the condition $\left|u_{d}(t)\right| \geq u_{d \min }, \forall t \geq 0$, is much less restrictive than a persistently exciting condition on the yaw reference velocity. The condition $\left|v_{d}(t)\right|<\left|u_{d}(t)\right|$ and $\left|w_{d}(t)\right|<\left|u_{d}(t)\right|$ implies that the underactuated underwater vehicle cannot track a helix with an arbitrarily large curvature and twist due to the vehicle's high inertia and underactuation in the sway and heave directions.

### 12.2 Coordinate Transformations

Since designing the control inputs $\tau_{u}, \tau_{q}$, and $\tau_{r}$ to achieve the control directly from (12.1) and (12.2) is difficult, we interpret the tracking errors in a frame attached to the vehicle body as follows:

$$
\left[\begin{array}{c}
x_{e}  \tag{12.4}\\
y_{e} \\
z_{e} \\
\theta_{e} \\
\psi_{e}
\end{array}\right]=\left[\begin{array}{ccccc}
\cos (\psi) \cos (\theta) & \sin (\psi) \cos (\theta) & -\sin (\theta) & 0 & 0 \\
-\sin (\psi) & \cos (\psi) & 0 & 0 & 0 \\
\sin (\theta) \cos (\psi) & \sin (\theta) \sin (\psi) & \cos (\theta) & 0 & 0 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 1
\end{array}\right]\left[\begin{array}{c}
x-x_{d} \\
y-y_{d} \\
z-z_{d} \\
\theta-\theta_{d} \\
\psi-\psi_{d}
\end{array}\right]
$$

Indeed, convergence of $\left(x_{e}, y_{e}, z_{e}, \theta_{e}, \psi_{e}\right)$ to the origin implies that of $\left(x-x_{d}, y-\right.$ $\left.y_{d}, z-z_{d}, \theta-\theta_{d}, \psi-\psi_{d}\right)$ since the matrix

$$
\left[\begin{array}{ccccc}
\cos (\psi) \cos (\theta) & \sin (\psi) \cos (\theta) & -\sin (\theta) & 0 & 0 \\
-\sin (\psi) & \cos (\psi) & 0 & 0 & 0 \\
\sin (\theta) \cos (\psi) & \sin (\theta) \sin (\psi) & \cos (\theta) & 0 & 0 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 1
\end{array}\right]
$$

is nonsingular for all $(\psi, \theta) \in \mathbb{R}^{2}$. We also define the velocity tracking errors as

$$
\begin{align*}
& u_{e}=u-u_{d}, \\
& v_{e}=v-v_{d}, \\
& w_{e}=w-w_{d}  \tag{12.5}\\
& q_{e}=q-q_{d}, \\
& r_{e}=r-r_{d}
\end{align*}
$$

Differentiating (12.4) along the solutions of (12.1) and (12.3) yields:

$$
\begin{align*}
\dot{x}_{e}= & u_{e}-\left(\cos \left(\theta_{e}\right)-1+\cos (\theta) \cos \left(\theta_{d}\right)\left(\cos \left(\psi_{e}\right)-1\right)\right) u_{d}- \\
& \cos (\theta) \sin \left(\psi_{e}\right) v_{d}+\left(\sin \left(\theta_{e}\right)-\cos (\theta) \sin \left(\theta_{d}\right)\left(\cos \left(\psi_{e}\right)-1\right)\right) w_{d}+ \\
& \left(r_{d}+r_{e}\right) y_{e}-\left(q_{d}+q_{e}\right) z_{e}, \\
\dot{y}_{e}= & v_{e}+\cos \left(\theta_{d}\right) \sin \left(\psi_{e}\right) u_{d}-\left(\cos \left(\psi_{e}\right)-1\right) v_{d}-\sin \left(\theta_{d}\right) \sin \left(\psi_{e}\right) w_{d}- \\
& \left(x_{e}+\tan (\theta) z_{e}\right)\left(r_{d}+r_{e}\right), \\
\dot{z}_{e}= & w_{e}-\left(\sin \left(\theta_{e}\right)+\sin (\theta) \cos \left(\theta_{d}\right)\left(\cos \left(\psi_{e}\right)-1\right)\right) u_{d}-\sin (\theta) \sin \left(\psi_{e}\right) v_{d}- \\
& \left(\cos \left(\theta_{e}\right)-1+\sin (\theta) \sin \left(\theta_{d}\right)\left(\cos \left(\psi_{e}\right)-1\right)\right) w_{d}+\tan (\theta)\left(r_{d}+r_{e}\right) y_{e}+ \\
& \left(q_{d}+q_{e}\right) x_{e}, \\
\dot{\theta}_{e}= & q_{e}, \\
\dot{\psi}_{e}= & \frac{r_{e}}{\cos (\theta)}+\frac{r_{d}}{\cos (\theta) \cos \left(\theta_{d}\right)}\left(\cos \left(\theta_{d}\right)\left(1-\cos \left(\theta_{e}\right)\right)+\sin \left(\theta_{d}\right) \sin \left(\theta_{e}\right)\right) . \tag{12.6}
\end{align*}
$$

From (12.6), we can directly see that $x_{e}, \theta_{e}$, and $\psi_{e}$ can be stabilized by $u_{e}, q_{e}$, and $r_{e}$. There are several options to stabilize $y_{e}$ and $z_{e}$. We can either use $q_{e}, r_{e}$, $v_{e}, w_{e}, x_{e}, \theta_{e}$, or $\psi_{e}$. If $q_{e}$ and $r_{e}$ are used, the control design will be extremely
complicated since $q_{e}$ and $r_{e}$ enter all of the first three equations of (12.6). On the other hand, the use of $v_{e}$ and $w_{e}$ to stabilize $y_{e}$ and $z_{e}$ will result in an undesired feature of marine vehicle control practice, namely the vessel will slide in the sway and heave directions. If we use $x_{e}$ and $z_{e}$, the reference yaw and heave velocities must then satisfy persistently exciting conditions. Hence we will choose $\theta_{e}$ and $\psi_{e}$ to stabilize the sway error $y_{e}$ and the heave error $z_{e}$, respectively. This choice also coincides with the ship control practice. Toward this end, we define the following coordinates

$$
\begin{align*}
& z_{1}=\psi_{e}+\arcsin \left(\frac{k_{1} y_{e}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}+z_{e}^{2}}}\right), \\
& z_{2}=\theta_{e}-\arcsin \left(\frac{k_{2} z_{e}}{\sqrt{1+x_{e}^{2}+y_{e}^{2}+z_{e}^{2}}}\right), \tag{12.7}
\end{align*}
$$

where the constants $k_{i}, i=1,2$, are such that $\left|k_{i}\right|<1$. These constants will be specified later. It can be seen that (12.7) is well defined and that convergence of $\left(z_{1}, y_{e}\right)$ and $\left(z_{2}, z_{e}\right)$ implies that of $\theta_{e}$ and $\psi_{e}$.

Remark 12.2. Using the nonlinear coordinate transformations (12.7) instead of $z_{1}=$ $\psi_{e}+k_{1} y_{e}$ and $z_{2}=\theta_{e}-k_{2} z_{e}$, we avoid the vessel whirling around when $y_{e}$ and $z_{e}$ are large. If one uses the transformations $z_{1}=\psi_{e}+\arcsin \left(k_{1} y_{e} / \sqrt{1+y_{e}^{2}}\right)$ and $z_{2}=\theta_{e}-\arcsin \left(k_{2} z_{e} / \sqrt{1+z_{e}^{2}}\right)$, the problem of the vessel whirling around is avoided. Indeed using these coordinate transformations will result in a much simpler tracking error system than using (12.7). However, these coordinate transformations make the design of the control inputs $\tau_{u}, \tau_{q}$, and $\tau_{r}$ difficult because the terms $\left(x_{e}+\tan (\theta) z_{e}\right) r_{e}$ and $\left(\tan (\theta) r_{e} y_{e}+x_{e} q_{e}\right)$ appear in the $y_{e}$ - and $z_{e}$-dynamics.

Using the coordinate changes (12.7), the tracking error system in the ( $x_{e}, y_{e}, z_{e}$, $\left.z_{1}, z_{2}, u_{e}, v_{e}, w_{e}, q_{e}, r_{e}\right)$ coordinates can be written as:

$$
\begin{aligned}
\dot{x}_{e}= & u_{e}-\left(\varpi_{2}-\varpi+\cos (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} u_{d}+ \\
& \cos (\theta) k_{1} v_{d} \varpi^{-1} y_{e}+\left(k_{2} z_{e}-\cos (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} w_{d}+ \\
& \left(r_{d}+r_{e}\right) y_{e}-\left(q_{d}+q_{e}\right) z_{e}+p_{x}, \\
\dot{y}_{e}= & v_{e}-\cos \left(\theta_{d}\right) k_{1} u_{d} \varpi^{-1} y_{e}-\left(\varpi_{1}-\varpi\right) \varpi^{-1} v_{d}+\sin \left(\theta_{d}\right) k_{1} w_{d} \varpi^{-1} y_{e}- \\
& \left(x_{e}+\tan (\theta) z_{e}\right)\left(r_{d}+r_{e}\right)+p_{y}, \\
\dot{z}_{e}= & w_{e}-\left(k_{2} z_{e}+\sin (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} u_{d}+\sin (\theta) k_{1} v_{d} \varpi^{-1} y_{e}- \\
& \left(\varpi_{2}-\varpi+\sin (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} w_{d}+\tan (\theta)\left(r_{d}+r_{e}\right) y_{e}+ \\
& \left(q_{d}+q_{e}\right) x_{e}+p_{z}, \\
\dot{z}_{1}= & \left(1-k_{1} \varpi_{1}^{-1}\left(\cos (\theta) x_{e}+\sin (\theta) z_{e}\right)\right) \cos (\theta)^{-1} r_{e}-k_{1} \varpi_{1}^{-1} \varpi^{-2} x_{e} y_{e} u_{e}+ \\
& f_{z 1}+p_{z 1}, \\
\dot{z}_{2}= & \left(1-k_{2} \varpi_{2}^{-1} x_{e}\right) q_{e}-k_{2} \tan (\theta) \varpi_{2}^{-1} y_{e} r_{e}+k_{2} \varpi_{2}^{-1} \varpi^{-2} x_{e} z_{e} u_{e}+f_{z 2}+p_{z 2},
\end{aligned}
$$

$$
\begin{align*}
\dot{u}_{e} & =\frac{m_{22}}{m_{11}} v r-\frac{m_{33}}{m_{11}} w q-\frac{d_{11}}{m_{11}} u+\frac{1}{m_{11}} \tau_{u}-\dot{u}_{d}, \\
\dot{v}_{e} & =-\frac{m_{11}}{m_{22}}\left(u_{e} r_{e}+u_{e} r_{d}+u_{d} r_{e}\right)-\frac{d_{22}}{m_{22}} v_{e}, \\
\dot{w}_{e} & =\frac{m_{11}}{m_{33}}\left(u_{e} q_{e}+u_{e} q_{d}+u_{d} q_{e}\right)-\frac{d_{33}}{m_{33}} w_{e}, \\
\dot{q}_{e} & =\frac{m_{33}-m_{11}}{m_{55}} u w-\frac{d_{55}}{m_{55}} q-\frac{\rho g \nabla \overline{G M_{L}} \sin (\theta)}{m_{55}}+\frac{1}{m_{55}} \tau_{q}-\dot{q}_{d}, \\
\dot{r}_{e} & =\frac{m_{11}-m_{22}}{m_{66}} u v-\frac{d_{66}}{m_{66}} r+\frac{1}{m_{66}} \tau_{r}-\dot{r}_{d}, \tag{12.8}
\end{align*}
$$

where, for simplicity, we have defined the following notations:

1. The terms $\varpi, \varpi_{1}$ and $\varpi_{2}$ :

$$
\begin{align*}
& \varpi=\sqrt{1+x_{e}^{2}+y_{e}^{2}+z_{e}^{2}}, \\
& \varpi_{1}=\sqrt{1+x_{e}^{2}+\left(1-k_{1}^{2}\right) y_{e}^{2}+z_{e}^{2}},  \tag{12.9}\\
& \varpi_{2}=\sqrt{1+x_{e}^{2}+y_{e}^{2}+\left(1-k_{2}^{2}\right) z_{e}^{2}} .
\end{align*}
$$

2. The terms exponentially tend to zero when $z_{1}$ and $z_{2}$ do:

$$
\begin{align*}
p_{x}= & -\left(\left(\cos \left(z_{2}\right)-1\right) \varpi_{2}-\sin \left(z_{2}\right) k_{2} z_{e}+\cos (\theta) \cos \left(\theta_{d}\right)\left(\left(\cos \left(z_{1}\right)-\right.\right.\right. \\
& \left.\left.1) \varpi_{1}+\sin \left(z_{1}\right) k_{1} y_{e}\right)\right) \varpi^{-1} u_{d}+\left(\sin \left(z_{2}\right) \varpi_{2}+\left(\cos \left(z_{2}\right)-1\right) \times\right. \\
& \left.k_{2} z_{e}-\cos (\theta) \sin \left(\theta_{d}\right)\left(\left(\cos \left(z_{1}\right)-1\right) \varpi_{1}+\sin \left(z_{1}\right) k_{1} y_{e}\right)\right) \varpi^{-1} \times \\
& w_{d}-\cos (\theta)\left(\sin \left(z_{1}\right) \varpi_{1}-\left(\cos \left(z_{1}\right)-1\right) k_{1} y_{e}\right) \varpi^{-1} v_{d}, \\
p_{y}= & \cos \left(\theta_{d}\right)\left(\sin \left(z_{1}\right) \varpi_{1}-\left(\cos \left(z_{1}\right)-1\right) k_{1} y_{e}\right) \varpi^{-1} u_{d}-\left(\left(\cos \left(z_{1}\right)-\right.\right. \\
& \left.1) \varpi_{1}+\sin \left(z_{1}\right) k_{1} y_{e}\right) \varpi^{-1} v_{d}-\sin \left(\theta_{d}\right)\left(\sin \left(z_{1}\right) \varpi_{1}-\left(\cos \left(z_{1}\right)-\right.\right. \\
& \left.1) k_{1} y_{e}\right) \varpi^{-1} w_{d}, \\
p_{z}= & -\left(\sin \left(z_{2}\right) \varpi_{2}+\left(\cos \left(z_{2}\right)-1\right) k_{2} z_{e}+\sin (\theta) \cos \left(\theta_{d}\right)\left(\left(\cos \left(z_{1}\right)-\right.\right.\right. \\
& \left.\left.1) \varpi_{1}+\sin \left(z_{1}\right) k_{1} y_{e}\right)\right) \varpi^{-1} u_{d}-\left(\left(\cos \left(z_{2}\right)-1\right) \varpi_{2}-\sin \left(z_{2}\right) \times\right. \\
& \left.k_{2} z_{e}+\sin (\theta) \sin \left(\theta_{d}\right)\left(\left(\cos \left(z_{1}\right)-1\right) \varpi_{1}+\sin \left(z_{1}\right) k_{1} y_{e}\right)\right) \varpi^{-1} \times \\
& w_{d}-\sin (\theta)\left(\sin \left(z_{1}\right) \varpi_{1}-\left(\cos \left(z_{1}\right)-1\right) k_{1} y_{e}\right) \varpi^{-1} v_{d}, \\
p_{z 1}= & k_{1} \varpi_{1}^{-1}\left(p_{y}-y_{e} \varpi^{-2}\left(x_{e} p_{x}+y_{e} p_{y}+z_{e} p_{z}\right)\right)+\cos (\theta)^{-1} \varpi^{-1} \times \\
& \left(\left(1-\cos \left(z_{2}\right)\right) \varpi_{2}+\sin \left(z_{2}\right) k_{2} z_{e}+\tan \left(\theta_{d}\right)\left(\sin \left(z_{2}\right)+\right.\right. \\
& \left(\cos \left(z_{2}\right)-1\right) k_{2} z_{e}, \\
p_{z 2}= & -k_{2} \varpi_{2}^{-1}\left(p_{z}-z_{e} \varpi^{-2}\left(x_{e} p_{x}+y_{e} p_{y}+z_{e} p_{z}\right)\right) . \tag{12.10}
\end{align*}
$$

3. The terms $f_{z 1}$ and $f_{z 2}$ :

$$
\begin{align*}
& f_{z 1}=\left(\varpi-\varpi_{2}+\tan \left(\theta_{d}\right) k_{2} z_{e}\right) \cos (\theta)^{-1} \varpi^{-1}+k_{1} \varpi_{1}^{-1}\left(v_{e}-\cos \left(\theta_{d}\right) \times\right. \\
&\left.k_{1} u_{d} \varpi^{-1} y_{e}-\varpi_{1}-\varpi\right) \varpi^{-1} v_{d}+\sin \left(\theta_{d}\right) k_{1} w_{d} \varpi^{-1} y_{e}-\left(x_{e}+\right. \\
&\left.\tan (\theta) z_{e}\right) r_{d}-y_{e} \varpi^{-2}\left(x _ { e } \left(-\left(\varpi_{2}-\varpi+\cos (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \times\right.\right. \\
& \varpi^{-1} u_{d}+\cos (\theta) k_{1} v_{d} \varpi^{-1} y_{e}+\left(k_{2} z_{e}-\cos (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \times \\
&\left.\varpi^{-1} w_{d}\right)+y_{e}\left(v_{e}-\cos \left(\theta_{d}\right) k_{1} u_{d} \varpi^{-1} y_{e}-\left(\varpi_{1}-\varpi\right) \varpi^{-1} v_{d}+\right. \\
&\left.\sin \left(\theta_{d}\right) k_{1} w_{d} \varpi^{-1} y_{e}\right)+z_{e}\left(w_{e}-\left(k_{2} z_{e}+\sin (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \times\right. \\
& \varpi^{-1} u_{d}+\sin (\theta) k_{1} v_{d} \varpi^{-1} y_{e}-\left(\varpi_{2}-\varpi+\sin (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \times \\
&\left.\left.\left.\left.\varpi^{-1} w_{d}\right)\right)\right)\right), \\
& f_{z 2}==-k_{2} \varpi_{2}^{-1}\left(w_{e}-\left(k_{2} z_{e}+\sin (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} u_{d}+\sin (\theta) \times\right. \\
& k_{1} v_{d} \varpi^{-1} y_{e}-\left(\varpi_{2}-\varpi+\sin (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} w_{d}+x_{e} q_{d}+ \\
& \tan (\theta) y_{e} r_{d}-z_{e} \varpi^{-2}\left(x _ { e } \left(-\left(\varpi_{2}-\varpi^{\left.\cos (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \times}\right.\right.\right. \\
& \varpi^{-1} u_{d}+\cos (\theta) k_{1} v_{d} \varpi^{-1} y_{e}+\left(k_{2} z_{e}-\cos (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \times \\
&\left.\varpi^{-1} w_{d}\right)+y_{e}\left(v_{e}-\cos \left(\theta_{d}\right) k_{1} u_{d} \varpi^{-1} y_{e}-\left(\varpi_{1}-\varpi\right) \varpi^{-1} v_{d}+\right. \\
&\left.\sin \left(\theta_{d}\right) k_{1} w_{d} \varpi^{-1} y_{e}\right)+z_{e}\left(w_{e}-\left(k_{2} z_{e}+\sin (\theta) \cos \left(\theta_{d}\right) \times\right.\right. \\
&\left(\varpi_{1}-\varpi\right) \varpi^{-1} u_{d}+\sin (\theta) k_{1} v_{d} \varpi^{-1} y_{e}-\left(\varpi_{2}-\varpi+\right. \\
&\left.\left.\left.\left.\left.\sin (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} w_{d}\right)\right)\right)\right) . \tag{12.11}
\end{align*}
$$

It is now clear that the problem of forcing the underactuated underwater vehicle given in (12.1) and (12.2) to track the virtual ship (12.3) becomes one of stabilizing the system (12.8). The efforts we have made so far are to put the tracking error dynamics in the triangular form of (12.8), and to have the terms $-\cos \left(\theta_{d}\right) k_{1} u_{d} \varpi^{-1} y_{e}$ and $-k_{2} u_{d} \varpi^{-1} z_{e}$ in the $y_{e}$ and $z_{e}$-dynamics, respectively. These terms play an important role in stabilizing the $y_{e}$ - and $z_{e}$-dynamics. In the next section we will design the control inputs $\tau_{u}, \tau_{q}$, and $\tau_{r}$ to asymptotically stabilize (12.8) at the origin.

### 12.3 Control Design

The triangular structure of (12.8) suggests that we design the actual controls $\tau_{u}, \tau_{q}$, and $\tau_{r}$ in two stages. First, we design the virtual velocity controls of $u_{e}, q_{e}$, and $r_{e}$ to asymptotically stabilize $x_{e}, y_{e}, z_{e}, z_{1}, z_{2}, w_{e}$, and $v_{e}$ at the origin. Based on the backstepping technique, the controls $\tau_{u}, \tau_{q}$, and $\tau_{r}$ will then be designed to force the errors between the virtual velocity controls and their actual values exponentially to zero. Since $u_{e}$ enters the $v_{e}$ - and $w_{e}$-dynamics, we will design a bounded virtual control of $u_{e}$ to simplify the stability analysis. The virtual controls of $q_{e}$ and $r_{e}$ are chosen to stabilize the $z_{1}$ - and $z_{2}$-dynamics.

### 12.3.1 Step 1

Define the virtual control errors as

$$
\begin{equation*}
\tilde{u}_{e}=u_{e}-u_{e}^{d}, \tilde{q}_{e}=q_{e}-q_{e}^{d}, \tilde{r}_{e}=r_{e}-r_{e}^{d}, \tag{12.12}
\end{equation*}
$$

where $u_{e}^{d}, q_{e}^{d}$, and $r_{e}^{d}$ are the virtual velocity controls of $u_{e}, q_{e}$, and $r_{e}$, respectively. The virtual controls $u_{e}^{d}, q_{e}^{d}$, and $r_{e}^{d}$ are chosen as follows:

$$
\begin{align*}
& u_{e}^{d}=-k_{0} \varpi^{-1} x_{e}+\left(\varpi_{2}-\varpi+\cos (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} u_{d}- \\
& \cos (\theta) k_{1} v_{d} \varpi^{-1} y_{e}-\left(k_{2} z_{e}-\cos (\theta) \sin \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} w_{d}, \\
& r_{e}^{d}=r_{1 e}^{d}+r_{2 e}^{d} \text {, } \\
& q_{e}^{d}=q_{1 e}^{d}+q_{2 e}^{d}, \tag{12.13}
\end{align*}
$$

where

$$
\begin{align*}
r_{1 e}^{d} & =\frac{\cos (\theta)}{1-k_{1} \varpi_{1}^{-1}\left(\cos (\theta) x_{e}+\sin (\theta) z_{e}\right)}\left(k_{1} \varpi_{1}^{-1} \varpi^{-2} x_{e} y_{e} u_{e}^{d}-f_{z 1}\right), \\
r_{2 e}^{d} & =\frac{\cos (\theta)}{1-k_{1} \varpi_{1}^{-1}\left(\cos (\theta) x_{e}+\sin (\theta) z_{e}\right)}\left(-c_{1} z_{1}-p_{z 1}\right), \\
q_{1 e}^{d} & =\frac{1}{1-k_{2} \varpi_{2}^{-1} x_{e}}\left(k_{2} \tan (\theta) \varpi_{2}^{-1} y_{e} r_{1 e}^{d}-k_{2} \varpi_{2}^{-1} \varpi^{-2} x_{e} z_{e} u_{e}^{d}-f_{z 2}\right), \\
q_{2 e}^{d} & =\frac{1}{1-k_{2} \varpi_{2}^{-1} x_{e}}\left(-c_{2} z_{2}+k_{2} \tan (\theta) \varpi_{2}^{-1} y_{e} r_{1 e}^{d}-p_{z 2}\right), \tag{12.14}
\end{align*}
$$

with $k_{0}$ being a positive design constant to be specified later, and $c_{1}$ and $c_{2}$ being positive constants. It is not hard to show that the virtual control $u_{e}^{d}$ is bounded as

$$
\begin{equation*}
\left|u_{e}^{d}\right| \leq k_{0}+\left(1+k_{1}^{2}+k_{2}^{2}\right)\left|u_{d}\right|+\left|k_{1} v_{d}\right|+\left(\left|k_{2}\right|+k_{1}^{2}\right)\left|w_{d}\right|:=u_{e m} \tag{12.15}
\end{equation*}
$$

Remark 12.3. Unlike the standard application of backstepping, in order to reduce complexity of the controller expressions, we have chosen a simple virtual control law $u_{e}^{d}$ without canceling some known terms. We have written $r_{e}^{d}=r_{1 e}^{d}+r_{2 e}^{d}$ and $q_{e}^{d}=q_{1 e}^{d}+q_{2 e}^{d}$ to simplify stability analysis, because $q_{2 e}^{d}$ and $r_{2 e}^{d}$ exponentially vanish when $z_{1}$ and $z_{2}$ do. From (12.14), we observe that $r_{e}^{d}$ and $q_{e}^{d}$ are Lipschitz in $\left(x_{e}, y_{e}, z_{e}, z_{1}, z_{2}, v_{e}, w_{e}\right)$. This observation plays a crucial role in the stability analysis of the closed loop system.

### 12.3.2 Step 2

By differentiating (12.12) along the solutions of (12.14) and (12.6), the actual controls $\tau_{u}, \tau_{q}$, and $\tau_{r}$ without canceling the useful damping terms are chosen as

$$
\begin{align*}
\tau_{u}= & m_{11}\left(-\rho_{1} \tilde{u}_{e}-\frac{m_{22}}{m_{11}} v r+\frac{m_{33}}{m_{11}} w q+\frac{d_{11}}{m_{11}}\left(u_{e}^{d}+u_{d}\right)+\dot{u}_{d}+\right. \\
& \left.\dot{u}_{e}^{d}+k_{1} \varpi_{1}^{-1} \varpi^{-2} x_{e} y_{e} z_{1}-k_{2} \varpi_{2}^{-1} \varpi^{-2} x_{e} z_{e} z_{2}\right) \\
\tau_{r}= & m_{66}\left(-\rho_{2} \tilde{r}_{e}-\frac{m_{11}-m_{22}}{m_{66}} u v+\frac{d_{66}}{m_{66}}\left(r_{e}^{d}+r_{d}\right)+\dot{r}_{d}+\dot{r}_{e}^{d}-\right. \\
& \left.\left(1-k_{1} \varpi_{1}^{-1}\left(\cos (\theta) x_{e}+\sin (\theta) z_{e}\right)\right) \frac{z_{1}}{\cos (\theta)}+k_{2} \tan (\theta) \varpi_{2}^{-1} y_{e} z_{2}\right) \\
\tau_{q}= & m_{55}\left(-\rho_{3} \tilde{q}_{e}-\frac{m_{33}-m_{11}}{m_{55}} u w+\frac{d_{55}}{m_{55}}\left(q_{e}^{d}+q_{d}\right)+\right. \\
& \left.\frac{\rho g \nabla \overline{G M_{L}} \sin (\theta)}{m_{55}}+\dot{q}_{d}+\dot{q}_{e}^{d}-\left(1-k_{2} \varpi_{2}^{-1} x_{e}\right) z_{2}\right) \tag{12.16}
\end{align*}
$$

where $\rho_{i}, i=1,2,3$, are positive constants. Substituting (12.16), (12.13), and (12.14) into (12.8) yields the closed loop system

$$
\begin{align*}
\dot{x}_{e}= & -k_{0} \varpi^{-1} x_{e}-\left(q_{d}+q_{e}\right) z_{e}+\left(r_{d}+r_{e}\right) y_{e}+p_{x}+\tilde{u}_{e}, \\
\dot{y}_{e}= & v_{e}-\cos \left(\theta_{d}\right) k_{1} u_{d} \varpi^{-1} y_{e}-\left(\varpi_{1}-\varpi\right) \varpi^{-1} v_{d}+ \\
& \sin \left(\theta_{d}\right) k_{1} w_{d} \varpi^{-1} y_{e}-\left(x_{e}+\tan (\theta) z_{e}\right)\left(r_{d}+r_{e}\right)+p_{y}, \\
\dot{z}_{e}= & w_{e}-\left(k_{2} z_{e}+\sin (\theta) \cos \left(\theta_{d}\right)\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} u_{d}+ \\
& \sin (\theta) k_{1} v_{d} \varpi^{-1} y_{e}-\left(\varpi_{2}-\varpi+\sin (\theta) \sin \left(\theta_{d}\right) \times\right. \\
& \left.\left(\varpi_{1}-\varpi\right)\right) \varpi^{-1} w_{d}+\tan (\theta)\left(r_{d}+r_{e}\right) y_{e}+ \\
& \left(q_{d}+q_{e}\right) x_{e}+p_{z}, \\
\dot{v}_{e}= & -m_{11} m_{22}^{-1}\left(u_{e}^{d} r_{1 e}^{d}+u_{e}^{d} r_{d}+u_{d} r_{1 e}^{d}\right)-\frac{d_{22}}{m_{22}} v_{e}-m_{11} m_{22}^{-1} \times \\
& \left(\left(u_{e}^{d}+u_{d}\right)\left(r_{2 e}^{d}+\tilde{r}_{e}\right)+\tilde{u}_{e}\left(r_{e}^{d}+r_{d}+\tilde{r}_{e}\right)\right), \\
\dot{w}_{e}= & m_{11} m_{33}^{-1}\left(u_{e}^{d} q_{1 e}^{d}+u_{e}^{d} r_{d}+u_{d} q_{1 e}^{d}\right)-\frac{d_{33}}{m_{33}} w_{e}+m_{11} m_{33}^{-1} \times \\
& \left(\left(u_{e}^{d}+u_{d}\right)\left(q_{2 e}^{d}+\tilde{q}_{e}\right)+\tilde{u}_{e}\left(q_{e}^{d}+q_{d}+\tilde{q}_{e}\right)\right), \\
\dot{z}_{1}= & -c_{1} z_{1}+\left(1-k_{1} \varpi_{1}^{-1}\left(\cos (\theta) x_{e}+\sin (\theta) z_{e}\right)\right) \times \\
& \cos (\theta)^{-1} \tilde{r}_{e}-k_{1} \varpi_{1}^{-1} \varpi^{-2} x_{e} y_{e} \tilde{u}_{e}, \\
\dot{z}_{2}= & -c_{2} z_{2}+\left(1-k_{2} \varpi_{2}^{-1} x_{e}\right) \tilde{q}_{e}-k_{2} \tan (\theta) \varpi_{2}^{-1} y_{e} \tilde{r}_{e}+k_{2} \varpi_{2}^{-1} \varpi^{-2} x_{e} z_{e} \tilde{u}_{e}, \\
\dot{\tilde{u}}_{e}= & -\left(\rho_{1}+d_{11} m_{11}^{-1}\right) \tilde{u}_{e}+k_{1} \varpi_{1}^{-1} \varpi^{-2} x_{e} y_{e} z_{1}-k_{2} \varpi_{2}^{-1} \varpi^{-2} x_{e} z_{e} z_{2}, \\
\dot{\tilde{q}}_{e}= & -\left(\rho_{2}+d_{55} m_{55}^{-1}\right) \tilde{q}_{e}-\left(1-k_{2} \varpi_{2}^{-1} x_{e}\right) z_{2} \\
\dot{\tilde{r}}_{e}= & -\left(\rho_{3}+d_{66} m_{66}^{-1}\right) \tilde{r}_{e}+k_{2} \tan (\theta) \varpi_{2}^{-1} y_{e} z_{2}- \\
& \left(1-k_{1} \varpi_{1}^{-1}\left(\cos (\theta) x_{e}+\sin (\theta) z_{e}\right)\right) \cos (\theta)^{-1} z_{1} . \tag{12.17}
\end{align*}
$$

We now state the main result of this section, the proof of which is given in the next section.

Theorem 12.1. Assume that the reference signals $\left(x_{d}, y_{d}, z_{d}, \theta_{d}, \psi_{d}, v_{d}, w_{d}\right)$ generated by the virtual vessel model (12.3), and Assumption 12.1 holds. If the state feedback control law (12.16) is applied to the vessel system (12.1) and (12.2), then the tracking errors $\left(x(t)-x_{d}(t), y(t)-y_{d}(t), z(t)-z_{d}(t), \theta(t)-\theta_{d}(t)\right.$, $\psi(t)-\psi_{d}(t), v(t)-v_{d}(t)$, and $\left.w(t)-w_{d}(t)\right)$ asymptotically converge to zero with an appropriate choice of the design constants $k_{0}, k_{1}$, and $k_{2}$, i.e., the closed loop system (12.17) is locally asymptotically stable at the origin.

### 12.4 Stability Analysis

To prove Theorem 12.1, we just need to show that the closed loop system (12.17) is asymptotically stable at the origin. To simplify stability analysis of this closed loop system, we observe that (12.17) consists of two subsystems ( $x_{e}, y_{e}, z_{e}, v_{e}, w_{e}$ ) and $\left(z_{1}, z_{2}, \tilde{u}_{e}, \tilde{q}_{e}, \tilde{r}_{e}\right)$ in an interconnected structure. Therefore we first consider the $\left(z_{1}, z_{2}, \tilde{u}_{e}, \tilde{q}_{e}, \tilde{r}_{e}\right)$-subsystem then move to $\left(x_{e}, y_{e}, z_{e}, v_{e}, w_{e}\right)$-subsystem.

## $\left(z_{1}, z_{2}, \tilde{u}_{e}, \tilde{q}_{e}, \tilde{r}_{e}\right)$-subsystem

From the last five equations of (12.17), it is direct to show that this subsystem is exponentially stable at the origin by taking the following Lyapunov function

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(z_{1}^{2}+z_{2}^{2}+\tilde{u}_{e}^{2}+\tilde{q}_{e}^{2}+\tilde{r}_{e}^{2}\right) \tag{12.18}
\end{equation*}
$$

whose time derivative along the solutions of (12.17) satisfies
$\dot{V}_{1}=-c_{1} z_{1}^{2}-c_{2} z_{2}^{2}-\left(\rho_{1}+d_{11} m_{11}^{-1}\right) \tilde{u}_{e}^{2}-\left(\rho_{2}+d_{55} m_{55}^{-1}\right) \tilde{q}_{e}^{2}-\left(\rho_{3}+d_{66} m_{66}^{-1}\right) \tilde{r}_{e}^{2}$,
which in turn implies that

$$
\begin{align*}
& \left\|\left(z_{1}(t), z_{2}(t), \tilde{u}_{e}(t), \tilde{q}_{e}(t), \tilde{r}_{e}(t)\right)\right\| \leq \\
& \left\|\left(z_{1}\left(t_{0}\right), z_{2}\left(t_{0}\right), \tilde{u}_{e}\left(t_{0}\right), \tilde{q}_{e}\left(t_{0}\right), \tilde{r}_{e}\left(t_{0}\right)\right)\right\| e^{-\sigma_{1}\left(t-t_{0}\right)}, \tag{12.20}
\end{align*}
$$

where

$$
\sigma_{1}=\min \left(c_{1}, c_{2},\left(\rho_{1}+d_{11} m_{11}^{-1}\right),\left(\rho_{2}+d_{55} m_{55}^{-1}\right),\left(\rho_{3}+d_{66} m_{66}^{-1}\right)\right) .
$$

$\left(x_{e}, y_{e}, z_{e}, v_{e}, w_{e}\right)$-subsystem
To analyze stability of this subsystem, we consider the following Lyapunov function:

$$
\begin{equation*}
V_{2}=\sqrt{1+x_{e}^{2}+y_{e}^{2}+y_{e}^{2}}-1+\frac{1}{2} k_{3}\left(v_{e}^{2}+w_{e}^{2}\right), \tag{12.21}
\end{equation*}
$$

where $k_{3}$ is a positive constant to be specified later. The time derivative of (12.21) along the solutions of the first five equations of (12.17), after a lengthy but simple calculation using completed squares, satisfies

$$
\begin{align*}
\dot{V}_{2} \leq & -\mu_{x}(t) \varpi^{-2} x_{e}^{2}-\mu_{y}(t) \varpi^{-2} y_{e}^{2}-\mu_{z}(t) \varpi^{-2} z_{e}^{2}-\mu_{v}(t) v_{e}^{2}-\mu_{w}(t) w_{e}^{2}+ \\
& \left(\chi_{1}(\cdot) V_{2}+\chi_{2}(\cdot)\right) e^{-\sigma_{1}\left(t-t_{0}\right)}, \tag{12.22}
\end{align*}
$$

where $\chi_{i}(\cdot), i=1,2$ are some nondecreasing functions of $\|\left(z_{1}\left(t_{0}\right), z_{2}\left(t_{0}\right), \tilde{u}_{e}\left(t_{0}\right)\right.$, $\left.\tilde{q}_{e}\left(t_{0}\right), \tilde{r}_{e}\left(t_{0}\right)\right) \|$,

$$
\begin{align*}
\mu_{x}(t)= & k_{0}-\frac{k_{3} m_{11}\left|r_{d}\right| \varepsilon_{1}}{m_{22}}-\frac{k_{3} m_{11} u_{e m} \varepsilon_{1}}{m_{22}\left(1-2\left|k_{1}\right|\right)}\left(\frac{\left|k_{1} r_{d}\right|}{1-k_{1}^{2}}+\left|k_{1}^{3} u_{d}\right|+k_{1}^{2}\left|v_{d}\right|+\right. \\
& \left.\left(k_{1}^{2}+k_{2}^{2}\right)\left|k_{1} w_{d}\right|\right)-\frac{k_{3} m_{11}\left|q_{d}\right| \varepsilon_{1}}{m_{33}}-\frac{k_{3} m_{11} u_{e m}\left|k_{2} \tan (\theta)\right| \varepsilon_{1}}{m_{33}\left(1-\left|k_{2}\right|\right)} \times \\
& \left(\left|k_{1} r_{d}\right|+\left|k_{1}^{3} u_{d}\right|+k_{1}^{2}\left|v_{d}\right|+\left(k_{1}^{2}+k_{2}^{2}\right)\left|k_{1} w_{d}\right|+u_{e m}\right)- \\
& \frac{k_{3} m_{11} u_{e m} \varepsilon_{1}}{m_{33}\left(1-\left|k_{2}\right|\right)}\left(\frac{\left|k_{2} q_{d}\right|}{1-k_{1}^{2}}+\left|k_{1}^{3} u_{d}\right|+k_{1}^{2}\left|k_{2} u_{d}\right|\right), \tag{12.23}
\end{align*}
$$

$$
\begin{align*}
\mu_{y}(t)= & k_{1} u_{d} \cos \left(\theta_{d}\right)-\left|k_{1} w_{d}\right|-\varepsilon_{3}-k_{1}^{2}\left(\left|v_{d}\right|+\left|u_{d}\right|+\left|w_{d}\right|\right)- \\
& \frac{k_{3} m_{11} u_{e m} \varepsilon_{1}}{m_{22}\left(1-\left|k_{2}\right|\right)}\left(k_{1}^{2}\left(\left|u_{d}\right|+\left|w_{d}\right|\right)+\left|k_{1} v_{d}\right|\right)-\frac{k_{3} m_{11} u_{e m} \varepsilon_{1}}{m_{22}\left(1-2\left|k_{1}\right|\right)} \times \\
& \left(u_{e m}\left|k_{1}\right|+k_{1}^{2}\left(\left|v_{d}\right|+2\left|u_{d}\right|+2\left|w_{d}\right|\right)+\left|k_{1}^{3} w_{d}\right|+2\left|k_{1}^{3} v_{d}\right|+\right. \\
& \left.\left(k_{1}^{2}+k_{2}^{2}\right)\left|k_{1} u_{d}\right|\right)-\frac{k_{3} m_{11}\left|q_{d}\right| \varepsilon_{1}}{m_{33}}\left(k_{1}^{2}\left(\left|u_{d}\right|+\left|w_{d}\right|\right)+\left|k_{1} v_{d}\right|\right)- \\
& \frac{k_{3} m_{11} u_{e m}\left|k_{2} \tan (\theta)\right| \varepsilon_{1}}{m_{33}\left(1-\left|k_{2}\right|\right)}\left(u_{e m}\left|k_{1}\right|+k_{1}^{2}\left(\left|v_{d}\right|+2\left|u_{d}\right|+2\left|w_{d}\right|\right)+\right. \\
& \left.\left|k_{1}^{3}\right|\left(2\left|v_{d}\right|+\left|w_{d}\right|\right)+\left(k_{1}^{2}+k_{2}^{2}\right)\left|k_{1} u_{d}\right|\right)-\frac{k_{3} m_{11} u_{e m} \varepsilon_{1}}{m_{33}\left(1-\left|k_{2}\right|\right)}\left(k_{1}^{2}\left|k_{2}\right| \times\right. \\
& \left.\left(\left|u_{d}\right|+\left|w_{d}\right|\right)+\left|k_{1} k_{2}\right|\left(\left|v_{d}\right|+\left|u_{d}\right|\right)+\frac{\left|k_{2} r_{d}\right|}{1-k_{2}^{2}}\right)-\frac{\left|k_{1} v_{d}\right|}{4 \varepsilon_{2}}, \tag{12.24}
\end{align*}
$$

$$
\mu_{z}(t)=k_{2} u_{d}-\varepsilon_{3}-\left|k_{1} v_{d}\right|-k_{2}^{2}\left|w_{d}\right|-\frac{k_{3} m_{11}\left|r_{d}\right| \varepsilon_{1}}{m_{22}}\left(k_{2}^{2}\left|u_{d}\right|+\left|k_{2} w_{d}\right|\right)-
$$

$$
\left(\frac{k_{3} m_{11} u_{e m}\left|k_{2} \tan (\theta)\right| \varepsilon_{1}}{m_{33}\left(1-\left|k_{2}\right|\right)}+\frac{k_{3} m_{11} u_{e m} \varepsilon_{1}}{m_{22}\left(1-2\left|k_{1}\right|\right)}\right)\left(k_{2}^{2}+\frac{\left|k_{1} r_{d}\right|}{1-k_{1}^{2}}+\right.
$$

$$
\left.\left|k_{2}\right|+\left|k_{1} k_{2}\right|\left(\left|u_{d}\right|+\left|w_{d}\right|\right)\right)-\frac{k_{3} m_{11}\left|q_{d}\right| \varepsilon_{1}}{m_{33}}\left(\left|k_{2} w_{d}\right|+k_{2}^{2}\left|u_{d}\right|\right)-
$$

$$
\begin{align*}
& \frac{k_{3} m_{11} u_{e m} \varepsilon_{1}}{m_{33}\left(1-\left|k_{2}\right|\right)}\left(2\left|k_{2}^{3} w_{d}\right|+\left(k_{2}^{2}+k_{1}^{2}\left|k_{2}\right|\right)\left(\left|u_{d}\right|+\left|v_{d}\right|+\left|w_{d}\right|\right)+\right. \\
& \left.\left|k_{1} k_{2}\right|\left(\left|u_{d}\right|+\left|v_{d}\right|+2\left|w_{d}\right|\right),\right) \tag{12.25}
\end{align*}
$$

$$
\begin{align*}
\mu_{v}(t)= & \frac{k_{3} d_{22}}{m_{22}}-\frac{1}{4 \varepsilon_{3}}-\frac{k_{3} m_{11}\left|r_{d}\right|}{m_{22} 4 \varepsilon_{1}}\left(k_{0}^{2}+\left(k_{1}^{2}+k_{2}^{2}\right)\left|u_{d}\right|+\left|k_{1} v_{d}\right|+\right. \\
& \left.\left|k_{2} w_{d}\right|+k_{1}^{2}\left|w_{d}\right|\right)-\frac{k_{3} m_{11} u_{e m}}{m_{22}\left(1-2\left|k_{1}\right|\right) 4 \varepsilon_{1}}\left(| k _ { 1 } | \left(u_{e m}+2.5+\right.\right. \\
& \left.\left|r_{d}\right|\left(1+\tan ^{2}(\theta)\right)\right)+k_{2}^{2}+k_{1}^{2}\left(\left|u_{d}\right|+3\left|v_{d}\right|+2\left|w_{d}\right|\right)+\left(k_{1}^{2}+\right. \\
& \left.\left.k_{2}^{2}\right)\left|k_{1}\right|\left(\left|u_{d}\right|+\left|w_{d}\right|\right)+\left|k_{1} k_{2}\right|\left(\left|u_{d}\right|+\left|w_{d}\right|\right)\right)+\left|k_{1}^{3}\right| \times\left(\left|u_{d}\right|+2 \times\right. \\
& \left.\left.\left|v_{d}\right|+\left|w_{d}\right|\right)+\left|k_{2}\right|\right)-\frac{k_{3} k_{2}^{2} m_{11} u_{e m}|\tan (\theta)|}{m_{33}\left(1-\left|k_{2}\right|\right) 8 \varepsilon_{1}}-\frac{k_{3} k_{2}^{2} m_{11} u_{e m}}{m_{33}\left(1-\left|k_{2}\right|\right) 8 \varepsilon_{1}} \tag{12.26}
\end{align*}
$$

$$
\begin{align*}
\mu_{w}(t) & =\frac{k_{3} d_{33}}{m_{33}}-\frac{1}{4 \varepsilon_{3}}-\frac{k_{3} m_{11}\left|q_{d}\right|}{m_{33} 4 \varepsilon_{1}}\left(k_{0}^{2}+\left(k_{1}^{2}+k_{2}^{2}\right)\left|u_{d}\right|+\left|k_{1} v_{d}\right|+\right. \\
& \left.\left|k_{2} w_{d}\right|+k_{1}^{2}\left|w_{d}\right|\right)-\frac{k_{3} m_{11} u_{e m}\left|k_{2} \tan (\theta)\right|}{m_{33}\left(1-\left|k_{2}\right|\right) 8 \varepsilon_{1}}\left(| k _ { 1 } | \left(u_{e m}+2.5 \varepsilon_{1}+\right.\right. \\
& \left.\left|r_{d}\right|\left(1+\tan ^{2}(\theta)\right)\right)+k_{2}^{2}+k_{1}^{2}\left(\left|u_{d}\right|+3\left|v_{d}\right|+2\left|w_{d}\right|\right)+\left(k_{1}^{2}+\right. \\
& \left.k_{2}^{2}\right)\left|k_{1}\right|\left(\left|u_{d}\right|+\left|w_{d}\right|\right)+\left|k_{1} k_{2}\right|\left(\left|u_{d}\right|+\left|w_{d}\right|\right)+\left|k_{1}^{3}\right|\left(\left|u_{d}\right|+\left|v_{d}\right|+\right. \\
& \left.\left.\left|w_{d}\right|\right)+\left|k_{2}\right|\right)-\frac{k_{3}\left|k_{2}\right| m_{11} u_{e m}}{m_{33}\left(1-\left|k_{2}\right|\right) 4 \varepsilon_{1}}\left(k_{0}+\left(k_{1}^{2}+k_{2}^{2}\right)\left|u_{d}\right|+\left|k_{1} v_{d}\right|+\right. \\
& \left.\left(\left|k_{2}\right|+k_{2}^{2}\right)\left|w_{d}\right|\right)-\frac{k_{3} m_{11} u_{e m}}{m_{33}\left(1-\left|k_{2}\right|\right) 4 \varepsilon_{1}}\left(| k _ { 2 } | \left(2.5 \varepsilon_{1}+\tan ^{2}(\theta)+\right.\right. \\
& \left.\left|q_{d}\right|\right)+k_{1}^{2}\left|k_{2}\right|\left(3\left|u_{d}\right|+2\left|v_{d}\right|+2\left|w_{d}\right|\right)+k_{2}^{2}\left(\left|u_{d}\right|+2\left|w_{d}\right|\right)+ \\
& \left.\left|k_{2}^{3}\right|\left(\left|u_{d}\right|+2\left|w_{d}\right|\right)+2\left|k_{1} k_{2}\right|\left(\left|u_{d}\right|+\left|v_{d}\right|+\left|w_{d}\right|\right)\right), \tag{12.27}
\end{align*}
$$

with $\varepsilon_{i}, i=1,2,3$ being some positive constants, and

$$
\begin{equation*}
\theta(t)=\theta_{d}(t)+z_{2}(t)+\arcsin \left(\frac{k_{2} z_{e}(t)}{\sqrt{1+x_{e}^{2}(t)+y_{e}^{2}(t)+z_{e}^{2}(t)}}\right) \tag{12.28}
\end{equation*}
$$

We now choose the design constants $k_{i}, 0 \leq i \leq 3$ such that

$$
\begin{align*}
& \mu_{x} \geq \mu_{x}^{*} \\
& \mu_{y}(t) \geq \mu_{y}^{*} \\
& \mu_{z}(t) \geq \mu_{z}^{*} \\
& \mu_{v}(t) \geq \mu_{v}^{*}  \tag{12.29}\\
& \mu_{w}(t) \geq \mu_{w}^{*}
\end{align*}
$$

for all $t \geq t_{0} \geq 0$ for some positive constants $\mu_{x}^{*}, \mu_{y}^{*}, \mu_{z}^{*}, \mu_{v}^{*}, \mu_{w}^{*}$, and $\theta$ being replaced by

$$
\begin{equation*}
\theta\left(t_{0}\right)=\theta_{d}\left(t_{0}\right)+z_{2}\left(t_{0}\right)+\arcsin \left(\frac{k_{2} z_{e}\left(t_{0}\right)}{\sqrt{1+x_{e}^{2}\left(t_{0}\right)+z_{e}^{2}\left(t_{0}\right)+z_{e}^{2}\left(t_{0}\right)}}\right) \tag{12.30}
\end{equation*}
$$

Substituting (12.29) into (12.26) yields

$$
\begin{align*}
\dot{V}_{2} \leq & -\mu_{x}^{*} \varpi^{-2} x_{e}^{2}-\mu_{y}^{*} \varpi^{-2} y_{e}^{2}-\mu_{z}^{*} \varpi^{-2} z_{e}^{2}-\mu_{v}^{*} v_{e}^{2}-\mu_{w}^{*} w_{e}^{2}+ \\
& \left(\chi_{1}(\cdot) V_{2}+\chi_{2}(\cdot)\right) e^{-\sigma_{1}\left(t-t_{0}\right)} . \tag{12.31}
\end{align*}
$$

From (12.31) we have $\dot{V}_{2} \leq\left(\chi_{1}(\cdot) V_{2}+\chi_{2}(\cdot)\right) e^{-\sigma_{1}\left(t-t_{0}\right)}$, which implies that $V_{2}(t) \leq$ $\chi_{3}(\cdot)$ with $\chi_{3}(\cdot)$ being a class- $K$ function of $\|\left(x_{e}\left(t_{0}\right), y_{e}\left(t_{0}\right), z_{e}\left(t_{0}\right), v_{e}\left(t_{0}\right), w_{e}\left(t_{0}\right)\right.$, $\left.z_{1}\left(t_{0}\right), z_{2}\left(t_{0}\right), \tilde{u}_{e}\left(t_{0}\right), \tilde{q}_{e}\left(t_{0}\right), \tilde{r}_{e}\left(t_{0}\right)\right) \|$. Substituting $V_{2}(t) \leq \chi_{3}(\cdot)$ into (12.31) results in

$$
\begin{align*}
\dot{V}_{2} \leq & -\mu_{x}^{*} \varpi^{-2} x_{e}^{2}-\mu_{y}^{*} \varpi^{-2} y_{e}^{2}-\mu_{z}^{*} \varpi^{-2} z_{e}^{2}-\mu_{v}^{*} v_{e}^{2}-\mu_{w}^{*} w_{e}^{2}+ \\
& \left(\chi_{1}(\cdot) \chi_{3}(\cdot)+\chi_{2}(\cdot)\right) e^{-\sigma_{1}\left(t-t_{0}\right)} . \tag{12.32}
\end{align*}
$$

From (12.21) and (12.32) it is not hard to show that there exists a nonnegative constant $\sigma_{2}$ and a class- $K$ function $\gamma_{2}(\cdot)$ depending on the initial conditions such that

$$
\begin{equation*}
\left\|\left(x_{e}(t), y_{e}(t), z_{e}(t), v_{e}(t), w_{e}(t)\right)\right\| \leq \gamma_{2}(\cdot) e^{-\sigma_{2}\left(t-t_{0}\right)} \tag{12.33}
\end{equation*}
$$

The dependence of $\sigma_{2}>0$ on the initial conditions implies that the closed loop system (12.17) is asymptotically stable at the origin. However one can straightforwardly show that (12.17) is also locally exponentially stable at the origin. To complete the proof of Theorem 12.1, we need to show that there exist the design constants $k_{i}, 0 \leq i \leq 3$ such that $\left|k_{1}\right|<1,\left|k_{2}\right|<1$, condition (12.29) holds, and that $|\theta(t)|<0.5 \pi$. Before discussing these conditions, we note the following observations.

1. Under Assumption 12.1, the reference surge velocity $u_{d}$ is always nonzero, and the magnitudes of the reference sway and heave velocities are always less than that of the reference surge velocity.
2. The mass, including added masses in the sway and heave dynamics, $m_{22}$ and $m_{33}$, is often larger than that in the surge dynamics, $m_{11}$, for underwater vehicles, i.e., $m_{11} m_{22}^{-1}<1$ and $m_{11} m_{33}^{-1}<1$.

The condition of $\left|k_{1}\right|<1$ and $\left|k_{2}\right|<1$ can be satisfied easily by picking small enough $k_{1}$ and $k_{2}$. A close look at (12.29) shows that it always holds if we pick $k_{1}$ and $k_{2}$ such that they are small and have the same sign with the surge reference velocity $u_{d}$, and pick large enough $k_{3}$ for some small positive constants $\varepsilon_{i}, i=$ $1,2,3$. For example, one can choose $\varepsilon_{i}=\left|k_{1}\right|=\left|k_{2}\right|$. We can see from the above choice that the value of $\left|k_{1}\right|$ and $\left|k_{2}\right|$ should be decreased if $\left|u_{d}\right|$ is large. This physically means that the distance from the vessel to the point it aims to track should
be increased if the surge velocity is large, otherwise the vessel will miss that point. Furthermore when $d_{22} m_{22}^{-1}$ and $d_{33} m_{33}^{-1}$ are small, $\left|k_{1}\right|$ and $\left|k_{2}\right|$ should also be decreased. Small $\left|k_{1}\right|$ and $\left|k_{2}\right|$ also imply a small value of $k_{0}$. This can be physically interpreted as follows: If the damping in the sway and heave dynamics is small, the control gain in the surge dynamics should also be small otherwise the vessel will slide in the sway and heave directions.

The condition $|\theta(t)|<0.5 \pi$ can be written as

$$
|\theta(t)| \leq\left|\theta_{d}(t)\right|+\gamma_{2}(\cdot)+\left|\arcsin \left(k_{2} \gamma_{2}(\cdot)\right)\right|<0.5 \pi .
$$

Hence there always exist initial conditions such that this condition holds under Assumption 12.1. Due to complicated expressions of $\mu_{x}(t), \mu_{y}(t), \mu_{z}(t), \mu_{v}(t)$ and $\mu_{w}(t)$, we provide some general guidelines to choose the design constants rather than present their extremely complex explicit expressions.

Select small values for $\left|k_{1}\right|$ and $\left|k_{2}\right|$, set $\varepsilon_{i}=\left|k_{1}\right|, i=1,2,3$, large enough value for $k_{3}$. Then increase $k_{3}$ and/or decrease $\left|k_{1}\right|$ and $\left|k_{2}\right|$ until (12.29) holds.

Finally, we note that for the ease of choosing the design constants, one can replace the absolute values of $u_{d}, v_{d}, w_{d}$ and $r_{d}$ in all of the negative terms in $\mu_{x}(t), \mu_{y}(t), \mu_{z}(t), \mu_{v}(t)$ and $\mu_{w}(t)$ by their maximum values. The trade-off is that the control gains $k_{i}, 0 \leq i \leq 2$ may be very small, which results in slow convergence of the tracking errors, if the surge reference velocity $u_{d}(t)$ varies largely, i.e. $u_{d \text { max }} \gg u_{d \text { min }}$ with $u_{d \text { max }}$ and $u_{d \text { min }}$ being the maximum and minimum values of $u_{d}(t)$, respectively.

### 12.5 Simulations

This section illustrates the effectiveness of the control law (12.16) by simulating it on an underwater vehicle with a length of 5.56 m , a mass of 1089.8 kg , and other parameters taken from [139] as follows:

$$
\begin{aligned}
& m_{11}=1116 \mathrm{~kg}, m_{22}=2133 \mathrm{~kg}, m_{33}=2133 \mathrm{~kg}, m_{44}=36.7 \mathrm{kgm}^{2}, \\
& m_{55}=4061 \mathrm{kgm}^{2}, m_{66}=4061 \mathrm{kgm}^{2}, d_{11}=25.5 \mathrm{kgs}^{-1}, d_{22}=138 \mathrm{kgs}^{-1}, \\
& d_{33}=138 \mathrm{kgs}^{-1}, d_{44}=10 \mathrm{kgm}^{2} \mathrm{~s}^{-1}, d_{55}=490 \mathrm{kgm}^{2} \mathrm{~s}^{-1}, d_{66}=490 \mathrm{kgm}^{2} \mathrm{~s}^{-1}, \\
& d_{u 2}=0, d_{u 3}=0, d_{p 2}=0, d_{p 3}=0, d_{q 2}=0, d_{q 3}=0, d_{r 2}=0, d_{r 3}=0, \\
& d_{v 2}=920.1 \mathrm{kgm}^{-2} \mathrm{~s}, d_{v 3}=750 \mathrm{kgm}^{-3} \mathrm{~s}^{2}, d_{w 2}=920.1 \mathrm{kgm}^{-2} \mathrm{~s}, \\
& d_{w 3}=750 \mathrm{kgm}^{-3} \mathrm{~s}^{2} .
\end{aligned}
$$

This vehicle has a minimum turning circle with a radius of 75 m , a maximum surge force of $2 \times 10^{4} \mathrm{~N}$, a maximum yaw moment of $1.5 \times 10^{4} \mathrm{Nm}$, a maximum pitch moment of $1.5 \times 10^{4} \mathrm{Nm}$, and the maximum roll moment of 120 Nm .

The reference trajectory is generated by (12.3) with

$$
\begin{align*}
& \tau_{u d}=5 d_{11}-\left(m_{22} v_{d} r_{d}-m_{33} w_{d} q_{d}\right) \\
& \tau_{q d}=\left(-\theta_{d}+0.2-\dot{\theta}_{d}\right) m_{55}-\left(m_{33}-m_{11}\right) u_{d} w_{d}-d_{55} q_{d}-\rho g \nabla \overline{G M_{L}} \sin \left(\theta_{d}\right), \\
& \tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} v_{d} \tag{12.34}
\end{align*}
$$

for the first 200 seconds, and

$$
\begin{align*}
& \tau_{u d}=5 d_{11}-\left(m_{22} v_{d} r_{d}-m_{33} w_{d} q_{d}\right), \\
& \tau_{q d}=\left(-\theta_{d}+0.2-\dot{\theta}_{d}\right) m_{55}-\left(m_{33}-m_{11}\right) u_{d} w_{d}-d_{55} q_{d}-\rho g \nabla \overline{G M_{L}} \sin \left(\theta_{d}\right), \\
& \tau_{r d}=-\left(m_{11}-m_{22}\right) u_{d} v_{d}+0.02 d_{66} \tag{12.35}
\end{align*}
$$

for the rest of simulation time. This choice means that the reference trajectory is a straight line for the first 200 seconds followed by a helix with constant curvature and torsion. The initial conditions are picked as follows:

$$
\begin{align*}
& \left(x\left(t_{0}\right), y\left(t_{0}\right), z\left(t_{0}\right), \theta\left(t_{0}\right), \psi\left(t_{0}\right), u\left(t_{0}\right), v\left(t_{0}\right), w\left(t_{0}\right), q\left(t_{0}\right), r\left(t_{0}\right)\right)= \\
& \quad(-50,-50,0,0,0,0,0,0,0,0), \\
& \left(x_{d}\left(t_{0}\right), y_{d}\left(t_{0}\right), z_{d}\left(t_{0}\right), \theta_{d}\left(t_{0}\right), \psi_{d}\left(t_{0}\right), u_{d}\left(t_{0}\right), v_{d}\left(t_{0}\right), w_{d}\left(t_{0}\right), q_{d}\left(t_{0}\right), r_{d}\left(t_{0}\right)\right)= \\
& \quad(0,0,20,0,0,0,0,10,0,0,0) . \tag{12.36}
\end{align*}
$$

Based on the proof of Theorem 12.1, the design constants are chosen as $k_{0}=0.8$, $k_{1}=k_{2}=0.4, c_{i}=2$, and $\rho_{i}=5$ with $i=1,2,3$. The reference and real trajectories in three dimensions are plotted in Figure 12.2. The vessel position and orientation are plotted in Figure 12.3a while the tracking errors are plotted in Figures 12.3b-c. The control inputs plotted in Figures 12.4a-c. As proven in Theorem 12.1, the tracking errors asymptotically converge to the origin. Moreover, the control inputs have not reached their limits. This means that we still can further shorten the transient time by increasing the control gains.

### 12.6 Conclusions

The key to the control development is the coordinate transformations (12.7) to transform the tracking error system, in which the tracking errors are interpreted in the frame attached to the vehicle body, to a triangular form, to which the backstepping technique can be applied. The control development in this chapter is based on the approach developed for underactuated surface ships in Chapter 5.


Figure 12.2 Reference and real trajectories in three dimensions


Figure 12.3 a. Position and orientation in the ( $\mathrm{x}, \mathrm{y}$ )-plane; $\mathbf{b}$. Tracking errors: $x-x_{d}$ (solid line), $y-y_{d}$ (dash-dotted line), $z-z_{d}$ (dash); $\mathbf{c}$. Tracking errors: $\psi-\psi_{d}$ (solid line), $\theta-\theta_{d}$ (dashdotted line)


Figure 12.4 a. Control input: $\tau_{u} ; \mathbf{b}$. Control input $\tau_{q}$; $\mathbf{c}$. Control input $\tau_{r}$

## Chapter 13

## Path-following of Underactuated Underwater Vehicles

This chapter extends the approach proposed for underactuated surface ships in Chapter 11 to design a path-following system for six degrees of freedom underactuated underwater vehicles. Although the control design is much more involved in comparison with that for underactuated surface ships in Chapter 11, it still guarantees that path-following errors asymptotically converge to a ball, with an adjustable radius, centered on a desired path, and covers both parking and point-to-point navigation problems.

### 13.1 Control Objective

For the reader's convenience, we rewrite the mathematical model of an underactuated underwater vehicle, which is described in detail in Section 3.4.2.1, moving in six degrees of freedom as follows:

$$
\begin{align*}
& \dot{\eta}_{1}=J_{1}\left(\eta_{2}\right) v_{1}, \\
& M_{1} \dot{v}_{1}=-C_{1}\left(v_{1}\right) v_{2}-D_{1} v_{1}-D_{n 1}\left(v_{1}\right) v_{1}+\tau_{1}+\tau_{1 E}, \\
& \dot{\eta}_{2}=J_{2}\left(\eta_{2}\right) v_{2},  \tag{13.1}\\
& M_{2} \dot{v}_{2}=-C_{1}\left(v_{1}\right) v_{1}-C_{2}\left(v_{2}\right) v_{2}-D_{2} v_{2}-D_{n 2}\left(v_{2}\right) v_{2}- \\
& \quad g_{2}\left(\eta_{2}\right)+\tau_{2}+\tau_{2 E}
\end{align*}
$$

where $\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ and $\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ are given by

$$
\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)=\left[\begin{array}{cc}
\cos (\psi) \cos (\theta) & -\sin (\psi) \cos (\phi)+\sin (\phi) \sin (\theta) \cos (\psi) \\
\sin (\psi) \cos (\theta) & \cos (\psi) \cos (\phi)+\sin (\phi) \sin (\theta) \sin (\psi) \\
-\sin (\theta) & \sin (\phi) \cos (\theta) \\
\sin (\psi) \sin (\phi)+\sin (\theta) \cos (\psi) \cos (\phi) \\
-\cos (\psi) \sin (\phi)+\sin (\theta) \sin (\psi) \cos (\phi) \\
\cos (\phi) \cos (\theta)
\end{array}\right],
$$

$$
\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)=\left[\begin{array}{ccc}
1 & \sin (\phi) \tan (\theta) & \cos (\phi) \tan (\theta)  \tag{13.2}\\
0 & \cos (\phi) & -\sin (\phi) \\
0 & \sin (\phi) / \cos (\theta) & \cos (\phi) / \cos (\theta)
\end{array}\right] .
$$

The matrices $\boldsymbol{M}_{\mathbf{1}}$ and $\boldsymbol{M}_{\mathbf{2}}$ are

$$
\boldsymbol{M}_{\mathbf{1}}=\left[\begin{array}{ccc}
m_{11} & 0 & 0  \tag{13.3}\\
0 & m_{22} & 0 \\
0 & 0 & m_{33}
\end{array}\right], \boldsymbol{M}_{\mathbf{2}}=\left[\begin{array}{ccc}
m_{44} & 0 & 0 \\
0 & m_{55} & 0 \\
0 & 0 & m_{66}
\end{array}\right]
$$

The matrices $\boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right)$ and $\boldsymbol{C}_{\mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right)$ are

$$
\begin{align*}
& \boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{1}\right)=\left[\begin{array}{ccc}
0 & m_{33} w & -m_{22} v \\
-m_{33} w & 0 & m_{11} u \\
m_{22} v & -m_{11} u & 0
\end{array}\right], \\
& \boldsymbol{C}_{\mathbf{2}}\left(\boldsymbol{v}_{2}\right)=\left[\begin{array}{ccc}
0 & m_{66} r & -m_{55} q \\
-m_{66} r & 0 & m_{44} p \\
m_{55} q & -m_{44} p & 0
\end{array}\right] . \tag{13.4}
\end{align*}
$$

The linear and nonlinear damping matrices $\boldsymbol{D}_{\mathbf{1}}, \boldsymbol{D}_{\mathbf{2}}, \boldsymbol{D}_{\boldsymbol{n} \mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right)$, and $\boldsymbol{D}_{\boldsymbol{n} 2}\left(\boldsymbol{v}_{\mathbf{2}}\right)$ are

$$
\begin{align*}
& \boldsymbol{D}_{\mathbf{1}}=\left[\begin{array}{ccc}
d_{11} & 0 & 0 \\
0 & d_{22} & 0 \\
0 & 0 & d_{33}
\end{array}\right], \boldsymbol{D}_{\mathbf{2}}=\left[\begin{array}{ccc}
d_{44} & 0 & 0 \\
0 & d_{55} & 0 \\
0 & 0 & d_{66}
\end{array}\right], \\
& \boldsymbol{D}_{\boldsymbol{n} \mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right)=\left[\begin{array}{ccc}
\sum_{i=2}^{3} d_{u i}|u|^{i-1} & 0 & 0 \\
0 & \sum_{i=2}^{3} d_{v i}|v|^{i-1} & 0 \\
0 & 0 & \sum_{i=2}^{3} d_{w i}|w|^{i-1}
\end{array}\right], \\
& \boldsymbol{D}_{\boldsymbol{n} \mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right)=\left[\begin{array}{ccc}
\sum_{i=2}^{3} d_{p i}|p|^{i-1} & 0 & 0 \\
0 & \sum_{i=2}^{3} d_{q i}|q|^{i-1} & 0 \\
0 & 0 & \sum_{i=2}^{3} d_{r i}|r|^{i-1}
\end{array}\right] . \tag{13.5}
\end{align*}
$$

The restoring force and moment vector $g_{2}\left(\eta_{2}\right)$ is given by

$$
\boldsymbol{g}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)=\left[\begin{array}{c}
\rho g \nabla \overline{G M_{T}} \sin (\phi) \cos (\theta)  \tag{13.6}\\
\rho g \nabla \overline{G M_{L}} \sin (\theta) \\
0
\end{array}\right]
$$

The propulsion force and moment vectors $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{\mathbf{2}}$ are

$$
\boldsymbol{\tau}_{\mathbf{1}}=\left[\begin{array}{c}
\tau_{u}  \tag{13.7}\\
0 \\
0
\end{array}\right], \quad \boldsymbol{\tau}_{\mathbf{2}}=\left[\begin{array}{c}
\tau_{p} \\
\tau_{q} \\
\tau_{r}
\end{array}\right],
$$

which imply that the vehicle under consideration does not have independent actuators in the sway and heave. The environmental disturbance vectors $\boldsymbol{\tau}_{\mathbf{1 E}}$ and $\boldsymbol{\tau}_{\mathbf{2} \boldsymbol{E}}$
are given by

$$
\boldsymbol{\tau}_{\mathbf{1 E}}=\left[\begin{array}{c}
\tau_{E u}(t)  \tag{13.8}\\
\tau_{E v}(t) \\
\tau_{E w}(t)
\end{array}\right], \boldsymbol{\tau}_{\mathbf{2 E}}=\left[\begin{array}{c}
\tau_{E p}(t) \\
\tau_{E q}(t) \\
\tau_{E r}(t)
\end{array}\right],
$$

where $\tau_{E u}(t), \tau_{E v}(t), \tau_{E w}(t), \tau_{E p}(t), \tau_{E q}(t)$, and $\tau_{E r}(t)$ are the environmental disturbance forces or moments acting on the surge, sway, heave, roll, pitch, and yaw axes, respectively.

We assume that these disturbances are bounded as follows:

$$
\begin{align*}
& \left|\tau_{E u}(t)\right| \leq \tau_{E u}^{\max }<\infty,\left|\tau_{E v}(t)\right| \leq \tau_{E v}^{\max }<\infty,\left|\tau_{E w}(t)\right| \leq \tau_{E w}^{\max }<\infty, \\
& \left|\tau_{E p}(t)\right| \leq \tau_{E p}^{\max }<\infty,\left|\tau_{E q}(t)\right| \leq \tau_{E q}^{\max }<\infty,\left|\tau_{E r}(t)\right| \leq \tau_{E r}^{\max }<\infty . \tag{13.9}
\end{align*}
$$

Since the sway and heave control forces are not available in the sway and heave dynamics, the vehicle model (13.1) is underactuated.

In this chapter, we consider a control objective of designing the control inputs $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{\mathbf{2}}$ to force the underactuated vehicle (13.1) to follow a specified path $\boldsymbol{\Omega}$, see Figure 13.1. If we are able to drive the vehicle to follow closely a virtual vessel that moves along the path with a desired speed $u_{0}$, then the control objective is fulfilled, i.e., the vessel is in a tube of nonzero diameter centered on the reference path and moves along the specified path at the speed $u_{0}$. Roughly speaking, the approach is to steer the vessel such that it heads toward the virtual one and diminishes the distance between the real and the virtual vessels.


Figure 13.1 General framework of underwater vehicle path-following

In Figure 13.1, $A$ is the center of the real vehicle and $A_{s}$ is a point on the reference path attached to the virtual vehicle. Define the following path-following errors

$$
\begin{align*}
& x_{e}=x_{d}-x \\
& y_{e}=y_{d}-y \\
& z_{e}=z_{d}-z \\
& d_{e}=\sqrt{x_{e}^{2}+y_{e}^{2}+z_{e}^{2}} \tag{13.10}
\end{align*}
$$

where $x_{d}, y_{d}$ and $z_{d}$ are the coordinates of $A_{s}$. Then the terms $a_{i}, 1 \leq i \leq 3$ in Figure 13.1 are obtained from $x_{e}, y_{e}$ and $z_{e}$ by rotating the body frame around the earth-fixed frame $O_{E} X_{E} Y_{E} Z_{E}$ the roll, pitch, and yaw angles, i.e.,

$$
\left[\begin{array}{l}
a_{1}  \tag{13.11}\\
a_{2} \\
a_{3}
\end{array}\right]=\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\left[\begin{array}{l}
x_{e} \\
y_{e} \\
z_{e}
\end{array}\right]
$$

Expanding (13.11) yields

$$
\begin{align*}
& a_{1}=x_{e} J_{1}^{11}\left(\eta_{2}\right)+y_{e} J_{1}^{21}\left(\eta_{2}\right)+z_{e} J_{1}^{31}\left(\eta_{\mathbf{2}}\right), \\
& a_{2}=x_{e} J_{1}^{12}\left(\eta_{2}\right)+y_{e} J_{1}^{22}\left(\eta_{2}\right)+z_{e} J_{1}^{32}\left(\eta_{2}\right),  \tag{13.12}\\
& a_{3}=x_{e} J_{1}^{13}\left(\eta_{\mathbf{2}}\right)+y_{e} J_{1}^{23}\left(\eta_{2}\right)+z_{e} J_{1}^{33}\left(\boldsymbol{\eta}_{\mathbf{2}}\right),
\end{align*}
$$

where $J_{1}^{i j}\left(\eta_{2}\right)$ is the element of $\boldsymbol{J}_{\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ at the $i$ th row and $j$ th column. Therefore the path-following orientation errors are defined by the angles $\alpha$ and $\beta$. It is noted that the angles $\gamma, \alpha$, and $\beta$ are not defined at $d_{e}=0$ but with the aid of a desired controller, $\lim _{\left\|\left(d_{e}, \alpha, \beta\right)\right\| \rightarrow 0}(\theta, \psi)=\left(\theta_{s}, \psi_{s}\right)$ with $\theta_{s}$ and $\psi_{s}$ being the orientation angles of the virtual vessel. Hence in this chapter, we will design a controller such that it guarantees $d_{e} \geq d_{e}^{*}$ with $d_{e}^{*}$ being an arbitrarily small positive constant to avoid chattering caused by $d_{e}=0$. With the above definitions, our control objective can be mathematically stated as follows:

Path-following Objective. Under Assumption 13.1, design the control inputs $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{\mathbf{2}}$ to force the underactuated vehicle (13.1) to follow the path $\boldsymbol{\Omega}$ given by

$$
\begin{align*}
x_{d} & =x_{d}(s), \\
y_{d} & =y_{d}(s),  \tag{13.13}\\
z_{d} & =z_{d}(s),
\end{align*}
$$

where $s$ is the path parameter variable, such that

$$
\begin{aligned}
& \lim _{t \rightarrow \infty} d_{e}(t) \leq \overline{d_{e}}, \quad \lim _{t \rightarrow \infty}|\alpha(t)| \leq \bar{\alpha}, \\
& \lim _{t \rightarrow \infty}|\beta(t)| \leq \bar{\beta}, \quad \lim _{t \rightarrow \infty}|\phi(t)| \leq \bar{\phi},
\end{aligned}
$$

with $\overline{d_{e}}, \bar{\phi}, \bar{\alpha}$, and $\bar{\beta}$ being arbitrarily small positive constants.
Assumption 13.1. The reference path is regular, i.e. there exist strictly positive constants $a_{3 \text { min }}, a_{3 \text { max }}, a_{2 \text { min }}$ and $a_{2 \text { max }}$ such that

$$
\begin{align*}
& a_{3 \text { min }} \leq \sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}+\left(\frac{\partial z_{d}}{\partial s}\right)^{2}} \leq a_{3 \max } \\
& a_{2 \min } \leq \sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2} \leq a_{2 \max } .} \tag{13.14}
\end{align*}
$$

## Remark 13.1.

1. We might refer to the above objective as a path-tracking one. However, we use the term "path-following" since our approach is to make the real vehicle follow the virtual one, see Figure 13.1.
2. Assumption 13.1 ensures that the path is feasible for the vessel to follow, see Section 13.3. The condition (13.14) implies that the reference trajectory cannot contain a vertical straight line to avoid singularity of $\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{2}\right)$ at the pitch angle $\theta= \pm 0.5 \pi$.
3. If the reference path is not regular, then we can often split it into regular pieces and consider each of them separately. This is the case of point-to-point navigation, which will be addressed in Section 13.6.
4. The path parameter, $s$, is not the arc length of the path in general. For example, a circle with radius $R$ centered at the origin can be described as $x_{d}=R \cos (s)$ and $y_{d}=R \sin (s)$.

### 13.2 Coordinate Transformations

From (13.10) and (13.12), we have the position kinematic error dynamics as follows:

$$
\begin{equation*}
\dot{d}_{e}=\frac{1}{d_{e}}\left(x_{e} \frac{\partial x_{d}}{\partial s}+y_{e} \frac{\partial y_{d}}{\partial s}+z_{e} \frac{\partial z_{d}}{\partial s}\right) \dot{s}-\frac{a_{1}}{d_{e}} u-\frac{a_{2}}{d_{e}} v-\frac{a_{3}}{d_{e}} w . \tag{13.15}
\end{equation*}
$$

For the path-following orientation errors, referring to Figure 13.1 and the control objective stated in the previous section, one can see that the following holds

$$
\begin{equation*}
\frac{a_{1}}{d_{e}}=\cos (\gamma)=\cos (\alpha) \cos (\beta) \tag{13.16}
\end{equation*}
$$

which in turn implies that

$$
\left\{\begin{array}{l}
\lim _{t \rightarrow \infty} \alpha(t)=0  \tag{13.17}\\
\lim _{t \rightarrow \infty} \beta(t)=0
\end{array} \Leftrightarrow \lim _{t \rightarrow \infty} \gamma(t)=0 \Rightarrow \lim _{t \rightarrow \infty}\left(\frac{a_{1}(t)}{d_{e}(t)}\right)=1 .\right.
$$

Hence we can either choose the angles $\alpha$ and $\beta$, or the angle $\gamma$, or the term $a_{1} / d_{e}$ as the orientation coordinates for the control design. We now discuss the above options and then choose one that results in a simple control design and enhance feasible initial conditions.

Using Angles $\alpha$ and $\beta$. In this case, the path-following orientation errors are defined as follows, see Figure 13.1:

$$
\begin{align*}
& \alpha=e_{\alpha}-2 \pi n_{\alpha}\left(e_{\alpha}\right), \\
& \beta=e_{\beta}-2 \pi n_{\beta}\left(e_{\beta}\right), \tag{13.18}
\end{align*}
$$

where

$$
\begin{align*}
& e_{\alpha}=\left\{\begin{array}{l}
2 \arctan \left(\frac{a_{3}}{a_{1}}\right) ;\left(a_{3}, a_{1}\right) \neq(0,0), \\
0 ;\left(a_{3}, a_{1}\right)=(0,0),
\end{array}\right. \\
& e_{\beta}=\left\{\begin{array}{l}
2 \arctan \left(\frac{a_{2}}{\sqrt{a_{1}^{2}+a_{3}^{2}}}\right) ;\left(a_{2}, \sqrt{a_{1}^{2}+a_{3}^{2}}\right) \neq(0,0), \\
0 ;\left(a_{2}, \sqrt{a_{1}^{2}+a_{3}^{2}}\right)=(0,0) .
\end{array}\right. \tag{13.19}
\end{align*}
$$

The functions $n_{\alpha}\left(e_{\alpha}\right)$ and $n_{\beta}\left(e_{\beta}\right)$ take values in $(0, \pm 1, \pm 2, \ldots)$ such that $\alpha$ and $\beta$ belong to $(-\pi, \pi]$. Hence $\alpha$ and $\beta$ are periodic and piecewise continuous functions with respect to $e_{\alpha}$ and $e_{\beta}$. The reason for introducing (13.18) is to convert all equilibrium points of $\alpha$ and $\beta$ to the origin. It is seen from (13.19) that $e_{\alpha}$ and $e_{\beta}$ are discontinuous on the following surfaces:

$$
\begin{align*}
& D_{\alpha}=\left\{\left(a_{1}, a_{3}\right): a_{3} \neq 0, a_{1}=0\right\} \\
& D_{\beta}=\left\{\left(a_{1}, a_{2}, a_{3}\right): a_{2} \neq 0, \sqrt{a_{1}^{2}+a_{3}^{2}}=0\right\} \tag{13.20}
\end{align*}
$$

It is also seen from (13.18) that $\alpha$ and $\beta$ are discontinuous on the surfaces:

$$
\begin{align*}
& C_{\alpha}=\left\{\left(a_{1}, a_{3}\right): \alpha=\pi\right\}  \tag{13.21}\\
& C_{\beta}=\left\{\left(a_{1}, a_{2}, a_{3}\right): \beta=\pi\right\} .
\end{align*}
$$

We now use (13.18) to transform the kinematic part of (13.1) to

$$
\begin{equation*}
\dot{\eta}_{\alpha \beta}=\boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\boldsymbol{s}}(\cdot) \dot{s}+\boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\boldsymbol{u}}(\cdot) u+\boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{v}(\cdot) v+\boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\boldsymbol{w}}(\cdot) w+\boldsymbol{A}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot) \boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) \boldsymbol{v}_{\mathbf{2}} \tag{13.22}
\end{equation*}
$$

where

$$
\begin{aligned}
& \eta_{\boldsymbol{\alpha} \boldsymbol{\beta}}=[\phi \alpha \beta]^{T} \\
& \boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{s}(\cdot)=\boldsymbol{F}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot)\left[\frac{\partial x_{d}}{\partial s} \frac{\partial y_{d}}{\partial s} \frac{\partial z_{d}}{\partial s}\right]^{T},
\end{aligned}
$$

$$
\begin{align*}
& \boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\boldsymbol{u}}(\cdot)=-\boldsymbol{F}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot)\left[J_{1}^{11}\left(\boldsymbol{\eta}_{2}\right) J_{1}^{21}\left(\boldsymbol{\eta}_{2}\right) J_{1}^{31}\left(\boldsymbol{\eta}_{2}\right)\right]^{T}, \\
& \boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\boldsymbol{v}}(\cdot)=-\boldsymbol{F}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot)\left[J_{1}^{12}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) J_{1}^{22}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) J_{1}^{32}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right]^{T} \text {, } \\
& f_{\alpha \beta}^{w}(\cdot)=-\boldsymbol{F}_{\alpha \beta}(\cdot)\left[J_{1}^{13}\left(\eta_{2}\right) J_{1}^{23}\left(\eta_{2}\right) J_{1}^{33}\left(\eta_{2}\right)\right]^{T} \text {, } \\
& \boldsymbol{F}_{\boldsymbol{\alpha} \beta}(\cdot)=\left[0 \boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\mathbf{1}}(\cdot) \boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\mathbf{2}}(\cdot)\right]^{T},  \tag{13.23}\\
& \boldsymbol{A}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot)=\left[\begin{array}{cc}
1 & 0 \\
-\cos (\alpha) \tan (\beta) & \cos (\phi)+\sin (\alpha) \tan (\beta) \sin (\phi) \\
\sin (\alpha) & \sin (\phi) \cos (\alpha)
\end{array}\right. \\
& \begin{array}{c}
0 \\
\cos (\theta) \sin (\phi)+\tan (\beta)(\sin (\theta) \cos (\alpha)-\cos (\theta) \sin (\alpha) \cos (\phi) \\
-(\cos (\alpha) \cos (\theta) \cos (\phi)+\sin (\alpha) \sin (\theta)
\end{array},
\end{align*}
$$

with

$$
\begin{align*}
\boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\mathbf{1}(\cdot)=} & \frac{1}{d_{e} \cos (\beta)}\left[\cos (\alpha) J_{1}^{13}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)-\sin (\alpha) J_{1}^{11}\left(\boldsymbol{\eta}_{\mathbf{2}}\right),\right. \\
& \cos (\alpha) J_{1}^{23}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)-\sin (\alpha) J_{1}^{21}\left(\boldsymbol{\eta}_{\mathbf{2}}\right), \\
& \left.\cos (\alpha) J_{1}^{33}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)-\sin (\alpha) J_{1}^{31}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right] \\
\boldsymbol{f}_{\boldsymbol{\alpha} \boldsymbol{\beta}}^{\mathbf{2}}(\cdot)= & {\left[-\frac{\sin (\beta) \cos (\alpha)}{d_{e}} J_{1}^{11}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\frac{\cos (\beta)}{d_{e}} J_{1}^{12}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)-\frac{\sin (\beta) \sin (\alpha)}{d_{e}} J_{1}^{13}\left(\boldsymbol{\eta}_{\mathbf{2}}\right),\right.} \\
& -\frac{\sin (\beta) \cos (\alpha)}{d_{e}} J_{1}^{21}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\frac{\cos (\beta)}{d_{e}} J_{1}^{22}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)-\frac{\sin (\beta) \sin (\alpha)}{d_{e}} J_{1}^{23}\left(\boldsymbol{\eta}_{\mathbf{2}}\right), \\
& \left.-\frac{\sin (\beta) \cos (\alpha)}{d_{e}} J_{1}^{31}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\frac{\cos (\beta)}{d_{e}} J_{1}^{32}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)-\frac{\sin (\beta) \sin (\alpha)}{d_{e}} J_{1}^{33}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right] . \tag{13.24}
\end{align*}
$$

From (13.23), we calculate the determination of the matrix $\boldsymbol{A}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot)$ as follows:

$$
\begin{align*}
\operatorname{det}\left(\boldsymbol{A}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot)\right)= & -\cos (\alpha) \cos (\theta)-\sin (\alpha) \sin (\theta) \cos (\phi)- \\
& \sin (\theta) \sin (\phi) \tan (\beta) . \tag{13.25}
\end{align*}
$$

It can be seen from (13.25) that the matrix $\boldsymbol{A}_{\boldsymbol{\alpha} \boldsymbol{\beta}}(\cdot)$ is not globally invertible even when $\phi=0$ and $\theta \neq \pm 0.5 \pi$. It is also observed that if we choose the angles $\alpha$ and $\beta$ as the orientation coordinates for control design, there are a number of discontinuous surfaces, see (13.19), (13.20), and (13.21), which make the stability analysis difficult.

Using Angle $\gamma$. In this case, the path-following orientation error is defined as follows, see Figure 13.1:

$$
\begin{equation*}
\gamma=e_{\gamma}-2 \pi n_{\gamma}\left(e_{\gamma}\right) \tag{13.26}
\end{equation*}
$$

where
$e_{\gamma}=\left\{\begin{array}{l}2 \arctan \left(\frac{\sqrt{a_{2}^{2}+a_{3}^{2}}}{a_{1}}\right) ;\left(a_{1}, \sqrt{a_{2}^{2}+a_{3}^{2}}\right) \neq(0,0), \\ 0 ;\left(a_{1}, \sqrt{a_{2}^{2}+a_{3}^{2}}\right)=(0,0) .\end{array}\right.$
The interpretation of the above expressions is similar to that of (13.18) and (13.19). It can be seen from (13.27) that $e_{\gamma}$ is discontinuous on the following surface:

$$
\begin{equation*}
D_{\gamma}=\left\{\left(a_{1}, a_{2}, a_{3}\right): \sqrt{a_{2}^{2}+a_{3}^{2}} \neq 0, a_{1}=0\right\} \tag{13.28}
\end{equation*}
$$

It can also be seen from (13.26) that $\gamma$ is discontinuous on the surface:

$$
\begin{equation*}
C_{\gamma}=\left\{\left(a_{1}, a_{2}, a_{3}\right): \gamma=\pi\right\} . \tag{13.29}
\end{equation*}
$$

Differentiating both sides of (13.26) results in

$$
\begin{align*}
\dot{\gamma}= & \frac{1}{e^{2} \sqrt{a_{2}^{2}+a_{3}^{2}}}\left(a_{1} a_{2}\left(\frac{\partial a_{2}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial a_{2}}{\partial y_{e}} \dot{y}_{e}+\frac{\partial a_{2}}{\partial z_{e}} \dot{z}_{e} \frac{\partial a_{2}}{\partial \phi} \dot{\phi}\right)+\right. \\
& a_{1} a_{3}\left(\frac{\partial a_{3}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial a_{3}}{\partial y_{e}} \dot{y}_{e}+\frac{\partial a_{3}}{\partial z_{e}} \dot{z}_{e} \frac{\partial a_{3}}{\partial \phi} \dot{\phi}\right)- \\
& \left(a_{2}^{2}+a_{3}^{2}\right)\left(\frac{\partial a_{1}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial a_{1}}{\partial y_{e}} \dot{y}_{e}+\frac{\partial a_{1}}{\partial z_{e}} \dot{z}_{e}+\frac{\partial a_{1}}{\partial \phi} \dot{\phi}\right)+ \\
& \left(a_{1}\left(a_{2} \frac{\partial a_{2}}{\partial \theta}+a_{3} \frac{\partial a_{3}}{\partial \theta}\right)-\left(a_{2}^{2}+a_{3}^{2}\right) \frac{\partial a_{1}}{\partial \theta}\right) \dot{\theta}+ \\
& \left.\left(a_{1}\left(a_{2} \frac{\partial a_{2}}{\partial \psi}+a_{3} \frac{\partial a_{3}}{\partial \psi}\right)-\left(a_{2}^{2}+a_{3}^{2}\right) \frac{\partial a_{1}}{\partial \psi}\right) \dot{\psi}\right) . \tag{13.30}
\end{align*}
$$

Since

$$
\begin{align*}
& {\left[\frac{\left(a_{1}\left(a_{2} \frac{\partial a_{2}}{\partial \theta}+a_{3} \frac{\partial a_{3}}{\partial \theta}\right)-\left(a_{2}^{2}+a_{3}^{2}\right) \frac{\partial a_{1}}{\partial \theta}\right)}{d_{e}^{2} \sqrt{a_{2}^{2}+a_{3}^{2}}}\right]^{2}+} \\
& {\left[\frac{\left(a_{1}\left(a_{2} \frac{\partial a_{2}}{\partial \psi}+a_{3} \frac{\partial a_{3}}{\partial \psi}\right)-\left(a_{2}^{2}+a_{3}^{2}\right) \frac{\partial a_{1}}{\partial}\right)}{d_{e}^{2} \sqrt{a_{2}^{2}+a_{3}^{2}}}\right]^{2} \neq 0} \tag{13.31}
\end{align*}
$$

for all $\left(x_{e}, y_{e}, z_{e}, \phi, \psi\right) \in \mathbb{R}^{5}$ and $\theta \in \mathbb{R} \backslash \pm 0.5 \pi$, the angle $\gamma$ can be chosen as the orientation coordinate for the control design. However (13.30) will result in a very
complicated control law.
Using $a_{1} / d_{e}$. By defining

$$
\begin{align*}
& a_{e}=\sqrt{x_{e}^{2}+y_{e}^{2}} \\
& \cos \left(\gamma_{1}\right)=\frac{a_{e}}{d_{e}} \cos (\theta)-\frac{z_{e}}{d_{e}} \sin (\theta)  \tag{13.32}\\
& \cos \left(\gamma_{2}\right)=\frac{x_{e}}{a_{e}} \cos (\psi)+\frac{y_{e}}{a_{e}} \sin (\psi)
\end{align*}
$$

we can write $a_{1} / d_{e}$ as follows

$$
\begin{equation*}
\frac{a_{1}}{d_{e}}=\cos \left(\gamma_{1}\right)+\frac{a_{e}}{d_{e}} \cos (\theta)\left(\cos \left(\gamma_{2}\right)-1\right) \tag{13.33}
\end{equation*}
$$

Since $\left|\frac{a_{e}}{d_{e}} \cos (\theta)\right| \leq 1, \forall\left(x_{e}, y_{e}, z_{e}, \theta\right) \in \mathbb{R}^{4}$, the conditions

$$
\begin{align*}
& \lim _{t \rightarrow \infty} \gamma_{1}=0 \\
& \lim _{t \rightarrow \infty} \gamma_{2}=0 \tag{13.34}
\end{align*}
$$

imply that $\lim _{t \rightarrow \infty}\left(a_{1} / d_{e}\right)=1$. Furthermore, from (13.32) we can write

$$
\begin{align*}
& \gamma_{1}=\theta+\theta_{d}  \tag{13.35}\\
& \gamma_{2}=\psi-\psi_{d}
\end{align*}
$$

where $\theta_{d}=\arccos \left(a_{e} / d_{e}\right)$ and $\psi_{d}=\arccos \left(x_{e} / a_{e}\right)$ are the desired orientation angles of the vessel in the horizontal and vertical planes, respectively. Hence (13.34) also implies the orientation control objective. If one differentiates both sides of (13.35) to obtain $\dot{\gamma}_{1}$ and $\dot{\gamma}_{2}$, there will be discontinuity in the $\gamma_{1}$ - and $\gamma_{2}$ - dynamics at $z_{e}=0$ and/or $y_{e}=0$, i.e., on the $a_{e}$ and/or $x_{e}$ axes. This discontinuity will cause difficulties in applying the backstepping technique. To get around this problem, we compute $\dot{\gamma}_{1}$ and $\dot{\gamma}_{2}$ based on (13.32) as follows:

$$
\begin{align*}
& \dot{\gamma}_{1}=\dot{\theta}+\frac{\dot{a}_{e} \sin (\theta)+\dot{z}_{e} \cos (\theta)}{d_{e} \cos \left(\gamma_{1}\right)}-\frac{\dot{d}_{e} \sin \left(\gamma_{1}\right)}{d_{e} \cos \left(\gamma_{1}\right)}, \\
& \dot{\gamma}_{2}=\dot{\psi}+\frac{\dot{x}_{e} \sin (\psi)-\dot{y}_{e} \cos (\psi)}{a_{e} \cos \left(\gamma_{2}\right)}-\frac{\dot{a}_{e} \sin \left(\gamma_{2}\right)}{a_{e} \cos \left(\gamma_{2}\right)} . \tag{13.36}
\end{align*}
$$

It can be seen that (13.36) is not defined at $\gamma_{i}= \pm 0.5 \pi, i=1,2$, and $a_{e}=0, d_{e}=0$. However, our controller will guarantee that $\left|\gamma_{i}(t)\right|<0.5 \pi$ and $d_{e}(t) \geq d_{e}^{*}, a_{e}(t) \geq$ $a_{e}^{*}, \forall t \geq t_{0} \geq 0$ with arbitrarily small positive constants $d_{e}^{*}$ and $a_{e}^{*}$ and for feasible initial conditions. From (13.36), we can see that $\dot{\theta}$ and $\dot{\psi}$ are decoupled. Hence, designing a controller to achieve the control objective posed in the previous section by using the orientation coordinate $a_{1} / d_{e}$ would be much simpler than using the angles $\gamma, \alpha$, and $\beta$. For convenience of control design, we rewrite the transformed
system dynamics (13.15) and (13.36) as follows:

$$
\begin{align*}
& \dot{d}_{e}=\frac{1}{d_{e}}\left(x_{e} \frac{\partial x_{d}}{\partial s}+y_{e} \frac{\partial y_{d}}{\partial s}+z_{e} \frac{\partial z_{d}}{\partial s}\right) \dot{s}-\frac{a_{1}}{d_{e}} u-\frac{a_{2}}{d_{e}} v-\frac{a_{3}}{d_{e}} w, \\
& \dot{\eta}_{2 \boldsymbol{r}}=\boldsymbol{f}_{\mathbf{2} \boldsymbol{\gamma}}(\cdot)+\boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) \boldsymbol{v}_{\mathbf{2}} \\
& \boldsymbol{M}_{\mathbf{1}} \dot{v}_{\mathbf{1}}=-\boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right) \boldsymbol{v}_{\mathbf{2}}-\boldsymbol{D}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right) \boldsymbol{v}_{\mathbf{1}}+\boldsymbol{\tau}_{\mathbf{1}}+\boldsymbol{\tau}_{\mathbf{1}}(t), \\
& \boldsymbol{M}_{\mathbf{2}} \dot{v}_{\mathbf{2}}=-\boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right) \boldsymbol{v}_{\mathbf{1}}-\boldsymbol{C}_{\mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right) \boldsymbol{v}_{\mathbf{2}}-\boldsymbol{D}_{\mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right) \boldsymbol{v}_{\mathbf{2}}-\boldsymbol{g}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\boldsymbol{\tau}_{\mathbf{2}}+\boldsymbol{\tau}_{\mathbf{2} E}(t), \tag{13.37}
\end{align*}
$$

where

$$
\begin{align*}
& \eta_{\mathbf{2 \gamma}}=\left[\phi \gamma_{1} \gamma_{2}\right]^{T}, \\
& \boldsymbol{f}_{\mathbf{2 \gamma}}(\cdot)=\left[\begin{array}{c}
0 \\
f_{1}^{s} \dot{s}+f_{1}^{u} u+f_{1}^{v} v+f_{1}^{w} w \\
f_{2}^{s} \dot{s}+f_{2}^{u} u+f_{2}^{v} v+f_{2}^{w} w
\end{array}\right], \\
& f_{1}^{s}=\varpi_{11} \frac{\partial x_{d}}{\partial s}+\varpi_{12} \frac{\partial y_{d}}{\partial s}+\varpi_{13} \frac{\partial z_{d}}{\partial s}, \\
& f_{1}^{u}=-\left[\varpi_{11} J_{1}^{11}\left(\eta_{2}\right)+\varpi_{12} J_{1}^{21}\left(\eta_{2}\right)+\varpi_{13} J_{1}^{31}\left(\eta_{\mathbf{2}}\right)\right], \\
& f_{1}^{v}=-\left[\varpi_{11} J_{1}^{12}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\varpi_{12} J_{1}^{22}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\varpi_{13} J_{1}^{32}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right],  \tag{13.38}\\
& f_{1}^{w}=-\left[\varpi_{11} J_{1}^{13}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\varpi_{12} J_{1}^{23}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\varpi_{13} J_{1}^{33}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right], \\
& f_{2}^{s}=\varpi_{21} \frac{\partial x_{d}}{\partial s}+\varpi_{22} \frac{\partial y_{d}}{\partial s}, \\
& f_{2}^{u}=-\left[\varpi_{21} J_{1}^{11}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\varpi_{22} J_{1}^{21}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right], \\
& f_{2}^{v}=-\left[\varpi_{21} J_{1}^{12}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)+\varpi_{22} J_{1}^{22}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right], \\
& f_{2}^{w}=-\left[\varpi_{21} J_{1}^{13}\left(\boldsymbol{\eta}_{2}\right)+\varpi_{22} J_{1}^{23}\left(\boldsymbol{\eta}_{2}\right)\right],
\end{align*}
$$

with

$$
\begin{align*}
\varpi_{11} & =\left(\frac{x_{e} \sin (\theta)}{a_{e} d_{e} \cos \left(\gamma_{1}\right)}-\frac{x_{e} \sin \left(\gamma_{1}\right)}{d_{e}^{2} \cos \left(\gamma_{1}\right)}\right), \\
\varpi_{12} & =\left(\frac{y_{e} \sin (\theta)}{a_{e} d_{e} \cos \left(\gamma_{1}\right)}-\frac{y_{e} \sin \left(\gamma_{1}\right)}{d_{e}^{2} \cos \left(\gamma_{1}\right)}\right) \\
\varpi_{13} & =\left(\frac{\cos (\theta)}{d_{e} \cos \left(\gamma_{1}\right)}-\frac{z_{e} \sin \left(\gamma_{1}\right)}{d_{e}^{2} \cos \left(\gamma_{1}\right)}\right)  \tag{13.39}\\
\varpi_{21} & =\left(\frac{\sin (\psi)}{a_{e} \cos \left(\gamma_{2}\right)}-\frac{x_{e} \sin \left(\gamma_{2}\right)}{a_{e}^{2} \cos \left(\gamma_{2}\right)}\right) \\
\varpi_{22} & =\left(-\frac{\cos (\psi)}{a_{e} \cos \left(\gamma_{2}\right)}-\frac{y_{e} \sin \left(\gamma_{2}\right)}{a_{e}^{2} \cos \left(\gamma_{2}\right)}\right)
\end{align*}
$$

Therefore, we will design the control inputs $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{\mathbf{2}}$ for (13.37) to yield the control objective. In Section 13.3, a procedure to design a stabilizer for the path-
following error system (13.37) is presented in detail. The structure of (13.37) suggests that we design the actual controls $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{\mathbf{2}}$ in two stages. First, we design the virtual velocity controls for $u$ and $\boldsymbol{v}_{2}$ and choose $\dot{s}$ to ultimately stabilize $d_{e}$ and $\gamma_{i}, i=1,2$ at the origin. Based on the backstepping technique, the controls $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{2}$ will then be designed to make the errors between the virtual velocity controls and their actual values asymptotically tend to a small ball centered at the origin. Since the vessel parameters are unknown, an adaptation scheme is also introduced in this step to estimate their values used in the control laws. We split the control design procedure into two steps. The first step is to design $\boldsymbol{\tau}_{\mathbf{1}}$ while the second step takes care of $\boldsymbol{\tau}_{\mathbf{2}}$. This allows us to simplify the choice of feasible initial conditions.

Since the transformed system (13.37) is not defined at $d_{e}(t)=0, a_{e}(t)=0$, $\gamma_{i}(t)= \pm 0.5 \pi, i=1,2$, we first assume that

$$
\begin{equation*}
d_{e}(t) \geq d_{e}^{*}, a_{e}(t) \geq a_{e}^{*},\left|\gamma_{i}(t)\right|<0.5 \pi, i=1,2, \forall t \geq t_{0} \geq 0 \tag{13.40}
\end{equation*}
$$

for some positive constants $d_{e}^{*}$ and $a_{e}^{*}$. Our controller design will guarantee (13.40) for feasible initial conditions.

### 13.3 Control Design

The $d_{e}$-dynamics have two inputs that can be chosen to stabilize $d_{e}$, namely $\dot{s}$ and $u$. We can either choose the input $u$ or $\dot{s}$ and then design the remaining input. If we fix $\dot{s}$, then the virtual vessel is allowed to move at a desired speed. The real vessel will follow the virtual one on the path by the controller, and vice versa. In this chapter, we choose to fix $\dot{s}$. This allows us to adjust the initial conditions in most cases without moving the real vessel, see Section 13.4.

Define

$$
\begin{equation*}
\tilde{u}=u-u_{d} \tag{13.41}
\end{equation*}
$$

where $u_{d}$ is the intermediate control of $u$. As discussed above, we choose the intermediate control $u_{d}$ and $\dot{s}$ as follows:

$$
\begin{gather*}
u_{d}=k_{1}\left(d_{e}-\delta_{e}\right)-\frac{1}{a_{1}}\left(a_{2} v+a_{3} w\right)+\frac{1}{d_{e}} \frac{u_{0}\left(t, d_{e}\right)\left(x_{e} \frac{\partial x_{d}}{\partial s}+y_{e} \frac{\partial y_{d}}{\partial s}+z_{e} \frac{\partial z_{d}}{\partial s}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}+\left(\frac{\partial z_{d}}{\partial s}\right)^{2}}}, \\
\dot{s}=\frac{a_{1}}{d_{e}} \frac{u_{0}\left(t, d_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}+\left(\frac{\partial z_{d}}{\partial s}\right)^{2}}}, \tag{13.42}
\end{gather*}
$$

where $k_{1}$ and $\delta_{e}$ are positive constants to be selected later, and $u_{0}\left(t, d_{e}\right) \neq 0, \forall t \geq$ $t_{0} \geq 0, d_{e}(t) \in \mathbb{R}$, is the speed of the virtual vessel on the path. Indeed, one can
choose this speed to be a constant. However, the time-varying speed and position path-following error dependence of the virtual vessel on the path is more desirable, especially when the underwater vehicle starts to follow the path. For example, one might choose

$$
\begin{equation*}
u_{0}\left(t, d_{e}\right)=u_{0}^{*}\left(1-\chi_{1} e^{-\chi_{2}\left(t-t_{0}\right)}\right) e^{-\chi_{3} d_{e}} \tag{13.44}
\end{equation*}
$$

where $u_{0}^{*} \neq 0, \chi_{i}>0, i=1,2,3, \chi_{1}<1$. The choice of $u_{0}\left(t, d_{e}\right)$ in (13.44) has the following desired feature: When the path-following error, $d_{e}$, is large, the virtual vessel will wait for the real one; when $d_{e}$ is small the virtual vessel will move along the path at the speed closed to $u_{0}^{*}$ and the real one follows it within the specified look ahead distance. This feature is suitable in practice because it avoids using a highgain control for large signal $d_{e}$. It is noted that $u_{d}$ is not defined at $a_{1}=0$. Since the terms $a_{2} / a_{1}$ and $a_{3} / a_{1}$ can be written as $\left(a_{2} / d_{e}\right) /\left(a_{1} / d_{e}\right)$ and $\left(a_{3} / d_{e}\right) /\left(a_{1} / d_{e}\right)$, and recalling that $\cos (\gamma)=a_{1} / d_{e}$, the intermediate control $u_{d}$ is well defined if (13.40) holds and

$$
\begin{equation*}
|\gamma(t)|<0.5 \pi, \forall t \geq t_{0} \geq 0 . \tag{13.45}
\end{equation*}
$$

We will come back to this issue in Section 13.4.

## Remark 13.2.

1. If we design the virtual control $u_{d}$ without canceling the terms $a_{2} v$ and $a_{3} w$ in the $d_{e}$-dynamics, then the condition (13.45) is not required for $u_{d}$ being well defined. However, an assumption of the sway and heave velocities being bounded is needed in advance in the stability analysis, i.e. assume stability to prove stability.
2. If the sway and heave velocities are assumed to be bounded by the surge velocity, the terms $a_{2} v$ and $a_{3} w$ are not required to be canceled either. This controller can be designed similarly to the one in this chapter. It is noted that the sway velocity does not require to be bounded by the surge velocity with a relatively small constant as in [129] for the case of path-following in the horizontal plane.

Substituting (13.42) and (13.43) into the first equation of (13.37) results in

$$
\begin{equation*}
\dot{d}_{e}=-k_{1} \frac{a_{1}}{d_{e}}\left(d_{e}-\delta_{e}\right)-\frac{a_{1}}{d_{e}} \tilde{u} . \tag{13.46}
\end{equation*}
$$

By noticing that under Assumption 13.1, see (13.40) and (13.45), the intermediate control $u_{d}$ is a smooth function of $x_{e}, y_{e}, z_{e}, s, u_{0}, \eta_{2}, v$, and $w$, differentiating both sides of (13.41) with (13.42) and (13.43) yields

$$
\begin{aligned}
\dot{\tilde{u}}= & \frac{m_{22}}{m_{11}} v r-\frac{m_{33}}{m_{11}} w q-\frac{d_{11}}{m_{11}} u-\sum_{i=2}^{3} \frac{d_{u i}}{m_{11}}|u|^{i-1} u+\frac{1}{m_{11}} \tau_{u}+ \\
& \frac{1}{m_{11}} \tau_{w u}(t)-\frac{\partial u_{d}}{\partial x_{e}} \dot{x}_{e}-\frac{\partial u_{d}}{\partial y_{e}} \dot{y}_{e}-\frac{\partial u_{d}}{\partial z_{e}} \dot{z}_{e}-\frac{\partial u_{d}}{\partial s} \dot{s}-\frac{\partial u_{d}}{\partial u_{0}} \dot{u}_{0}- \\
& \frac{\partial u_{d}}{\partial \eta_{2}} \dot{\eta}_{2}-\frac{\partial u_{d}}{\partial v}\left(\frac{m_{33}}{m_{22}} w p-\frac{m_{11}}{m_{22}} u r-\frac{d_{22}}{m_{22}} v-\sum_{i=2}^{3} \frac{d_{v i}}{m_{22}}|v|^{i-1} v+\right.
\end{aligned}
$$

$$
\begin{align*}
& \left.\frac{1}{m_{22}} \tau_{w v}(t)\right)-\frac{\partial u_{d}}{\partial w}\left(\frac{m_{11}}{m_{33}} u q-\frac{m_{22}}{m_{33}} v p-\frac{d_{33}}{m_{33}} w-\right. \\
& \left.\sum_{i=2}^{3} \frac{d_{w i}}{m_{33}}|w|^{i-1} w+\frac{1}{m_{33}} \tau_{w w}(t)\right) \tag{13.47}
\end{align*}
$$

where for convenience of choosing $u_{0}$, the terms $\frac{\partial u_{d}}{\partial x_{e}}, \frac{\partial u_{d}}{\partial y_{e}}$, and $\frac{\partial u_{d}}{\partial z_{e}}$ do not include $\frac{\partial u_{0}}{\partial x_{e}}, \frac{\partial u_{0}}{\partial y_{e}}$ and $\frac{\partial u_{0}}{\partial z_{e}}$, which are lumped into $\dot{u}_{0}$. From (13.47), we choose the actual control $\tau_{1}$ or $\tau_{u}$ without canceling useful nonlinear damping terms as

$$
\begin{align*}
\tau_{u}= & -c_{1} \tilde{u}-\hat{\theta}_{1}^{T} f_{1}(\cdot)-\hat{\theta}_{21} \tanh \left(\frac{\tilde{u} \hat{\theta}_{21}}{\varepsilon_{21}}\right)-\hat{\theta}_{22} \frac{\partial u_{d}}{\partial v} \tanh \left(\frac{\partial u_{d}}{\partial v} \frac{\tilde{u} \hat{\theta}_{22}}{\varepsilon_{22}}\right)- \\
& \hat{\theta}_{23} \frac{\partial u_{d}}{\partial w} \tanh \left(\frac{\partial u_{d}}{\partial w} \frac{\tilde{u} \hat{\theta}_{23}}{\varepsilon_{23}}\right), \tag{13.48}
\end{align*}
$$

and the update law as

$$
\begin{align*}
& \dot{\hat{\theta}}_{1 j}=\gamma_{1 j} \operatorname{proj}\left(\tilde{u} f_{1 j}(\cdot), \hat{\theta}_{1 j}\right), 1 \leq j \leq 16, \\
& \dot{\hat{\theta}}_{21}=\gamma_{21} \operatorname{proj}\left(|\tilde{u}|, \hat{\theta}_{21}\right), \\
& \dot{\hat{\theta}}_{22}=\gamma_{22} \operatorname{proj}\left(\left|\tilde{u} \frac{\partial u_{d}}{\partial v}\right|, \hat{\theta}_{22}\right),  \tag{13.49}\\
& \dot{\hat{\theta}}_{23}=\gamma_{23} \operatorname{proj}\left(\left|\tilde{u} \frac{\partial u_{d}}{\partial w}\right|, \hat{\theta}_{23}\right),
\end{align*}
$$

where $c_{1}, \varepsilon_{2 i}, \gamma_{1 j}, \gamma_{2 i}, 1 \leq i \leq 3$ and $1 \leq j \leq 16$, are positive constants to be selected later, and $f_{1 j}(\cdot)$ are the $j$ th elements of $f_{1}(\cdot)$, respectively, with

$$
\begin{align*}
f_{1}(\cdot)= & {\left[v r,-w q,-u_{d},-|u| u_{d},-u^{2} u_{d},-\left(\frac{\partial u_{d}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial u_{d}}{\partial y_{e}} \dot{y}_{e}+\frac{\partial u_{d}}{\partial z_{e}} \dot{z}_{e}+\right.\right.} \\
& \left.\frac{\partial u_{d}}{\partial s} \dot{s}+\frac{\partial u_{d}}{\partial u_{0}} \dot{u}_{0}+\frac{\partial u_{d}}{\partial \eta_{\mathbf{2}}} \dot{\eta}_{2}\right),-\frac{\partial u_{d}}{\partial v} w p, \frac{\partial u_{d}}{\partial v} u r, \frac{\partial u_{d}}{\partial v} v, \frac{\partial u_{d}}{\partial v}|v| v, \\
& \left.\frac{\partial u_{d}}{\partial v} v^{3},-\frac{\partial u_{d}}{\partial w} u q, \frac{\partial u_{d}}{\partial w} v p, \frac{\partial u_{d}}{\partial w} w, \frac{\partial u_{d}}{\partial w}|w| w, \frac{\partial u_{d}}{\partial w} w^{3}\right]^{T}, \tag{13.50}
\end{align*}
$$

$\hat{\theta}_{i j}, 1 \leq i \leq 2$ is the $j$ th element of $\hat{\theta}_{i}$, which is an estimate of $\theta_{i}$ with

$$
\begin{aligned}
\theta_{1}= & {\left[m_{22}, m_{33}, d_{11}, d_{u 2}, d_{u 3}, m_{11}, \frac{m_{11} m_{33}}{m_{22}}, \frac{m_{11}^{2}}{m_{22}}, \frac{m_{11} d_{22}}{m_{22}}, \frac{m_{11} d_{v 2}}{m_{22}},\right.} \\
& \left.\frac{m_{11} d_{v 3}}{m_{22}}, \frac{m_{11}^{2}}{m_{33}}, \frac{m_{11} m_{22}}{m_{33}}, \frac{m_{11} d_{33}}{m_{33}}, \frac{m_{11} d_{w 2}}{m_{33}}, \frac{m_{11} d_{w 3}}{m_{33}}\right]^{T},
\end{aligned}
$$

$$
\begin{equation*}
\theta_{2}=\left[\tau_{w u}^{\max }, \frac{m_{11}}{m_{22}} \tau_{w v}^{\max }, \frac{m_{11}}{m_{33}} \tau_{w w}^{\max }\right]^{T} \tag{13.51}
\end{equation*}
$$

The operator, proj, is the Lipschitz continuous projection algorithm repeated here for the convenience of the reader as follows:

$$
\begin{align*}
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \quad \text { if } \quad \Xi(\hat{\omega}) \leq 0, \\
& \operatorname{proj}(\varpi, \hat{\omega})=\varpi \quad \text { if } \quad \Xi(\hat{\omega}) \geq 0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi \leq 0,  \tag{13.52}\\
& \operatorname{proj}(\varpi, \hat{\omega})=(1-\Xi(\hat{\omega})) \varpi \quad \text { if } \Xi(\hat{\omega})>0 \text { and } \Xi_{\hat{\omega}}(\hat{\omega}) \varpi>0,
\end{align*}
$$

where $\Xi(\hat{\omega})=\frac{\hat{\omega}^{2}-\omega_{M}^{2}}{\mu^{2}+2 \mu \omega_{M}}, \Xi_{\hat{\omega}}(\hat{\omega})=\frac{\partial \Xi(\hat{\omega})}{\partial \hat{\omega}}, \mu$ is an arbitrarily small positive constant, $\hat{\omega}$ is an estimate of $\omega$ and $|\omega| \leq \omega_{M}$.

The projection algorithm is such that if $\dot{\hat{\omega}}=\operatorname{proj}(\varpi, \hat{\omega})$ and $\hat{\omega}\left(t_{0}\right) \leq \omega_{M}$ and then

1. $\hat{\omega}(t) \leq \omega_{M}+\xi, \forall 0 \leq t_{0} \leq t<\infty$,
2. $\operatorname{proj}(\varpi, \hat{\omega})$ is Lipschitz continuous,
3. $|\operatorname{proj}(\varpi, \hat{\omega})| \leq|\varpi|$,
4. $\tilde{\omega} \operatorname{proj}(\varpi, \hat{\omega}) \geq \tilde{\omega} \varpi$ with $\tilde{\omega}=\omega-\hat{\omega}$.

Substituting (13.48) into (13.47) yields the error dynamics

$$
\left.\begin{array}{rl}
\dot{\tilde{u}}= & -\frac{1}{m_{11}}\left(c_{1}+d_{11}+\sum_{i=2}^{3} d_{u i}|u|^{i-1}\right) \tilde{u}+\frac{1}{m_{11}} \theta_{1}^{T} f_{1}(\cdot)-\frac{1}{m_{11}} \hat{\theta}_{1}^{T} f_{1}(\cdot)+ \\
& \frac{1}{m_{11}} \tau_{w u}(t)-\frac{1}{m_{11}} \hat{\theta}_{21} \tanh \left(\frac{\tilde{u} \hat{\theta}_{21}}{\varepsilon_{21}}\right)-\frac{\partial u_{d}}{\partial v} \frac{1}{m_{22}} \tau_{w v}(t)-\frac{1}{m_{11}} \hat{\theta}_{22} \frac{\partial u_{d}}{\partial v} \times \\
& \tanh \left(\frac{\partial u_{d}}{\partial v} \frac{\tilde{u}}{\hat{\theta}_{22}}\right.  \tag{13.53}\\
\varepsilon_{22}
\end{array}\right)-\frac{\partial u_{d}}{\partial w} \frac{1}{m_{33}} \tau_{w w}(t)-\frac{1}{m_{11}} \hat{\theta}_{23} \frac{\partial u_{d}}{\partial w} \tanh \left(\frac{\partial u_{d}}{\partial w} \frac{\tilde{u} \hat{\theta}_{23}}{\varepsilon_{23}}\right) . ~ \$
$$

Define

$$
\begin{equation*}
\tilde{v}_{2}=v_{2}-v_{2 d}, \tag{13.54}
\end{equation*}
$$

where $\boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}=\left[p_{d}, q_{d}, r_{d}\right]^{T}$ is the intermediate control of $\boldsymbol{v}_{\mathbf{2}}$. Recalling that our goal is to ultimately stabilize $\eta_{\mathbf{2} \gamma}=\left[\phi, \gamma_{1}, \gamma_{2}\right]^{T}$ at the origin, the second equation of (13.37) suggests that we choose this intermediate control as follows:

$$
\begin{equation*}
v_{2 d}=J_{2}^{-1}\left(\eta_{2}\right)\left(-f_{2 \gamma}(\cdot)-K_{2} \eta_{2 \gamma}\right), \tag{13.55}
\end{equation*}
$$

where $\boldsymbol{K}_{\mathbf{2}}=\operatorname{diag}\left(k_{21}, k_{22}, k_{23}\right)$ is a positive definite diagonal matrix. Substituting (13.54) and (13.55) into the second equation of (13.37) yields

$$
\begin{equation*}
\dot{\eta}_{2 \gamma}=-K_{2} \eta_{2 \gamma}+J_{2}\left(\eta_{2}\right) \tilde{v}_{2} . \tag{13.56}
\end{equation*}
$$

To design the actual control $\boldsymbol{\tau}_{\mathbf{2}}$, we first note that under Assumption 13.1, (13.40) and (13.45), the intermediate control $\boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}$ is a smooth function of $x_{e}, y_{e}, z_{e}, s, u_{0}, \boldsymbol{\eta}_{\mathbf{2}}$, and $\boldsymbol{v}_{\mathbf{1}}$. Differentiating both sides of (13.54) and multiplying by $\boldsymbol{M}_{\mathbf{2}}$, along the solutions of the last two equations of (13.37) results in

$$
\begin{equation*}
M_{2} \dot{\tilde{v}}_{2}=-C_{2}\left(v_{2}\right) \tilde{v}_{2}-D_{2}\left(v_{2}\right) \tilde{v}_{2}+F(\cdot) \theta_{3}+G(\cdot) \theta_{4}(t)+\tau_{2} \tag{13.57}
\end{equation*}
$$

where

$$
\begin{align*}
& \boldsymbol{F}(\cdot) \boldsymbol{\theta}_{\mathbf{3}}=-\boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right) \boldsymbol{v}_{\mathbf{1}}-\boldsymbol{C}_{\mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right) \boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}-\boldsymbol{D}_{\mathbf{2}}\left(\boldsymbol{v}_{\mathbf{2}}\right) \boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}-\boldsymbol{g}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)- \\
& \boldsymbol{M}_{\mathbf{2}}\left(\frac{\partial \boldsymbol{v}_{\mathbf{2}} \boldsymbol{d}}{\partial x_{e}} \dot{x}_{e}+\frac{\partial v_{2 d}}{\partial y_{e}} \dot{y}_{e}+\frac{\partial \boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}}{\partial z_{e}} \dot{z}_{e}+\frac{\partial \boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}}{\partial s} \dot{s}+\frac{\partial \boldsymbol{v}_{\mathbf{2}} \boldsymbol{d}}{\partial u_{0}} \dot{u}_{0}+\right. \\
& \left.\frac{\partial \boldsymbol{v}_{\mathbf{2}} \boldsymbol{d}}{\partial \boldsymbol{\eta}_{\mathbf{2}}} \dot{\boldsymbol{\eta}}_{\mathbf{2}}\right)-\boldsymbol{M}_{\mathbf{2}} \frac{\partial \boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}}{\partial \boldsymbol{v}_{\mathbf{1}}} \boldsymbol{M}_{\mathbf{1}}^{-\mathbf{1}}\left(-\boldsymbol{C}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right) \boldsymbol{v}_{\mathbf{2}}-\boldsymbol{D}_{\mathbf{1}}\left(\boldsymbol{v}_{\mathbf{1}}\right) \boldsymbol{v}_{\mathbf{1}}+\boldsymbol{\tau}_{\mathbf{1}}\right) \\
& \boldsymbol{G}(\cdot) \boldsymbol{\theta}_{\mathbf{4}}(t)=\boldsymbol{\tau}_{\boldsymbol{w} \mathbf{2}}(t)-\boldsymbol{M}_{\mathbf{2}} \frac{\partial \boldsymbol{v}_{\mathbf{2} \boldsymbol{d}}}{\partial \boldsymbol{v}_{\mathbf{1}}} \boldsymbol{M}_{\mathbf{1}}^{\mathbf{1}} \boldsymbol{\tau}_{\boldsymbol{w} \mathbf{1}}(t), \tag{13.58}
\end{align*}
$$

with $\boldsymbol{F}(\cdot) \in \mathbb{R}^{3 \times m_{3}}$ and $\boldsymbol{G}(\cdot) \in \mathbb{R}^{3 \times m_{4}}$ being the regression matrices, $\boldsymbol{\theta}_{\mathbf{3}} \in \mathbb{R}^{3 \times m_{3}}$ and $\boldsymbol{\theta}_{\mathbf{4}}(t) \in \mathbb{R}^{m_{4}}$ being the vectors of unknown vessel and environmental disturbance parameters. For the sake of simplicity, the regression matrices $\boldsymbol{F}(\cdot)$ and $\boldsymbol{G}(\cdot)$, and the vectors $\boldsymbol{\theta}_{\mathbf{3}}$ and $\boldsymbol{\theta}_{\mathbf{4}}(t)$ are not written down explicitly. From (13.56) and (13.57), we choose the actual control $\boldsymbol{\tau}_{\mathbf{2}}$ and update laws as follows:

$$
\begin{gather*}
\boldsymbol{\tau}_{\mathbf{2}}=-\boldsymbol{K}_{\mathbf{3}} \tilde{\boldsymbol{v}}_{\mathbf{2}}-\left(\boldsymbol{\eta}_{\mathbf{2} \boldsymbol{\gamma}}^{\boldsymbol{T}} \boldsymbol{J}_{\mathbf{2}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)\right)^{\boldsymbol{T}}-\boldsymbol{F}(\cdot) \hat{\boldsymbol{\theta}}_{\mathbf{3}}-\boldsymbol{G}(\cdot) \hat{\boldsymbol{\theta}}_{\mathbf{4}}  \tag{13.59}\\
\dot{\hat{\theta}}_{3 i}=\gamma_{3 i} \operatorname{proj}\left(\sum_{j=1}^{3} \tilde{v}_{2 j} f_{j i}, \hat{\theta}_{3 i}\right), 1 \leq i \leq m_{3} \\
\dot{\hat{\theta}}_{4 i}=\gamma_{4 i} \operatorname{proj}\left(\sum_{j=1}^{3}\left|\tilde{v}_{2 j} g_{j i}\right|, \hat{\theta}_{4 i}\right), 1 \leq i \leq m_{4} \tag{13.60}
\end{gather*}
$$

where $\varepsilon_{4 i}>0,1 \leq i \leq 3, \gamma_{3 j}>0,1 \leq j \leq m_{3}, \gamma_{4 l}>0,1 \leq l \leq m_{4}, \boldsymbol{K}_{\mathbf{3}}=\operatorname{diag}\left(k_{31}\right.$, $k_{32}, k_{33}$ ) is a positive definite diagonal matrix, $\hat{\theta}_{3 i}$ is an estimate of the $i$ th element of $\boldsymbol{\theta}_{\mathbf{3}}$, and $\hat{\theta}_{4 i}$ is an estimate of the maximum value of the $i$ th element of $\boldsymbol{\theta}_{\mathbf{4}}(t)$. For simplicity of notation, we have defined

$$
\boldsymbol{F}(\cdot) \hat{\boldsymbol{\theta}}_{\mathbf{3}}:=\left[\begin{array}{l}
\sum_{i=1}^{m_{3}} f_{1 i} \hat{\theta}_{3 i}  \tag{13.61}\\
\sum_{i=1}^{m_{3}} f_{2 i} \hat{\theta}_{3 i} \\
\sum_{i=1}^{m_{3}} f_{3 i} \hat{\theta}_{3 i}
\end{array}\right], \boldsymbol{G}(\cdot) \hat{\boldsymbol{\theta}}_{\mathbf{4}}:=\left[\begin{array}{l}
\sum_{i=1}^{m_{4}} g_{1 i} \hat{\theta}_{4 i} \tanh \left(\varepsilon_{41}^{-1} \tilde{v}_{21} \sum_{i=1}^{m_{2}} g_{1 i} \hat{\theta}_{4 i}\right) \\
\sum_{i=1}^{m_{4}} g_{2 i} \hat{\theta}_{4 i} \tanh \left(\varepsilon_{42}^{-1} \tilde{v}_{22} \sum_{i=1}^{m_{2}} g_{2 i} \hat{\theta}_{4 i}\right) \\
\sum_{i=1}^{m_{4}} g_{3 i} \hat{\theta}_{4 i} \tanh \left(\varepsilon_{43}^{-1} \tilde{v}_{23} \sum_{i=1}^{m_{2}} g_{3 i} \hat{\theta}_{4 i}\right)
\end{array}\right]
$$

with $f_{j i}, 1 \leq j \leq 3,1 \leq i \leq m_{3}$, being the element in $j$ th column and $i$ th row of the regression matrix $\boldsymbol{F}(\cdot)$. Similarly $g_{j i}, 1 \leq j \leq 3,1 \leq i \leq m_{4}$, is the element in $j$ th column and $i$ th row of the regression matrix $\boldsymbol{G}(\cdot)$.

Substituting (13.59) into (13.57) yields the error dynamics

$$
\begin{align*}
M_{2} \dot{v}_{2}= & -C_{2}\left(v_{2}\right) \tilde{v}_{2}-\left(K_{3}+D_{2}\left(v_{2}\right)\right) \tilde{v}_{2}-\left(\eta_{2 \gamma}^{T} J_{2}\left(\eta_{2}\right)\right)^{T}+F(\cdot) \theta_{3}+ \\
& G(\cdot) \theta_{4}(t)-F(\cdot) \hat{\theta}_{3}-G(\cdot) \hat{\theta}_{4} . \tag{13.62}
\end{align*}
$$

We now present the main result of this chapter, the proof of which is given in the next section.

Theorem 13.1. Assume that

1. the vessel inertia, added mass and damping matrices are diagonal;
2. the environmental disturbances are bounded;
3. the vessel parameters are unknown but constant;
4. the reference path satisfies Assumption 13.1.

If the state feedback control laws (13.48) and (13.59), and the update laws (13.49) and (13.60) are applied to the vessel system (13.1) then there exist feasible initial conditions such that the path-following errors $\left(x(t)-x_{d}(t), y(t)-y_{d}(t), z(t)-\right.$ $\left.z_{d}(t), \gamma_{1}(t), \gamma_{2}(t)\right)$ converge to a ball centered on the desired path $\boldsymbol{\Omega}$ asymptotically. Furthermore, the radius of this ball can be made arbitrarily small by adjusting the control gains.

### 13.4 Stability Analysis

To prove Theorem 13.1, we first consider the $\left(\boldsymbol{\eta}_{\boldsymbol{2} \boldsymbol{\gamma}}, \tilde{\boldsymbol{v}}_{\mathbf{2}}\right)$-subsystem and then the ( $\left.\boldsymbol{d}_{\boldsymbol{e}}, \tilde{\boldsymbol{u}}\right)$-subsystem.

## $\left(\eta_{2 \gamma}, \tilde{v}_{2}\right)$-subsystem

To investigate stability of this subsystem, we consider the following Lyapunov function

$$
\begin{equation*}
V_{1}=\frac{1}{2} \eta_{\mathbf{2} \gamma}^{\boldsymbol{T}} \eta_{\mathbf{2} \gamma}+\frac{1}{2} \tilde{\boldsymbol{v}}_{\mathbf{2}}^{\boldsymbol{T}} \boldsymbol{M}_{\mathbf{2}} \tilde{\boldsymbol{v}}_{\mathbf{2}}+\frac{1}{2} \sum_{i=3}^{4} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}^{\boldsymbol{T}} \boldsymbol{\Gamma}_{\boldsymbol{i}}^{-\mathbf{1}} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}} \tag{13.63}
\end{equation*}
$$

where $\tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}=\boldsymbol{\theta}_{\boldsymbol{i}}-\hat{\boldsymbol{\theta}}_{\boldsymbol{i}}$ and $\boldsymbol{\Gamma}_{\boldsymbol{i}}=\operatorname{diag}\left(\gamma_{i j}\right), 1 \leq j \leq m_{3}$ for $i=3,1 \leq j \leq m_{4}$ for $i=4$. Differentiating both sides of (13.63) along (13.56), (13.60), and (13.62) yields

$$
\begin{equation*}
\dot{V}_{1} \leq-\eta_{\mathbf{2} \gamma}^{\boldsymbol{T}} \boldsymbol{K}_{\mathbf{2}} \eta_{\mathbf{2} \gamma}-\tilde{\boldsymbol{v}}_{\mathbf{2}}^{\boldsymbol{T}}\left(\boldsymbol{D}_{\mathbf{2 0}}+\boldsymbol{K}_{\mathbf{3}}\right) \tilde{\boldsymbol{v}}_{\mathbf{2}}+0.2785 \sum_{i=1}^{3} \varepsilon_{4 i} \tag{13.64}
\end{equation*}
$$

where $\boldsymbol{D}_{\mathbf{2 0}}=\operatorname{diag}\left(d_{44}, d_{55}, d_{66}\right)$ and we have used $|x|-x \tanh (x / \lambda) \leq 0.2785 \lambda$ for all $x \in \mathbb{R}$ and $\lambda>0$. From (13.64), we conclude that $\boldsymbol{\eta}_{\mathbf{2} \gamma}$ and $\tilde{\boldsymbol{v}}_{\mathbf{2}}$ are ultimately asymptotically stable at the origin. To estimate the upper bound of $\boldsymbol{\eta}_{\mathbf{2} \gamma}$ and $\tilde{\boldsymbol{v}}_{\mathbf{2}}$, we subtract and add $\frac{1}{2} \sum_{i=3}^{4} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}^{\boldsymbol{T}} \boldsymbol{\Gamma}_{\boldsymbol{i}}^{-\mathbf{1}} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}$ to the right-hand side of (13.64) to obtain

$$
\begin{equation*}
\dot{V}_{1} \leq-\sigma_{1} V_{1}+\rho_{1}, \tag{13.65}
\end{equation*}
$$

where

$$
\begin{align*}
& \sigma_{1}=\min \left(1,2 \lambda_{\min }\left(\boldsymbol{K}_{\mathbf{2}}\right), \frac{2 \lambda_{\min }\left(\boldsymbol{D}_{\mathbf{2 0}}+\boldsymbol{K}_{\mathbf{3}}\right)}{\lambda_{\max }\left(\boldsymbol{M}_{\mathbf{2}}\right)}\right), \\
& \rho_{1}=\frac{1}{2} \sum_{i=3}^{4} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}^{\boldsymbol{T}} \boldsymbol{\Gamma}_{\boldsymbol{i}}^{-\mathbf{1}} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}+0.2785 \sum_{i=1}^{3} \varepsilon_{4 i} . \tag{13.66}
\end{align*}
$$

From (13.65), it is direct to show that

$$
\begin{equation*}
V_{1}(t) \leq V_{1}\left(t_{0}\right) e^{-\sigma_{1}\left(t-t_{0}\right)}+\frac{\rho_{1}}{\sigma_{1}} \tag{13.67}
\end{equation*}
$$

which further yields

$$
\begin{align*}
& \left\|\boldsymbol{\eta}_{\boldsymbol{2} \boldsymbol{\gamma}}(\boldsymbol{t})\right\| \leq \sqrt{2 V_{1}\left(t_{0}\right)} e^{-\frac{\sigma_{1}}{2}\left(t-t_{0}\right)}+\sqrt{\frac{2 \rho_{1}}{\sigma_{1}}:=\alpha_{\eta}(\cdot) e^{-\frac{\sigma_{1}}{2}\left(t-t_{0}\right)}+\rho_{\eta}} \\
& \left\|\tilde{\boldsymbol{v}}_{\mathbf{2}}(\boldsymbol{t})\right\| \leq \sqrt{\frac{2 V_{1}\left(t_{0}\right)}{\lambda_{\min }\left(\boldsymbol{M}_{\mathbf{2}}\right)}} e^{-\frac{\sigma_{1}}{2}\left(t-t_{0}\right)}+\sqrt{\frac{2 \rho_{1}}{\sigma_{1} \lambda_{\min }\left(\boldsymbol{M}_{\mathbf{2}}\right)}}:=\alpha_{v}(\cdot) e^{-\frac{\sigma_{1}}{2}\left(t-t_{0}\right)}+\rho_{v} \tag{13.68}
\end{align*}
$$

## $\left(d_{e}, \tilde{u}\right)$-subsystem

To analyze the stability of this subsystem more easily, we first consider the $\tilde{u}$ dynamics and then the $d_{e}$-dynamics.
$\tilde{u}$-dynamics. Consider the following Lyapunov function

$$
\begin{equation*}
V_{2}=\frac{m_{11}}{2} \tilde{u}^{2}+\frac{1}{2} \sum_{i=1}^{2} \tilde{\theta}_{i}^{T} \Gamma_{i}^{-1} \tilde{\theta}_{i} \tag{13.69}
\end{equation*}
$$

where $\tilde{\theta}_{i}=\theta_{i}-\hat{\theta}_{i}$ and $\Gamma_{i}=\operatorname{diag}\left(\gamma_{i j}\right), 1 \leq j \leq 16$ for $i=1,1 \leq j \leq 3$ for $i=2$. Differentiating both sides of (13.69) along (13.47), (13.48), and (13.49) yields

$$
\begin{equation*}
\dot{V}_{2} \leq-\left(c_{1}+d_{11}\right) \tilde{u}^{2}+0.2785 \sum_{i=1}^{3} \varepsilon_{2 i} . \tag{13.70}
\end{equation*}
$$

Subtracting and adding $\frac{1}{2} \sum_{i=1}^{2} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}^{\boldsymbol{T}} \boldsymbol{\Gamma}_{\boldsymbol{i}}^{\mathbf{- 1}} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}$ to the right-hand side of (13.70) yields

$$
\begin{equation*}
\dot{V}_{2} \leq-\sigma_{2} V_{2}+\rho_{2} \tag{13.71}
\end{equation*}
$$

where

$$
\begin{align*}
& \sigma_{2}=\min \left(1, \frac{2\left(c_{1}+d_{11}\right)}{m_{11}^{\min }}\right), \\
& \rho_{2}=\frac{1}{2} \sum_{i=1}^{2} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}^{\boldsymbol{T}} \boldsymbol{\Gamma}_{\boldsymbol{i}}^{-\mathbf{1}} \tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}+0.2785 \sum_{i=1}^{3} \varepsilon_{2 i}, \tag{13.72}
\end{align*}
$$

with $m_{11}^{\min }$ being the minimum value of $m_{11}$. From (13.71), it is direct to show that

$$
\begin{equation*}
V_{2}(t) \leq V_{2}\left(t_{0}\right) e^{-\sigma_{2}\left(t-t_{0}\right)}+\frac{\rho_{2}}{\sigma_{2}} \tag{13.73}
\end{equation*}
$$

which further yields

$$
\begin{equation*}
|\tilde{u}(t)| \leq \sqrt{2 V_{2}\left(t_{0}\right)} e^{-\sigma_{2} / 2\left(t-t_{0}\right)}+\sqrt{\frac{2 \rho_{2}}{\sigma_{2}}}:=\alpha_{u}(\cdot) e^{-\sigma_{2} / 2\left(t-t_{0}\right)}+\rho_{u} \tag{13.74}
\end{equation*}
$$

Remark 13.3. It is noted that, due to the use of the projection algorithm, by adjusting $\boldsymbol{K}_{\mathbf{2}}, \boldsymbol{K}_{\mathbf{3}}, c_{1}, \varepsilon_{2 i}, \varepsilon_{4 i}, \gamma_{1 j}, \gamma_{2 i}, \gamma_{3 n}, \gamma_{4 l}, 1 \leq i \leq 3,1 \leq j \leq 16,1 \leq n \leq m_{3}$, and $1 \leq l \leq m_{4}$, we can make $\rho_{u}, \rho_{\eta}$ and $\rho_{v}$ arbitrarily small. This observation plays an important role in the stability analysis of the $d_{e}$-dynamics.
$d_{e}$-dynamics. We first calculate the lower-bound of $d_{e}$. We now show that there exist initial conditions such that $d_{e}(t) \geq d_{e}^{*}>0$. From (13.46) and (13.16), we have

$$
\begin{equation*}
\dot{\tilde{d}}_{e} \geq-k_{1} \tilde{d}_{e}-\left(\alpha_{u} e^{-\sigma_{2} / 2\left(t-t_{0}\right)}+\rho_{u}\right) \tag{13.75}
\end{equation*}
$$

where $\tilde{d}_{e}=d_{e}-\delta_{e}$, which with $\sigma_{2}>2 k_{1}$ further yields

$$
\begin{align*}
\tilde{d}_{e}(t) \geq & \tilde{d}_{e}\left(t_{0}\right) e^{-k_{1}\left(t-t_{0}\right)}+\frac{\alpha_{u}(\cdot) e^{-k_{1}\left(t-t_{0}\right)}}{\sigma_{2} / 2-k_{1}}\left(-1+e^{-\left(\sigma_{2} / 2-k_{1}\right)\left(t-t_{0}\right)}\right)- \\
& \frac{\rho_{u}}{k_{1}}\left(1-e^{-k_{1}\left(t-t_{0}\right)}\right) \tag{13.76}
\end{align*}
$$

Therefore, the condition $d_{e}(t) \geq d_{e}^{*}>0$ holds when

$$
\begin{equation*}
\sigma_{2}>2 k_{1}, \delta_{e} \geq d_{e}^{*}+\frac{\rho_{u}}{k_{1}}, d_{e}\left(t_{0}\right) \geq \frac{\alpha_{u}(\cdot)}{\sigma_{2} / 2-k_{1}}+\delta_{e}-\rho_{u} \tag{13.77}
\end{equation*}
$$

We will come back to this issue in the next section. We now calculate the upperbound of $d_{e}$. We rewrite (13.46) as

$$
\begin{equation*}
\left.\dot{\tilde{d}}_{e}=-k_{1} \tilde{d}_{e}-\cos (\gamma) \tilde{u}-k_{1}\left(\cos \left(\gamma_{1}\right)-1\right)+\cos (\gamma)\left(\cos \left(\gamma_{2}\right)-1\right)\right) \tilde{d}_{e} . \tag{13.78}
\end{equation*}
$$

Since (13.78) and the ( $\tilde{u}, \gamma_{1}, \gamma_{2}$ )-subsystem are in the cascade form, one might think that the stability results developed for cascade systems in [17] and [69] can be applied. However, the stability results in those papers were developed for cascade systems without nonvanishing disturbances. In fact, nonvanishing disturbances may destroy the stability of a cascade system that satisfies all conditions stated in the above papers. Therefore, we will use Lemma 10.1 to investigate the stability of the system (13.78) by verifying all conditions $\mathrm{C} 1-\mathrm{C} 4$.

Verifying Condition C1. Take the Lyapunov function

$$
\begin{equation*}
V_{3}=\frac{1}{2} d_{e}^{2} \tag{13.79}
\end{equation*}
$$

It is direct to show that C 1 holds with

$$
\begin{equation*}
c_{0}=0, c_{1}=c_{2}=0.5, c_{3}=1, c_{4}=k_{1} \tag{13.80}
\end{equation*}
$$

Verifying Condition C2. By noting that

$$
\begin{gather*}
\left.\mid \cos (\gamma) \tilde{u}+k_{1}\left(\cos \left(\gamma_{1}\right)-1\right)+\cos (\gamma)\left(\cos \left(\gamma_{2}\right)-1\right)\right) \tilde{d}_{e} \mid \leq \\
|\tilde{u}|+k_{1}\left(\left|\gamma_{1}\right|+\left|\gamma_{2}\right|\right)\left|\tilde{d}_{e}\right|, \tag{13.81}
\end{gather*}
$$

we have

$$
\begin{equation*}
\lambda_{1}=1, \lambda_{2}=k_{1} \tag{13.82}
\end{equation*}
$$

Verifying Condition C3. This condition directly holds from (13.74) and (13.68).
Verifying Condition C4. From (13.80) and (13.82), condition C4 becomes

$$
\begin{equation*}
k_{1}-\max \left(\rho_{\eta}, \rho_{u}\right)\left(k_{1}+0.25 \mu_{0}\right)>0 . \tag{13.83}
\end{equation*}
$$

From Remark 13.3 and noting that $\mu_{0}$ is an arbitrarily positive constant, we can see that there always exists $k_{1}$ such that (13.83) holds. All conditions of Lemma 10.1 have been verified, and we therefore have

$$
\begin{equation*}
\left|d_{e}(t)\right| \leq \alpha_{d e}(\cdot) e^{-\sigma_{d e}\left(t-t_{0}\right)}+\rho_{d e} \tag{13.84}
\end{equation*}
$$

where $\alpha_{d e}, \sigma_{d e}$ and $\rho_{d e}$ are calculated as in Lemma 10.1.
$(v, w)$-dynamics. Expanding (13.55) gives

$$
\begin{aligned}
q_{d}= & -\left(J_{2}^{21}\left(\boldsymbol{\eta}_{2}\right) k_{21} \phi+J_{2}^{22}\left(\boldsymbol{\eta}_{2}\right) k_{22} \gamma_{1}+J_{2}^{23}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) k_{23} \gamma_{2}\right)- \\
& J_{2}^{22}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{s} \dot{s}+f_{1}^{u} u\right)-J_{2}^{23}\left(\boldsymbol{\eta}_{2}\right)\left(f_{2}^{s} \dot{s}+f_{2}^{u} u\right)- \\
& \left(J_{2}^{22}\left(\boldsymbol{\eta}_{2}\right) f_{1}^{v}+J_{2}^{23}\left(\boldsymbol{\eta}_{2}\right) f_{2}^{v}\right) v-\left(J_{2}^{22}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) f_{1}^{w}+J_{2}^{23}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) f_{2}^{w}\right) w,
\end{aligned}
$$

$$
\begin{align*}
r_{d}= & -\left(J_{2}^{31}\left(\eta_{\mathbf{2}}\right) k_{21} \phi+J_{2}^{32}\left(\eta_{\mathbf{2}}\right) k_{22} \gamma_{1}+J_{2}^{33}\left(\eta_{\mathbf{2}}\right) k_{23} \gamma_{2}\right)- \\
& J_{2}^{32}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{s} \dot{s}+f_{1}^{u} u\right)-J_{2}^{33}\left(\eta_{2}\right)\left(f_{2}^{s} \dot{s}+f_{2}^{u} u\right)- \\
& \left(J_{2}^{32}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) f_{1}^{v}+J_{2}^{33}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) f_{2}^{v}\right) v-\left(J_{2}^{33}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) f_{1}^{w}+J_{2}^{33}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) f_{2}^{w}\right) w, \tag{13.85}
\end{align*}
$$

where $J_{2}^{i j}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$ is the element at the $i$ th row and $j$ th of $\boldsymbol{J}_{\mathbf{2}}^{-\mathbf{1}}\left(\boldsymbol{\eta}_{\mathbf{2}}\right)$. To show that the sway and heave velocities are bounded, we take the following quadratic function

$$
\begin{equation*}
V_{4}=\frac{1}{2} m_{22}^{2} v^{2}+\frac{1}{2} m_{33}^{2} w^{2}, \tag{13.86}
\end{equation*}
$$

whose derivative along (13.37), (13.54), and (13.85) satisfies

$$
\begin{align*}
\dot{V}_{4} \leq & -m_{22} d_{v 3} v^{4}-m_{33} d_{w 2}|w| w^{2}-m_{33} d_{w 3} w^{4}+A_{3}^{\max } v^{2}+A_{4}^{\max } w^{2}+ \\
& \frac{1}{4 \varepsilon_{5}} A_{1}^{\max }+\frac{1}{4 \varepsilon_{6}} A_{2}^{\max }, \tag{13.87}
\end{align*}
$$

where $\varepsilon_{i}>0, i=1,2, A_{j}^{\max }, 1 \leq j \leq 4$ is the maximum value of $A_{j}$ with

$$
\begin{align*}
A_{1}= & -m_{11} m_{22} u\left(\tilde{r}-J_{2}^{31}\left(\boldsymbol{\eta}_{2}\right) k_{21} \phi-J_{2}^{32}\left(\boldsymbol{\eta}_{2}\right) k_{22} \gamma_{1}-J_{2}^{33}\left(\boldsymbol{\eta}_{2}\right) k_{23} \gamma_{2}-\right. \\
& \left.J_{2}^{32}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{s} \dot{s}+f_{1}^{u} u\right)-J_{2}^{33}\left(\boldsymbol{\eta}_{2}\right)\left(f_{2}^{s} \dot{s}+f_{2}^{u} u\right)\right)+m_{22} \tau_{w v(t)}, \\
A_{2}= & m_{11} m_{33} u\left(\tilde{q}-J_{2}^{21}\left(\boldsymbol{\eta}_{2}\right) k_{21} \phi-J_{2}^{22}\left(\boldsymbol{\eta}_{2}\right) k_{22} \gamma_{1}-J_{2}^{23}\left(\boldsymbol{\eta}_{\mathbf{2}}\right) k_{23} \gamma_{2}-\right. \\
& \left.J_{2}^{22}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{s} \dot{s}+f_{1}^{u} u\right)-J_{2}^{23}\left(\boldsymbol{\eta}_{2}\right)\left(f_{2}^{s} \dot{s}+f_{2}^{u} u\right)\right)+m_{33} \tau_{w w(t),} \\
A_{3}= & \left(-m_{22} d_{22}-m_{22} d_{v 2}|v|+\varepsilon_{5}\left|A_{1}\right|+\frac{m_{11} m_{22}\left|u J_{2}^{33}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{w}+f_{2}^{w}\right)\right|}{2}+\right. \\
& \left.\frac{m_{11} m_{33}\left|u J_{2}^{22}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{v}+f_{2}^{v}\right)\right|}{2}\right), \\
A_{4}= & \left(-m_{33} d_{33}-m_{33} d_{w 2}|w|+\varepsilon_{6}\left|A_{2}\right|+\frac{m_{11} m_{22}\left|u J_{2}^{33}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{w}+f_{2}^{w}\right)\right|}{2}+\right. \\
& \left.\frac{m_{11} m_{33}\left|u J_{2}^{22}\left(\boldsymbol{\eta}_{2}\right)\left(f_{1}^{v}+f_{2}^{v}\right)\right|}{2}\right) . \tag{13.88}
\end{align*}
$$

It can be seen from (13.88) that $A_{i}^{\max }$ exist and are finite since their arguments are bounded as shown above. Hence (13.87) and (13.86) guarantee a finite upper bound of the sway and heave velocities.

Initial Conditions for $|\gamma(t)|<\frac{\pi}{2},\left|\gamma_{i}(t)\right|<\frac{\pi}{2}, i=1,2, \forall t \geq t_{0} \geq 0$. Since $\left|\gamma_{i}(t)\right| \leq$ $\left\|\eta_{2 \gamma}(t)\right\|, i=1,2, \forall t \geq t_{0} \geq 0$, from (13.68), it is direct to show that the condition $\left|\gamma_{i}(t)\right|<0.5 \pi, i=1,2, \forall t \geq t_{0} \geq 0$ holds if the initial conditions are such that

$$
\begin{equation*}
\sqrt{2 V_{1}\left(t_{0}\right)}+\sqrt{2 \rho_{1} / \sigma_{1}}<0.5 \pi \tag{13.89}
\end{equation*}
$$

which is further equivalent to

$$
\begin{equation*}
\sqrt{\left\|\eta_{2 \gamma}\left(t_{0}\right)\right\|^{2}+\lambda_{\max }\left(M_{2}\right)\left\|\tilde{v}_{2}\left(t_{0}\right)\right\|^{2}+\sum_{i=1}^{4} \Gamma_{i}^{-1}\left\|\tilde{\theta}_{i}\left(t_{0}\right)\right\|^{2}}+\sqrt{2 \frac{\rho_{1}}{\sigma_{1}}}<\frac{\pi}{2} \tag{13.90}
\end{equation*}
$$

for all $t \geq t_{0} \geq 0$. It is noted that the terms $\sum_{i=1}^{4} \Gamma_{i}^{-1}\left\|\tilde{\boldsymbol{\theta}}_{i}\left(t_{0}\right)\right\|^{2}$ and $\rho_{1}$ can be made arbitrarily small, see Remark 13.3.

From (13.33) and (13.16), the condition $|\gamma(t)|<0.5 \pi, \forall t \geq t_{0} \geq 0$ holds if the initial conditions are such that

$$
\begin{equation*}
\cos \left(\gamma_{1}(t)\right)+\frac{a_{e}(t)}{d_{e}(t)} \cos (\theta(t))\left(\cos \left(\gamma_{2}(t)-1\right)>0\right. \tag{13.91}
\end{equation*}
$$

Under the assumption that $|\theta(t)|<0.5 \pi, \forall t \geq t_{0} \geq 0$, the above condition is equivalent to

$$
\begin{equation*}
\cos \left(\gamma_{1}(t)\right)+\cos \left(\gamma_{2}(t)\right)>1 . \tag{13.92}
\end{equation*}
$$

From (13.68), the condition (13.92) holds if the initial conditions are such that

$$
\begin{align*}
& \sqrt{\left\|\boldsymbol{\eta}_{\boldsymbol{2} \boldsymbol{\gamma}}\left(t_{0}\right)\right\|^{2}+\lambda_{\max }\left(\boldsymbol{M}_{\mathbf{2}}\right)\left\|\tilde{\boldsymbol{v}}_{\boldsymbol{2}}\left(t_{0}\right)\right\|^{2}+\sum_{i=1}^{4} \boldsymbol{\Gamma}_{\boldsymbol{i}}^{-\mathbf{1}}\left\|\tilde{\boldsymbol{\theta}}_{\boldsymbol{i}}\left(t_{0}\right)\right\|^{2}}+ \\
& \sqrt{2 \frac{\rho_{1}}{\sigma_{1}}}<\arccos (0.5), \forall t \geq t_{0} \geq 0 \tag{13.93}
\end{align*}
$$

Since $\arccos (0.5)<0.5 \pi$, the condition (13.93) covers the condition (13.90).
Initial Conditions for $a_{e}(t) \geq a_{e}^{*}>0, \forall t \geq t_{0} \geq 0$. Since $a_{e}^{2}=d_{e}^{2}-z_{e}^{2}$, we have

$$
\begin{align*}
\overbrace{a_{e}^{2}}= & 2\left(d_{e}\left(-k_{1} \cos (\gamma) d_{e}+k_{1} \cos (\gamma) \delta_{e}-\cos (\gamma) \tilde{u}\right)-\right. \\
& \left.z_{e}\left(\frac{\partial z_{d}}{\partial s} \dot{s}+\sin (\theta) u-\cos (\theta) \sin (\phi) v-\cos (\theta) \cos (\phi) w\right)\right) \\
= & -2 k_{1} \cos (\gamma) a_{e}^{2}+2 k_{1} \cos (\gamma) z_{e}^{2}+2 k_{1} \cos (\gamma) d_{e} \delta_{e}-2 \cos (\gamma) d_{e} \tilde{u}- \\
& 2 z_{e}\left(\frac{\partial z_{d}}{\partial s} \dot{s}+\sin (\theta) u-\cos (\theta) \sin (\phi) v-\cos (\theta) \cos (\phi) w\right) \tag{13.94}
\end{align*}
$$

From (13.94), it is not hard to see that under Assumption 13.1 and $d_{e}(t) \geq d_{e}^{*}>0$, if there exists a strictly positive constant $a_{e}^{0}$ such that

$$
\begin{equation*}
a_{e}\left(t_{0}\right) \geq a_{e}^{0} \tag{13.95}
\end{equation*}
$$

then there exists a strictly positive constant $a_{e}^{*}$ such that the condition $a_{e}(t) \geq a_{e}^{*}>$ $0, \forall t \geq t_{0} \geq 0$ holds.

In summary, the feasible initial conditions are such that the conditions (13.77), (13.93) and (13.95) hold. Roughly speaking, with the above initial conditions, due to the underactuated configuration in the sway and heave, the sway and heave velocities are not able to push the vehicle to the point $a_{e}=0$ and $d_{e}=0$.

### 13.5 Discussion of the Initial Condition

We now discuss how to obtain the initial conditions such that (13.77), (13.93), and (13.95) hold. A close look at these conditions shows that they are always satisfied by selecting the initial value, $s\left(t_{0}\right)$, if the vessel heads toward the conical space containing the initial path to be followed, see Figure 13.2. If the vessel does not, the surge control should be turned off and the yaw and pitch controls should make the vessel turn until (13.77), (13.93), and (13.95) hold before applying the proposed path-following controller. The angle $\delta_{0}$ (see Figure 13.2) should be increased if the initial velocities $\boldsymbol{v}_{\mathbf{1}}\left(t_{0}\right)$ and $\boldsymbol{v}_{\mathbf{2}}\left(t_{0}\right)$ are large. Otherwise the vessel might cross the edge-line of the subspace in question, which might result in $\gamma_{i}= \pm 0.5 \pi$ and/or $\gamma= \pm 0.5 \pi$.

### 13.6 Parking and Point-to-point Navigation

### 13.6.1 Parking

Parking Objective. Design the controls $\boldsymbol{\tau}_{\mathbf{1}}$ and $\boldsymbol{\tau}_{\mathbf{2}}$ to park the underactuated underwater vehicle (13.1) from the initial position and orientation $\left(x\left(t_{0}\right), y\left(t_{0}\right), z\left(t_{0}\right)\right.$, $\left.\phi\left(t_{0}\right), \theta\left(t_{0}\right), \psi\left(t_{0}\right)\right)$ to the desired parking position and orientation of $\left(x_{p}, y_{p}, z_{p}\right.$, $\left.\phi_{p}, \theta_{p}, \psi_{p}\right)$ under the following conditions:

1. There exists a large enough positive constant $\varpi_{p}$ such that

$$
\sqrt{\left(x\left(t_{0}\right)-x_{p}\right)^{2}+\left(y\left(t_{0}\right)-y_{p}\right)^{2}+\left(z\left(t_{0}\right)-z_{p}\right)^{2}} \geq \varpi_{p} .
$$

2. The vessel heads toward the feasible cone containing the desired parking orientation, see Section 13.5.
3. At the desired parking position and orientation, the environmental disturbances are negligible.

The above conditions normally hold for parking practice. However, if the first two conditions do not hold, one can apply the strategy in Section 13.5 to move the vessel until they do hold. Having formulated the parking problem as above, one might claim that the path-following controller proposed in Section 13.3 can be applied by setting $u_{0}$ equal zero. However, this will result in an orientation that may be very different from the desired parking one, at the desired parking position, since


Figure 13.2 Feasible initial conditions
our proposed path-following controller is designed to drive $d_{e}$ to a small ball, not to zero for reasons of robustness. To resolve this problem, we first generate a regular curve, $\boldsymbol{\Omega}_{\boldsymbol{p}}\left(x_{d}, y_{d}, z_{d}\right)$, which goes via the parking position and its orientation at the parking position is equal to the desired parking condition. For simplicity of calculation, the curve can be taken as a straight line in almost all cases of the vessel's initial conditions. Then the proposed path-following controller can be used to make the vessel follow $\boldsymbol{\Omega}_{\boldsymbol{p}}\left(x_{d}, y_{d}, z_{d}\right)$. In this case, the velocity $u_{0}$ should be chosen such that it goes to zero when the virtual vessel tends to the desired parking position, i.e., $\lim _{d_{e p} \rightarrow 0} u_{0}=0$ with $\mathrm{d}_{e p}=\sqrt{\left(x_{d}-x_{p}\right)^{2}+\left(y_{d}-y_{p}\right)^{2}+\left(z_{d}-z_{p}\right)^{2}}$. A simple choice can be taken as

$$
\begin{equation*}
u_{0}=u_{0}^{*}\left(1-e^{-\chi_{1} d_{e p}}\right) e^{-\chi_{2} d_{e}}, \tag{13.96}
\end{equation*}
$$

where $\chi_{i}>0, i=1,2$. Special care should be taken to choose the initial values of $\left(x_{d}\left(t_{0}\right), y_{d}\left(t_{0}\right), z_{d}\left(t_{0}\right)\right)$, see Section 13.5, and the sign of $u_{0}$ such that it results in a short parking time.

Remark 13.4. Once at the desired parking position and orientation, if there are large environmental disturbances, there will be an oscillatory behavior in the yaw and pitch dynamics, and the vessel might diverge from its desired position. This phenomenon is well known in ship dynamic positioning.

### 13.6.2 Point-to-point Navigation

As seen in Section 13.1, the requirement that the reference path be a regular curve might be too cumbersome in practice, since this curve has to go via desired points generated by the helmsman and its derivatives are needed in the path-following controller. These restrictions motivate us to consider the point-to-point navigation problem as follows.

Point-to-point Navigation Objective. Design the surge force $\boldsymbol{\tau}_{\boldsymbol{u}}$ and the yaw moment $\boldsymbol{\tau}_{\boldsymbol{r}}$ to force the underactuated underwater vehicle (13.1) from the initial position and orientation, $\left(x\left(t_{0}\right), y\left(t_{0}\right), z\left(t_{0}\right), \phi\left(t_{0}\right), \theta\left(t_{0}\right), \psi\left(t_{0}\right)\right)$, to go via the desired points generated by a path planner.

To achieve this control objective, we first assume that the path planner generates desired points, which are feasible for the vehicle to be navigated through. We then apply the path-following controller proposed in Section 13.3 to each regular curve segments connecting desired points in sequence. The regular curve segments can be straight line, arc, or known regular curve ones. It is, however, noted that a fundamental difference between point-to-point navigation and the proposed smooth path-following is that there are a finite number of "peaks", equal to the number of points, in the orientation errors, $\gamma_{1}$ and $\gamma_{2}$. This phenomenon is because the path is non-smooth in the orientation at the points.

### 13.7 Numerical Simulations

This section validates the control laws (13.48) and (13.59) by simulating them on a 5.56 m long underwater vehicle whose parameters are given in Section 12.5. The values of the vehicle parameters are assumed to be of the real vessel and are estimated on-line by the adaptation laws (13.49) and (13.60). We assume that these parameters fluctuate around the above values $\pm 15 \%$. This fluctuation is chosen here for the purpose of calculating the maximum and minimum values used in the choice of the design constants. Indeed, a different fluctuation of the vehicle parameters results in different maximum and minimum values used in the choice of the design constants. In the simulation, we assume that the environmental disturbances are

$$
\begin{aligned}
& \tau_{w u}=0.2 m_{11} d(t), \tau_{w v}=0.2 m_{22} d(t), \tau_{w w}=0.2 m_{33} d(t), \tau_{w p}=0.2 m_{44} d(t), \\
& \tau_{w q}=0.2 m_{55} d(t), \tau_{w r}=0.2 m_{66} d(t)
\end{aligned}
$$

where $d(t)=1+0.1 \sin (0.2 t)$. This choice results in nonzero-mean disturbances. In practice, the environmental disturbances may be different. We take the above disturbances for an illustration of the robustness properties of our proposed controller. It should be noted that only upper bounds of the environmental disturbances are needed in our proposed controller.

In the simulation, based on Section 13.3 the control parameters and initial conditions are taken as

$$
\begin{aligned}
& k_{1}=0.5, c_{1}=2, \boldsymbol{K}_{\mathbf{2}}=\operatorname{diag}(0.05), \boldsymbol{K}_{\mathbf{3}}=\operatorname{diag}(2), \boldsymbol{\Gamma}_{\boldsymbol{i}}=\operatorname{diag}(10), \delta_{e}=0.2, \\
& {\left[\boldsymbol{\eta}_{\mathbf{1}}{ }^{T}\left(t_{0}\right), \boldsymbol{\eta}_{2}^{T}\left(t_{0}\right), \boldsymbol{v}_{\mathbf{1}}^{T}\left(t_{0}\right), \boldsymbol{v}_{2}^{T}\left(t_{0}\right), s\left(t_{0}\right)\right]^{T}=} \\
& {[-145,-15,-5,0,0.2,0.5,0,0,0,0,0,0,0]^{T},}
\end{aligned}
$$

and all initial values of parameter estimates are taken to be $70 \%$ of their assumed true ones. The virtual vessel velocity on the path is taken as $u_{0}\left(t, d_{e}\right)=$ $5\left(1-0.8 e^{-2 t}\right) e^{-0.5 d_{e}}$. The reference path is given by $\left(x_{d}=-90 \cos (s), y_{d}=\right.$
$90 \sin (s), z_{d}=3 s$ ), i.e., a helix with constant curvature and torsion. Figure 13.3 plots the trajectory of the vessel and the path (dotted line) to be followed in three dimensions. The trajectory of the vessel in the horizontal plane and the path following error are plotted in Figure 13.4. Figure 13.5 plots the control inputs, $\tau_{u}, \tau_{r}, \tau_{p}$, and $\tau_{q}$. With nonvanishing environmental disturbances, our proposed controller is able to force the vehicle to follow a predefined path as expected in the control design. As can be seen in Figure 13.5, $d_{e}$ converges to a nonzero small value, i.e., the sway and heave velocities cannot push the vessel to the point where $d_{e}=0$. From Figure 13.5, it can be seen that the control inputs are below their limits. Therefore, we can still further shorten the transient time by increasing the control gains.

### 13.8 Conclusions

The control scheme developed for path-following of underactuated surface ships in Chapter 11 was extended to design a path-following system for six degrees of freedom underactuated underwater vehicles. The key to the development of the proposed path-following system is the proper selection of the coordinate transformations in Section 13.2. The work presented in this chapter is based on [140, 141].


Figure 13.3 Simulation results: Path-following real and reference trajectories in three-dimensional space


Figure 13.4 Simulation results: a. Path-following trajectory in the horizontal plane; b. Pathfollowing error $d_{e}$


Figure 13.5 Simulation results (control inputs): a. Surge force $\tau_{u}[\mathrm{~N}]$; b. Yaw moment $\tau_{r}[\mathrm{Nm}]$; c. Pitch moment $\tau_{p}[\mathrm{Nm}]$; d. Roll moment $\tau_{q}[\mathrm{Nm}]$

# Part V <br> Control of Other Underactuated <br> Mechanical Systems 

## Chapter 14 <br> Control of Other Underactuated Mechanical Systems

This chapter presents several applications of the observer and control design techniques developed in the previous chapters to the control of other underactuated mechanical systems including mobile robots and VTOL aircraft. For mobile robots, a global exponential observer is first designed based on the observer design for underactuated ships in Chapter 7. Output feedback simultaneous stabilization and trajectory-tracking, and path-following controllers are then developed using the control design techniques proposed for underactuated ships in Chapters 6 and 11. For VTOL aircraft, the observer and control design strategies used for underactuated ships in Chapters 5 and 6 are utilized to design a global output feedback trajectorytracking controller.

### 14.1 Mobile Robots

### 14.1.1 Basic Motion Tasks

In order to derive the most suitable feedback controllers for each case, it is convenient to classify the possible motion tasks as follows:

- Point-to-point motion: The robot must reach a desired goal configuration starting from a given initial configuration, see Figure 14.1a.
- Path-following: The robot must reach and follow a geometric path in the Cartesian space starting from a given initial configuration (on or off the path), see Figure 14.1b.
- Trajectory-tracking: The robot must reach and follow a trajectory in the Cartesian space (i.e., a geometric path with an associated timing law) starting from a given initial configuration (on or off the trajectory), see Figure 14.1c.

The three tasks are sketched in Figure 14.1 with reference to a three wheel car-like robot. Execution of these tasks can be achieved using either feedforward commands, or feedback control, or a combination of the two. Indeed, feedback solutions exhibit


Figure 14.1 Basic motion tasks for a mobile robot
an intrinsic degree of robustness. Explanations of the above motion tasks are similar to those of ocean vehicles, see Section 3.2.

### 14.1.2 Modeling and Control Properties

### 14.1.2.1 Modeling

We consider a unicycle-type mobile robot, which under the assumption of no wheel slips has the following dynamics [142]:

$$
\begin{align*}
& \dot{\eta}=J(\eta) \omega  \tag{14.1}\\
& M \dot{\omega}=-C(\dot{\eta}) \omega-D \omega+\tau
\end{align*}
$$

where $\eta=\left[\begin{array}{ll}x & y\end{array} \phi\right]^{T}$ denotes the position $(x, y)$, the coordinates of the middle point, $P_{0}$, between the left and right driving wheels, and heading angle $\phi$ of the robot coordinated in the earth-fixed frame $O X Y$, see Figure 14.2, $\omega=\left[\omega_{1} \omega_{2}\right]^{T}$, with $\omega_{1}$ and $\omega_{2}$ being the angular velocities of the wheels of the robot, $\boldsymbol{\tau}=\left[\begin{array}{ll}\tau_{1} & \tau_{2}\end{array}\right]^{T}$, with $\tau_{1}$ and $\tau_{2}$ being the control torques applied to the wheels of the robot. The rotation matrix $\boldsymbol{J}(\boldsymbol{\eta})$, mass matrix $\boldsymbol{M}$, Coriolis matrix $\boldsymbol{C}(\dot{\boldsymbol{\eta}})$, and damping matrix $\boldsymbol{D}$ in (14.1) are given by

$$
\begin{align*}
& \boldsymbol{J}(\boldsymbol{\eta})=\frac{r}{2}\left[\begin{array}{cc}
\cos (\phi) & \cos (\phi) \\
\sin (\phi) & \sin (\phi) \\
b^{-1} & -b^{-1}
\end{array}\right], \boldsymbol{M}=\left[\begin{array}{ll}
m_{11} & m_{12} \\
m_{12} & m_{11}
\end{array}\right], \\
& \boldsymbol{C}(\dot{\boldsymbol{\eta}})=\left[\begin{array}{cc}
0 & c \dot{\phi} \\
-c \dot{\phi} & 0
\end{array}\right], \boldsymbol{D}=\left[\begin{array}{cc}
d_{11} & 0 \\
0 & d_{22}
\end{array}\right], \tag{14.2}
\end{align*}
$$

with

$$
\begin{align*}
& c=\frac{1}{2 b} r^{2} m_{c} a, m_{11}=\frac{1}{4 b^{2}} r^{2}\left(m b^{2}+I\right)+I_{w}, m_{12}=\frac{1}{4 b^{2}} r^{2}\left(m b^{2}-I\right) \\
& m=m_{c}+2 m_{w}, I=m_{c} a^{2}+2 m_{w} b^{2}+I_{c}+2 I_{m} \tag{14.3}
\end{align*}
$$

where $m_{c}$ and $m_{w}$ are the masses of the body and wheel with a motor; $I_{c}, I_{w}$, and $I_{m}$ are the moments of inertia of the body about the vertical axis through $P_{c}$ (center of mass), the wheel with the rotor of a motor about the wheel axis, and the wheel with the rotor of a motor about the wheel diameter, respectively; $r, a$, and $b$ are defined in Figure 14.2; the nonnegative constants $d_{11}$ and $d_{22}$ are the damping coefficients. If these damping coefficients are zero, we have an undamped case. On the other hand, if the damping coefficients are positive, we have a damped case. We take the physical parameters from [142]: $b=0.75 \mathrm{~m}, a=0.3 \mathrm{~m}, r=0.15$ $\mathrm{m}, m_{c}=30 \mathrm{~kg}, m_{w}=1 \mathrm{~kg}, I_{c}=15.625 \mathrm{kgm}^{2}, I_{w}=0.005 \mathrm{kgm}^{2}, I_{m}=0.0025$ $\mathrm{kgm}^{2}, d_{11}=d_{22}=5 \mathrm{kgs}^{-1}$ for numerical simulations. For convenience, we convert the wheel angular velocities $\left(\omega_{1}, \omega_{2}\right)$ of the robot to its linear, $v$, and angular, $w$, velocities by:

$$
\varpi=\boldsymbol{B}^{-1} \boldsymbol{\omega}, \boldsymbol{B}=\frac{1}{r}\left[\begin{array}{cc}
1 & b  \tag{14.4}\\
1 & -b
\end{array}\right],
$$

where $\boldsymbol{\varpi}=\left[\begin{array}{ll}v & w\end{array}\right]^{T}$ and $\boldsymbol{B}$ is invertible since $\operatorname{det}(\boldsymbol{B})=-2 b / r$. With (14.4), we can write the robot dynamics (14.1) as follows:

$$
\begin{align*}
& \dot{x}=v \cos (\phi), \\
& \dot{y}=v \sin (\phi), \\
& \dot{\phi}=w,  \tag{14.5}\\
& \overline{\boldsymbol{M}} \dot{\boldsymbol{\omega}}=-\overline{\boldsymbol{C}}(w) \boldsymbol{\varpi}-\overline{\boldsymbol{D}} \boldsymbol{\varpi}+\overline{\boldsymbol{B}} \boldsymbol{\tau}
\end{align*}
$$



Figure 14.2 Robot parameters
where

$$
\begin{align*}
& \overline{\boldsymbol{M}}=\boldsymbol{B}^{-\mathbf{1}} \boldsymbol{M} \boldsymbol{B}=\left[\begin{array}{cc}
\bar{m}_{11} & 0 \\
0 & \bar{m}_{22}
\end{array}\right] \\
& \overline{\boldsymbol{C}}(w)=\boldsymbol{B}^{-1} \boldsymbol{C}(\dot{\boldsymbol{\eta}}) \boldsymbol{B}=\left[\begin{array}{cc}
0 & -b c w \\
\frac{c}{b} w & 0
\end{array}\right] \\
& \overline{\boldsymbol{D}}=\boldsymbol{B}^{-\mathbf{1}} \boldsymbol{D} \boldsymbol{B}=\left[\begin{array}{l}
\bar{d}_{11} \\
\bar{d}_{21} \\
\bar{d}_{22}
\end{array}\right], \overline{\boldsymbol{B}}=\boldsymbol{B}^{-1}  \tag{14.6}\\
& \bar{m}_{11}=m_{11}+m_{12}, \bar{m}_{22}=m_{11}-m_{12} \\
& \bar{d}_{11}=\frac{1}{2}\left(d_{11}+d_{22}\right), \bar{d}_{12}=\frac{b}{2}\left(d_{11}-d_{22}\right), \\
& \bar{d}_{21}=\frac{1}{2 b}\left(d_{11}-d_{22}\right), \bar{d}_{22}=\frac{1}{2}\left(d_{11}+d_{22}\right)
\end{align*}
$$

### 14.1.2.2 Control Properties

Since the last equation of (14.5) is a square system if we consider the robot velocities $v$ and $w$ as outputs and the torques $\tau_{1}$ and $\tau_{2}$ as inputs, we only need to investigate control properties of the first three equations of (14.5), i.e., we investigate the control properties of the robot kinematic model:

$$
\begin{align*}
\dot{x} & =v \cos (\phi), \\
\dot{y} & =v \sin (\phi), \\
\dot{\phi} & =w . \tag{14.7}
\end{align*}
$$

From the first two equations of (14.7), the nonholonomic constraint is

$$
\begin{equation*}
\dot{x} \sin (\phi)-\dot{y} \cos (\phi)=0 . \tag{14.8}
\end{equation*}
$$

Controllability at a Point. The tangent linearization of (14.7) at any point $\boldsymbol{\eta}_{\boldsymbol{e}}$ is the linear system

$$
\dot{\tilde{\eta}}=\left[\begin{array}{c}
\cos \left(\phi_{e}\right)  \tag{14.9}\\
\sin \left(\phi_{e}\right) \\
0
\end{array}\right] v+\left[\begin{array}{l}
0 \\
0 \\
1
\end{array}\right] w, \tilde{\boldsymbol{\eta}}=\boldsymbol{\eta}-\boldsymbol{\eta}_{\boldsymbol{e}}
$$

which is clearly not controllable. This implies that a linear controller will never achieve posture stabilization, not even in a local sense. In order to study the controllability of the unicycle, we need to use tools from nonlinear control theory [4]. Let us define

$$
\boldsymbol{g}_{\mathbf{1}}=\left[\begin{array}{c}
\cos (\phi)  \tag{14.10}\\
\sin (\phi) \\
0
\end{array}\right], \boldsymbol{g}_{\mathbf{2}}=\left[\begin{array}{l}
0 \\
0 \\
1
\end{array}\right]
$$

It is easy to check that the accessibility rank condition is satisfied globally (at any $\eta_{e}$ ), since

$$
\begin{equation*}
\operatorname{rank}\left[g_{1} \quad g_{2} \quad\left[g_{1}, g_{2}\right]\right]=3 \tag{14.11}
\end{equation*}
$$

where the Lie bracket $\left[\boldsymbol{g}_{\mathbf{1}}, \boldsymbol{g}_{\mathbf{2}}\right]$ of the two input vector fields $\boldsymbol{g}_{\mathbf{1}}$ and $\boldsymbol{g}_{\mathbf{2}}$ is

$$
\left[\boldsymbol{g}_{\mathbf{1}}, \boldsymbol{g}_{\mathbf{2}}\right]=\frac{\partial \boldsymbol{g}_{\mathbf{2}}}{\partial \boldsymbol{\eta}} \boldsymbol{g}_{\mathbf{1}}-\frac{\partial \boldsymbol{g}_{\mathbf{1}}}{\partial \boldsymbol{\eta}} \boldsymbol{g}_{\mathbf{2}}=\left[\begin{array}{c}
\sin (\phi)  \tag{14.12}\\
-\cos (\phi) \\
0
\end{array}\right]
$$

Since the system is driftless, condition (14.11) implies its controllability. Controllability can also be shown constructively, i.e., by providing an explicit sequence of maneuvers bringing the robot from any start configuration ( $x_{s}, y_{s}, \phi_{s}$ ) to any desired goal configuration $\left(x_{g}, y_{g}, \phi_{g}\right)$. Since the unicycle can rotate on itself, this task is simply achieved by an initial rotation on ( $x_{s}, y_{s}$ ) until the unicycle is oriented toward $\left(x_{g}, y_{g}\right)$, followed by a translation to the goal position, and by a final rotation on $\left(x_{g}, y_{g}\right)$ so as to align $\phi$ with $\phi_{g}$. As for the stabilizability of system (14.9) to a point, the failure of the previous linear analysis indicates that exponential stability cannot be achieved by smooth feedback [5]. Things turn out to be even worse: If smooth (in fact, even continuous) time-invariant feedback laws are used, Lyapunov stability is not applicable. This negative result is established on the basis of a necessary condition due to Brockett [21]: Smooth stabilizability of a driftless regular system (i.e., such that the input vector fields are well defined and linearly independent at $\eta_{\boldsymbol{e}}$ ) requires a number of inputs equal to the number of states. The above difficulty has a deep impact on the control design. In fact, to obtain a posture stabilizing controller it is either necessary to give up the continuity requirement and/or
to resort to time-varying control laws.
Controllability About a Trajectory. Given a desired Cartesian motion for the unicycle mobile robot, it may be convenient to generate a corresponding state trajectory $\eta_{\boldsymbol{d}}(t)=\left(x_{d}(t), y_{d}(t), \phi_{d}(t)\right)$. In order to be feasible, the latter must satisfy the nonholonomic constraint on the vehicle motion or, equivalently, be consistent with (14.8). The generation of $\boldsymbol{\eta}_{\boldsymbol{d}}(t)$ and of the corresponding reference velocity inputs $v_{d}(t)$ and $w_{d}(t)$ will be addressed properly.

Defining the state tracking error as $\tilde{\boldsymbol{\eta}}=\boldsymbol{\eta}-\boldsymbol{\eta}_{\boldsymbol{d}}$ and the input variations as $\tilde{v}=v-$ $v_{d}$ and $\tilde{w}=w-w_{d}$, the tangent linearization of system (14.7) about the reference trajectory is

$$
\dot{\tilde{\boldsymbol{\eta}}}=\left[\begin{array}{ccc}
0 & 0 & -v_{d} \sin \left(\phi_{d}\right)  \tag{14.13}\\
0 & 0 & v_{d} \cos \left(\phi_{d}\right) \\
0 & 0 & 0
\end{array}\right] \tilde{\boldsymbol{\eta}}+\left[\begin{array}{cc}
\cos \left(\phi_{d}\right) & 0 \\
\sin \left(\phi_{d}\right) & 0 \\
0 & 1
\end{array}\right]\left[\begin{array}{c}
\tilde{v} \\
\tilde{w}
\end{array}\right]=\boldsymbol{A}(t) \tilde{\boldsymbol{\eta}}+\boldsymbol{B}(t)\left[\begin{array}{c}
\tilde{v} \\
\tilde{w}
\end{array}\right] .
$$

Since the linearized system is time-varying, a necessary and sufficient controllability condition is that the controllability Grammian is nonsingular. However, a simpler analysis can be conducted by defining the state tracking error as

$$
\tilde{\eta}_{\boldsymbol{R}}=\left[\begin{array}{ccc}
\cos \left(\phi_{d}\right) & \sin \left(\phi_{d}\right) & 0  \tag{14.14}\\
-\sin \left(\phi_{d}\right) & \cos \left(\phi_{d}\right) & 0 \\
0 & 0 & 1
\end{array}\right] \tilde{\boldsymbol{\eta}} .
$$

Using (14.7), we obtain

$$
\dot{\tilde{\eta}}_{\boldsymbol{R}}=\left[\begin{array}{ccc}
0 & w_{d} & 0  \tag{14.15}\\
-w_{d} & 0 & v_{d} \\
0 & 0 & 0
\end{array}\right] \tilde{\eta}_{\boldsymbol{R}}+\left[\begin{array}{ll}
1 & 0 \\
0 & 0 \\
0 & 1
\end{array}\right]\left[\begin{array}{c}
\tilde{v} \\
\tilde{w}
\end{array}\right] .
$$

When $v_{d}$ and $w_{d}$ are constant, the above linear system becomes time-invariant and controllable, since matrix

$$
\boldsymbol{C}=\left[\begin{array}{lll}
\boldsymbol{B} & \boldsymbol{A} \boldsymbol{B} & \boldsymbol{A}^{\mathbf{2}} \boldsymbol{B}
\end{array}\right]=\left[\begin{array}{ccccccc}
1 & 0 & 0 & 0 & -w_{d}^{2} & v_{d} w_{d}  \tag{14.16}\\
0 & 0 & -w_{d} & v_{d} & 0 & 0 \\
0 & 1 & 0 & 0 & 0 & 0
\end{array}\right]
$$

has rank 3 provided that either $v_{d}$ or $w_{d}$ are nonzero. Therefore, we conclude that the kinematic system (14.7) can be locally stabilized by linear feedback about trajectories that consist of linear or circular paths, executed with constant velocity.

### 14.1.2.3 Feedback Linearizability

Based on the previous discussion, it is easy to see that the driftless nonholonomic system (14.7) cannot be transformed into a linear controllable one using static state feedback. In particular, the controllability condition (14.11) implies that the distribution generated by vector fields $\boldsymbol{g}_{\mathbf{1}}$ and $\boldsymbol{g}_{\mathbf{2}}$ is not involutive, thus violating the necessary condition for full-state feedback linearizability [4]. However, when matrix

$$
\boldsymbol{G}(\boldsymbol{\eta})=\left[\begin{array}{cc}
\cos (\phi) & 0  \tag{14.17}\\
\sin (\phi) & 0 \\
0 & 1
\end{array}\right]
$$

has full column rank, two equations can always be transformed via feedback into simple integrators (input-output linearization and decoupling). The choice of the linearizing outputs is not unique and can be accommodated for special purposes as illustrated in the following example. Define the two outputs as

$$
\begin{align*}
& y_{1}=x+d \cos (\phi) \\
& y_{2}=y+d \sin (\phi) \tag{14.18}
\end{align*}
$$

with $d \neq 0$, i.e., the Cartesian coordinates of a point, B , displaced at a distance $d$ along the main axis of the unicycle.

Using the globally defined state feedback

$$
\left[\begin{array}{c}
v  \tag{14.19}\\
w
\end{array}\right]=\left[\begin{array}{cc}
\cos (\phi) & \sin (\phi) \\
-\sin (\phi) / d & \cos (\phi) / d
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
u_{2}
\end{array}\right],
$$

the unicycle kinematic is equivalent to

$$
\begin{align*}
& \dot{y}_{1}=u_{1} \\
& \dot{y}_{1}=u_{1} \\
& \dot{\phi}=\frac{u_{2} \cos (\phi)-u_{1} \sin (\phi)}{d} \tag{14.20}
\end{align*}
$$

As a consequence, a linear feedback controller for $u=\left(u_{1}, u_{2}\right)$ will make the point B track any reference trajectory, even with discontinuous tangent to the path (e.g., a square without stopping at corners). Moreover, it is easy to show that the internal state evolution $\phi(t)$ is bounded. This approach, however, will not be pursued in this chapter because of its limited interest, since the robot orientation $\phi(t)$ is not controlled.

### 14.1.2.4 Chained Forms

The existence of canonical forms for kinematic models of nonholonomic robots allows a general and systematic development of both open loop and closed loop control strategies. The most useful canonical structure is the chained form, which in the case of two-input systems is

$$
\begin{gather*}
\dot{z}_{1}=u_{1} \\
\dot{z}_{2}=u_{2} \\
\dot{z}_{3}=z_{2} u_{1}  \tag{14.21}\\
\quad \vdots \\
\dot{z}_{n}=z_{n-1} u_{1}
\end{gather*}
$$

It has been shown that a two-input driftless nonholonomic system with up to $n=4$ generalized coordinates can always be transformed into a chained form by static feedback transformation [143]. As a matter of fact, most (but not all) wheeled mobile robots can be transformed in a chained form. For the kinematic model (14.7) of the unicycle, we introduce the following globally defined coordinate transformation

$$
\begin{align*}
& z_{1}=\phi, \\
& z_{2}=x \cos (\phi)+y \sin (\phi),  \tag{14.22}\\
& z_{3}=x \sin (\phi)-y \cos (\phi),
\end{align*}
$$

and a static state feedback

$$
\begin{align*}
& v=u_{2}+z_{3} u_{1}, \\
& w=u_{1} \tag{14.23}
\end{align*}
$$

to give

$$
\begin{align*}
& \dot{z}_{1}=u_{1}, \\
& \dot{z}_{2}=u_{2}  \tag{14.24}\\
& \dot{z}_{3}=z_{2} u_{1} .
\end{align*}
$$

Note that $\left(z_{2}, z_{3}\right)$ is the position of the unicycle in a rotating left-hand frame having the $z_{2}$-axis aligned with the vehicle orientation. Equation (14.24) is another example of static input output linearization, with $z_{1}$ and $z_{2}$ as linearizing outputs. We note also that the transformation in a chained form is not unique. Mathematical model and control properties of other types of mobile robots are given in [15, 130]. The reader is referred to $[22,33-40,46,51,52,144-146]$ for various control design methods ranging from discontinuous to time-varying ones on stabilization and trajectory-tracking control of systems that cover the chained form (14.21).

### 14.1.3 Output Feedback Simultaneous Stabilization and Trajectory-tracking

### 14.1.3.1 Control Objective

For the reader's convenience, we rewrite the equations of motion here, see (14.1):

$$
\begin{align*}
& \dot{\eta}=J(\eta) \omega \\
& M \dot{\omega}+C(\dot{\eta}) \omega+D \omega=\tau \tag{14.25}
\end{align*}
$$

We assume that the reference trajectory is generated by the following virtual robot:

$$
\begin{align*}
\dot{x}_{d} & =\cos \left(\phi_{d}\right) u_{1 d}, \\
\dot{y}_{d} & =\sin \left(\phi_{d}\right) u_{1 d},  \tag{14.26}\\
\dot{\phi}_{d} & =u_{2 d},
\end{align*}
$$

where $\left(x_{d}, y_{d}, \phi_{d}\right)$ are the position and orientation of the virtual robot, and $u_{1 d}$ and $u_{2 d}$ are the linear and angular velocities of the virtual robot, respectively.

Control Objective. Under Assumption 14.1, design the control input vector $\boldsymbol{\tau}$ to force the position and orientation, $(x, y, \phi)$ of the real robot (14.25) to globally asymptotically track $\left(x_{d}, y_{d}, \phi_{d}\right)$ generated by (14.26) with only ( $x, y, \phi$ ) available for feedback.

## Assumption 14.1.

1. The reference signals $u_{1 d}, \dot{u}_{1 d}, \ddot{u}_{1 d}, u_{2 d}$, and $\dot{u}_{2 d}$ are bounded. In addition, one of the following conditions holds:

$$
\begin{align*}
& \text { C1. } \int_{0}^{\infty}\left(\left|u_{1 d}(t)\right|+\left|u_{2 d}(t)\right|\right) \mathrm{d} t \leq \mu_{11} \\
& \text { C2. } \int_{0}^{\infty}\left|u_{1 d}(t)\right| \mathrm{d} t \leq \mu_{21} \text { and }\left|u_{2 d}(t)\right| \geq \mu_{22}  \tag{14.27}\\
& \text { C3. } \int_{t_{0}}^{t} u_{1 d}^{2}(\tau) \mathrm{d} \tau \geq \mu_{31}\left(t-t_{0}\right)-\mu_{32}, \forall t \geq t_{0} \geq 0
\end{align*}
$$

where $\mu_{11}, \mu_{21}$, and $\mu_{32}$ are nonnegative constants, and $\mu_{22}$ and $\mu_{31}$ are strictly positive constants.
2. The robot wheel velocities $\boldsymbol{\omega}=\left[\omega_{1} \omega_{2}\right]^{T}$ are not available for feedback.

Remark 14.1. The problem of set-point regulation/stabilization, tracking a path approaching a set-point is included in Condition C1. Tracking linear and circular paths
belongs to Condition C3. Condition C2 implies that the case, where the robot linear velocity is zero or approaches zero and its angular velocity is of sinusoidal type, is excluded. The reason is that our control approach introduces a sinusoid signal in the robot angular velocity virtual control to handle set-point stabilization/regulation. Therefore, this case is excluded to avoid two signals canceling each other. If the reference velocity $u_{2 d}$ is known completely in advance, the above case can be included. Moreover, item (2) of Assumption 14.1 implies that we need to design an output feedback controller.

Remark 14.2. The problem of simultaneous stabilization and tracking is not only of theoretical interest but also possesses some advantages over the use of separate stabilization and tracking controllers such as only one controller and transient improvement because of the lack of switching. Moreover, if the switching time is unknown, a separate stabilization and tracking control approach cannot be used.

### 14.1.3.2 Observer Design

We first remove the quadratic velocity terms in the mobile robot dynamics by introducing the following coordinate change:

$$
\begin{equation*}
X=\boldsymbol{Q}(\eta) \omega \tag{14.28}
\end{equation*}
$$

where $\boldsymbol{Q}(\boldsymbol{\eta})$ is a globally invertible matrix with bounded elements to be determined. Using (14.28), we write the second equation of (14.25) as follows:

$$
\begin{equation*}
\dot{X}=\left[\dot{Q}(\eta) \omega-Q(\eta) M^{-\mathbf{1}} C(\dot{\eta}) \omega\right]+Q(\eta) M^{-1}(-D \omega+\tau) \tag{14.29}
\end{equation*}
$$

In [101], $\boldsymbol{Q}(\boldsymbol{\eta})$ is required with the above properties such that

$$
\dot{Q}(\eta)=Q(\eta) M^{-1} C(\dot{\eta}), \forall \eta \in \mathbb{R}^{3},
$$

which does not exist, as a simple calculation shows.
Our method is to cancel the square bracket on the right-hand side of (14.29) for all $(\boldsymbol{\eta}, \boldsymbol{\omega}) \in \mathbb{R}^{5}$. We assume that $q_{i j}(\boldsymbol{\eta}), i=1,2, j=1,2$, are the elements of $\boldsymbol{Q}(\boldsymbol{\eta})$. Using the first equation of (14.25), it can readily be shown that the above square bracket is zero for all $(\boldsymbol{\eta}, \boldsymbol{\omega}) \in \mathbb{R}^{5}$ if

$$
\begin{align*}
& \frac{\partial q_{i 1}}{\partial x} \cos (\phi)+\frac{\partial q_{i 1}}{\partial y} \sin (\phi)+\frac{\partial q_{i 1}}{\partial \phi} \frac{1}{b}+\frac{n_{12} c}{b} q_{i 1}+\frac{n_{11} c}{b} q_{i 2}=0 \\
& \frac{\partial q_{i 2}}{\partial x} \cos (\phi)+\frac{\partial q_{i 2}}{\partial y} \sin (\phi)-\frac{\partial q_{i 2}}{\partial \phi} \frac{1}{b}+\frac{n_{11} c}{b} q_{i 1}+\frac{n_{12} c}{b} q_{i 2}=0 \\
& \left(\frac{\partial q_{i 1}}{\partial x}+\frac{\partial q_{i 2}}{\partial x}\right) \cos (\phi)+\left(\frac{\partial q_{i 1}}{\partial y}+\frac{\partial q_{i 2}}{\partial y}\right) \sin (\phi)+  \tag{14.30}\\
& \left(\frac{\partial q_{i 2}}{\partial \phi}-\frac{\partial q_{i 1}}{\partial \phi}\right) \frac{1}{b}-\left(n_{11}+n_{12}\right) \frac{c}{b}\left(q_{i 1}+q_{i 2}\right)=0 .
\end{align*}
$$

Using the characteristic method to solve the above partial differential equations gives a family of solutions with $i=1,2$

$$
\begin{align*}
& q_{i 1}=C_{i 1} \sin (c \Delta \phi)+C_{i 2} \cos (c \Delta \phi) \\
& q_{i 2}=n_{11}^{-1}\left(\left(C_{i 2} \Delta-C_{i 1} n_{12}\right) \sin (c \Delta \phi)-\left(C_{i 1} \Delta+C_{i 2} n_{12}\right) \cos (c \Delta \phi)\right), \tag{14.31}
\end{align*}
$$

where $n_{11}=m_{11}\left(m_{11}^{2}-m_{12}^{2}\right)^{-1}, n_{12}=-m_{12}\left(m_{11}^{2}-m_{12}^{2}\right)^{-1}, \Delta=\sqrt{n_{11}^{2}-n_{12}^{2}}$, and $C_{i 1}$ and $C_{i 2}$ are arbitrary constants. A choice of $C_{11}=C_{22}=0, C_{12}=C_{21}=$ $n_{11}$ results in

$$
\boldsymbol{Q}(\boldsymbol{\eta})=\left[\begin{array}{ll}
n_{11} \cos (a \Delta \phi) & \Delta \sin (a \Delta \phi)-n_{12} \cos (a \Delta \phi)  \tag{14.32}\\
n_{11} \sin (a \Delta \phi) & -n_{12} \sin (a \Delta \phi)-\Delta \cos (a \Delta \phi)
\end{array}\right]
$$

This matrix is globally invertible and its elements are bounded. Now we write (14.25) in the ( $\boldsymbol{\eta}, \boldsymbol{X}$ ) coordinates as

$$
\begin{align*}
& \dot{\eta}=J(\eta) Q^{-\mathbf{1}}(\eta) X \\
& \dot{X}=-D_{\eta}(\eta) X+Q(\eta) M^{-1} \tau \tag{14.33}
\end{align*}
$$

where $\boldsymbol{D}_{\boldsymbol{\eta}}(\boldsymbol{\eta})=\boldsymbol{Q}(\boldsymbol{\eta}) \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{D} \boldsymbol{Q}^{-\mathbf{1}}(\boldsymbol{\eta})$. It can be seen that (14.33) is linear in the unmeasured states. We here use the following passive observer:

$$
\begin{align*}
& \dot{\hat{\eta}}=J(\eta) Q^{-1}(\eta) \hat{X}+K_{01}(\eta-\hat{\eta}), \\
& \dot{\hat{X}}=-D_{\eta}(\eta) \hat{X}+Q(\eta) M^{-1} \tau+K_{02}(\eta-\hat{\eta}), \tag{14.34}
\end{align*}
$$

where $\hat{\boldsymbol{\eta}}$ and $\hat{\boldsymbol{X}}$ are the estimates of $\boldsymbol{\eta}$ and $\boldsymbol{X}$, respectively. The observer gain matrices $\boldsymbol{K}_{\mathbf{0 1}}$ and $\boldsymbol{K}_{\mathbf{0 2}}$ are chosen such that

$$
\begin{aligned}
& Q_{01}=K_{01}^{T} P_{01}+P_{01} K_{01} \\
& Q_{02}=D_{\eta}^{T}(\eta) P_{02}+P_{02} D_{\eta}(\eta)
\end{aligned}
$$

are positive definite and

$$
\begin{equation*}
\left(J(\eta) Q^{-1}(\eta)\right)^{T} P_{01}-P_{02} K_{02}=0 \tag{14.35}
\end{equation*}
$$

with $\boldsymbol{P}_{\mathbf{0 1}}$ and $\boldsymbol{P}_{\mathbf{0 2}}$ being positive definite matrices. Since $\boldsymbol{D}_{\boldsymbol{\eta}}(\boldsymbol{\eta})$ is positive definite, $\boldsymbol{K}_{\mathbf{0 1}}$ and $\boldsymbol{K}_{\mathbf{0 2}}$ always exist. From (14.34) and (14.33), we have

$$
\begin{align*}
& \dot{\tilde{\eta}}=J(\eta) Q^{-1}(\eta) \tilde{X}-K_{01} \tilde{\eta}, \\
& \dot{\tilde{X}}=-D_{\eta}(\eta) \tilde{X}-K_{02} \tilde{\eta}, \tag{14.36}
\end{align*}
$$

where $\tilde{\eta}=\eta-\hat{\boldsymbol{\eta}}$ and $\tilde{\boldsymbol{X}}=\boldsymbol{X}-\hat{\boldsymbol{X}}$. It can now be seen that (14.36) is GES at the origin by taking the Lyapunov function

$$
V_{0}=\tilde{\eta}^{\boldsymbol{T}} \boldsymbol{P}_{\mathbf{0 1}} \tilde{\eta}+\tilde{\boldsymbol{X}}^{\boldsymbol{T}} \boldsymbol{P}_{\mathbf{0 2}} \tilde{\boldsymbol{X}}
$$

whose derivative along the solutions of (14.36) and using (14.35) satisfies

$$
\dot{V}_{0}=-\tilde{\eta}^{\boldsymbol{T}} Q_{\mathbf{0 1}} \tilde{\eta}-\tilde{X}^{\boldsymbol{T}} Q_{\mathbf{0 2}} \tilde{X}
$$

which in turn implies that there exists a strictly positive constant $\sigma_{0}$ such that

$$
\begin{equation*}
\|(\tilde{\boldsymbol{\eta}}(t), \tilde{\boldsymbol{X}}(t))\| \leq\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\| e^{-\sigma_{0}\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0 \tag{14.37}
\end{equation*}
$$

Define $\hat{\boldsymbol{\omega}}=\left[\hat{\omega}_{1} \hat{\omega}_{2}\right]^{T}$ as an estimator of the velocity vector $\boldsymbol{\omega}$ as

$$
\begin{equation*}
\hat{\omega}=Q^{-1}(\eta) \hat{X} . \tag{14.38}
\end{equation*}
$$

The velocity estimate error vector, $\tilde{\boldsymbol{\omega}}=\boldsymbol{\omega}-\hat{\boldsymbol{\omega}}$, satisfies

$$
\begin{equation*}
\tilde{\omega}=Q^{-1}(\eta) \tilde{X} \tag{14.39}
\end{equation*}
$$

To prepare for the control design in the next section, we convert the wheel velocities $\omega_{1}$ and $\omega_{2}$ to the linear, $v$, and angular, $w$, velocities of the robot by the relationship:

$$
\left[\begin{array}{c}
v  \tag{14.40}\\
w
\end{array}\right]=\boldsymbol{B}^{-1}\left[\begin{array}{l}
\omega_{1} \\
\omega_{2}
\end{array}\right] \text { with } \boldsymbol{B}=\frac{1}{r}\left[\begin{array}{cc}
1 & b \\
1 & -b
\end{array}\right]
$$

By defining $\tilde{v}=v-\hat{v}, \tilde{w}=w-\hat{w}$, with $\hat{v}$ and $\hat{w}$ being estimates of $v$ and $w$, we can see from (14.39) and (14.40) that

$$
\begin{equation*}
\|(\tilde{v}(t), \tilde{w}(t))\| \leq \gamma_{0}\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\| e^{-\sigma_{0}\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0 \tag{14.41}
\end{equation*}
$$

where $\gamma_{0}$ is a positive constant. We now write (14.25) in conjunction with (14.38) and (14.40) as

$$
\begin{align*}
\dot{x} & =\cos (\phi) \hat{v}+\cos (\phi) \tilde{v} \\
\dot{y} & =\sin (\phi) \hat{v}+\sin (\phi) \tilde{v} \\
\dot{\phi} & =\hat{w}+\tilde{w}  \tag{14.42}\\
\dot{\hat{v}} & =\tau_{v c}+\Omega_{v} \\
\dot{\hat{w}} & =\tau_{w c}+\Omega_{w}
\end{align*}
$$

where $\Omega_{v}$ and $\Omega_{w}$ are the first and second rows of $\boldsymbol{\Omega}$

$$
\begin{align*}
& \boldsymbol{\Omega}=\boldsymbol{B}^{-\mathbf{1}} \boldsymbol{N}_{\boldsymbol{c}} \boldsymbol{B}\left[\begin{array}{l}
\hat{v} \\
\hat{w}
\end{array}\right] \tilde{w}+\boldsymbol{B}^{-\mathbf{1}} \boldsymbol{Q}^{-\mathbf{1}}(\eta) \boldsymbol{K}_{\mathbf{0 2}} \tilde{\eta} \boldsymbol{\eta} \\
& \boldsymbol{N}_{\boldsymbol{c}}=c\left[\begin{array}{l}
n_{12}-n_{11} \\
n_{11}-n_{12}
\end{array}\right] \tag{14.43}
\end{align*}
$$

and we have chosen the control torque

$$
\boldsymbol{\tau}=\boldsymbol{M} \boldsymbol{B}\left(\boldsymbol{B}^{-\mathbf{1}} \boldsymbol{N}_{\boldsymbol{c}} \boldsymbol{B}\left[\begin{array}{l}
\hat{v}  \tag{14.44}\\
\hat{w}
\end{array}\right] \hat{\boldsymbol{w}}-\boldsymbol{B}^{-\mathbf{1}} \boldsymbol{M}^{-\mathbf{1}} \boldsymbol{D} \boldsymbol{B}\left[\begin{array}{l}
\hat{v} \\
\hat{w}
\end{array}\right]+\left[\begin{array}{l}
\tau_{v c} \\
\tau_{w c}
\end{array}\right]\right),
$$

with $\tau_{v c}$ and $\tau_{w c}$ being the new control inputs to be designed in the next section.

### 14.1.3.3 Control Design

We first interpret the tracking errors as

$$
\left[\begin{array}{l}
x_{e}  \tag{14.45}\\
y_{e} \\
\phi_{e}
\end{array}\right]=\left[\begin{array}{lll}
\cos (\phi) & \sin (\phi) & 0 \\
-\sin (\phi) & \cos (\phi) & 0 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{l}
x-x_{d} \\
y-y_{d} \\
\phi-\phi_{d}
\end{array}\right] .
$$

Using (14.45), (14.26), and the kinematic part of (14.42) results in

$$
\begin{align*}
& \dot{x}_{e}=\hat{v}-u_{1 d} \cos \left(\phi_{e}\right)+y_{e}(\hat{w}+\tilde{w})+\tilde{v}, \\
& \dot{y}_{e}=u_{1 d} \sin \left(\phi_{e}\right)-x_{e}(\hat{w}+\tilde{w}),  \tag{14.46}\\
& \dot{\phi}_{e}=\hat{w}-u_{2 d}+\tilde{w} .
\end{align*}
$$

Since (14.46) and the last two equations of (14.42) are of a lower triangular structure, we use the backstepping technique [3] to design $\tau_{v c}$ and $\tau_{w c}$ in two steps.

## Step 1

In this step, we consider $\hat{v}$ and $\hat{w}$ as the controls. From (14.46), it can be seen that $\hat{v}$ and $\hat{w}$ can be directly used to stabilize $x_{e}$ and $\phi_{e}$-dynamics. To stabilize $y_{e}$-dynamics, $\phi_{e}$ can be used when $u_{1 d}$ is persistently exciting. When $u_{1 d}$ is not persistently exciting (stabilization/regulation case), we need some persistently exciting signal in $\hat{w}$ to stabilize $y_{e}$-dynamics via $x_{e}$. With these observations in mind, we define

$$
\begin{align*}
& \bar{v}=\hat{v}-\alpha_{v} \\
& \bar{w}=\hat{w}-\alpha_{w}  \tag{14.47}\\
& \bar{\phi}_{e}=\phi_{e}-\alpha_{\phi e}
\end{align*}
$$

where $\alpha_{v}, \alpha_{w}$, and $\alpha_{\phi e}$ are the virtual controls of $\hat{v}, \hat{w}$, and $\phi_{e}$, respectively. From the above discussion, we first choose the virtual controls $\alpha_{v}$ and $\alpha_{\phi e}$ as

$$
\begin{align*}
& \alpha_{v}=-c_{1} \Omega_{1}^{-1} x_{e}+u_{1 d} \cos \left(\phi_{e}\right) \\
& \alpha_{\phi e}=-\arcsin \left(\frac{k(t)}{\Omega_{1}} y_{e}\right),  \tag{14.48}\\
& k(t)=\lambda_{1} u_{1 d}+\lambda_{2} \cos \left(\lambda_{3} t\right)
\end{align*}
$$

where $\Omega_{1}=\sqrt{1+x_{e}^{2}+y_{e}^{2}}, c_{1}$ is a positive constant, and $\lambda_{i}, i=1,2,3$ are positive constants such that $|k(t)| \leq k_{*}<1, \forall t$. They will be specified later. For simplicity, the virtual control $\alpha_{v}$ does not cancel a known term $y_{e} \hat{w}$ in the $x_{e}$-dynamics. It is of interest to note that the choice of (14.48) will result in the global result and bounded virtual velocity controls.

To design $\alpha_{w}$, differentiating $\bar{\phi}_{e}=\phi_{e}-\alpha_{\phi e}$ along the solutions of (14.46) together with (14.48) yields

$$
\begin{align*}
\dot{\bar{\phi}}_{e}= & \left(1-\frac{k}{\Omega_{2}} x_{e}\right)\left(\alpha_{w}+\bar{w}+\tilde{w}\right)-u_{2 d}-\frac{k}{\Omega_{2} \Omega_{1}^{2}} x_{e} y_{e}(\bar{v}+\tilde{v})+ \\
& \frac{1}{\Omega_{2}}\left(\dot{k} y_{e}+\frac{k c_{1}}{\Omega_{1}^{3}} x_{e}^{2} y_{e}+\frac{k u_{1 d}}{\Omega_{1}^{2}}\left(1+x_{e}^{2}\right) \sin \left(\phi_{e}\right)\right) \tag{14.49}
\end{align*}
$$

which suggests that we choose

$$
\begin{align*}
\alpha_{w}= & \frac{1}{1-k \Omega_{2}^{-1} x_{e}}\left(-\frac{c_{2} \bar{\phi}_{e}}{\sqrt{1+\phi_{e}^{2}}}+u_{2 d}-\frac{1}{\Omega_{2}}\left(\dot{k} y_{e}+\frac{k c_{1}}{\Omega_{1}^{3}} x_{e}^{2} y_{e}+\right.\right. \\
& \left.\frac{k u_{1 d}}{\Omega_{1}^{2}}\left(1+x_{e}^{2}\right) \sin \left(\phi_{e}\right)\right) \tag{14.50}
\end{align*}
$$

where $\Omega_{2}=\sqrt{1+x_{e}^{2}+\left(1-k^{2}\right) y_{e}^{2}}$ and $c_{2}$ is a positive constant.
Remark 14.3. From (14.48) and (14.50), the virtual controls $\alpha_{v}$ and $\alpha_{w}$, as a simple calculation shows, are bounded by some constants depending on the upper bound of $u_{1 d}, \dot{u}_{1 d}$, and $u_{2 d}$.

Substituting (14.48) and (14.50) into (14.46) and (14.49) results in

$$
\begin{align*}
& \dot{x}_{e}=-\frac{c_{1}}{\Omega_{1}} x_{e}+y_{e}(\hat{w}+\tilde{w})+\bar{v}+\tilde{v} \\
& \dot{y}_{e}=-\frac{k u_{1 d}}{\Omega_{1}} y_{e}-x_{e}(\hat{w}+\tilde{w})+\frac{u_{1 d}}{\Omega_{1}}\left(\sin \left(\bar{\phi}_{e}\right) \Omega_{2}-\left(\cos \left(\bar{\phi}_{e}\right)-1\right) k y_{e}\right),  \tag{14.51}\\
& \dot{\bar{\phi}}_{e}=-\frac{c_{2} \bar{\phi}_{e}}{\sqrt{1+\phi_{e}^{2}}}+\left(1-\frac{k}{\Omega_{2}} x_{e}\right)(\bar{w}+\tilde{w})-\frac{k}{\Omega_{2} \Omega_{1}^{2}} x_{e} y_{e}(\bar{v}+\tilde{v})
\end{align*}
$$

## Step 2

At this step, the control inputs $\tau_{v c}$ and $\tau_{w c}$ are designed. We note that $\alpha_{v}$ is a smooth function of $x_{e}, y_{e}, \phi_{e}$, and $u_{1 d}$, and that $\alpha_{w}$ is a smooth function of $x_{e}, y_{e}, \phi_{e}, u_{1 d}$, $\dot{u}_{1 d}, u_{2 d}$, and $t$. By differentiating $\bar{v}=\hat{v}-\alpha_{v}$ and $\bar{w}=\hat{w}-\alpha_{w}$ along the solutions of (14.46) and the last two equations of (14.42), and noting the last equation of (14.51), we choose $\tau_{v c}$ and $\tau_{w c}$ as

$$
\begin{align*}
\tau_{v c}= & -c_{3} \bar{v}+\frac{\partial \alpha_{v}}{\partial x_{e}}\left(\hat{v}-u_{1 d} \cos \left(\phi_{e}\right)+y_{e} \hat{w}\right)+\frac{\partial \alpha_{v}}{\partial \phi_{e}}\left(\hat{w}-u_{2 d}\right)+ \\
& \frac{\partial \alpha_{v}}{\partial y_{e}}\left(u_{1 d} \sin \left(\phi_{e}\right)-x_{e} \hat{w}\right)+\frac{\partial \alpha_{v}}{\partial u_{1 d}} \dot{u}_{1 d}-\delta_{v}\left(\hat{v}^{2}+\hat{w}^{2}\right) \bar{v}+\frac{k}{\Omega_{2} \Omega_{1}^{2}} x_{e} y_{e} \bar{\phi}_{e}, \\
\tau_{w c}= & -c_{4} \bar{w}+\frac{\partial \alpha_{w}}{\partial x_{e}}\left(\hat{v}-u_{1 d} \cos \left(\phi_{e}\right)+y_{e} \hat{w}\right)+\frac{\partial \alpha_{w}}{\partial \phi_{e}}\left(\hat{w}-u_{2 d}\right)+ \\
& \frac{\partial \alpha_{w}}{\partial y_{e}}\left(u_{1 d} \sin \left(\phi_{e}\right)-x_{e} \hat{w}\right)+\frac{\partial \alpha_{w}}{\partial u_{1 d}} \dot{u}_{1 d}+\frac{\partial \alpha_{w}}{\partial \dot{u}_{1 d}} \ddot{u}_{1 d}+\frac{\partial \alpha_{w}}{\partial u_{2 d}} \dot{u}_{2 d}- \\
& \left(1-\frac{k}{\Omega_{2}} x_{e}\right) \bar{\phi}_{e}-\delta_{w}\left(\hat{v}^{2}+\hat{w}^{2}\right) \bar{w}, \tag{14.52}
\end{align*}
$$

where $c_{3}, c_{4}, \delta_{v}$, and $\delta_{w}$ are positive constants. The terms multiplied by $\delta_{v}$ and $\delta_{w}$ are the nonlinear damping terms to overcome the effect of observer errors, see (14.43). The choice of (14.52) results in

$$
\begin{align*}
\dot{\bar{v}}= & -c_{3} \bar{v}-\frac{\partial \alpha_{v}}{\partial x_{e}}\left(y_{e} \tilde{w}+\tilde{v}\right)-\frac{\partial \alpha_{v}}{\partial \phi_{e}} \tilde{w}+\frac{\partial \alpha_{v}}{\partial y_{e}} x_{e} \tilde{w}+ \\
& \Omega_{v}-\delta_{v}\left(\hat{v}^{2}+\hat{w}^{2}\right) \bar{v}+\frac{k}{\Omega_{2} \Omega_{1}^{2}} x_{e} y_{e} \bar{\phi} \\
\dot{\bar{w}}= & -c_{4} \bar{w}-\frac{\partial \alpha_{w}}{\partial x_{e}}\left(y_{e} \tilde{w}+\tilde{v}\right)-\frac{\partial \alpha_{w}}{\partial \phi_{e}} \tilde{w}+\frac{\partial \alpha_{w}}{\partial y_{e}} x_{e} \tilde{w}+  \tag{14.53}\\
& \Omega_{w}-\left(1-\frac{k}{\Omega_{2}} x_{e}\right) \bar{\phi}_{e}-\delta_{w}\left(\hat{v}^{2}+\hat{w}^{2}\right) \bar{w} .
\end{align*}
$$

### 14.1.3.4 Stability Analysis

To analyze the closed loop system consisting of (14.51) and (14.53), we first consider the $\left(\bar{\phi}_{e}, \bar{v}, \bar{w}\right)$-subsystem, then move to the $\left(x_{e}, y_{e}\right)$-subsystem.
$\left(\bar{\phi}_{e}, \bar{v}, \bar{w}\right)$-subsystem. For this subsystem, consider the Lyapunov function

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(\bar{\phi}_{e}^{2}+\bar{v}^{2}+\bar{w}^{2}\right), \tag{14.54}
\end{equation*}
$$

whose derivative along the solutions of the last equation of (14.51) and (14.53) satisfies

$$
\begin{align*}
\dot{V}_{1} & \leq-\frac{c_{2}}{\sqrt{1+\bar{\phi}_{e}^{2}}} \bar{\phi}_{e}^{2}-c_{3} \bar{v}^{2}-c_{4} \bar{w}^{2}+\left(\chi_{11}+\chi_{12} V_{1}\right) e^{-\sigma_{0}\left(t-t_{0}\right)} \\
& \leq\left(\chi_{11} V_{1}+\chi_{12}\right) e^{-\sigma_{0}\left(t-t_{0}\right)} \tag{14.55}
\end{align*}
$$

where $\chi_{11}$ and $\chi_{12}$ are class- $K$ functions of $\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{X}\left(t_{0}\right)\right)\right\|$. The second line of (14.55) implies that $V_{1}(t) \leq \chi_{13}$, with $\chi_{13}$ being a class- $K$ function of $\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right)\right.$,
$\left.\tilde{\boldsymbol{X}}\left(t_{0}\right), \bar{X}\left(t_{0}\right)\right) \|$ with $\overline{\boldsymbol{X}}(t)=\left[\bar{\phi}_{e}(t) \bar{v}(t) \bar{w}(t)\right]^{T}$. Substituting this bound into the first line of (14.55) yields

$$
\begin{equation*}
\dot{V}_{1} \leq-2 \min \left(\frac{c_{2}}{\sqrt{1+2 \chi_{13}}}, c_{3}, c_{4}\right) V_{1}+\left(\chi_{11}+\chi_{12} \chi_{13}\right) e^{-\sigma_{0}\left(t-t_{0}\right)} \tag{14.56}
\end{equation*}
$$

which implies that there exist $\sigma_{1}>0$ and a class- $K$ function $\chi_{1}$ depending on $\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right), \overline{\boldsymbol{X}}\left(t_{0}\right)\right)\right\|$ such that $\|\overline{\boldsymbol{X}}(t)\| \leq \chi_{1} e^{-\sigma_{1}\left(t-t_{0}\right)}$, that is, the $\left(\bar{\phi}_{e}, \bar{v}, \bar{w}\right)$ subsystem is GAS.
$\left(x_{e}, y_{e}\right)$-subsystem. We first prove that the trajectories $\left(x_{e}, y_{e}\right)$ are bounded by taking the Lyapunov function

$$
\begin{equation*}
V_{2}=\sqrt{1+x_{e}^{2}+y_{e}^{2}}-1 \tag{14.57}
\end{equation*}
$$

whose derivative along the solutions of the first two equations of (14.51) satisfies

$$
\begin{align*}
\dot{V}_{2} & \leq-\frac{c_{1}}{\Omega_{1}^{2}} x_{e}^{2}-\frac{k u_{1 d}}{\Omega_{1} 2} y_{e}^{2}+\chi_{21} e^{-\sigma_{21}\left(t-t_{0}\right)} \\
& \leq \frac{\lambda_{2} u_{1 d} \cos \left(\lambda_{3} t\right)}{\Omega_{1}^{2}} y_{e}^{2}+\chi_{21} e^{-\sigma_{21}\left(t-t_{0}\right)} \tag{14.58}
\end{align*}
$$

where $\sigma_{21}=\min \left(\sigma_{0}, \sigma_{1}\right)$ and $\chi_{21}$ is a class- $K$ function of $\left\|\left(\tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right), \overline{\boldsymbol{X}}\left(t_{0}\right)\right)\right\|$. Integrating both sides of the second line of (14.58) yields

$$
\begin{align*}
V_{2}(t) & \leq V_{2}\left(t_{0}\right)+2 \lambda_{2} u_{1 d}^{\max }+\frac{\chi_{21}}{\sigma_{21}} \\
& \leq \chi_{22} \tag{14.59}
\end{align*}
$$

where $u_{1 d}^{\max }$ is the upper bound of $\left|u_{1 d}(t)\right|$. Therefore, the trajectories $\left(x_{e}, y_{e}\right)$ are bounded on $[0, \infty)$. To prove convergence of $\left(x_{e}, y_{e}\right)$ to zero, we consider each case of Assumption 14.1.

Cases C1 and C2. From the first line of (14.58) and noting (14.48), we have

$$
\begin{equation*}
\dot{V}_{2} \leq-\frac{c_{1}}{\Omega_{1}^{2}} x_{e}^{2}+\left|\lambda_{2} u_{1 d}\right|+\chi_{21}(\cdot) e^{-\sigma_{21}\left(t-t_{0}\right)} \tag{14.60}
\end{equation*}
$$

Integrating both sides of (14.60) and using Barbalat's lemma, it is straightforward to show that $\lim _{t \rightarrow \infty} x_{e}(t)=0$. To prove that $\lim _{t \rightarrow \infty} y_{e}(t)=0$, applying Lemma 2.6 to the first equation of (14.51) yields:

$$
\begin{equation*}
\lim _{t \rightarrow \infty}\left(y_{e}\left(\alpha_{w}+\bar{w}+\tilde{w}\right)+\tilde{v}+\bar{v}\right)=0 \tag{14.61}
\end{equation*}
$$

which is equivalent to:

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \Xi(t)=0 \tag{14.62}
\end{equation*}
$$

where

$$
\begin{equation*}
\Xi(t)=y_{e}(t)\left(\frac{\dot{k}(t) y_{e}(t)}{\sqrt{1+\left(1-k^{2}(t)\right) y_{e}^{2}(t)}}-u_{2 d}(t)\right) \tag{14.63}
\end{equation*}
$$

On the other hand, from (14.60), we have

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t}\left(V_{2}-\int_{0}^{t}\left|\lambda_{2} u_{1 d}(\tau)\right| \mathrm{d} \tau+\frac{1}{\sigma_{21}} \chi_{21}(\cdot) e^{-\sigma_{21}\left(t-t_{0}\right)}\right) \leq 0 \tag{14.64}
\end{equation*}
$$

which means that $V_{2}-\int_{0}^{t}\left|\lambda_{2} u_{1 d}(\tau)\right| \mathrm{d} \tau+\sigma_{21}^{-1} \chi_{21}(\cdot) e^{-\sigma_{21}\left(t-t_{0}\right)}$ is nonincreasing. Since $V_{2}$ is bounded from below by zero, $V_{2}$ tends to a finite nonnegative constant depending on $\left\|\overline{\boldsymbol{X}}_{e}\left(t_{0}\right)\right\|$ with $\overline{\boldsymbol{X}}_{e}=\left(x_{e}, y_{e}, \overline{\boldsymbol{X}}, \tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)$. This implies that the limit of $\left|y_{e}(t)\right|$ exists and is finite, say $l_{y_{e}}$. If $l_{y_{e}}$ were not zero, there would exist a sequence of increasing time instants $\left\{t_{i}\right\}_{i=1}^{\infty}$ with $t_{i} \rightarrow \infty$, such that both of the limits of $\dot{k}\left(t_{i}\right)$ and $\Xi\left(t_{i}\right)$ would not be zero. With this in mind, if we choose $\lambda_{i} \neq 0$ and such that

$$
\begin{equation*}
\frac{\lambda_{2} \lambda_{3}}{\sqrt{1-k_{*}^{2}}}<\mu_{22} \tag{14.65}
\end{equation*}
$$

then under conditions (14.27), $\dot{k}\left(t_{i}\right)$ and $\Xi\left(t_{i}\right)$ cannot be nonzero simultaneously for any $t_{i}$. Hence $l_{y_{e}}$ must be zero, which allows us to conclude from (14.62) that $\lim _{t \rightarrow \infty} y_{e}(t)=0$, i.e., the $\left(x_{e}, y_{e}\right)$-subsystem is asymptotically stable.

Case C3. In this case, from (14.58), we have

$$
\begin{equation*}
\dot{V}_{2} \leq-\frac{1}{\Omega_{1}^{2}}\left(c_{1} x_{e}^{2}+\left(\lambda_{1} u_{1 d}^{2}-\lambda_{2}\left|u_{1 d}\right|\right) y_{e}^{2}\right)+\chi_{21} e^{-\sigma_{21}\left(t-t_{0}\right)}, \tag{14.66}
\end{equation*}
$$

which means that there exist $\sigma_{2}>0$ and a class- $K$ function $\chi_{2}$ depending on $\left\|\bar{X}\left(t_{0}\right)\right\|$ such that

$$
\begin{equation*}
\left\|\left(x_{e}(t), y_{e}(t)\right)\right\| \leq \chi_{2} \sigma^{-\sigma_{2}\left(t-t_{0}\right)} \tag{14.67}
\end{equation*}
$$

as long as

$$
\begin{equation*}
\lambda_{1} \mu_{31}-\lambda_{2} u_{1 d}^{\max } \geq \mu_{31}^{*} \tag{14.68}
\end{equation*}
$$

where $\mu_{31}^{*}$ is a positive constant. In addition, it can be shown that in this case the closed loop of (14.51) and (14.53) is also locally exponentially stable. Under Assumption 14.1 there always exist $\lambda_{i}$ such that (14.65) and (14.68) hold. We have thus proven the following result.

Theorem 14.1. Under Assumption 14.1 the output feedback control laws consisting of (14.44) and (14.52) force the mobile robot (14.25) to globally asymptotically track the virtual vehicle (14.26) if the constants $\lambda_{i}, i=1,2,3$ are chosen such that $\lambda_{i} \neq 0$, (14.65) and (14.68) hold.

### 14.1.3.5 Simulations

The physical parameters of the robot are given in Section 14.1.2. We perform two simulations. For the first simulation, the reference velocities are chosen as: $u_{1 d}=0.5\left(\tanh \left(t_{s}-t\right)+1\right)$, and $u_{2 d}=0$, where $t_{s}$ is a positive constant. A switching combination of a tracking controller and a stabilization one in the literature cannot be used to fulfill this task if $t_{s}$ is unknown in advance. A calculation shows that for $t \leq t_{s}$ (tracking a curve), condition C3 holds with $\mu_{31}=0.25$ and for $t>t_{s}$ (parking) Case C 1 holds. Hence, our proposed controller can be applied. We also assume that due to some sudden impact at the time $t_{m}>t_{s}$, the robot position is perturbed to $y=y_{m} \neq 0$ to illustrate the regulation ability of our proposed controller. For the second simulation, the reference velocities are $u_{1 d}=0, u_{2 d}=0.2$, i.e., Case C2 holds with $\mu_{22}=0.2$. The initial conditions are: $\left(\boldsymbol{\eta}^{\boldsymbol{T}}, \boldsymbol{\omega}^{\boldsymbol{T}}\right)=((-2,2,-0.5),(0,0))$, $\left(\hat{\boldsymbol{\eta}}^{\boldsymbol{T}}, \hat{\boldsymbol{X}}^{\boldsymbol{T}}\right)=((0,0,0),(0,0)),\left(x_{d}, y_{d}, \phi_{d}\right)=(0,0,0)$, and we take $t_{s}=20, t_{m}=30$, $y_{m}=1.5$. The control and observer gains are chosen as $c_{i}=2,1 \leq i \leq 4, \delta_{v}=$ $\delta_{w}=0.1, \boldsymbol{P}_{\mathbf{0 1}}=\boldsymbol{P}_{\mathbf{0 2}}=\operatorname{diag}(1,1), \lambda_{1}=\lambda_{3}=0.5, \lambda_{2}=0.1, \boldsymbol{K}_{\mathbf{0 1}}=\operatorname{diag}(1,1)$, $\boldsymbol{K}_{\mathbf{0 2}}=\left(\boldsymbol{J}(\eta) \boldsymbol{Q}^{-\mathbf{1}}(\eta)\right)^{T}$. The above choice satisfies the requirements in Theorem 14.1. The robot trajectories in the $(x, y)$-plane are plotted in Figures 14.3 and 14.4. The tracking errors in the form of $\sqrt{x_{e}^{2}+y_{e}^{2}+\phi_{e}^{2}}$ are plotted in Figure 14.5. This figure indicates that convergence of the tracking errors for the case of regulation to zero is much slower than those for the other cases, which is a quite well-known effect when using the smooth time-varying controllers. Convergence of tracking errors in Case C2 is slower than that in Case C3, since Case C3 yields local exponential stability but only asymptotic stability for Case C2 (see the proof of Theorem 14.1).


Figure 14.3 First simulation: Robot position in the $(x, y)$-plane


Figure 14.4 Second simulation: Robot position in the $(x, y)$-plane


Figure 14.5 Tracking errors with respect to the first and second simulations

### 14.1.4 Output Feedback Path-following

### 14.1.4.1 Control Objective

For the reader's convenience, we rewrite the equations of motion here, see (14.1):

$$
\begin{align*}
& \dot{\eta}=J(\eta) \omega,  \tag{14.69}\\
& M \dot{\omega}+C(\dot{\eta}) \omega+D \omega=\tau .
\end{align*}
$$

After an observer and a primary control design as in Section 14.1.3.2, we only need to consider the robot model (i.e., the dynamics (14.42)):

$$
\begin{align*}
& \dot{x}=\cos (\phi) \hat{v}+\cos (\phi) \tilde{v}, \\
& \dot{y}=\sin (\phi) \hat{v}+\sin (\phi) \tilde{v}, \\
& \dot{\phi}=\hat{w}+\tilde{w},  \tag{14.70}\\
& \dot{\hat{v}}=\tau_{v c}+\Omega_{v}, \\
& \dot{\hat{w}}=\tau_{w c}+\Omega_{w} .
\end{align*}
$$

In this section, we consider a control objective of designing the control vector $\boldsymbol{\tau}$ to force the mobile robot to follow a specified path $\Gamma$, see Figure 14.6. If we are able to drive the robot to closely follow a virtual robot that moves along the path with a desired speed $v_{0}$, which is tangential to the path, then the control objective is fulfilled, i.e., the robot is in a tube of nonzero diameter centered on the reference path and moves along the specified path at the speed $v_{0}$. Roughly speaking, the approach is to steer the robot such that it heads toward the virtual robot and eliminates the distance between itself and the virtual robot. We define the following variables to mathematically formulate the control objective:

$$
\begin{align*}
& x_{e}=x_{d}-x, \\
& y_{e}=y_{d}-x,  \tag{14.71}\\
& \phi_{e}=\phi-\phi_{d}, \\
& z_{e}=\sqrt{x_{e}^{2}+y_{e}^{2}}
\end{align*}
$$

where

$$
\begin{equation*}
\phi_{d}=\arcsin \left(\frac{y_{e}}{z_{e}}\right) . \tag{14.72}
\end{equation*}
$$

Control Objective. Under Assumption 14.2, design the controls $\tau_{1}$ and $\tau_{2}$ to force the mobile robot (14.70) to follow the path $\Gamma$ given by

$$
\begin{align*}
x_{d} & =x_{d}(s), \\
y_{d} & =y_{d}(s) \tag{14.73}
\end{align*}
$$



Figure 14.6 General framework of mobile robot path-following
where $s$ is the path parameter variable, such that

$$
\begin{align*}
& \lim _{t \rightarrow \infty} z_{e}(t) \leq \bar{z}_{e} \\
& \lim _{t \rightarrow \infty}\left|\phi_{e}(t)\right|=0 \tag{14.74}
\end{align*}
$$

with $\bar{z}_{e}$ being an arbitrarily small positive constant.

## Assumption 14.2.

1. The reference path is regular, i.e., there exist positive constants $R_{\min }$ and $R_{\max }$ such that

$$
0<R_{\min } \leq\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2} \leq R_{\max }<\infty
$$

2. The minimum radius of the osculating circle of the path is larger than or equal to the minimum possible turning radius of the robot.

## Remark 14.4.

1. Assumption 14.2 ensures that the path is feasible for the robot to follow.
2. If the reference path is not regular, then we can often split it into regular sections and consider each of them separately.
3. The path parameter, $s$, is not the arc length of the path in general. For example, a circle with radii of $R$ centered at the origin can be described as $x_{d}=R \cos (s)$ and $y_{d}=R \sin (s)$, see [136] for more details.

If one differentiates both sides of $\phi_{e}=\phi-\phi_{d}$ to obtain the $\dot{\phi}_{e}$-dynamics, there will be discontinuity in the $\dot{\phi}_{e}$-dynamics when $x_{e}$ changes its sign. This discontinu-
ity will cause difficulties in applying the backstepping technique. To get around this problem, we compute the $\dot{\phi}_{e}$-dynamics based on

$$
\begin{align*}
& \sin \left(\phi_{e}\right)=\frac{x_{e} \sin (\phi)-y_{e} \cos (\phi)}{z_{e}} \\
& \cos \left(\phi_{e}\right)=\frac{x_{e} \cos (\phi)+y_{e} \sin (\phi)}{z_{e}} \tag{14.75}
\end{align*}
$$

We now use (14.71) and (14.75) to transform (14.71) to

$$
\begin{align*}
\dot{z}_{e}= & -\cos \left(\phi_{e}\right) \hat{v}+\left(\frac{x_{e}}{z_{e}} \frac{\partial x_{d}}{\partial s}+\frac{y_{e}}{z_{e}} \frac{\partial y_{d}}{\partial s}\right) \dot{s}-\cos \left(\phi_{e}\right) \tilde{v} \\
\dot{\phi}_{e}= & \hat{w}+\left[\left(\frac{\sin (\phi)}{z_{e}}-\frac{x_{e} \sin \left(\phi_{e}\right)}{z_{e}^{2}}\right) \frac{\partial x_{d}}{\partial s}-\left(\frac{\cos (\phi)}{z_{e}}+\frac{y_{e} \sin \left(\phi_{e}\right)}{z_{e}^{2}}\right) \frac{\partial y_{d}}{\partial s}\right] \times \\
& \frac{\dot{s}}{\cos \left(\phi_{e}\right)}+\frac{\sin \left(\phi_{e}\right)}{z_{e}}(\hat{v}+\tilde{v})+\tilde{w} \\
\dot{\hat{v}}= & \tau_{v c}+\Omega_{v} \\
\dot{\hat{w}}= & \tau_{w c}+\Omega_{w} . \tag{14.76}
\end{align*}
$$

It is noted that (14.75) is not defined at $z_{e}=0$. However, our controller will guarantee that $z_{e}(t) \geq z_{e}^{*}>0$, for all $0 \leq t<\infty$ for feasible initial conditions. The second equation of (14.76) is not defined at $\phi_{e}(t)= \pm 0.5 \pi$ but we will design $\dot{s}$ to overcome this problem. Therefore, we will design the controls $\tau_{v c}$ and $\tau_{w c}$ for (14.76) to yield the control objective. In the next section, a procedure to design a stabilizer for the path-following error system (14.76) is presented in detail. The triangular structure of (14.76) suggests that we design the controls $\tau_{v c}$ and $\tau_{w c}$ in two stages. First, we design the virtual velocity controls for $\hat{v}$ and $\hat{w}$ and choose $\dot{s}$ to ultimately stabilize $z_{e}$ and $\phi_{e}$ at the origin. Based on the backstepping technique, the controls $\tau_{v c}$ and $\tau_{w c}$ will be then designed.

### 14.1.4.2 Control Design

## Step 1

The $z_{e}$ - and $\phi_{e}$-dynamics have three inputs that can be chosen to stabilize $z_{e}$ and $\phi_{e}$, namely $\dot{s}, \hat{v}$, and $\hat{w}$. The input $\hat{w}$ should be designed to stabilize the $\phi_{e}$-dynamics at the origin. Therefore, two inputs, $\dot{s}$ and $\hat{v}$, can be used to ultimately stabilize $z_{e}$ at the origin. We can either choose the input $\hat{v}$ and $\dot{s}$ and then design the remaining input. If we fix $\dot{s}$, then the virtual robot is allowed to move at a desired speed. The real robot will follow the virtual one on the path by the controller, and vice versa. We here choose to fix $\dot{s}$. This allows us to adjust the initial conditions in most cases without moving the robot. Since the transformed system (14.76) is not defined at
$z_{e}=0$, we first assume in the following that $z_{e}(t) \geq z_{e}^{*}>0, \forall 0 \leq t<\infty$. We will then show that there exist initial conditions such that this hypothesis holds.

Define

$$
\begin{align*}
& v_{e}=\hat{v}-\hat{v}_{r},  \tag{14.77}\\
& w_{e}=\hat{w}-\hat{w}_{r},
\end{align*}
$$

where $\hat{v}_{r}$ and $\hat{w}_{r}$ are the virtual controls of $\hat{v}$ and $\hat{w}$, respectively. As discussed above, we choose the virtual controls and $\dot{s}$ as follows:

$$
\begin{align*}
\hat{v}_{r}= & k_{1}\left(z_{e}-\delta_{e}\right)+\left(\frac{x_{e}}{z_{e}} \frac{\partial x_{d}}{\partial s}+\frac{y_{e}}{z_{e}} \frac{\partial y_{d}}{\partial s}\right) \frac{v_{0}\left(t, z_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}}, \\
\hat{w}_{r}= & -k_{2} \phi_{e}-\left[\left(\frac{\sin (\phi)}{z_{e}}-\frac{x_{e} \sin \left(\phi_{e}\right)}{z_{e}^{2}}\right) \frac{\partial x_{d}}{\partial s}-\left(\frac{\cos (\phi)}{z_{e}}+\frac{y_{e} \sin \left(\phi_{e}\right)}{z_{e}^{2}}\right) \frac{\partial y_{d}}{\partial s}\right] \\
& \times \frac{v_{0}\left(t, z_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}}-\frac{\sin \left(\phi_{e}\right)}{z_{e}} \hat{v}_{r}-\delta_{1}\left(\frac{\sin \left(\phi_{e}\right)}{z_{e}}\right)^{2} \phi_{e}, \\
\dot{s}= & \frac{\cos \left(\phi_{e}\right) v_{0}\left(t, z_{e}\right)}{\sqrt{\left(\frac{\partial x_{d}}{\partial s}\right)^{2}+\left(\frac{\partial y_{d}}{\partial s}\right)^{2}}}, \tag{14.78}
\end{align*}
$$

where $k_{1}>0, k_{2}>0$, and $\delta_{1}>0$. The term multiplied by $\delta_{1}$ is a nonlinear damping term to overcome the observer error effect. $v_{0}\left(t, z_{e}\right) \neq 0$ for all $t \geq t_{0} \geq 0$ and $z_{e} \in \mathbb{R}$, is the speed of the virtual robot on the path. Indeed, one can choose this speed to be a constant. However, the time-varying speed and position path-following dependence of the virtual robot on the path is more desirable, especially when the robot starts to follow the path. For example, one might choose

$$
\begin{equation*}
v_{0}\left(t, z_{e}\right)=v_{0}^{*}\left(1-\chi_{1} e^{-\chi_{2}\left(t-t_{0}\right)}\right) e^{-\chi_{3} z_{e}}, \tag{14.79}
\end{equation*}
$$

where $v_{0}^{*} \neq 0, \chi_{i}>0, i=1,2,3$, and $\chi_{1}<1$. The choice of $v_{0}\left(t, z_{e}\right)$ in (14.79) has the following desired feature: When the path-following error, $z_{e}$, is large, the virtual robot will wait for the real one; when $z_{e}$ is small, the virtual robot will move along the path at the speed closed to $v_{0}^{*}$ and the real one follows it within the specified look ahead distance. This feature is suitable in practice because it avoids using a high gain control for a large signal $z_{e}$.

Substituting (14.78) into the first two equations of (14.76) results in

$$
\begin{align*}
& \dot{z}_{e}=-k_{1} \cos \left(\phi_{e}\right)\left(z_{e}-\delta_{e}\right)-\cos \left(\phi_{e}\right)\left(\tilde{v}+v_{e}\right), \\
& \dot{\phi}_{e}=-k_{2} \phi_{e}-\delta_{1}\left(\frac{\sin \left(\phi_{e}\right)}{z_{e}}\right)^{2} \phi_{e}+\frac{\sin \left(\phi_{e}\right)}{z_{e}}\left(\tilde{v}+v_{e}\right)+\tilde{w}+w_{e} \tag{14.80}
\end{align*}
$$

## Step 2

By noting that the virtual control $\hat{v}_{r}$ is a function of $t, x_{e}, y_{e}$, and $s$, and the virtual control $\hat{w}_{r}$ is a function of $t, x_{e}, y_{e}, s$, and $\phi$, differentiating both sides of (14.77) along the solutions of the last two equations of (14.76) results in

$$
\begin{align*}
& {\left[\begin{array}{c}
\dot{v}_{e} \\
\dot{w}_{e}
\end{array}\right]=\left[\begin{array}{l}
\tau_{v c}+\Omega_{v} \\
\tau_{w c}+\Omega_{w}
\end{array}\right]+\left[\begin{array}{l}
\left(\frac{\partial \hat{v}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{v}_{r}}{\partial y_{e}} \sin (\phi)\right) \tilde{v} \\
\left(\frac{\partial \hat{w}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{w}_{r}}{\partial y_{e}} \sin (\phi)\right) \tilde{v}-\frac{\partial \hat{w}_{r}}{\partial \phi} \tilde{w}-\frac{\partial \hat{w}_{r}}{\partial t}
\end{array}\right]-} \\
& {\left[\begin{array}{l}
\frac{\partial \hat{v}_{r}}{\partial t}+\frac{\partial \hat{v}_{r}}{\partial x_{e}}\left(\frac{\partial x_{d}}{\partial s} \dot{s}-\hat{v} \cos (\phi)\right)+\frac{\partial \hat{v}_{r}}{\partial y_{e}}\left(\frac{\partial y_{d}}{\partial s} \dot{s}-\hat{v} \sin (\phi)\right)+\frac{\partial \hat{v}_{r}}{\partial s} \dot{s} \\
\left.\frac{\partial \hat{w}_{r}}{\partial x_{e}}\left(\frac{\partial x_{d}}{\partial s} \dot{s}-\hat{v} \cos (\phi)\right)+\frac{\partial \hat{w}_{r}}{\partial y_{e}}\left(\frac{\partial y_{d}}{\partial s} \dot{s}-\hat{v} \sin (\phi)\right)+\frac{\partial \hat{w}_{r}}{\partial s} \dot{s}+\frac{\partial \hat{w}_{r}}{\partial \phi} \hat{w}\right] .
\end{array} . .\right.}
\end{align*}
$$

From (14.80) and (14.81), we choose the control inputs with nonlinear damping terms to overcome the observer error effect as follows:

$$
\begin{align*}
& {\left[\begin{array}{c}
\tau_{v c} \\
\tau_{w c}
\end{array}\right]=-\left[\begin{array}{c}
c_{21} v_{e} \\
c_{22} w_{e}
\end{array}\right]-\delta_{2}\left[\begin{array}{c}
\left(\frac{\partial \hat{v}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{v}_{r}}{\partial y_{e}} \sin (\phi)\right)^{2} v_{e} \\
\left(\frac{\partial \hat{w}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{w}_{r}}{\partial y_{e}} \sin (\phi)\right)^{2} w_{e}+\left(\frac{\partial \hat{w}_{r}}{\partial \phi}\right)^{2}
\end{array}\right]+} \\
& {\left[\begin{array}{c}
\frac{\partial \hat{v}_{r}}{\partial t}+\frac{\partial \hat{v}_{r}}{\partial x_{e}}\left(\frac{\partial x_{d}}{\partial s} \dot{s}-\hat{v} \cos (\phi)\right)+\frac{\partial \hat{v}_{r}}{\partial y_{e}}\left(\frac{\partial y_{d}}{\partial s} \dot{s}-\hat{v} \sin (\phi)\right)+\frac{\partial \hat{v}_{r}}{\partial s} \dot{s} \\
\frac{\partial \hat{w}_{r}}{\partial x_{e}}\left(\frac{\partial x_{d}}{\partial s} \dot{s}-\hat{v} \cos (\phi)\right)+\frac{\partial \hat{w}_{r}}{\partial y_{e}}\left(\frac{\partial y_{d}}{\partial s} \dot{s}-\hat{v} \sin (\phi)\right)+\frac{\partial \hat{w}_{r}}{\partial s} \dot{s}+\frac{\partial \hat{w}_{r}}{\partial \phi} \hat{w}
\end{array}\right]-} \\
& \quad \delta_{3}\left[\begin{array}{c}
\left(\hat{v}^{2}+\hat{w}^{2}\right) v_{e} \\
\left(\hat{v}^{2}+\hat{w}^{2}\right) w_{e}
\end{array}\right]-\left[\begin{array}{c}
\frac{\sin \left(\phi_{e}\right)}{z_{e}} \phi_{e} \\
\phi_{e}-\frac{\partial \hat{w}_{r}}{\partial t}
\end{array}\right], \tag{14.82}
\end{align*}
$$

where $c_{21}, c_{22}, \delta_{2}$, and $\delta_{3}$ are positive constants. Substituting (14.82) into (14.81) results in

$$
\begin{align*}
& {\left[\begin{array}{c}
\dot{v}_{e} \\
\dot{w}_{e}
\end{array}\right]=}-\left[\begin{array}{c}
c_{21} v_{e} \\
c_{22} w_{e}
\end{array}\right]+\left[\begin{array}{c}
\Omega_{v} \\
\Omega_{w}
\end{array}\right]-\delta_{3}\left[\begin{array}{c}
\left(\hat{v}^{2}+\hat{w}^{2}\right) v_{e} \\
\left(\hat{v}^{2}+\hat{w}^{2}\right) w_{e}
\end{array}\right]-\left[\begin{array}{c}
\sin \left(\phi_{e}\right) \phi_{e} / z_{e} \\
\phi_{e}
\end{array}\right]+ \\
& {\left[\begin{array}{c}
\left(\frac{\partial \hat{v}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{v}_{r}}{\partial y_{e}} \sin (\phi)\right) \tilde{v} \\
\left(\frac{\partial \hat{w}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{w}_{r}}{\partial y_{e}} \sin (\phi)\right) \tilde{v}-\frac{\partial \hat{w}_{r}}{\partial \phi} \tilde{w}
\end{array}\right]-} \\
& \delta_{2}\left[\begin{array}{c}
\left(\frac{\partial \hat{v}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{v}_{r}}{\partial y_{e}} \sin (\phi)\right)^{2} v_{e} \\
\left.\left(\frac{\partial \hat{w}_{r}}{\partial x_{e}} \cos (\phi)+\frac{\partial \hat{w}_{r}}{\partial y_{e}} \sin (\phi)\right)^{2} w_{e}+\left(\frac{\partial \hat{w}_{r}}{\partial \phi}\right)^{2}\right] .
\end{array} \$ .\right. \tag{14.83}
\end{align*}
$$

### 14.1.4.3 Stability Analysis

To analyze the closed loop system consisting of (14.80) and (14.83), we first consider the $\left(\phi_{e}, v_{e}, w_{e}\right)$-subsystem, then move to the $z_{e}$-dynamics.
( $\phi_{e}, v_{e}, w_{e}$ )-subsystem. For this subsystem, we take the following Lyapunov function:

$$
\begin{equation*}
V_{1}=\frac{1}{2}\left(\phi_{e}^{2}+v_{e}^{2}+w_{e}^{2}\right) \tag{14.84}
\end{equation*}
$$

whose derivative along the solutions of the last equation of (14.80) and (14.83), after some manipulation, satisfies

$$
\begin{equation*}
\dot{V}_{1} \leq-\rho_{1} V_{1}+\chi_{1}(\cdot) e^{-\sigma_{0}\left(t-t_{0}\right)} \tag{14.85}
\end{equation*}
$$

where $\rho_{1}$ is a positive constant and can be made arbitrarily large by increasing the design constants $k_{1}, k_{2}, c_{21}$ and $c_{22}$, and $\chi_{1}(\cdot)$ is a class- $K$ function of $\left\|\left(z_{e}\left(t_{0}\right), \phi_{e}\left(t_{0}\right), \tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\|$. From (14.85), it is not hard to show that

$$
\begin{equation*}
\left\|\left(\phi_{e}(t), v_{e}(t), w_{e}(t)\right)\right\| \leq \alpha_{1}(\cdot) e^{-\sigma_{1}\left(t-t_{0}\right)} \tag{14.86}
\end{equation*}
$$

where $\alpha_{1}(\cdot)$ is a class- $K$ function of $\left\|\left(z_{e}\left(t_{0}\right), \phi_{e}\left(t_{0}\right), \tilde{\eta}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\|$, and $\sigma_{1}$ is a positive constant. Hence (14.86) implies that the $\left(\phi_{e}, v_{e}, w_{e}\right)$-subsystem is $K$ exponentially stable at the origin.
$z_{e}$-subsystem: Lower Bound of $z_{e}$. Defining $\tilde{z}_{e}=z_{e}-\delta_{e}$, the first equation of (14.80) can be written as

$$
\begin{align*}
\dot{\tilde{z}}_{e} & =-k_{1} \cos \left(\phi_{e}\right) \tilde{z}_{e}-\cos \left(\phi_{e}\right)\left(\tilde{v}+v_{e}\right) \\
& \geq-k_{1} \tilde{z}_{e}-\left|\tilde{v}+v_{e}\right| \\
& \geq-k_{1} \tilde{z}_{e}-\alpha_{2}(\cdot) e^{\sigma_{2}\left(t-t_{0}\right)} \tag{14.87}
\end{align*}
$$

where $\alpha_{2}(\cdot)$ is a class- $K$ function of $\left\|\left(z_{e}\left(t_{0}\right), \phi_{e}\left(t_{0}\right), \tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\|$, and $\sigma_{2}$ is a positive constant. From (14.87) and the comparison principle, we have

$$
\begin{equation*}
\tilde{z}_{e}(t) \geq \tilde{z}_{e}\left(t_{0}\right) e^{-k_{1}\left(t-t_{0}\right)}+\frac{\alpha_{2}(\cdot)}{\sigma_{2}-k_{1}}\left(e^{-\sigma_{2}\left(t-t_{0}\right)}-e^{-k_{1}\left(t-t_{0}\right)}\right) \tag{14.88}
\end{equation*}
$$

Therefore the condition $z_{e}(t) \geq z_{e}^{*}$ holds if

$$
\begin{equation*}
\sigma_{2}>k_{1}, z_{e}\left(t_{0}\right) \geq-2 \delta_{e}+\frac{\alpha_{2}(\cdot)}{\sigma_{2}-k_{1}}+z_{e}^{*} \tag{14.89}
\end{equation*}
$$

$z_{e}$-subsystem: Upper Bound of $z_{e}$. To estimate the upper bound of $z_{e}$, we write the first equation of (14.80) as

$$
\begin{equation*}
\dot{z}_{e}=-k_{1} z_{e}-k_{1}\left(\cos \left(\phi_{e}\right)-1\right) z_{e}+k_{1} \cos \left(\phi_{e}\right) \delta_{e}-\cos \left(\phi_{e}\right)\left(\tilde{v}+v_{e}\right) \tag{14.90}
\end{equation*}
$$

By taking the Lyapunov function $V_{2}=0.5 z_{e}^{2}$ and noting that $\phi_{e}, v_{e}$, and $\tilde{v}$ exponentially converge to zero, it is not hard to show that

$$
\begin{equation*}
\left|z_{e}(t)\right| \leq \alpha_{3}(\cdot) e^{-\sigma_{3}\left(t-t_{0}\right)}+\rho_{3} \tag{14.91}
\end{equation*}
$$

where $\alpha_{3}(\cdot)$ is a class- $K$ function of $\left\|\left(z_{e}\left(t_{0}\right), \phi_{e}\left(t_{0}\right), \tilde{\boldsymbol{\eta}}\left(t_{0}\right), \tilde{\boldsymbol{X}}\left(t_{0}\right)\right)\right\|$, and $\sigma$ and $\rho_{3}$ are positive constants. The constant $\rho_{3}$ can be made arbitrarily small by reducing $\delta_{e}$.

### 14.1.4.4 Simulations

To illustrate the effectiveness of the proposed output feedback path-following controller, we perform a numerical simulation. The physical parameters are given in Section 14.1.2. The design constants are:

$$
k_{1}=0.5 ; k_{2}=5, c_{21}=c_{22}=2, \delta_{1}=\delta_{2}=\delta_{3}=0.05, \delta_{e}=0.2
$$

The initial conditions are:

$$
\left(\boldsymbol{\eta}^{\boldsymbol{T}}, \boldsymbol{\omega}^{\boldsymbol{T}}\right)=((-5,0,0.5),(0,1)),\left(\hat{\boldsymbol{\eta}}^{\boldsymbol{T}}, \hat{\boldsymbol{X}}^{\boldsymbol{T}}\right)=((0,0,0),(0,0)), s(0)=0
$$

The reference speed of the virtual mobile robot is $v_{0}=2.5 \mathrm{~m} / \mathrm{s}$. The path $\Gamma$ is chosen to be a sinusoidal path specified by $x_{d}=s$ and $y_{d}=10 \sin (0.15 s)$. This path is regular and satisfies all requirements in Assumption 14.2. Simulation results are plotted in Figure 14.7. From this figure, it can be clearly seen that the proposed output feedback controller is able to force the mobile robot in question to follow the reference path accurately.

### 14.1.5 Notes and References

This section has shown that the observer design and control design techniques developed for underactuated ships in Chapters 5, 6, 7, and 11 can be successfully applied to solve the challenging problems of output feedback simultaneous stabilization and trajectory-tracking, and path-following for mobile robots. Indeed, the main difficulty in solving stabilization and trajectory-tracking control of mobile robots is due to the fact that the motion of the systems to be controlled has more degrees of freedom to be controlled than the number of control inputs under nonholonomic constraints. In addition, the cross-terms of the robot velocities due to the Coriolis matrix prevents the success of the traditional observer designs. Stabilization and trajectory-tracking control problems of mobile robots are usually addressed separately in the literature. Discontinuous and time-varying approaches are often used


Figure 14.7 a. Robot position and orientation in the $(x, y)$-plane, $\mathbf{b}$. Control torques
to solve the stabilization problem where the robot dynamics is transformed to a chained system. Various nonlinear controllers have been developed based on Lyapunov's direct method and the backstepping technique to solve the tracking control problem, see [18, 41-44, 47, 131, 142, 147-150]. The simultaneous stabilization and trajectory-tracking control problem of mobile robots has been solved using the time-varying approach, see [133, 151, 152] and the high-gain (using an oscillator) approach, see $[148,150]$. It is also mentioned that all proofs of the main results in this section are based on Lyapunov's direct method for the reader's convenience. Indeed, we can use the stability result for cascade systems given in Section 2.1.3 to analyze stability of the closed loop systems. Moreover, if nonlinear damping terms are included in the robot dynamics, Lemma 2.2 can be used with the proposed coordinate transformation to design an exponential/asymptotic observer.

The work in this section is based on $[131,151]$.

### 14.2 Vertical Take-off and Landing Aircraft

### 14.2.1 Control Objective

A scaled mathematical model of a VTOL aircraft can be described as, see [153]:

$$
\begin{aligned}
& \dot{x}_{1}=x_{2} \\
& \dot{x}_{2}=-u_{1} \sin (\theta)+\varepsilon u_{2} \cos (\theta), \\
& \dot{y}_{1}=y_{2} \\
& \dot{y}_{2}=u_{1} \cos (\theta)+\varepsilon u_{2} \sin (\theta)-g, \\
& \dot{\theta}=\omega \\
& \dot{\omega}=u_{2}
\end{aligned}
$$

where $x_{1}, y_{1}$, and $\theta$ denote position of the aircraft center of mass and roll angle, $x_{2}$, $y_{2}$, and $\omega$ denote linear and roll angular velocities of the aircraft, respectively, $u_{1}$ and $u_{2}$ are the vertical control force and rotational moment, $g>0$ is the gravitational acceleration, and $\varepsilon$ is the constant coupling between the roll moment and the lateral force, see Figure 14.8.


Figure 14.8 The VTOL aircraft

The aircraft model (14.92) is underactuated if we want to use two available control inputs ( $u_{1}$ and $u_{2}$ ) to control three outputs ( $x_{1}, y_{1}$, and $\theta$ ). Furthermore, the
zero-dynamics of the aircraft model (14.92) is nonminimum phase for $\varepsilon \neq 0$ at the steady state when considering $\left(x_{1}, y_{1}\right)$ as the output and $\theta$ as an internal state. This phenomenon can be seen from (14.92) by setting $x_{1}=y_{1}=x_{2}=y_{2}=0$.

We assume that the reference trajectory to be tracked is generated by

$$
\begin{align*}
& \dot{x}_{1 r}=x_{2 r}, \\
& \dot{x}_{2 r}=-u_{1 r} \sin \left(\theta_{r}\right)+\varepsilon u_{2 r} \cos \left(\theta_{r}\right), \\
& \dot{y}_{1 r}=y_{2 r},  \tag{14.93}\\
& \dot{y}_{2 r}=u_{1 r} \cos \left(\theta_{r}\right)+\varepsilon u_{2 r} \sin \left(\theta_{r}\right)-g, \\
& \dot{\theta}_{r}=\omega_{r}, \\
& \dot{\omega}_{r}=u_{2 r},
\end{align*}
$$

where all of the variables in (14.93) have the same meaning as in (14.92). In this section, we are interested in designing the control inputs $u_{1}$ and $u_{2}$ to force the aircraft model (14.92) to globally asymptotically track the reference model (14.93) without measurements of the velocities $x_{2}, y_{2}$, and $\omega$ under the following assumption:

## Assumption 14.3.

1. The reference signals $\omega_{r}, u_{1 r}, u_{2 r}, \dot{u}_{1 r}, \ddot{u}_{1 r}$ and $\dot{u}_{2 r}$ are bounded.
2. There exists a strictly positive constant $u_{1 r}^{*}$ such that

$$
\begin{equation*}
u_{1 r}-\varepsilon \omega_{r}^{2} \geq u_{1 r}^{*} \tag{14.94}
\end{equation*}
$$

Remark 14.5. The condition (14.94) covers the stabilization/regulation of the VTOL aircraft and implies that the aircraft is not allowed to land faster than it freely falls under the gravitational force.

### 14.2.2 Observer Design

From (14.92), a reduced-order observer can indeed be designed. However, such an observer is often noise-sensitive. Here we use the following full-order observer:

$$
\begin{align*}
& \dot{\hat{x}}_{1}=\hat{x}_{2}+k_{11}\left(x_{1}-\hat{x}_{1}\right), \\
& \dot{\hat{x}}_{2}=-u_{1} \sin (\theta)+\varepsilon u_{2} \cos (\theta)+k_{12}\left(x_{1}-\hat{x}_{1}\right), \\
& \dot{\hat{y}}_{1}=\hat{y}_{2}+k_{21}\left(y_{1}-\hat{y}_{1}\right), \\
& \dot{\hat{y}}_{2}=u_{1} \cos (\theta)+\varepsilon u_{2} \sin (\theta)-g+k_{22}\left(y_{1}-\hat{y}_{1}\right),  \tag{14.95}\\
& \dot{\hat{\theta}}=\hat{\omega}+k_{31}(\theta-\hat{\theta}), \\
& \dot{\hat{\omega}}=u_{2}+k_{32}(\theta-\hat{\theta}),
\end{align*}
$$

where $k_{i j}, 1 \leq i \leq 3$ and $1 \leq j \leq 2$ are positive constant observer gains. By defining the observer error as

$$
\tilde{X}=\left[\begin{array}{c}
\tilde{x}_{1}  \tag{14.96}\\
\tilde{x}_{2} \\
\tilde{y}_{1} \\
\tilde{y}_{2} \\
\tilde{\theta} \\
\tilde{\omega}
\end{array}\right]=\left[\begin{array}{c}
x_{1}-\hat{x}_{1} \\
x_{2}-\hat{x}_{2} \\
y_{1}-\hat{y}_{1} \\
y_{2}-\hat{y}_{2} \\
\theta-\hat{\theta} \\
\omega-\hat{\omega}
\end{array}\right]
$$

and subtracting (14.95) from (14.92), we have

$$
\dot{\tilde{X}}=A \tilde{X}, \quad A=\operatorname{diag}\left(A_{i}\right), A_{i}=\left[\begin{array}{ll}
-k_{i 1} & 1  \tag{14.97}\\
-k_{i 2} & 0
\end{array}\right], i=1,2,3 .
$$

It is easy to show that

$$
\begin{equation*}
\|\tilde{X}(t)\| \leq \varphi_{0}\left\|\tilde{X}\left(t_{0}\right)\right\| e^{-\sigma_{0}\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0 \tag{14.98}
\end{equation*}
$$

for some positive constants $\varphi_{0}$ and $\sigma_{0}$, which implies that (14.95) is a global exponential observer of (14.92). Therefore, in the following we will design the desired tracking controllers $u_{1}$ and $u_{2}$ based on the following transformed system:

$$
\begin{align*}
& \dot{x}_{1}=\hat{x}_{2}+\tilde{x}_{2}, \\
& \dot{\hat{x}}_{2}=-u_{1} \sin (\theta)+\varepsilon u_{2} \cos (\theta)+k_{12} \tilde{x}_{1}, \\
& \dot{y}_{1}=\hat{y}_{2}+\tilde{y}_{2},  \tag{14.99}\\
& \dot{\hat{y}}_{2}=u_{1} \cos (\theta)+\varepsilon u_{2} \sin (\theta)-g+k_{22} \tilde{y}_{1}, \\
& \dot{\theta}=\hat{\omega}+\tilde{\omega}, \\
& \dot{\hat{\omega}}=u_{2}+k_{32} \tilde{\theta} .
\end{align*}
$$

### 14.2.3 Coordinate Transformations

Solving $u_{2}$ from the last two equations of (14.92) and substituting it into the second and fourth equations of (14.92), we have

$$
\begin{align*}
& \frac{\mathrm{d}}{\mathrm{~d} t}\left(x_{1}-\varepsilon \sin (\theta)\right)=x_{2}-\varepsilon \cos (\theta) \omega \\
& \frac{\mathrm{d}}{\mathrm{~d} t}\left(x_{2}-\varepsilon \cos (\theta) \omega\right)=-\sin (\theta)\left(u_{1}-\varepsilon \omega^{2}\right) \\
& \frac{\mathrm{d}}{\mathrm{~d} t}\left(y_{1}+\varepsilon \cos (\theta)\right)=y_{2}-\varepsilon \sin (\theta) \omega  \tag{14.100}\\
& \frac{\mathrm{d}}{\mathrm{~d} t}\left(y_{2}-\varepsilon \sin (\theta) \omega\right)=\cos (\theta)\left(u_{1}-\varepsilon \omega^{2}\right)-g
\end{align*}
$$

If we consider $x_{1}-\varepsilon \sin (\theta)$ and $y_{1}+\varepsilon \cos (\theta)$ as outputs, it can be seen that (14.100) is of a triangular form and does not depend on $u_{2}$. It is of interest to note that the above outputs coincide with the aircraft center of oscillation and they are the flat outputs [154]. Motivated by the above discussion, we define the following coordinate changes:

$$
\begin{align*}
& z_{1}=x_{1}-\varepsilon \sin (\theta), \\
& \hat{z}_{2}=\hat{x}_{2}-\varepsilon \cos (\theta) \hat{\omega},  \tag{14.101}\\
& w_{1}=y_{1}+\varepsilon \cos (\theta) \\
& \hat{w}_{2}=\hat{y}_{2}-\varepsilon \sin (\theta) \hat{\omega}
\end{align*}
$$

Applying the above coordinate changes to (14.99) results in

$$
\begin{align*}
& \dot{z}_{1}=\hat{z}_{2}+\tilde{x}_{2}-\varepsilon \cos (\theta) \tilde{\omega}, \\
& \dot{w}_{1}=\hat{w}_{2}+\tilde{y}_{2}-\varepsilon \sin (\theta) \tilde{\omega}, \\
& \dot{z}_{2}=-\bar{u}_{1} \sin (\theta)+\varepsilon \sin (\theta) \hat{\omega} \tilde{\omega}+k_{12} \tilde{x}_{1}-\varepsilon k_{32} \cos (\theta) \tilde{\theta},  \tag{14.102}\\
& \dot{\hat{w}}_{2}=\bar{u}_{1} \cos (\theta)-g-\varepsilon \cos (\theta) \hat{\omega} \tilde{\omega}+k_{22} \tilde{y}_{1}-\varepsilon k_{32} \sin (\theta) \tilde{\theta}, \\
& \dot{\theta}=\hat{\omega}+\tilde{\omega}, \\
& \dot{\hat{\omega}}=u_{2}+k_{32} \tilde{\theta},
\end{align*}
$$

where $\bar{u}_{1}=u_{1}-\varepsilon \hat{\omega}^{2}$. Similarly, applying the coordinate changes

$$
\begin{aligned}
& z_{1 r}=x_{1 r}-\varepsilon \sin \left(\theta_{r}\right), \\
& z_{2 r}=x_{2 r}-\varepsilon \cos \left(\theta_{r}\right) \omega_{r} \\
& w_{1 r}=y_{1 r}+\varepsilon \cos \left(\theta_{r}\right), \\
& w_{2 r}=y_{2 r}-\varepsilon \sin \left(\theta_{r}\right) \omega_{r}
\end{aligned}
$$

to (14.93) yields

$$
\begin{align*}
& \dot{z}_{1 r}=z_{2 r}, \\
& \dot{w}_{1 r}=w_{2 r}, \\
& \dot{z}_{2 r}=-\bar{u}_{1 r} \sin \left(\theta_{r}\right),  \tag{14.104}\\
& \dot{w}_{2 r}=\bar{u}_{1 r} \cos \left(\theta_{r}\right)-g, \\
& \dot{\theta}_{r}=\omega_{r}, \\
& \dot{\omega}_{r}=u_{2 r},
\end{align*}
$$

where $\bar{u}_{1 r}=u_{1 r}-\varepsilon \omega_{r}^{2}$.
If we directly design the control inputs $\bar{u}_{1}$ and $u_{2}$ to force (14.102) to track (14.104), it will be very complicated since the angle $\theta$ enters in the third and fourth equations of (14.102). Motivated by controlling underactuated ships, we interpret the tracking errors in a frame attached to the virtual aircraft as

$$
\begin{align*}
& {\left[\begin{array}{l}
z_{1 e} \\
w_{1 e}
\end{array}\right]=J\left(\theta_{r}\right)\left[\begin{array}{l}
z_{1}-z_{1 r} \\
w_{1}-w_{1 r}
\end{array}\right],} \\
& {\left[\begin{array}{l}
z_{2 e} \\
w_{2 e}
\end{array}\right]=J\left(\theta_{r}\right)\left[\begin{array}{l}
\hat{z}_{2}-z_{2 r} \\
\hat{w}_{2}-w_{2 r}
\end{array}\right],}  \tag{14.105}\\
& {\left[\begin{array}{l}
\theta_{e} \\
\omega_{e}
\end{array}\right]=\left[\begin{array}{l}
\theta-\theta_{r} \\
\hat{\omega}-\omega_{r}
\end{array}\right],}
\end{align*}
$$

where

$$
J\left(\theta_{r}\right)=\left[\begin{array}{cc}
-\sin \left(\theta_{r}\right) & \cos \left(\theta_{r}\right)  \tag{14.106}\\
\cos \left(\theta_{r}\right) & \sin \left(\theta_{r}\right)
\end{array}\right]
$$

It is seen that the above coordinate changes are globally invertible and that convergence of $\left(z_{1 e}, w_{1 e}, z_{2 e}, w_{2 e}\right)$ to the origin implies that of $\left(z_{1}-z_{1 r}, w_{1}-w_{1 r}, \hat{z}_{2}-\right.$ $\left.z_{2 r}, \hat{w}_{2}-w_{2 r}\right)$.

Remark 14.6. If $J(\theta)$ instead of $J\left(\theta_{r}\right)$ is used in the coordinate transformation (14.105), it will be extremely difficult to design the control inputs for the resulting system.

Differentiating both sides of (14.105) along the solutions of (14.102) and (14.104) yields

$$
\begin{align*}
& \dot{z}_{1 e}=z_{2 e}-w_{1 e} \omega_{r}+O_{z 1}, \\
& \dot{w}_{1 e}=w_{2 e}+z_{1 e} \omega_{r}+O_{w 1}, \\
& \dot{z}_{2 e}=\bar{u}_{1} \cos \left(\theta_{e}\right)-\bar{u}_{1 r}-w_{2 e} \omega_{r}-\varepsilon \cos \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{z 2},  \tag{14.107}\\
& \dot{w}_{2 e}=-\bar{u}_{1} \sin \left(\theta_{e}\right)+z_{2 e} \omega_{r}+\varepsilon \sin \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{w 2}, \\
& \dot{\theta}_{e}=\omega_{e}+\tilde{\omega}, \\
& \dot{\omega}_{e}=u_{2}-u_{2 r}+k_{32} \tilde{\theta},
\end{align*}
$$

where for simplicity of presentation, we have defined the following terms which exponentially converge to zero when $\tilde{X}$ does:

$$
\begin{align*}
& O_{z 1}=-\sin \left(\theta_{r}\right)\left(\tilde{x}_{2}-\varepsilon \cos (\theta) \tilde{\omega}\right)+\cos \left(\theta_{r}\right)\left(\tilde{y}_{2}-\varepsilon \sin (\theta) \tilde{\omega}\right), \\
& O_{w 1}=\cos \left(\theta_{r}\right)\left(\tilde{x}_{2}-\varepsilon \cos (\theta) \tilde{\omega}\right)+\sin \left(\theta_{r}\right)\left(\tilde{y}_{2}-\varepsilon \sin (\theta) \tilde{\omega}\right), \\
& O_{z 2}=-\sin \left(\theta_{r}\right)\left(k_{12} \tilde{x}_{1}-\varepsilon k_{32} \cos (\theta) \tilde{\theta}\right)+ \\
& \quad \cos \left(\theta_{r}\right)\left(k_{22} \tilde{y}_{1}-\varepsilon k_{32} \sin (\theta) \tilde{\theta}\right)-\varepsilon \cos \left(\theta_{e}\right) \omega_{r} \tilde{\omega},  \tag{14.108}\\
& O_{w 2}=\cos \left(\theta_{r}\right)\left(k_{12} \tilde{x}_{1}-\varepsilon k_{32} \cos (\theta) \tilde{\theta}\right)+ \\
& \quad \sin \left(\theta_{r}\right)\left(k_{22} \tilde{y}_{1}-\varepsilon k_{32} \sin (\theta) \tilde{\theta}\right)+\varepsilon \sin \left(\theta_{e}\right) \omega_{r} \tilde{\omega}
\end{align*}
$$

We have therefore converted the tracking control problem to a problem of stabilizing (14.107) at the origin.

### 14.2.4 Control Design

It is observed that (14.107) is of a triangular form. Although there are connection terms between the $\left(z_{1 e}, z_{2 e}\right)$ - and $\left(w_{1 e}, w_{2 e}\right)$-subsystems, they do not prevent us from applying the popular backstepping technique [10] to (14.107), see below. To make our control design clear, we divide the control design into two steps. The first step designs $\bar{u}_{1}$. The control $u_{2}$ is designed in the second step.

## Step 1

At this step, we design the control input $\bar{u}_{1}$ and as an application of our approach used in underactuated ship tracking, the angle error $\theta_{e}$ is also used as a "control" to stabilize the $\left(z_{1 e}, z_{2 e}, w_{1 e}, w_{2 e}\right)$-subsystem. We proceed in two substeps.

Substep 1
Define

$$
\begin{align*}
& \bar{z}_{2 e}=z_{2 e}-\alpha_{z}, \\
& \bar{w}_{2 e}=w_{2 e}-\alpha_{w}, \tag{14.109}
\end{align*}
$$

where $\alpha_{z}$ and $\alpha_{w}$ are virtual controls of $z_{2 e}$ and $w_{2 e}$, respectively. Consider the Lyapunov function

$$
\begin{equation*}
V_{11}=\sqrt{1+z_{1 e}^{2}+w_{1 e}^{2}}-1 \tag{14.110}
\end{equation*}
$$

whose time derivative along the solutions of (14.109) and the first two equations of (14.107) is

$$
\begin{equation*}
\dot{V}_{11}=\frac{z_{1 e}}{\Delta_{1}}\left(\bar{z}_{2 e}+\alpha_{z}+O_{z 1}\right)+\frac{w_{1 e}}{\Delta_{1}}\left(\bar{w}_{2 e}+\alpha_{w}+O_{w 1}\right), \tag{14.111}
\end{equation*}
$$

where $\Delta_{1}=\sqrt{1+z_{1 e}^{2}+w_{1 e}^{2}}$. From (14.111), we choose

$$
\begin{equation*}
\alpha_{z}=-\frac{k_{1} z_{1 e}}{\Delta_{2}}, \alpha_{w}=-\frac{k_{2} w_{1 e}}{\Delta_{2}}, \tag{14.112}
\end{equation*}
$$

where $\Delta_{2}=\sqrt{1+z_{1 e}^{2}+w_{1 e}^{2}+z_{2 e}^{2}+w_{2 e}^{2}}$, and $k_{1}$ and $k_{2}$ are positive constants to be specified later.

Remark 14.7. The choice of the virtual controls $\alpha_{z}$ and $\alpha_{w}$ in (14.112) is different from a standard application of the backstepping technique in the sense that $\alpha_{z}$ and $\alpha_{w}$ depend on $z_{2 e}$ and $w_{2 e}$.This choice together with the Lyapunov function (14.110) instead of a quadratic function will result in a bounded control input $\bar{u}_{1}$ (see Substep 1.2). This bounded control is crucial to obtain a global result, see Step
2. In addition, the choice of (14.110) and (14.112) can also be considered as an interesting application of the bounded backstepping approaches recently proposed in [155] and [156] to the VTOL aircraft system.

Substituting (14.112) into (14.111) gives

$$
\begin{equation*}
\dot{V}_{11}=-\frac{k_{1} z_{1 e}^{2}+k_{2} w_{1 e}^{2}}{\Delta_{1} \Delta_{2}}+\frac{z_{1 e} \bar{z}_{2 e}}{\Delta_{1}}+\frac{w_{1 e} \bar{w}_{2 e}}{\Delta_{1}}+O_{11} \tag{14.113}
\end{equation*}
$$

where the term $O_{11}$ containing the observer errors as a factor is defined as

$$
\begin{equation*}
O_{11}=\frac{z_{1 e} O_{z 1}}{\Delta_{1}}+\frac{w_{1 e} O_{w 1}}{\Delta_{1}} \tag{14.114}
\end{equation*}
$$

## Substep 2

Define

$$
\begin{equation*}
\bar{\theta}_{e}=\theta_{e}-\alpha_{\theta} \tag{14.115}
\end{equation*}
$$

where $\alpha_{\theta}$ is a virtual control of $\theta_{e}$. To design $\bar{u}_{1}$ and $\alpha_{\theta}$, we take the following Lyapunov function:

$$
\begin{equation*}
V_{12}=V_{11}+\frac{1}{2} k_{3}\left(\bar{z}_{2 e}^{2}+\bar{w}_{2 e}^{2}\right), \tag{14.116}
\end{equation*}
$$

where $k_{3}$ is a positive constant, which is introduced to enhance the flexibility of choosing the design constants. Differentiating both sides of (14.116) along the solutions of (14.112) and (14.113), the third and fourth equations of (14.107) give

$$
\begin{gather*}
\dot{V}_{12}=-\frac{k_{1} z_{1 e}^{2}+k_{2} w_{1 e}^{2}}{\Delta_{1} \Delta_{2}}+k_{3} \bar{z}_{2 e}\left[b_{11}\left(\bar{u}_{1} \cos \left(\alpha_{\theta}\right)-\bar{u}_{1 r}\right)+b_{12} \bar{u}_{1} \sin \left(\alpha_{\theta}\right)+\right. \\
\left.f_{z}\right]+k_{3} \bar{w}_{2 e}\left[b_{21}\left(\bar{u}_{1} \cos \left(\alpha_{\theta}\right)-\bar{u}_{1 r}\right)+b_{22} \bar{u}_{1} \sin \left(\alpha_{\theta}\right)+f_{w}\right]+g_{\theta}+O_{12} \tag{14.117}
\end{gather*}
$$

where for simplicity of presentation, we have defined the following terms: The terms $b_{i j}, i=1,2, j=1,2$, are

$$
\begin{align*}
& b_{11}=1-\frac{k_{1} z_{1 e} z_{2 e}}{\Delta_{2}^{3}}, \\
& b_{12}=\frac{k_{1} z_{1 e} w_{2 e}}{\Delta_{2}^{3}}, \\
& b_{21}=-\frac{k_{2} w_{1 e} z_{2 e}}{\Delta_{2}^{3}},  \tag{14.118}\\
& b_{22}=-1+\frac{k_{2} w_{1 e} w_{2 e}}{\Delta_{2}^{3}} .
\end{align*}
$$

The terms $f_{z}$ and $f_{w}$, which will be canceled in this step, are

$$
\begin{align*}
f_{z} & =\frac{z_{1 e}}{k_{3} \Delta_{1}}-\alpha_{w} \omega_{r}+\frac{k_{1}}{\Delta_{2}}\left(z_{2 e}-w_{1 e} \omega_{r}\right)-\frac{k_{1} z_{1 e}}{\Delta_{2}^{3}}\left(z_{1 e} z_{2 e}+w_{1 e} w_{2 e}\right) \\
f_{w} & =\frac{w_{1 e}}{k_{3} \Delta_{1}}+\alpha_{z} \omega_{r}+\frac{k_{2}}{\Delta_{2}}\left(w_{2 e}+z_{1 e} \omega_{r}\right)-\frac{k_{2} w_{1 e}}{\Delta_{2}^{3}}\left(z_{1 e} z_{2 e}+w_{1 e} w_{2 e}\right) . \tag{14.119}
\end{align*}
$$

The term $g_{\theta}$, which will be canceled in the next step, is

$$
\begin{align*}
& g_{\theta}=k_{3} \bar{u}_{1}\left[\left(b_{11} \bar{z}_{2 e}+b_{21} \bar{w}_{2 e}\right)\left(\left(\cos \left(\bar{\theta}_{e}\right)-1\right) \cos \left(\alpha_{\theta}\right)-\sin \left(\bar{\theta}_{e}\right) \sin \left(\alpha_{\theta}\right)\right)+\right. \\
& \left.\quad\left(b_{12} \bar{z}_{2 e}+b_{22} \bar{w}_{2 e}\right)\left(\left(\cos \left(\bar{\theta}_{e}\right)-1\right) \sin \left(\alpha_{\theta}\right)+\sin \left(\bar{\theta}_{e}\right) \cos \left(\alpha_{\theta}\right)\right)\right] . \tag{14.120}
\end{align*}
$$

The term $O_{12}$, which contains the observer error as a factor, is defined as

$$
\begin{align*}
O_{12}= & k_{3} \bar{z}_{2 e}\left[-\varepsilon \cos \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{z 2}+\frac{k_{1} O_{z 1}}{\Delta_{2}}\right]+ \\
& k_{3} \bar{w}_{2 e}\left[-\varepsilon \sin \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{w 2}+\frac{k_{2} O_{w 1}}{\Delta_{2}}\right]-  \tag{14.121}\\
& \frac{k_{3}\left(k_{1} z_{1 e} \bar{z}_{2 e}+k_{2} w_{1 e} \bar{w}_{2 e}\right)}{\Delta_{2}^{3}}\left[z_{1 e} O_{z 1}+w_{1 e} O_{w 1}+\right. \\
& \left.z_{2 e}\left(-\varepsilon \cos \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{z 2}\right)+w_{2 e}\left(\varepsilon \sin \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{w 2}\right)\right]+O_{11} .
\end{align*}
$$

Since $O_{12}$ is Lipschitz in $\bar{z}_{2 e} \omega_{e}$ and $\bar{w}_{2 e} \omega_{e}$, from (14.117) we choose the control input $\bar{u}_{1}$ and $\alpha_{\theta}$ such that

$$
\begin{align*}
& b_{11}\left(\bar{u}_{1} \cos \left(\alpha_{\theta}\right)-\bar{u}_{1 r}\right)+b_{12} \bar{u}_{1} \sin \left(\alpha_{\theta}\right)+f_{z}=-k_{4} \bar{z}_{2 e} / \Delta_{3}, \\
& b_{21}\left(\bar{u}_{1} \cos \left(\alpha_{\theta}\right)-\bar{u}_{1 r}\right)+b_{22} \bar{u}_{1} \sin \left(\alpha_{\theta}\right)+f_{w}=-k_{5} \bar{w}_{2 e} / \Delta_{3}, \tag{14.122}
\end{align*}
$$

where $\Delta_{3}=\sqrt{1+\bar{z}_{2 e}^{2}+\bar{w}_{2 e}^{2}}, k_{4}$ and $k_{5}$ are positive constants to be selected later. Solving (14.122) results in

$$
\begin{align*}
\alpha_{\theta} & =\arctan \left(\bar{u}_{12} /\left(\bar{u}_{1 r}+\bar{u}_{11}\right)\right),  \tag{14.123}\\
\bar{u}_{1} & =\left(\bar{u}_{1 r}+\bar{u}_{11}\right) \cos \left(\alpha_{\theta}\right)+\bar{u}_{12} \sin \left(\alpha_{\theta}\right)
\end{align*}
$$

where

$$
\begin{align*}
& \bar{u}_{11}=\frac{1}{b_{11} b_{22}-b_{12} b_{21}}\left(-\left(\frac{k_{4} \bar{z}_{2 e}}{\Delta_{3}}+f_{z}\right) b_{22}+\left(\frac{k_{5} \bar{w}_{2 e}}{\Delta_{3}}+f_{w}\right) b_{12}\right),  \tag{14.124}\\
& \bar{u}_{12}=\frac{1}{b_{11} b_{22}-b_{12} b_{21}}\left(\left(\frac{k_{4} \bar{z}_{2 e}}{\Delta_{3}}+f_{z}\right) b_{21}-\left(\frac{k_{5} \bar{w}_{2 e}}{\Delta_{3}}+f_{w}\right) b_{11}\right) .
\end{align*}
$$

From expressions of $b_{i j}, i=1,2, j=1,2$, given in (14.118), we have

$$
\begin{equation*}
b_{11} b_{22}-b_{12} b_{21}=-1+\frac{k_{1} z_{1 e} z_{2 e}}{\Delta_{2}^{3}}+\frac{k_{2} w_{1 e} w_{2 e}}{\Delta_{2}^{3}}-\frac{2 k_{1} k_{2} z_{1 e} z_{2 e} w_{1 e} w_{2 e}}{\Delta_{2}^{6}} \tag{14.125}
\end{equation*}
$$

Therefore, there is no singularity in (14.124), if $k_{1}$ and $k_{2}$ are chosen such that

$$
\begin{equation*}
k_{1}+k_{2}+2 k_{1} k_{2}<1 \tag{14.126}
\end{equation*}
$$

Moreover, the virtual control $\alpha_{\theta}$ in (14.123) is well defined if

$$
\begin{equation*}
\bar{u}_{1 r}+\bar{u}_{11}>0 . \tag{14.127}
\end{equation*}
$$

This condition, as a simple calculation shows, is equivalent to

$$
\begin{align*}
& u_{1 r}^{*}-( \frac{\omega_{r}^{\max }\left(k_{1}+k_{2}\right)\left(1+k_{1}+k_{2}\right)+\left(1+k_{2}\right)\left(3 k_{1}+1 / k_{3}+k_{4}\right)}{1-\left(k_{1}+k_{2}+2 k_{1} k_{2}\right)}+ \\
&\left.\frac{k_{1}\left(k_{5}+1 / k_{3}+3 k_{2}\right)}{1-\left(k_{1}+k_{2}+2 k_{1} k_{2}\right)}\right)>0 \tag{14.128}
\end{align*}
$$

Indeed, there always exist positive constants $k_{i}, 1 \leq i \leq 5$ such that (14.128) holds. Also it can be shown from (14.123) and (14.124) that $\bar{u}_{1}$ is bounded by some positive constant, which can be easily calculated from (14.123) and (14.124). Note that under the conditions (14.126) and (14.128), $\alpha_{\theta}$ is a smooth function of $z_{1 e}, z_{2 e}$, $w_{1 e}, w_{2 e}, \omega_{r}$, and $\bar{u}_{1 r}$. Substituting (14.123) into (14.117) yields

$$
\begin{equation*}
\dot{V}_{12}=-\frac{k_{1} z_{1 e}^{2}+k_{2} w_{1 e}^{2}}{\Delta_{1} \Delta_{2}}-\frac{k_{3}\left(k_{4} \bar{z}_{2 e}^{2}+k_{5} \bar{w}_{2 e}^{2}\right)}{\Delta_{3}}+g_{\theta}+O_{12} \tag{14.129}
\end{equation*}
$$

## Step 2

In this step, we design $u_{2}$ to stabilize the $\left(\theta_{e}, \omega_{e}\right)$-system. Define

$$
\begin{equation*}
\bar{\omega}_{e}=\omega_{e}-\alpha_{\omega}, \tag{14.130}
\end{equation*}
$$

where $\alpha_{\omega}$ is a virtual control of $\omega_{e}$. To design $\alpha_{\omega}$, we take the Lyapunov function

$$
\begin{equation*}
V_{21}=V_{12}+0.5 \bar{\theta}_{e}^{2} . \tag{14.131}
\end{equation*}
$$

By differentiating both sides of (14.131) along the solutions of (14.129) and the fifth equation of (14.107), and using (14.130), the virtual control $\alpha_{\omega}$ is chosen as

$$
\begin{align*}
\alpha_{\omega}= & -k_{6} \bar{\theta}_{e}-\frac{g_{\theta}}{\bar{\theta}_{e}}+\frac{\partial \alpha_{\theta}}{\partial z_{1 e}}\left(z_{2 e}-w_{1 e} \omega_{r}\right)+\frac{\partial \alpha_{\theta}}{\partial z_{2 e}}\left(\bar{u}_{1} \cos \left(\theta_{e}\right)-\bar{u}_{1 r}-w_{2 e} \omega_{r}\right)+ \\
& \frac{\partial \alpha_{\theta}}{\partial w_{1 e}}\left(w_{2 e}+z_{1 e} \omega_{r}\right)+\frac{\partial \alpha_{\theta}}{\partial w_{2 e}}\left(-\bar{u}_{1} \sin \left(\theta_{e}\right)+z_{2 e} \omega_{r}\right)+\frac{\partial \alpha_{\theta}}{\partial \omega_{r}} \dot{\omega}_{r}+\frac{\partial \alpha_{\theta}}{\partial \bar{u}_{1 r}} \dot{\bar{u}}_{1 r}, \tag{14.132}
\end{align*}
$$

where $k_{6}$ is a positive constant. Note that $g_{\theta} / \bar{\theta}_{e}$ with $g_{\theta}$ being given in (14.120) is well defined because $\frac{\sin \left(\bar{\theta}_{e}\right)}{\bar{\theta}_{e}}=\int_{0}^{1} \cos \left(\bar{\theta}_{e} \lambda\right) \mathrm{d} \lambda$ and $\frac{\left(\cos \left(\bar{\theta}_{e}\right)-1\right)}{\bar{\theta}_{e}}=-\int_{0}^{1} \sin \left(\bar{\theta}_{e} \lambda\right) \mathrm{d} \lambda$ are smooth functions of $\bar{\theta}_{e}$. Hence $\alpha_{\omega}$ is a smooth function of $z_{1 e}, z_{2 e}, w_{1 e}, w_{2 e}, \theta_{e}$, $\omega_{r}, \dot{\omega}_{r}, \bar{u}_{1 r}$ and $\dot{\bar{u}}_{1 r}$. Note that at this stage, we do not need a nonlinear damping term for the observer effects since $O_{21}$, see (14.134), is Lipschitz in $\bar{\theta}_{e} \omega_{e}$ and we wish to design a weak nonlinear control law of $\alpha_{\omega}$.

Differentiating both sides of (14.131) along the solutions of (14.129) and the fifth equation of (14.107), and using (14.130), (14.132) yields

$$
\begin{equation*}
\dot{V}_{21}=-\frac{k_{1} z_{1 e}^{2}+k_{2} w_{1 e}^{2}}{\Delta_{1} \Delta_{2}}-\frac{k_{3}\left(k_{4} \bar{z}_{2 e}^{2}+k_{5} \bar{w}_{2 e}^{2}\right)}{\Delta_{3}}-k_{6} \bar{\theta}_{e}^{2}+\bar{\theta}_{e} \bar{\omega}_{e}+O_{21} \tag{14.133}
\end{equation*}
$$

where the term $O_{21}$, which contains the observer error as a factor, is defined as

$$
\begin{gather*}
O_{21}=\bar{\theta}_{e}\left[\tilde{\omega}+\frac{\partial \alpha_{\theta}}{\partial z_{1 e}} O_{z 1}+\frac{\partial \alpha_{\theta}}{\partial w_{1 e}} O_{w 1}+\frac{\partial \alpha_{\theta}}{\partial z_{2 e}}\left(-\varepsilon \cos \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{z 2}\right)+\right. \\
\left.\frac{\partial \alpha_{\theta}}{\partial w_{2 e}}\left(\varepsilon \sin \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{w 2}\right)\right]+O_{12} \tag{14.134}
\end{gather*}
$$

We are now ready to design the control input $u_{2}$ by considering the Lyapunov function

$$
\begin{equation*}
V_{22}=V_{21}+\frac{1}{2} \bar{\omega}_{e}^{2} \tag{14.135}
\end{equation*}
$$

By differentiating both sides of (14.135) along the solutions of (14.133) and the sixth equation of (14.107), and using (14.130) and (14.132), the control input $u_{2}$ is chosen as

$$
\begin{align*}
u_{2}= & -k_{7} \bar{\omega}_{e}-\bar{\theta}_{e}+u_{2 r}+\frac{\partial \alpha_{\omega}}{\partial \omega_{r}} \dot{\omega}_{r}+\frac{\partial \alpha_{\omega}}{\partial \dot{\omega}_{r}} \ddot{\omega}_{r}+\frac{\partial \alpha_{\omega}}{\partial \bar{u}_{1 r}} \dot{\bar{u}}_{1 r}+ \\
& \frac{\partial \alpha_{\omega}}{\partial \dot{\bar{u}}_{1 r}} \ddot{\bar{u}}_{1 r}+\frac{\partial \alpha_{\omega}}{\partial z_{1 e}}\left(z_{2 e}-w_{1 e} \omega_{r}\right)+\frac{\partial \alpha_{\omega}}{\partial z_{2 e}}\left(\bar{u}_{1} \cos \left(\theta_{e}\right)-\bar{u}_{1 r}-w_{2 e} \omega_{r}\right)+ \\
& \frac{\partial \alpha_{\omega}}{\partial w_{1 e}}\left(w_{2 e}+z_{1 e} \omega_{r}\right)+\frac{\partial \alpha_{\omega}}{\partial w_{2 e}}\left(-\bar{u}_{1} \sin \left(\theta_{e}\right)+z_{2 e} \omega_{r}\right)+\frac{\partial \alpha_{\omega}}{\partial \theta_{e}} \omega_{e}+u_{2 \mathrm{dam}}, \tag{14.136}
\end{align*}
$$

where $k_{7}$ is a positive constant and the damping term $u_{2 \text { dam }}$ to take care of observer error effects is given by

$$
\begin{align*}
u_{2 d a m}= & -\delta\left[\left(\frac{\partial \alpha_{\omega}}{\partial z_{1 e}}\right)^{2}+\left(\frac{\partial \alpha_{\omega}}{\partial w_{1 e}}\right)^{2}+\left(\frac{\partial \alpha_{\omega}}{\partial z_{2 e}}\right)^{2}\left(\omega_{e}^{2}+1\right)+\right. \\
& \left.\left(\frac{\partial \alpha_{\omega}}{\partial w_{2 e}}\right)^{2}\left(\omega_{e}^{2}+1\right)+\left(\frac{\partial \alpha_{\omega}}{\partial \theta_{e}}\right)^{2}\right] \bar{\omega}_{e} \tag{14.137}
\end{align*}
$$

where $\delta$ is a positive constant. Differentiating both sides of (14.135) along the solutions of (14.133) and the sixth equation of (14.107), and using (14.130), (14.132), and (14.136) yields

$$
\begin{align*}
\dot{V}_{22}= & -\frac{k_{1} z_{1 e}^{2}+k_{2} w_{1 e}^{2}}{\Delta_{1} \Delta_{2}}-\frac{k_{3}\left(k_{4} \bar{z}_{2 e}^{2}+k_{5} \bar{w}_{2 e}^{2}\right)}{\Delta_{3}}- \\
& k_{6} \bar{\theta}_{e}^{2}-k_{7} \bar{\omega}_{e}^{2}+u_{2 \operatorname{dam}} \bar{\omega}_{e}+O_{22}, \tag{14.138}
\end{align*}
$$

where the term $O_{22}$ containing the observer error as a factor, is defined as

$$
\begin{align*}
O_{22}= & \bar{\omega}_{e}\left[k_{32} \tilde{\theta}+\frac{\partial \alpha_{\omega}}{\partial z_{1 e}} O_{z 1}+\frac{\partial \alpha_{\omega}}{\partial w_{1 e}} O_{w 1}+\frac{\partial \alpha_{\omega}}{\partial z_{2 e}}\left(-\varepsilon \cos \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{z 2}\right)+\right. \\
& \left.\frac{\partial \alpha_{\omega}}{\partial w_{2 e}}\left(\varepsilon \sin \left(\theta_{e}\right) \omega_{e} \tilde{\omega}+O_{w 2}\right)+\frac{\partial \alpha_{\omega}}{\partial \theta_{e}} \tilde{\omega}\right]+O_{21} \tag{14.139}
\end{align*}
$$

By replacing $\omega_{e}$ by $\bar{\omega}_{e}+\alpha_{\omega}$ in all of the terms $O_{12}, O_{21}$, and $O_{22}$, and noting that $\bar{u}_{1}$ is bounded by some positive constant, after a lengthy but simple calculation by completing squares, it is readily shown that

$$
\begin{align*}
\dot{V}_{22} \leq & -\frac{k_{1} z_{1 e}^{2}+k_{2} w_{1 e}^{2}}{\Delta_{1} \Delta_{2}}-\frac{k_{3}\left(k_{4} \bar{z}_{2 e}^{2}+k_{5} \bar{w}_{2 e}^{2}\right)}{\Delta_{3}}-k_{6} \bar{\theta}_{e}^{2}-k_{7} \bar{\omega}_{e}^{2}+ \\
& \left(\chi_{1}(\cdot) V_{22}+\chi_{2}(\cdot)\right) e^{-\sigma_{0}\left(t-t_{0}\right)} \\
& \leq\left(\chi_{1}(\cdot) V_{22}+\chi_{2}(\cdot)\right) e^{-\sigma_{0}\left(t-t_{0}\right)}, \tag{14.140}
\end{align*}
$$

where $\chi_{1}(\cdot)$ and $\chi_{2}(\cdot)$ are some class- $K$ functions of $\left\|\tilde{X}\left(t_{0}\right)\right\|$. To prove convergence of $\left(z_{1 e}, w_{1 e}, \bar{z}_{2 e}, \bar{w}_{2 e}, \bar{\theta}_{e}, \bar{\omega}_{e}\right)$ to zero based on (14.140), we need the following lemma.

Lemma 14.1. Consider the following first-order scalar differential equation

$$
\begin{equation*}
\dot{x}=(a x+b) e^{-c\left(t-t_{0}\right)}, \forall t \geq t_{0} \geq 0 \tag{14.141}
\end{equation*}
$$

where $a \geq 0, b \geq 0$ and $c>0$ are constants. The solution of (14.141) is bounded and satisfies

$$
\begin{equation*}
|x(t)| \leq\left|x\left(t_{0}\right)\right| e^{a / c}+b a^{-1}\left(e^{a / c}-1\right):=\pi\left(\left|x\left(t_{0}\right)\right|\right), \forall t \geq t_{0} \geq 0 \tag{14.142}
\end{equation*}
$$

Proof. Solving (14.141) for $x$ then taking the norm of $x$ gives (14.142).
Applying the above lemma to the second inequality of (14.140), we have

$$
V_{22}(t) \leq \pi_{22}\left(\left\|\left(\tilde{X}\left(t_{0}\right), z_{1 e}\left(t_{0}\right), w_{1 e}\left(t_{0}\right), \bar{z}_{2 e}\left(t_{0}\right), \bar{w}_{2 e}\left(t_{0}\right), \bar{\theta}_{e}\left(t_{0}\right), \bar{\omega}_{e}\left(t_{0}\right)\right)\right\|\right)
$$

with $\pi_{22}(\cdot)$ being calculated from the above lemma. This implies that $z_{2 e}, w_{2 e}, \theta_{e}$, $\omega_{e}, \bar{u}_{1}$, and $u_{2}$ are also bounded for all $t \in[0, \infty)$. With this observation in mind, the first line of (14.140) yields

$$
\begin{align*}
\dot{V}_{22} \leq & -\frac{k_{1} z_{1 e}^{2}+k_{2} w_{1 e}^{2}}{\Delta_{1 m} \Delta_{2 m}}-\frac{k_{3}\left(k_{4} \bar{z}_{2 e}^{2}+k_{5} \bar{w}_{2 e}^{2}\right)}{\Delta_{3 m}}-k_{6} \bar{\theta}_{e}^{2}-k_{7} \bar{\omega}_{e}^{2}+ \\
& \left(\chi_{1}(\cdot) \pi_{22}(\cdot)+\chi_{2}(\cdot)\right) e^{-\sigma_{0}\left(t-t_{0}\right)}, \tag{14.143}
\end{align*}
$$

where $\Delta_{i m}, i=1,2,3$ are $\Delta_{i}$ with the arguments being replaced by their upper bounds. Applying Babarlat's lemma to (14.143), it is readily shown that

$$
\begin{equation*}
\lim _{t \rightarrow \infty}\left(z_{1 e}(t), w_{1 e}(t), \bar{z}_{2 e}(t), \bar{w}_{2 e}(t), \bar{\theta}_{e}(t), \bar{\omega}_{e}(t)\right)=0 \tag{14.144}
\end{equation*}
$$

By construction, see (14.109), (14.112), (14.115) and (14.123), the limit (14.144) implies that

$$
\lim _{t \rightarrow \infty}\left(z_{1 e}(t), w_{1 e}(t), z_{2 e}(t), w_{2 e}(t), \theta_{e}(t), \omega_{e}(t)\right)=0
$$

We now summarize the main result of this section in the following theorem.
Theorem 14.2. Under Assumption 14.3, the dynamic output feedback control law consisting of (14.99), (14.123), and (14.136) forces the aircraft (14.92) to globally asymptotically track the reference model (14.93) if the design constants $k_{i}, 1 \leq i \leq$ 5, are chosen such that (14.126) and (14.128) hold.

### 14.2.5 Simulations

In this section, we perform a numerical simulation to illustrate the effectiveness of the proposed controller with $\varepsilon=0.8$. The observer and control gains are chosen as: $k_{11}=k_{21}=k_{31}=1, k_{12}=k_{22}=k_{32}=2, k_{1}=k_{2}=0.15, k_{3}=5$, $k_{4}=k_{5}=0.4, k_{6}=k_{7}=5, \delta=0.1$, and initial conditions are $x_{1}(0)=y_{1}(0)=5$, $\theta(0)=0.2, x_{2}(0)=y_{2}(0)=\omega(0)=0.2, \hat{x}_{1}(0)=\hat{y}_{1}(0)=3, \hat{\theta}(0)=0.1, x_{1 r}(0)=$ $y_{1 r}(0)=\theta_{r}(0)=x_{2 r}(0)=y_{2 r}(0)=\omega_{r}(0)=0, \hat{x}_{2}(0)=\hat{y}_{2}(0)=\hat{\omega}(0)=0$. The goal is to force the aircraft to track a sinusoid signal of $5(\sin (0.1 t)+1.2)$ in the vertical plane generated by (14.93). It is shown that there exists $u_{1 r}$ for this case with $u_{1 r}^{*}>2$ (see Assumption 14.3) and that conditions (14.126) and (14.128) hold. The simulation results are plotted in Figures 14.9, 14.10, and 14.11. It is seen that the proposed controller forces the aircraft to track the reference trajectory very well. The oscillation in the control input $u_{2}$ is due to the initial conditions and the above selection of the control and observer gains. By further tuning these gains, we can improve the transient response. In general, higher control and observer gains result in a shorter transient response time but more overshoot in the velocities and control inputs of the aircraft.


Figure 14.9 Reference (dotted line) and real (solid line) position trajectories: a. Position in $x$ direction; $\mathbf{b}$. Position in $y$ direction; $\mathbf{c}$. Velocity in $x$ direction; $\mathbf{d}$. Velocity in $y$ direction


Figure 14.10 Reference (dotted line) and real (solid line) angular trajectories, and control inputs: a. Angles; b. Angular velocities; c. Control input $u_{1} ;$ d. Control input $u_{2}$


Figure 14.11 Observer errors: a. Errors $x_{1}-\hat{x}_{1}$ (solid line) and $y_{1}-\hat{y}_{1}$ (dash-dotted line); $\mathbf{b}$. Errors $x_{2}-\hat{x}_{2}$ (solid line) and $y_{2}-\hat{y}_{2}$ (dash-dotted line); c. Errors $\theta-\hat{\theta}$ (solid line) and $\omega-\hat{\omega}$ (dash-dotted line)

### 14.2.6 Notes and References

The main difficulty in controlling VTOL aircraft is that they are underactuated and nonminimum phase. An approximate input-output linearization approach was used in $[1,153,157-159]$ to develop a controller for stabilization and output tracking/regulation of a VTOL aircraft. In these papers, the controller was initially designed by ignoring the coupling between roll moment and thrust. The controller parameters were then selected to take into account the effects of the coupling. In [154], by noting that the output at a fixed point with respect to the aircraft body (Huygens center of oscillation) can be used, an interesting approach was introduced to design an output tracking controller. However, the proposed controller was not defined in the whole space. A simple approach was developed in [160] to provide a global controller for the stabilization of a VTOL aircraft. An optimal controller was provided in [6] for robust hovering control of a VTOL aircraft. In [161], dynamic inversion and robust control techniques were used to deal with the nonminimum phase dynamics. However, this approach imposed restrictions on the desired reference trajectories. Recently, a dynamic high-gain approach was used in [162] to design a controller to force the VTOL aircraft to globally practically track a reference trajectory generated by a reference model. In all of the aforementioned papers, all of the VTOL aircraft states are required for feedback. The output feedback tracking controller for a VTOL
aircraft in this section is based on [163]. Under this controller, the VTOL aircraft globally asymptotically tracks the reference trajectory generated by the reference model. Indeed, the tracking controller proposed in this section covered stabilization and output tracking/regulation problems studied in the above-mentioned papers.

### 14.3 Conclusions

This chapter has illustrated that a careful investigation of the observer design and control design techniques developed for the ocean vessels in this book can result in various solutions of many control problems for other underactuated mechanical systems that are common in practice.

## Chapter 15 <br> Conclusions and Perspectives

### 15.1 Summary of the Book

This monograph has concentrated on the control of underactuated ocean vessels including ships and underwater vehicles. These vessels have more degrees of freedom to be controlled than the number of independent control inputs. Ships do not have an independent sway actuator while for underwater vehicles there are no independent sway and heave actuators. The book started with a review of the necessary background on ocean vessel dynamics and nonlinear control theory. The authors then demonstrated a systematic approach based on various nontrivial coordinate transformations together with advanced nonlinear control design methodologies founded on the basis of Lyapunov's direct method, backstepping, and parameter projection techniques for the development and analysis of a number of ocean vessel control systems to achieve advanced motion control tasks. These tasks include stabilization, trajectory-tracking, path-tracking, and path-following. The book has offered new knowledge regarding the nonlinear control of underactuated ocean vessels, efficient controllers for practical control of underactuated ocean vessels, and general methods and strategies to solve nonlinear control problems of other underactuated systems, including underactuated land and aerial vehicles. Numerical simulations and real-time implementations of the designed control systems on a scaled model ship have been included to illustrate the effectiveness and practical guidance of the control systems developed in the book. As an illustration of the constructive feature of the book, the observer design, control design, and stability analysis techniques developed for underactuated ocean vessels have been applied to control of other underactuated mechanical systems including mobile robots and VTOL aircraft.

The book consists of 15 Chapters. Chapter 1 presented a brief review of the development in nonlinear control theory and its applications, and difficulties in the control of underactuated ocean vessels. The main contributions of the book are organized into five parts.

In Part I (Chapter 2), various mathematical tools such as Barbalat-like lemmas, Lyapunov stability theory, and the backstepping technique were presented for control design and stability analysis of controlled systems designed in the book.

Part II consists of Chapters 3 and 4. This part addressed modeling, motion control tasks, and control properties of ocean vessels. The existing literature on the control of underactuated ocean vessels was also reviewed in this part. This review motivated contributions of the book on new solutions for the motion control of underactuated ships and underwater vehicles.

Part III consists of Chapters 5, 6, 7, 8, 9, 10, and 11. This part focused on underactuated surface ships by proposing a number of solutions for various control problems including stabilization, trajectory-tracking, path-tracking and path-following. Chapter 5 addressed the problem of trajectory-tracking control of underactuated ships. These ships do not have independent actuators in the sway axis. The reference trajectory was generated by a suitable virtual ship. The control development was based on an elegant coordinate transformation, Lyapunov's direct method, and the backstepping technique, and utilized passivity properties of ship dynamics and their interconnected structure. Chapter 6 examined the problem of designing a single controller that achieved stabilization and trajectory-tracking simultaneously for underactuated ships. In comparison with the preceding chapter, a path approaching the origin and a set-point were included in the reference trajectory, that is, stabilization/regulation was also considered. The control development was based on several special coordinate transformations plus the techniques in the preceding chapter. Chapter 7 presented global partial-state feedback and output feedback control schemes for trajectory-tracking control of underactuated ships. For the case of partial-state feedback, measurements of the ship sway and surge velocities were not needed, while for the case of output feedback, no ship's velocities were required for feedback. Global nonlinear coordinate changes were introduced to transform the ship dynamics to a system affine in the ship velocities. This affine form allowed us to design observers that globally exponentially estimate unmeasured velocities. These observers plus the techniques in Chapters 4 and 5 facilitated the development of the controllers. Chapter 8 dealt with the problem of path-tracking control of underactuated ships. In comparison with Chapters 5, 6, and 7, the requirement of the reference trajectory generated by a suitable virtual ship was relaxed. Both full state feedback and output feedback cases were considered. The control development was based on a series of nontrivial coordinate transformations plus the techniques in the previous chapters. Chapter 9 addressed the problem of way-point tracking control of underactuated ships. Both full state feedback and output feedback controllers were designed. The controllers in this chapter could be regarded as an advanced version of the conventional course-keeping controllers in the sense that in addition to maintaining the desired heading, both the sway displacement (lateral distance) and sway velocity were controlled. Chapter 10 developed full state feedback and output feedback controllers that forced underactuated ships to follow a predefined path. The control development was motivated by an observation that it is practical to steer a vessel such that the vessel is on the reference path and its total velocity is tangent to the reference path, and that the vessel forward speed is controlled separately by the
main thruster control system. The techniques in the previous chapters plus the use of the Serret-Frenet frame facilitated the results. Chapter 11 presented a different approach from Chapter 10 to solve a path-following problem for underactuated ships. Unlike Chapter 10, the control development was based on the method of generating reference paths by the helmsman. The path-following errors were first interpreted in polar coordinates, then the techniques developed in the previous chapters were used to design path-following controllers. Interestingly, some situations of practical importance such as parking and point-to-point navigation were covered in this chapter as a by-product of the developed path-following system.

Part IV consists of Chapters 12 and 13. This part focused on underactuated underwater vehicles. Chapter 12 addressed the problem of trajectory-tracking control of underactuated underwater vehicles. These vehicles do not have independent actuators in the sway and heave axes. The controller development was built on the techniques developed for underactuated ships in Chapters 5, 6, and 7. Due to complex dynamics of the underwater vehicles in comparison with that of the ships, the control design and stability analysis required more complicated and coordinate transformations and control design than those already developed for underactuated ships in Chapters 5, 6, and 7. Chapter 13 extended the results of Chapter 11 to the design of a path-following control system for underactuated underwater vehicles. A series of path-following strategies for the vehicles was first discussed. A practical approach was then chosen to facilitate the control development. We also addressed the parking and point-to-point navigation problems for the vehicles in this chapter.

Part V (Chapter 14) presented several applications of the observer and control design techniques developed in the previous chapters to other underactuated mechanical systems including mobile robots, and VTOL aircraft.

Finally this chapter concludes the book by briefly summarizing the main results presented in the previous chapters and presenting related open problems for further investigation.

### 15.2 Perspectives and Open Problems

Control of underactuated ocean vessels in particular and of underactuated mechanical systems in general is far from being well completed. There are in fact numerous important problems, which are still open to further investigation. In the previous chapters, we provided a number of control objectives and their solutions. However, this book has not covered all aspects of dynamics and control of the ocean vessels. Here, we highlight some open problems with detailed discussion. The problems are mainly chosen based on the following two criteria. First, to our knowledge these problems are still open and have not been solved in the literature. Second, we think that these problems potentially have important impacts on the development of control theory for underactuated systems and significant applications in practice. These criteria are highly subjective and mainly reflect our personal experience and points
of view. However, the discussions are usually made in an elementary and intuitive manner.

### 15.2.1 Further Issues on Control of Single Underactuated Ocean Vessels

### 15.2.1.1 Theoretical issues

Robust and Adaptive Issues. Although the trajectory-tracking controllers in this book possess a certain robustness, more research is needed to design such a controller that explicitly takes the environmental disturbances into account. Disturbances entering the actuated dynamics (e.g., surge and yaw for ships) can be easily taken care of. However, disturbances that enter the unactuated dynamics (e.g., sway for ships) are more difficult to handle. A possible way to reject the constant part of these disturbances on the unactuated dynamics is to introduce a small angle to the heading error to compensate for the disturbances.

Since the controllers in the book require exact knowledge of the system parameters, it is of interest to examine the problem of entirely unknown vessel parameters. As opposed to the approach based on time-varying linear system stability, the proposed control designs in this book generate explicit Lyapunov functions. This feature can be further exploited in conjunction with adaptive control algorithms in the literature to address the adaptive tracking problem.

Output Feedback Issues. The global coordinate transformations in the book bring the ship and mobile robot systems to an affine form for which a global exponential nonlinear observer can be easily designed. It is worth investigating the possibility of extending this type of transformation to a certain class of Lagrange systems. Indeed, the most difficult task is to obtain a solution of a set of partial differential equations, which are in general not easy to solve. Furthermore, an adaptation rule should be included in the observers to estimate the constant components (bias) of the environmental disturbances.

### 15.2.1.2 Real Application Issues

Although several controllers presented in this book have been successfully implemented on a scaled model ship, the following problems of technology transfer to real application studies should be considered:

1. assessment of benefits obtainable from implementing nonlinear controllers on real ocean-going vessels;
2. the issues of controller implementation in practice;
3. water-tank trials of controllers with scaled models;
4. sea-going trials: a full set of application trials;
5. commercial implementation.

### 15.2.2 Coordination Control of Multiple Underactuated Ocean Vessels

Coordination control involves controlling positions of a group of ocean vehicles such that they perform desired tasks such as optimizing objective functions from measurements taken by each vessel, and stabilization/tracking desired locations relative to reference point(s). Therefore from a reference trajectory/path point of view, we can divide coordination control of a network of vessels into two main classes: (1) Known reference trajectories in advance and (2) unknown reference trajectories in advance, for the agents in the network. Three popular approaches to the known reference trajectory in advance class are leader-following (e.g. [164, 165]), behavioral (e.g., $[166,167]$ ), and use of virtual structures (e.g., $[168,169]$ ). Most research work investigating formation control utilize one or more of these approaches in either a centralized or a decentralized manner. Centralized schemes ( $[165,170,171])$ use a single controller that generates collision free trajectories in the workspace. Although these guarantee a complete solution, centralized schemes require high computational power and are not robust. Decentralized schemes (e.g., [57, 172]) require less computational effort, and are relatively more scalable to the team size. However, it is very difficult to predict and control the critical points. For the unknown reference trajectory in advance class, coordination control often involves optimizing objective functions to give reference trajectories for the agents in the network to track. An application of formation control based on the potential field method [165] and Lyapunov's direct method [173] to gradient climbing was addressed in [174]. Geometric formation based on Voronoi's partition optimization is given in [175]. Some other work belonging to this class includes [176] on geometric formation, [177,178] on pattern formation, [57] on flocking, [179] on swarm aggregation, and [180] on deployment and task allocation. The main problem with this class is that the final arrangement of the agents cannot be foretold due to the fact that the aforementioned results use local optimization methods to optimize nonconvex objective functions. This means that only local results can be obtained. For both classes, the aforementioned works assume that the vehicles have very simple dynamics such as single or double integrators. Vehicle dynamics are usually complex in the sense of being nonlinear and are subject to disturbances. Therefore, motion coordination control systems and deployment algorithms for networked vehicles should consider these complex features of the vehicle dynamics. It is suggested to combine the control developments for the single underactuated ocean vessels in this book with the available coordination control algorithms to achieve the aforementioned coordination objectives. Indeed, this combination is not a trivial task. It requires a careful examination of the algorithms developed for single vessels and coordination objectives. The reader is referred to [181-185] for several coordination control results based
on the control of single mobile robots in Chapter 14 of this book and the formation control results in [186].

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